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# Interlayer engineering-induced charge redistribution in Bi<sub>2</sub>Te<sub>3</sub> toward efficient Zn<sup>2+</sup> and NH<sub>4</sub><sup>+</sup> storage†

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Bismuth-based materials show promise for aqueous energy storage systems due to their unique layered structures and high storage capacity. Some bismuth-based materials have been applied to store  $Zn^{2+}$  or  $NH_4^+$ , indicating that one bismuth-based compound may be innovatively used in both zinc-ion and ammonium-ion batteries (ZIBs and AIBs). Herein, we successfully design a poly(3,4-ethylenedioxythiophene) (PEDOT) coated and embedded  $Bi_2Te_3$  ( $Bi_2Te_3$ @PEDOT). Theoretical calculations and experimental studies demonstrate that the PEDOT coating and its intercalation into the interlayer enhance the structural stability of  $Bi_2Te_3$  and significantly improve the storage capacities for  $Zn^{2+}$  and  $NH_4^+$ . The PEDOT intercalation results in an increased interlayer spacing and a charge redistribution in the interlayer, facilitating charge transfer. Additionally, the insertion-type mechanism of  $Zn^{2+}$  and  $NH_4^+$  in  $Bi_2Te_3$ @PEDOT is revealed through E0 E1 in ZIBs and 235 mA h E2 in AIBs at 0.2 A E3 E4 E5 whibits high discharge capacities of 385 mA h E3 in ZIBs and 235 mA h E4 in AIBs at 0.2 A E5 in AIBs at 0.2 A E6 in AIBs and E7 in AIBs at 0.2 A E8 in AIBs and E9. Bi2E9 and Bi2E9 and Bi2E9 and Bi2E9 and Bi2E9 and Bi2E9 and Bi2E9 full cells offer high reversible capacities. This work provides a reference for designing bifunctional energy storage materials.

### Introduction

Aqueous batteries (ABs) are a new type of promising energy storage device due to their abundant raw materials, inherent safety, and high cost-effectiveness. <sup>1-5</sup> Various types of ABs have been developed such as Na<sup>+</sup>, K<sup>+</sup>, Al<sup>3+</sup>, Zn<sup>2+</sup>, Ca<sup>2+</sup>, H<sup>+</sup> and NH<sub>4</sub><sup>+</sup> ion batteries. <sup>6-9</sup> Among them, zinc-ion and ammonium-ion batteries (ZIBs and AIBs) have gained significant attention owing to the high theoretical specific capacity (5855 mA h cm<sup>-3</sup>)

of Zn anodes and the small hydrated ion size (0.331 nm) of  $\mathrm{NH_4}^{+.10^{-13}}$  Though some materials have been reported to show good electrochemical performances in ZIBs and AIBs, such as  $\mathrm{MoO_3}$  and CuS, two different modification methods are essential to separately obtain high storage capacities for  $\mathrm{Zn^{2^+}}$  and  $\mathrm{NH_4}^{+.14^{-16}}$ 

Recently, the design of bifunctional materials has been drawing more and more attention. For instance, Yu *et al.* synthesized a heterostructural  $VS_2/VO_x$  by an *in situ* electrochemical method; the  $VS_2/VO_x$  cathode displayed a high specific capacity of 156 mA h g<sup>-1</sup> at 10 A g<sup>-1</sup> in ZIBs and maintained a reversible capacity of 150 mA h g<sup>-1</sup> over more than 1000 cycles in AIBs.<sup>17</sup> Li and coworkers designed a layered NH<sub>4</sub>V<sub>4</sub>O<sub>10</sub> cathode with an adjustable interlayer distance through an *in situ* electrochemical strategy, and the optimized electrode exhibited good electrochemical performances for both Zn<sup>2+</sup> and NH<sub>4</sub><sup>+</sup> storage.<sup>18</sup> However, there are few reports on bifunctional anodes for both ZIBs and AIBs.

In previous works, layered bismuth-based materials such as BiOXs, Bi<sub>2</sub>O<sub>2</sub>CO<sub>3</sub>, and Bi<sub>2</sub>SeO<sub>5</sub> exhibited good Zn<sup>2+</sup> or NH<sub>4</sub><sup>+</sup> storage properties due to their large interlayer spacings and robust crystal structures. <sup>19-21</sup> Therefore, bismuth-based materials show great potential for both Zn<sup>2+</sup> and NH<sub>4</sub><sup>+</sup> storage. To achieve excellent electrochemical properties in bismuth-based materials, it is very important to improve electronic/ionic

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conductivity.<sup>22,23</sup> Doping/defect engineering is often used to promote electronic rearrangement and therefore enhances ion/electron transfer, such as Co-doped BiOBr and Se vacancy-activated Bi<sub>2</sub>Se<sub>3</sub>.<sup>20,24</sup> However, localized or uneven doping and defects result in their restricted enhancement. For instance, excessive Co doping increases the diffusion barrier in a local area.<sup>20</sup> Moreover, the designs of the coating and composite can only improve electron or ion transfer. Fortunately, interlayer engineering is a good choice for layered materials. Benefiting from the existence of structural pillars (cations, graphene, polymers and so on), the interlamellar spacing is increased and the electrostatic interaction between layers is decreased, enabling fast reversible ion diffusion/intercalation and high structural stability.<sup>25-27</sup>

In this study, we prepare a PEDOT-coated/embedded  $Bi_2Te_3$  ( $Bi_2Te_3$ @PEDOT) as a bifunctional anode material for efficient  $Zn^{2+}$  and  $NH_4^+$  storage. The coating of PEDOT improves the stability and electronic conductivity of  $Bi_2Te_3$ , and the embedded PEDOT changes the interlayer spacing and charge distribution of  $Bi_2Te_3$  and thus enhances  $Zn^{2+}/NH_4^+$  transport kinetics. The  $Bi_2Te_3$ @PEDOT electrode shows superior capacity for  $Zn^{2+}$  and  $NH_4^+$  storage at mass loadings of 5 and 10 mg cm<sup>-2</sup>. Remarkably,  $Zn^{2+}$  and  $NH_4^+$  full cells also display satisfactory discharge capacity and cycling stability.

## Results and discussion

The preparation processes of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT are schematically displayed in Fig. 1a. Firstly, hexagonal Bi<sub>2</sub>Te<sub>3</sub> nanosheets are synthesized via a solvothermal method, and then PEDOT is coated on and embedded in Bi<sub>2</sub>Te<sub>3</sub> during the polymerization process of EDOT. The compositions and distinctions of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT are characterized by various tests. The X-ray diffraction (XRD) patterns reveal that both samples exhibit similar characteristic peaks which are consistent with Bi<sub>2</sub>Te<sub>3</sub> (JCPDS no. 15-0863), suggesting that the crystal structure of Bi<sub>2</sub>Te<sub>3</sub> remains unchanged after the incorporation of PEDOT (Fig. 1b). Significantly, the (006) crystal plane in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT (17.22°) shifts to a low angle compared to Bi<sub>2</sub>Te<sub>3</sub> (17.48°), indicating that the lattice spacing increases from 0.507 to 0.515 nm due to the insertion of PEDOT (inset of Fig. 1b). The pure PEDOT is characterized by means of multiple methods (Fig. S1†). Raman spectra further confirm the presence of both Bi<sub>2</sub>Te<sub>3</sub> and PEDOT in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT (Fig. 1c). The peaks at 62, 102, and 134 cm<sup>-1</sup> arise from the  $A_{1g}^1$ ,  $E_g^2$ , and  $A_{1g}^2$ modes of Bi<sub>2</sub>Te<sub>3</sub>, while the characteristic peaks of PEDOT are located at 990, 1263, 1365, 1435, 1510, and 1566 cm<sup>-1</sup>.27-29 Additionally, the peaks in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT exhibit a visible red shift (inset in Fig. 1c). X-ray photoelectron spectroscopy (XPS) is employed to analyze the chemical state of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>-Te<sub>3</sub>@PEDOT. In the Te 3d XPS spectra (Fig. 1d), the peaks at 571.8 and 582.2 eV in Bi<sub>2</sub>Te<sub>3</sub> originate from the Te-Bi bonds, and the Te-O bonds at 572.5 and 586.0 eV are attributed to surface oxidation.30 Similarly, the peaks at 157.2/162.5 eV and 158.6/164.0 eV in the Bi 4f XPS spectra correspond to the Te-Bi and O-Bi bonds, respectively (Fig. S2a†).31,32 However, those peaks in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT shift towards higher binding energies,

indicating the change in the chemical environment due to the introduction of PEDOT. Further evidence for the successful synthesis of  $\rm Bi_2Te_3$ @PEDOT is provided by the C 1s and S 2p XPS spectra and the Fourier-transform infrared (FT-IR) spectroscopy spectra (Fig. S2b-d†).<sup>33,34</sup> Nitrogen adsorption/desorption isotherms and pore size distribution curves reveal that the cumulative pore volume of  $\rm Bi_2Te_3$  rises from 0.0395 to 0.1201 cm<sup>3</sup> g<sup>-1</sup>, and the specific surface area grows from 7.54 to 19.66 m<sup>2</sup> g<sup>-1</sup> after the introduction of PEDOT (Fig. S3†).

Scanning electron microscopy (SEM) and transmission electron microscopy (TEM) images provide insights into the microstructure and morphology of the samples. The SEM image reveals that Bi2Te3 exhibits a smooth hexagonal nanosheet structure (Fig. 1e). Conversely, Bi<sub>2</sub>Te<sub>3</sub>@PEDOT displays a rough surface, due to the PEDOT coating layer (Fig. 1f). This explains the increased specific surface area in Bi2Te3@PEDOT and enhances the charge transfer between Bi<sub>2</sub>Te<sub>3</sub>@PEDOT and electrolyte. The TEM images reveal a PEDOT coating layer with a thickness of approximately 10 nm in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT (Fig. S4a and b†). The high-resolution TEM (HRTEM) image exhibits the lattice fringes (0.383 nm) which corresponds to the (101) crystal plane of Bi<sub>2</sub>Te<sub>3</sub> (greater than 0.377 nm for pure Bi<sub>2</sub>Te<sub>3</sub>) (Fig. 1g).35 This result corresponds to the shift in the XRD test. Simultaneously, the polycrystalline nature of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT is confirmed by the selected area electron diffraction (SAED) test (Fig. S4c†). Elemental mapping illustrates that Bi, Te, S, O, and C elements are evenly distributed across the nanosheet (Fig. 1h). The above results demonstrate the successful synthesis of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT.

Density functional theory (DFT) calculations are used to investigate the role of the embedded PEDOT in the Bi<sub>2</sub>Te<sub>3</sub> structure. The optimized crystal structures of  $Bi_2Te_3$  (8.87  $\times$  $15.38 \times 30.55 \text{ Å}, \alpha = 90.15^{\circ}, \beta = 90^{\circ}, \gamma = 90^{\circ})$  and PEDOT-Bi<sub>2</sub>Te<sub>3</sub>  $(8.94 \times 15.44 \times 37.96 \text{ Å}, \alpha = 90.09^{\circ}, \beta = 90.03^{\circ}, \gamma = 89.98^{\circ})$  are presented in Fig. S5.† The insertion of PEDOT results in an expanded interlayer distance and a slight lattice distortion, which is consistent with the experimental results. Differential charge distribution diagrams reveal significant interlamellar electron exchange in the Bi<sub>2</sub>Te<sub>3</sub>, reflecting a strong interaction in the interlayer (Fig. 1i). PEDOT intercalation reduces such electron exchange, weakening the interlayer interaction (Fig. 1j). Bader charge analysis further supports this finding. Specifically, the average electron count of Te atoms adjacent to PEDOT decreases from 0.37 (without PEDOT) to 0.06 e, indicating that more electron transfer occurs between PEDOT and its neighboring Te atoms. The density of states (DOS) is simulated for both Bi<sub>2</sub>Te<sub>3</sub> and PEDOT-Bi<sub>2</sub>Te<sub>3</sub> (Fig. S6†). The Bi<sub>2</sub>Te<sub>3</sub> model exhibits typical semiconductor properties. The downward shift of the DOS peaks in the PEDOT-Bi<sub>2</sub>Te<sub>3</sub> model along with the significant DOS at the Fermi level indicates a transformation from semiconductor to conductor, suggesting that the insertion of PEDOT significantly improves the electrical conductivity of PEDOT-Bi<sub>2</sub>Te<sub>3</sub>.36 The reduced interlayer electronic interactions and improved electrical conductivity are expected to facilitate ion diffusion in the interlayer, therefore enhancing the electrochemical performances Bi<sub>2</sub>Te<sub>3</sub>@PEDOT.

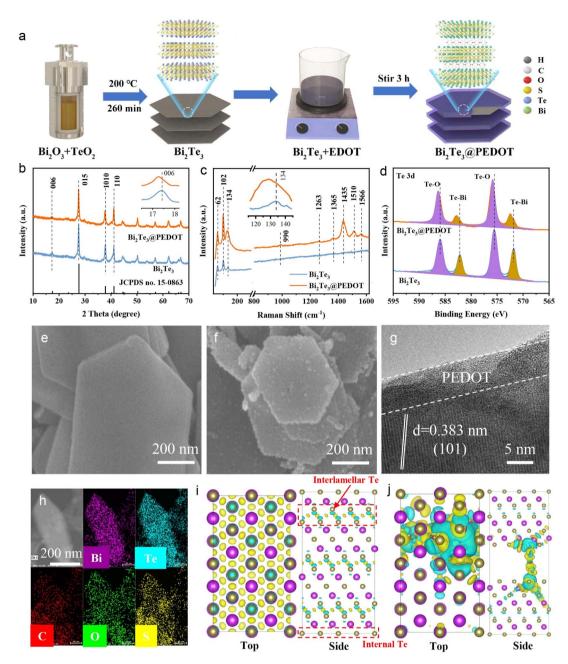


Fig. 1 (a) Schematic diagram of the preparation of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. (b) XRD patterns, (c) Raman spectra, and (d) Te 3d XPS spectra of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. SEM images of (e) Bi<sub>2</sub>Te<sub>3</sub> and (f) Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. (g) TEM images of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. (h) Elemental mapping of Bi<sub>2</sub>-Te<sub>3</sub>@PEDOT. Charge density distributions of (i) Bi<sub>2</sub>Te<sub>3</sub> and (j) PEDOT-Bi<sub>2</sub>Te<sub>3</sub> (the purple and brown colors represent Bi and Te atoms, and the yellow and blue regions reflect electron accumulation and electron depletion, respectively).

The electrochemical properties of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@-PEDOT in ZIBs are evaluated firstly. In cyclic voltammetry (CV) tests, Bi<sub>2</sub>Te<sub>3</sub> exhibits two redox peaks at 0.62, 0.68, 1.15, and 1.20 V, suggesting the multi-step insertion/deinsertion of Zn<sup>2+</sup> in Bi<sub>2</sub>Te<sub>3</sub>. Meanwhile, Bi<sub>2</sub>Te<sub>3</sub>@PEDOT shows a pair of redox peaks at ~0.6/1.2 V, reflecting a variation in the insertion/ deinsertion process of Zn2+ after the introduction of PEDOT (Fig. S7†).37 This change is explored in the subsequent studies. Meanwhile, the first three CV curves of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT overlap better than those of Bi<sub>2</sub>Te<sub>3</sub>, indicating its improved reversibility.

The first three galvanostatic charge-discharge (GCD) curves show an average discharge plateau at about 0.69 V in Bi<sub>2</sub>Te<sub>3</sub> and 0.67 V in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT, implying their suitability as anodes in full cells (Fig. S8†). The rate performances and cycling properties are further explored. At current densities of 0.2, 0.5, 1.0, 2.0, 3.0, 5.0, and 10 A  $g^{-1}$ , Bi<sub>2</sub>Te<sub>3</sub>@PEDOT offers specific capacities of 378, 340, 320, 312, 310, 300, and 253 mA h  $g^{-1}$ , significantly higher than those of Bi<sub>2</sub>Te<sub>3</sub> (195, 180, 152, 123, 110, 97, and 80 mA h  $g^{-1}$ ) and PEDOT (Fig. 2a and S9†). Fig. 2b and c shows the GCD curves at different current densities, in which

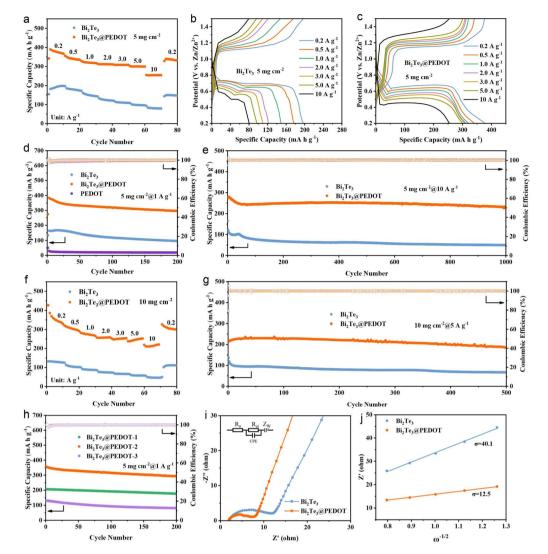


Fig. 2 (a) Rate performances and (b and c) corresponding GCD curves of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. (d) Cycle life of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT, and PEDOT at 1 A q<sup>-1</sup>. (e) Cyclic performances of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub> (PEDOT at 10 A q<sup>-1</sup> (5 mg cm<sup>-2</sup>). (f) Rate capabilities and (g) cyclic performances of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT at 5 A g<sup>-1</sup> (10 mg cm<sup>-2</sup>). (h) Cycle life of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT-x at 1 A g<sup>-1</sup>. (i) EIS spectra and (j) relationship between  $\omega^{-1/2}$ and Z' of  $Bi_2Te_3$  and  $Bi_2Te_3$ @PEDOT.

Bi<sub>2</sub>Te<sub>3</sub>@PEDOT exhibits more stable charging and discharging plateaus. In the cycle life test, Bi2Te3@PEDOT maintains a discharge capacity of 295 mA h g<sup>-1</sup> after 200 cycles at 1 A g<sup>-1</sup>, which is significantly higher than that of Bi<sub>2</sub>Te<sub>3</sub> (96 mA h g<sup>-1</sup>) (Fig. 2d). Significantly, PEDOT exhibits a negligible Zn<sup>2+</sup> storage capacity (18 mA h  $\mathrm{g}^{-1}$ ), reflecting its low capacity contribution to Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. After 1000 cycles at 10 A g<sup>-1</sup>, Bi<sub>2</sub>Te<sub>3</sub>@PEDOT retains 82% of its initial capacity, while Bi<sub>2</sub>Te<sub>3</sub> retains only 40% (Fig. 2e). Additionally, the high mass loading Bi<sub>2</sub>Te<sub>3</sub>@PEDOT electrodes (10 mg cm<sup>-2</sup>) are further tested. The discharge capacities of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT are 325, 290, 263, 257, 253, 251, and 219 mA h g<sup>-1</sup> at current densities from 0.2 to 10 A g<sup>-1</sup>, much higher than those of Bi<sub>2</sub>Te<sub>3</sub> (125, 99, 85, 75, 64, 56, and 49 mA h  $g^{-1}$ ) (Fig. 2f). After 500 cycles at 5 A  $g^{-1}$ , Bi<sub>2</sub>Te<sub>3</sub>@PEDOT maintains a discharge capacity of 187 mA h g<sup>-1</sup>, while Bi<sub>2</sub>Te<sub>3</sub> retains only 68 mA h g<sup>-1</sup> (Fig. 2g). The rate and cycling performances of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT-x are compared (Fig. 2h and S10†). It is clear that Bi<sub>2</sub>Te<sub>3</sub>@PEDOT-2 (Bi<sub>2</sub>Te<sub>3</sub>@PEDOT) outperforms Bi<sub>2</sub>Te<sub>3</sub>@PEDOT-1 and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT-3, reflecting that a suitable amount of PEDOT is critical to battery performance.

The electrochemical kinetics of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT is analyzed by CV tests at different sweep rates (Fig. S11a and b†). Fitting the slopes of the  $\log(i)$  versus  $\log(v)$  curves (Fig. S11c†), the b-values of the redox peaks R1, O1, R2, and O2 are 0.72, 0.98, 0.62, and 0.67, respectively. This suggests the enhanced diffusion-controlled behavior of Bi2Te3@PEDOT due to its fast ion diffusion.<sup>38,39</sup> The electrochemical impedance spectra (EIS) show that the ohmic resistance ( $R_0 = 2.13 \Omega$ ), charge transfer resistance ( $R_{\rm ct} = 9.9 \,\Omega$ ), and Warburg factor ( $\sigma = 40.1$ ) of Bi<sub>2</sub>Te<sub>3</sub> are larger than those of  $Bi_2Te_3$ @PEDOT ( $R_0 = 1.94 \Omega$ ,  $R_{ct} = 4.56$  $\Omega$ , and  $\sigma = 16.8$ ), thereby indicating an improved reaction kinetics in Bi2Te3@PEDOT (Fig. 2i and j). To further demonstrate the superiority of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT, the Zn<sup>2+</sup> diffusion coefficient is calculated by the galvanostatic intermittent

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titration technique (GITT) test and Fick's second law. As shown in Fig. S12,† the ion diffusion rate in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT is larger than that in Bi<sub>2</sub>Te<sub>3</sub>. This result is attributed to the interlayer expansion and reduced electrostatic interaction of Bi<sub>2</sub>Te<sub>3</sub> and thus the improved ion diffusion rate.<sup>40</sup>

To explore the storage mechanism in Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>(a)-PEDOT, ex situ XRD, XPS, and SEM are conducted. Ex situ XRD profiles of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT are displayed in Fig. 3a-d. Both materials maintain similar XRD patterns during the charge and discharge processes, indicating their intercalation/ deintercalation mechanism. The (006) crystal plane in Bi<sub>2</sub>Te<sub>3</sub> shifts to a higher angle after the electrode is discharged to 0.2 V, which results from strong electrostatic interactions between the embedded Zn2+ and Bi2Te3 layers (Fig. 3b).41 In contrast, the (006) crystal plane of Bi<sub>2</sub>Te<sub>3</sub>(a)PEDOT shifts to a lower angle at the fully discharged state. This opposite result is attributed to the insertion of PEDOT and the enlarged interlayer spacing which reduces the interaction forces between Zn2+ and Bi2Te3 layers (this is proved by DFT calculations).42 Meanwhile, the byproduct of Zn<sub>4</sub>SO<sub>4</sub>(OH)<sub>6</sub>·4H<sub>2</sub>O (ZSH) is generated in both Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT, though the intensity of ZSH peaks is weak in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. To assess the stability of the materials, Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT are soaked in 2 mol L<sup>-1</sup> ZnSO<sub>4</sub> electrolyte for 2 h. As shown in Fig. S13a,† numerous bubbles are clearly observed on the surface of Bi<sub>2</sub>Te<sub>3</sub>, and some white byproducts appear on Bi<sub>2</sub>Te<sub>3</sub> after drying. In contrast, there is no obvious change for Bi2Te3@PEDOT. In addition, the XRD pattern of Bi<sub>2</sub>Te<sub>3</sub> shows the peaks of Bi<sub>2</sub>TeO<sub>5</sub> and Zn(SO<sub>4</sub>)(H<sub>2</sub>-O)6, whereas no obvious new phase is detected in Bi2Te3@-PEDOT (Fig. S13b†).

The energy storage mechanism is further investigated using ex situ XPS spectra. In the Bi 4f spectra of Bi<sub>2</sub>Te<sub>3</sub> (Fig. 3e), the peaks of the Bi-O and Bi-Te bonds irreversibly shift to high binding energy during the first cycle, and the Bi-Te peaks almost disappear at the fully charged state. Moreover, Te-O and Te-Bi peaks in the Te 3d XPS spectra also exhibit similar change, implying the poor reversibility and stability of Bi<sub>2</sub>Te<sub>3</sub> (Fig. 3f). For comparison, those peaks in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT do not change significantly during the reaction, highlighting its enhanced stability due to the coating and insertion of PEDOT (Fig. 3g and h). In the Zn 2p XPS spectra, no Zn signal is observed at the original state (Fig. S14†). Zn 2p peaks emerge when the electrode is discharged to 0.2 V and become weak at the fully charged state, indicating the Zn2+ intercalation/deintercalation.37,43 The O 1s XPS spectra can be divided into two peaks: the O-H peak at 532.0 eV and the O-Bi/O-Te peak at 530.5 eV (Fig. S15†).44 The area ratios of O-H and O-Bi/O-Te peaks increase and decrease in the discharge and charge

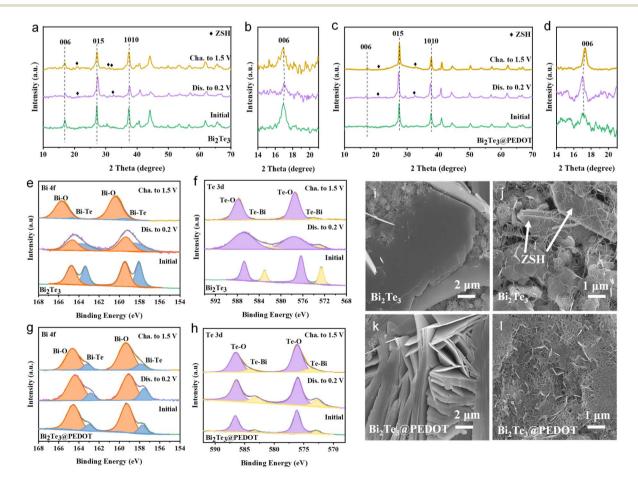


Fig. 3 (a-d) Ex situ XRD patterns, (e and g) Bi 4f XPS spectra, and (f and h) Te 3d XPS spectra of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. SEM images of (i and j) Bi<sub>2</sub>Te<sub>3</sub> and (k and l) Bi<sub>2</sub>Te<sub>3</sub>@PEDOT at the fully discharged/charged states.

processes, respectively, indicating the insertion/extraction of H<sup>+</sup>. The discharge capacities of H<sup>+</sup> and Zn<sup>2+</sup> in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT are studied in ZnSO<sub>4</sub>/DMSO and H<sub>2</sub>SO<sub>4</sub> (pH = 5) electrolytes, respectively (Fig. S16†). The specific discharge capacity of H<sup>+</sup> (194 mA h  $g^{-1}$ ) is higher than that of  $Zn^{2+}$  (103 mA h  $g^{-1}$ ), indicating dominant H<sup>+</sup> storage. Additionally, the morphology changes of the electrodes at different charge states are observed by ex situ SEM. Many large-sized ZSH flakes are observed in both Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT at 0.2 V (Fig. 3i and k). The ZSH flakes are still visible in Bi<sub>2</sub>Te<sub>3</sub> at the fully charged state (Fig. 3j). Meanwhile, the ZSH almost disappears in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT, implying its enhanced reaction reversibility (Fig. 31). The above study demonstrates the co-insertion mechanism of Zn<sup>2+</sup> and H<sup>+</sup> in both Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. Bi<sub>2</sub>Te<sub>3</sub>@PEDOT shows high stability and reversibility in the Zn<sup>2+</sup>/H<sup>+</sup> insertion/deinsertion processes.

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The remarkable performances of  $Bi_2Te_3$ @PEDOT in ZIBs prompt an exploration of its electrochemical properties in AIBs.  $Bi_2Te_3$ @PEDOT exhibits good rate performance with specific capacities of 266, 235, 200, 164, and 113 mA h g<sup>-1</sup> at 0.1, 0.2, 0.5, 1.0, and 2.0 A g<sup>-1</sup>, respectively, which is higher than those of  $Bi_2Te_3$  (245, 212, 168, 120, and 71 mA h g<sup>-1</sup>) (Fig. 4a). The corresponding GCD curves are displayed in Fig. S17,† and  $Bi_2Te_3$ @PEDOT exhibits more obvious discharge plateaus at all current densities.  $Bi_2Te_3$ @PEDOT and  $Bi_2Te_3$  deliver the

discharge capacities of 156 and 113 mA h g $^{-1}$  with capacity retentions of 72% and 67%, respectively, after 700 cycles at 1 A g $^{-1}$  (Fig. 4b). Notably, PEDOT shows a low NH $_4$ <sup>+</sup> storage capacity, suggesting its low capacity contribution to Bi $_2$ Te $_3$ @-PEDOT (Fig. S18 $^{\dagger}$ ). Additionally, Bi $_2$ Te $_3$  and Bi $_2$ Te $_3$ @PEDOT electrodes are tested under a high mass loading of 10 mg cm $^{-2}$ . The discharge capacities of Bi $_2$ Te $_3$ @PEDOT are 257, 219, 169, 112, and 50 mA h g $^{-1}$  at current densities from 0.1 to 2.0 A g $^{-1}$ , while Bi $_2$ Te $_3$  exhibits lower discharge capacities (Fig. 4c). After 300 cycles at 1.0 A g $^{-1}$ , Bi $_2$ Te $_3$ @PEDOT exhibits a significantly higher discharge capacity of 115 mA h g $^{-1}$  than that of Bi $_2$ Te $_3$  (66 mA h g $^{-1}$ ) (Fig. 4d). Evidently, Bi $_2$ Te $_3$ @PEDOT is well suited for use as a host material for NH $_4$ <sup>+</sup> storage.

The reaction kinetics of  $Bi_2Te_3$  and  $Bi_2Te_3$ @PEDOT is further investigated through CV, EIS, and GITT tests. In the first five CV curves, the redox peaks of  $Bi_2Te_3$ @PEDOT are more obvious, suggesting its higher capacity for  $NH_4^+$  storage (Fig. S19†). The corresponding *b*-values for peaks R1, O1, R2, and O2 are 0.67, 0.77, 0.83, and 0.74 respectively, suggesting an increased capacitive contribution of  $Bi_2Te_3$ @PEDOT (Fig. S20†). The  $R_{ct}$  and Warburg factor of  $Bi_2Te_3$ @PEDOT (68  $\Omega$  and 38.2) are significantly lower than those of  $Bi_2Te_3$  (108  $\Omega$  and 57.6), implying its enhanced charge transport (Fig. 4e and f). The GITT results also reveal the fast insertion and extraction rates of  $NH_4^+$  in  $Bi_2Te_3$ @PEDOT (Fig. S21†). The above results suggest the

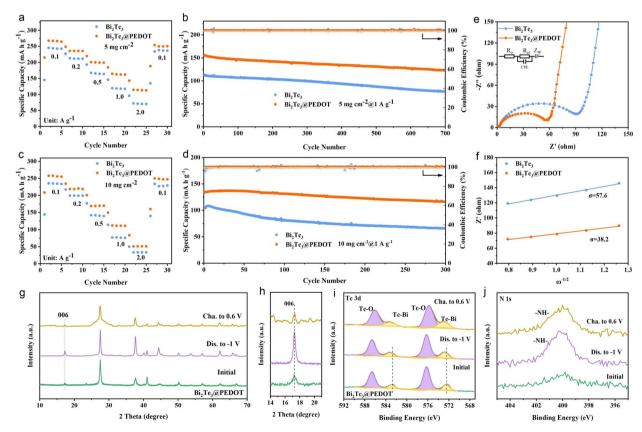


Fig. 4 (a) Rate performances and (b) cycle life of  $Bi_2Te_3$  and  $Bi_2Te_3@PEDOT$  (5 mg cm<sup>-2</sup>) at 1 A g<sup>-1</sup>. (c) Rate capabilities and (d) cycle life of  $Bi_2Te_3$  and  $Bi_2Te_3@PEDOT$  (10 mg cm<sup>-2</sup>) at 1 A g<sup>-1</sup>. (e) EIS spectra and (f) relationship between Z' and  $\omega^{-1/2}$  of  $Bi_2Te_3$  and  $Bi_2Te_3@PEDOT$ . (g and h) Ex situ XRD patterns, (i) Te 3d XPS spectra, and (j) N 1s XPS spectra of  $Bi_2Te_3@PEDOT$ .

improved charge transfer of Bi2Te3 after the introduction of PEDOT.

To study the NH<sub>4</sub><sup>+</sup> storage mechanism in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT, ex situ XRD and XPS tests are conducted. The XRD test reveals that no new phase appears after the electrode is fully discharged/ charged (Fig. 4g). The partial enlargement of the XRD pattern is shown in Fig. 4h, in which the (006) crystal plane remains

stable, implying that the insertion and extraction of NH<sub>4</sub><sup>+</sup> do not affect the structure of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. This good stability is due to the protective layer of PEDOT and the increased interlayer spacing in Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. As exhibited in Bi 4f and Te 3d XPS spectra, the Bi-Te peaks shift to higher binding energies after discharging, while the peak positions remain unchanged after charging, indicating a degree of irreversible reaction in the first

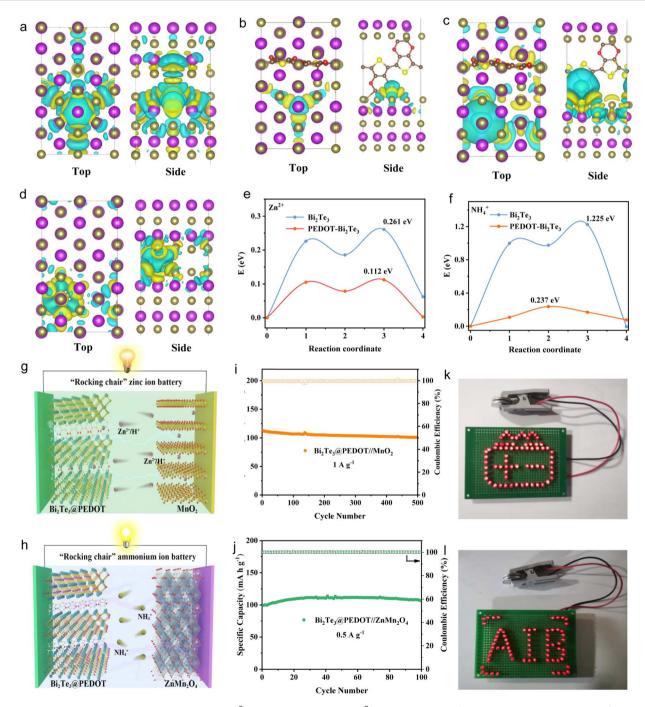


Fig. 5 Charge density distributions of (a)  $Bi_2Te_3 + Zn^{2+}$ , (b) PEDOT- $Bi_2Te_3 + Zn^{2+}$ , (c)  $Bi_2Te_3 + NH_4^+$ , and (d) PEDOT- $Bi_2Te_3 + NH_4^+$  (the purple and brown colors represent Bi and Te atoms, and the yellow and blue regions reflect electron accumulation and electron depletion, respectively). (e)  $Zn^{2+}$  and (f)  $NH_4^+$  diffusion barriers of  $Bi_2Te_3$  and  $PEDOT-Bi_2Te_3$ . (g and h) Structural illustration and (i and j) cycle life of  $Bi_2Te_3$ @PEDOT//  $MnO_2 \ and \ Bi_2Te_3@PEDOT//ZnMn_2O_4 \ full \ cells. \ Digital \ photos \ of \ two \ (k) \ Bi_2Te_3@PEDOT//MnO_2 \ and \ (l) \ Bi_2Te_3@PEDOT//ZnMn_2O_4 \ full \ cells.$ lighting up red LEDs.

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cycle (Fig. 4i and S22a†). In the N 1s spectra, a strong peak of NH– appears at the fully discharged state, providing evidence that NH<sub>4</sub><sup>+</sup> inserts into Bi<sub>2</sub>Te<sub>3</sub>@PEDOT (Fig. 4j). However, a weakened peak can still be seen after charging, reflecting the incomplete extraction of NH<sub>4</sub><sup>+</sup>. The intensity of O–H bonds in the O 1s spectra is enhanced after discharging and decreased after charging, demonstrating the co-insertion of H<sub>2</sub>O along with NH<sub>4</sub><sup>+</sup> (Fig. S22b†). The above results prove the insertion of NH<sub>4</sub><sup>+</sup> into Bi<sub>2</sub>Te<sub>3</sub>@PEDOT. Significantly, Bi<sub>2</sub>Te<sub>3</sub>@PEDOT outperforms most of the reported anodes in both ZIBs and AIBs (Tables S1 and S2†). Moreover, the high mass loading anode has been reported rarely.

The Zn<sup>2+</sup>and NH<sub>4</sub><sup>+</sup> storage properties of Bi<sub>2</sub>Te<sub>3</sub>, as well as the function of the embedded PEDOT, are further validated by DFT calculations. The ionic transport is affected by adsorption energy and lattice distortion. The charge density distributions are used to reveal the effect of PEDOT insertion on Zn2+ and NH<sub>4</sub><sup>+</sup> adsorption. In the PEDOT-Bi<sub>2</sub>Te<sub>3</sub> model, the charge transfer between Zn<sup>2+</sup> and surrounding Te atoms is reduced, leading to a change in adsorption energy from -0.547 (Bi<sub>2</sub>Te<sub>3</sub>) to -0.492 eV (PEDOT-Bi<sub>2</sub>Te<sub>3</sub>) (Fig. 5a and b). This is beneficial to Zn<sup>2+</sup> diffusion. According to Bader charge analysis, when Zn<sup>2+</sup> is inserted into Bi<sub>2</sub>Te<sub>3</sub>, the Zn transfers a charge of 0.428 e to Bi<sub>2</sub>Te<sub>3</sub>, and it drops to 0.089 e in PEDOT-Bi<sub>2</sub>Te<sub>3</sub>. Meanwhile, the three neighboring Te atoms obtain an average electron number of 0.046 e before PEDOT intercalation, and it decreases to 0.009 e after PEDOT intercalation. This is consistent with the charge density distributions. By contrast, the adsorption energy between  $NH_4^+$  and Te atoms increases from -0.735 (Bi<sub>2</sub>Te<sub>3</sub>) to -2.204 eV (PEDOT-Bi<sub>2</sub>Te<sub>3</sub>) (Fig. 5c and d). Bader charge analysis shows that NH<sub>4</sub><sup>+</sup> transfers 0.53 e before PEDOT insertion and 0.73 e after PEDOT insertion, while the average electron numbers gained by the three neighboring Te atoms are 0.028 and 0.131 e, respectively, indicating the formation of strong covalent interactions between NH<sub>4</sub><sup>+</sup> and Bi<sub>2</sub>Te<sub>3</sub> after the PEDOT insertion and the adverse NH<sub>4</sub><sup>+</sup> diffusion. On the other hand, the average bond length of Te-Bi near the adsorption site of Bi<sub>2</sub>Te<sub>3</sub> is 3.08 Å and increases to 3.15 and 3.27 Å after Zn<sup>2+</sup> and NH<sub>4</sub><sup>+</sup> insertion. This change shows a decrease from 3.07 Å to 3.07 and 3.10 Å with the help of interlayer engineering, facilitating Zn<sup>2+</sup> and NH<sub>4</sub> diffusion. Further analysis of the Zn<sup>2+</sup>/ NH<sub>4</sub><sup>+</sup> diffusion barriers in Bi<sub>2</sub>Te<sub>3</sub> and PEDOT-Bi<sub>2</sub>Te<sub>3</sub> reveals a significant reduction in diffusion barriers with the insertion of PEDOT. The diffusion pathways of Zn<sup>2+</sup>/NH<sub>4</sub><sup>+</sup> are illustrated in Fig. S23.† Specifically, the Zn<sup>2+</sup>/NH<sub>4</sub><sup>+</sup> diffusion barrier in Bi<sub>2</sub>Te<sub>3</sub> (1.225/0.237 eV) declines to 0.261/0.112 eV with the help of PEDOT (Fig. 5e and f). This reflects that the insertion of PEDOT improves the diffusion of Zn<sup>2+</sup>/NH<sub>4</sub><sup>+</sup> due to the combined influence of adsorption energy and lattice distortion.

Finally, full cells are assembled to investigate the practical application of  $\rm Bi_2Te_3@PEDOT$  in both "rocking-chair" ZIBs and AIBs, respectively. Commercial MnO $_2$  and ZnMn $_2O_4$  are used as cathodes due to their low cost and high operating voltages. The schematics of  $\rm Bi_2Te_3@PEDOT//MnO_2$  and  $\rm Bi_2Te_3@PEDOT//ZnMn_2O_4$  full cells are shown in Fig. 5g and h, where Zn $^{2+}/\rm H^+$  or NH $_4^+$  shuttles between the anode and cathode through the electrolyte. The zinc-ion full battery exhibits a high capacity of

110 mA h g $^{-1}$  at 1 A g $^{-1}$  and a capacity retention of 88% after 500 cycles (Fig. 5i). The ammonium-ion full battery delivers a discharge capacity of 110 mA h g $^{-1}$  at 0.5 A g $^{-1}$  and good cycling stability over 100 cycles (Fig. 5j). The GCD curves of the ZIB and AIB full cells are displayed in Fig. S24.† In addition, both full cells drive the self-assembled red LED logos (Fig. 5k and l).

#### Conclusions

In summary, a PEDOT-coated/embedded Bi<sub>2</sub>Te<sub>3</sub> is successfully fabricated through a hydrothermal method followed by the polymerization of EDOT. Various characterizations and DFT calculations demonstrate that the embedded PEDOT enlarges the interlayer spacing and reduces the interlamellar interaction, thus improving the reaction kinetics of Bi<sub>2</sub>Te<sub>3</sub>. Meanwhile, the coating of PEDOT enhances structural stability, suppresses side reactions with the aqueous solution, and increases electronic conductivity. Ex situ tests demonstrate the insertion-type mechanism of Bi<sub>2</sub>Te<sub>3</sub>@PEDOT in both ZIBs and AIBs. The Bi<sub>2</sub>-Te<sub>3</sub>@PEDOT electrode exhibits high discharge capacities of 385 and 235 mA h  $g^{-1}$  at 0.2 A  $g^{-1}$  in ZIBs and AIBs, respectively. Even at a high mass loading of 10 mg cm<sup>-2</sup>, Bi<sub>2</sub>Te<sub>3</sub>@PEDOT shows satisfactory rate performances and cycling stability. Furthermore, "rocking-chair" ZIBs and AIBs with the Bi<sub>2</sub>Te<sub>3</sub>@-PEDOT anode offer high reversible capacities of 110 mA h  $g^{-1}$  at 1.0 A  $g^{-1}$  (ZIBs) and 110 mA h  $g^{-1}$  at 0.5 A  $g^{-1}$  (AIBs), respectively. This work extends the utilizing scope of layered bismuthbased materials and serves as a reference for bifunctional energy storage electrodes.

# Data availability

The data supporting this article have been included in the main text and ESI.†

### **Author contributions**

X. L.: writing – original draft, data curation; F. L.: writing – original draft, data curation; H. Y.: data curation, investigation; Y. D.: formal analysis, resources, software; L. C.: investigation; T. S.: investigation; Y. P.: resources, software; X. W.: project administration, funding acquisition; B. L.: writing – review & editing, conceptualization, supervision; Y. X.: writing – review & editing, conceptualization, supervision; X. W.: writing – review & editing, project administration, funding acquisition. X. L. and F. L. contributed equally to this work.

# Conflicts of interest

There are no conflicts to declare.

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#### References

**Edge Article** 

- 1 M. Li, X. P. Wang, J. S. Meng, C. L. Zuo, B. K. Wu, C. Li, W. Sun and L. O. Mai, Adv. Mater., 2024, 36, 2308628.
- 2 Z. N. Ju, Q. Zhao, D. L. Chao, Y. Hou, H. G. Pan, W. P. Sun, Z. Y. Yuan, H. Li, T. Y. Ma, D. W. Su and B. H. Jia, Adv. Energy Mater., 2022, 12, 2201074.
- 3 J. Liu, K. Wang, Y. Sun, H. Li, X. Han, X. Duan, Z. Huang and T. Ma, Nano Energy, 2025, 136, 110764.
- 4 C. Luo, H. Lei, Y. Xiao, X. Nie, Y. Li, Q. Wang, W. Cai, C. Dai, M. Yao, Y. Zhang and D. Yuan, Energy Mater., 2024, 4,
- 5 X. Liu, B. Xu, S. Deng, J. Han, Y. An, J. Zhao, Q. Yang, Y. Xiao and C. Fang, Carbon Energy, 2024, 6, e603.
- 6 G. Zeng, Q. Sun, S. Horta, P. R. Martínez-Alanis, P. Wu, J. Li, S. Wang, M. Ibáñez, Y. Tian, L. Ci and A. Cabot, Energy Environ. Sci., 2025, 18, 1683-1695.
- 7 W. Jiang, K. Zhu, W. Xie, Z. Wang, Z. Ou and W. Yang, Chem. Sci., 2024, 15, 2601-2611.
- 8 Y. Zhang, Z. Huang, L. Lei, H. Fan, X. Han, H. Li and T. Ma, Adv. Energy Mater., 2025, 15, 2404732.
- 9 F. Xiankai, X. Kaixiong, Z. Wei, D. Weina, Z. Hai, C. Liang and C. Han, Carbon Energy, 2024, 6, e536.
- 10 B.-H. Xiao, K. Xiao, J.-X. Li, C.-F. Xiao, S. Cao and Z.-Q. Liu, Chem. Sci., 2024, 15, 11229-11266.
- 11 G. Zeng, Q. Sun, S. Horta, S. Wang, X. Lu, C. Y. Zhang, J. Li, J. Li, L. Ci, Y. Tian, M. Ibáñez and A. Cabot, Adv. Mater., 2024, 36, 2305128.
- 12 B. Long, X. Ma, L. Chen, T. Song, Y. Pei, X. Wang and X. Wu, Adv. Funct. Mater., 2024, 34, 2411430.
- 13 S. Gong, M. Zhu, Y. Zhou, R. Li, J. Zhang, X. Jia, D. Chao and C. Wang, Chem. Sci., 2024, 15, 19870-19885.
- 14 J. Shi, Y. Hou, Z. Liu, Y. Zheng, L. Wen, J. Su, L. Li, N. Liu, Z. Zhang and Y. Gao, Nano Energy, 2022, 91, 106651.
- 15 H. Liu, K. Zhang, S. Wang and X. Cai, Small, 2023, 20, 2310835.
- 16 Q. Lei, J. Zhang, Z. Liang, Y. Yue, Z. Ren, Y. Sun, Z. Yao, J. Li, Y. Zhao, Y. Yin, P. Huai, Z. Lv, J. Li, Z. Jiang, W. Wen, X. Li, X. Zhou and D. Zhu, Adv. Energy Mater., 2022, 12, 2200547.
- 17 D. Yu, Z. Wei, X. Zhang, Y. Zeng, C. Wang, G. Chen, Z. X. Shen and F. Du, Adv. Funct. Mater., 2020, 31, 2008743.
- 18 S. Li, D. Yu, J. Liu, N. Chen, Z. Shen, G. Chen, S. Yao and F. Du, Adv. Sci., 2023, 10, 2206836.
- 19 M. Han, Y. Qian, X. Li, N. Wang, T. Song, L. Liu, X. Wang, X. Wu, M.-K. Law and B. Long, J. Colloid Interface Sci., 2023, 645, 483-492.
- 20 B. Long, Q. Zhang, T. Duan, T. Song, Y. Pei, X. Wang, C. Zhi, X. Wu, Q. Zhang and Y. Wu, Adv. Sci., 2022, 9, 2204087.
- 21 D. Ling, Q. Wang, G. Tian, H. Yu, D. Zhang and Q. Wang, ACS Nano, 2023, 17, 25222-25233.

- 22 Q. Wang, M. Wang, L. Wen, W. Zeng, B. Ge, C. Zhang, Y. Yue and S. Wang, Adv. Funct. Mater., 2023, 34, 2214506.
- 23 X. Wang, F. Zhou, Q. Liang, Q. Zhang, Y. Zhu, Z. Xiao and L. Wang, Adv. Funct. Mater., 2024, 34, 2408203.
- 24 B. Long, X. Ma, L. Chen, T. Song, Y. Pei, X. Wang and X. Wu, Adv. Funct. Mater., 2024, 34, 2411430.
- 25 S. Li, Y. Liu, X. Zhao, Q. Shen, W. Zhao, Q. Tan, N. Zhang, P. Li, L. Jiao and X. Qu, Adv. Mater., 2021, 33, e2007480.
- 26 G. Zhang, T. Wu, H. Zhou, H. Jin, K. Liu, Y. Luo, H. Jiang, K. Huang, L. Huang and J. Zhou, ACS Energy Lett., 2021, 6, 2111-2120.
- 27 W. Bi, G. Gao, G. Wu, M. Atif, M. S. AlSalhi and G. Cao, Energy Storage Mater., 2021, 40, 209-218.
- 28 J. Feng, Z. Zhuang, Y. Zhou and C. Li, Adv. Funct. Mater., 2024, 34, 2315188.
- 29 G. Wang, F. Meng, Y. Chen, A. Lotnyk and X. Shen, Adv. Sci., 2024, 11, e2308056.
- 30 X. Liu, Y. Si, K. Li, Y. Xu, Z. Zhao, C. Li, Y. Fu and D. Li, Energy Storage Mater., 2021, 41, 255-263.
- 31 C. Zheng, Z. Guo, B. Jian, Z. Chen, J. Zhong, N. Li and S. Huang, Chem. Eng. J., 2023, 475, 146408.
- 32 S. Wang, Y. Qiao, X. Liu, S. Zhu, Y. Zheng, H. Jiang, Y. Zhang, J. Shen, Z. Li, Y. Liang, Z. Cui, P. K. Chu and S. Wu, Adv. Funct. Mater., 2022, 33, 2210098.
- 33 X. B. Liao, C. L. Pan, H. X. Yan, Y. Zhu, Y. S. Pan and C. J. Yin, Chem. Eng. J., 2022, 440, 135930.
- 34 K. Kim, M. Kim, H. Lee, D. W. Chung and J. Kim, Small, 2024, 20, e2402341.
- 35 S. Chong, L. Yuan, Q. Zhou, Y. Wang, S. Qiao, T. Li, M. Ma, B. Yuan and Z. Liu, Small, 2023, 19, 2303985.
- 36 Y. Wu, Q. Zong, Y. Zhuang, Q. Wang, C. Liu, Q. Zhang, D. Tao, J. Zhang, J. Wang and G. Cao, Chem. Eng. J., 2024, 491, 152064.
- 37 D. Bin, W. Huo, Y. Yuan, J. Huang, Y. Liu, Y. Zhang, F. Dong, Y. Wang and Y. Xia, Chem, 2020, 6, 968-984.
- 38 X. Tan, F. Zhang, D. Chen, P. Wang, Y. Liu, C. Meng and Y. Zhang, Chem. Eng. J., 2024, 489, 151119.
- 39 X. Tan, F. Zhang, D. Chen, J. n. Gong, J. Sun, C. Meng and Y. Zhang, J. Colloid Interface Sci., 2024, 669, 2-13.
- 40 K. Zhu, W. Jiang, Z. Wang, W. Li, W. Xie, H. Yang and W. Yang, Angew Chem. Int. Ed. Engl., 2023, 62, e202213368.
- 41 H. Geng, M. Cheng, B. Wang, Y. Yang, Y. Zhang and C. C. Li, Adv. Funct. Mater., 2019, 30, 1907684.
- 42 L. Chen, H. Nie, S. Zhou, G. Cao and A. Pan, Sci. China Mater., 2023, 66, 3453-3460.
- 43 Y. Hou, P. Ma, F. Long, M. Liu, Y. Zheng, L. Sun, J. Shi, K. Niu, J. Su, Y. Ma and Y. Gao, ACS Nano, 2024, 18, 27358-27371.
- 44 S. Hu, X. Huang, L. Zhang, G. Li, S. Chen, J. Zhang and X. Liu, Adv. Funct. Mater., 2023, 33, 2214161.
- 45 Y. Song, Q. Pan, H. Lv, D. Yang, Z. Qin, M. Y. Zhang, X. Sun and X. X. Liu, Angew. Chem., Int. Ed., 2021, 60, 5718-5722.
- 46 X. Mu, Y. Song, Z. Qin, J. Meng, Z. Wang and X.-X. Liu, Chem. Eng. J., 2023, 453, 139575.