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Excited-state proton-coupled electron transfer within ion pairs†

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The use of light to drive proton-coupled electron transfer (PCET) reactions has received growing interest, with recent focus on the direct use of excited states in PCET reactions (ES-PCET). Electrostatic ion pairs provide a scaffold to reduce reaction orders and have facilitated many discoveries in electron-transfer chemistry. Their use, however, has not translated to PCET. Herein, we show that ion pairs, formed solely through electrostatic interactions, provide a general, facile means to study an ES-PCET mechanism. These ion pairs formed readily between salicylate anions and tetracationic ruthenium complexes in acetonitrile solution. Upon light excitation, quenching of the ruthenium excited state occurred through ES-PCET oxidation of salicylate within the ion pair. Transient absorption spectroscopy identified the reduced ruthenium complex and oxidized salicylate radical as the primary photoproducts of this reaction. The reduced reaction order due to ion pairing allowed the first-order PCET rate constants to be directly measured through nanosecond photoluminescence spectroscopy. These PCET rate constants saturated at larger driving forces consistent with approaching the Marcus barrierless region. Surprisingly, a proton-transfer tautomer of salicylate, with the proton localized on the carboxylate functional group, was present in acetonitrile. A pre-equilibrium model based on this tautomerization provided non-adiabatic electron-transfer rate constants that were well described by Marcus theory. Electrostatic ion pairs were critical to our ability to investigate this PCET mechanism without the need to covalently link the donor and acceptor or introduce specific hydrogen bonding sites that could compete in alternate PCET pathways.

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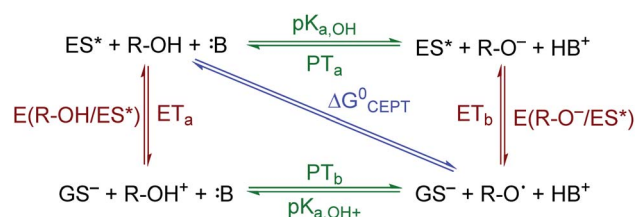
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Introduction

The creation of energy-rich fuels from small molecules is dependent upon the ability to effectively couple proton and electron transfer. There is growing interest in the use of solar energy to drive these proton-coupled electron transfer (PCET) reactions.^{1–4} To accomplish this feat, molecular systems that effectively couple light energy to proton and electron transfer are needed. Two approaches that have been utilized to couple light to PCET include the flash-quench technique, where photoexcitation of a sensitizer is followed by rapid electron-transfer quenching by a redox mediator to yield an oxidized or reduced sensitizer. This oxidized or reduced sensitizer then initiates a thermal PCET reaction with a secondary substrate.⁵

Alternatively, the excited sensitizer can directly participate in the PCET reaction.⁶ While excited-state PCET (ES-PCET) has been less explored than thermal PCET, there are a growing number of fundamental and application-based studies in the field.^{1,6–9}

ES-PCET reactions occur through three distinct mechanisms, Scheme 1,^{1,10} (1) stepwise electron transfer-proton transfer (ET_aPT_b), (2) stepwise proton transfer-electron transfer (PT_aET_b), and (3) concerted electron-proton transfer (CEPT), in which the proton and electron are transferred in the same step with a common transition state. The concerted mechanism is expected to be more valuable for selective catalysis as it avoids



Scheme 1 ES-PCET reaction diagram. ES* is the excited sensitizer, GS[−] is the reduced sensitizer, R-OH is the protonated substrate, and B is a base.

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100 mM TBAClO₄/CH₃CN was used. All potentials were externally referenced to Fc/Fc⁺ (630 mV vs. NHE).⁴⁷

UV-visible spectroscopy

UV-vis absorption spectra were acquired on a Varian Cary 50 UV-vis spectrophotometer in 1 cm path length spectrophotometric quartz cuvettes. Resolution of 1 nm was used.

Time-resolved photoluminescence and nanosecond transient absorption spectroscopy

Time-resolved photoluminescence and nanosecond transient absorption (TA) single wavelength kinetic data were collected on an Applied Photophysics spectrometer. Optical excitation was afforded by an OPO (opotek) pumped by a Q-switched, frequency tripled (355 nm) Nd:YAG laser (Quantel, Brilliant B). Pulses had an ~ 7 ns FWHM at 460 nm (*c.a.* 10 mJ per pulse). A pulsed Xenon lamp of an Applied Photophysics LKS60 setup provided probe light that was passed through 1 cm² quartz cuvette 90° to the laser and through a monochromator before hitting the P928 photomultiplier. For photoluminescence measurements the probe lamp was not used, and photoluminescence was detected at 90° to the incident laser through the same detection system. The PMT signal was converted and digitized using an HP Infiniium S5 digital oscilloscope (2G samples per s). Transient absorption traces were generated from the raw data using the LKS60 software.

Full transient absorption spectra were acquired on a Spectra Physics Quanta-Ray system with a frequency doubled (532 nm), Q-switch Nd:YAG laser. The pulse laser was connected to a transient absorption detection system (Edinburgh Instruments), equipped with a monochromator and a pulsed Xe arc lamp. A transient absorption spectrum of the sample was collected at 90° to the incident laser by a Tektronix 500 MHz digital oscilloscope coupled to a CCD camera. The output was processed with Edinburgh Instruments' L900 software. All data analysis was performed on OriginPro 2016 and 2017 software.

Titration

All samples were purged with nitrogen for 5–10 minutes prior to measurements and a flow of nitrogen was maintained over the samples during data collection. Ruthenium concentrations were held around 20–30 μM . Stock solutions of the salicylate derivatives were prepared at 5–10 mM and were titrated into the ruthenium samples in 10–100 μL amounts.

Cage escape quantum yields

Cage escape quantum yields were determined from the nanosecond transient absorption experiments through eqn (1). Ru(bpy)₃²⁺ was utilized as an actinometer assuming a unity yield of intersystem crossing. The $\Delta\epsilon_{450}$ between the ground-state Ru(bpy)₃²⁺ and the excited-state Ru(bpy)₃^{2+*} was $-1.5 \times 10^4 \text{ M}^{-1} \text{ cm}^{-1}$,⁴⁸ and $\Delta\epsilon_{510}$ between the ground-state **Ru-Bpz**⁴⁺ and the reduced **Ru-Bpz**³⁺ was $1.05 \times 10^4 \text{ M}^{-1} \text{ cm}^{-1}$. Salicylate concentrations of $\sim 75 \mu\text{M}$ were utilized as at this concentration >98% of the photoluminescence had been quenched.

$$\mathcal{O}_{\text{CE}} = \frac{\left(\frac{A_{j1}}{\varepsilon_{j1}}\right)_{\text{Ru-LL}}}{\left(\frac{A_{j2}}{\varepsilon_{j2}}\right)_{\text{Ru(bpy)}_3}} \frac{(1 - 10^{-A_{460}})_{\text{Ru(bpy)}_3}}{(1 - 10^{-A_{460}})_{\text{Ru-LL}}} \quad (1)$$

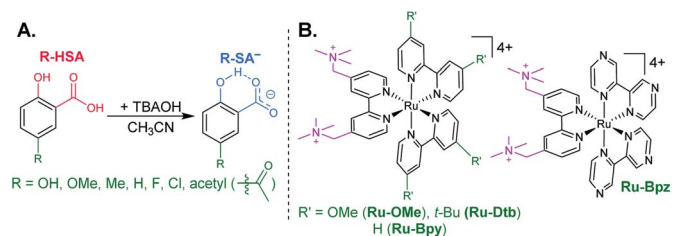
Results and discussion

Characterization of salicylate anions

The seven salicylic acids (**R-HSA**), where *R* is the functional group *para*- to the phenol (**OH**-, **OMe**-, **Me**-, **H**-, **F**-, **Cl**-, **acetyl**-), were readily deprotonated in acetonitrile (CH₃CN) through the *in situ* addition of tetrabutylammonium hydroxide 30-hydrate (TBAOH) to form the salicylate anions (**R-SA**⁻), Scheme 2. The deprotonation was monitored by UV-vis spectroscopy. Hypsochromic shifts of 10–20 nm (~0.1–0.18 eV) were accompanied by a slight decrease in the absorption intensity. This change in absorption was linear with respect to the TBAOH concentration up to one equivalent, upon which the spectral changes saturated. A set of isosbestic points was maintained throughout the titration indicative of clean conversion to the deprotonated anion. Fig. 1A–C shows the deprotonation for **H-HSA** to **H-SA**⁻, along with the extinction coefficient spectra of each entity. **Acetyl-HSA** showed different spectroscopic changes than the other six **R-HSA** derivatives. Upon the addition of TBAOH, a significant increase in the low-energy absorption intensity occurred, Fig. 1D.

The pK_a associated with this deprotonation in CH_3CN was determined through the spectrophotometric titration of 2,4-bis(tetramethylphenyl)-7-(dimethylamino)quinoline ($pK_a = 15.2$ in CH_3CN).⁴⁹ The measured pK_a of **H-HSA** (16.7) aligned well with that previously reported as 16.7.⁵⁰ For **OH-**, **OMe-**, and **Me-HSA**, the pK_a s were found in the range of 16.9–16.6. Whereas **F-**, **Cl-**, and **acetyl-HSA** were more acidic, in the range 15.8–15.4. The pK_a values are presented in Table 1.

Upon deprotonation, a small red-absorbing shoulder appeared at ~ 340 nm for all the **R-SA⁻** compounds, except **acetyl-SA⁻**. This absorption was previously identified for **H-SA⁻** in acetonitrile and ethanol as a proton-transfer tautomer, Fig. 2A.^{16,17} In this tautomeric form, the proton is localized on the carboxylate functional group instead of the phenolic oxygen. This tautomeric form was not expected at the outset of this study as previous PCET experiments in solely H₂O did not observe this tautomer.¹⁹ However, in non-polar organic solvents this tautomer has been characterized in both *intra*- and



Scheme 2 (A) Salicylate derivatives and (B) tetracationic ruthenium sensitizers.

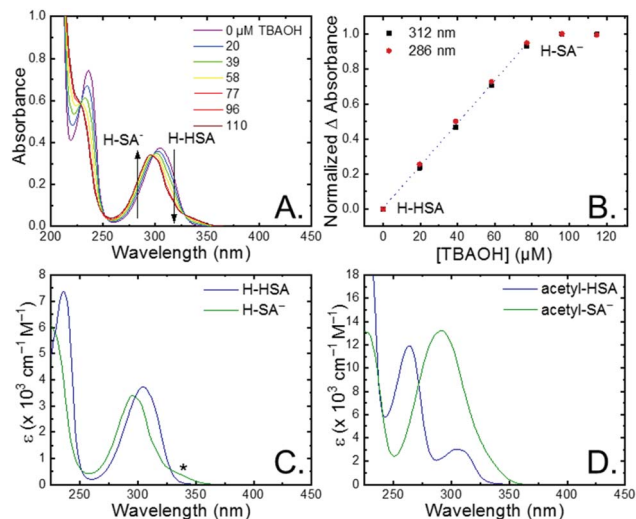


Fig. 1 (A) UV-vis absorption spectral changes upon the addition of tetrabutylammonium hydroxide to an $\sim 80 \mu\text{M}$ solution of H-HSA in CH_3CN . (B) Normalized change in absorbance monitored at the absorbance maxima for H-HSA and H-SA^- . Dotted line is a linear fit through the first five data points. (C and D) Extinction coefficient spectrum for R-HSA and R- SA^- when (C) R = H or (D) R = acetyl. Asterisk marks the tautomer absorbance.

intermolecular phenol-carboxylate systems.^{16,17,51} In the intermolecular studies, non-polar solvents were suggested to better stabilize larger, delocalized anions, such as expected for a phenolate, and the equilibrium between the normal and tautomeric form was found to shift towards the tautomer as solvent polarity decreased.⁵¹ As acetonitrile is a reasonably polar organic solvent, only a small amount of tautomer absorbance was found for the salicylate derivatives studied, barring acetyl- SA^- .

The addition of water to a CH_3CN solution of salicylate (H-SA^-) was shown to shift the equilibrium towards the normal isomer (protonated phenol).¹⁶ Therefore, to show that the low intensity absorbance measured for the expanded series of R- SA^- derivatives was due to this tautomer, titrations with deionized water were performed, Fig. 2. A loss of the tautomer absorbance was correlated to a growth of the normal isomer absorbance and confirmed our assignment of the low-energy absorbance to the tautomer.

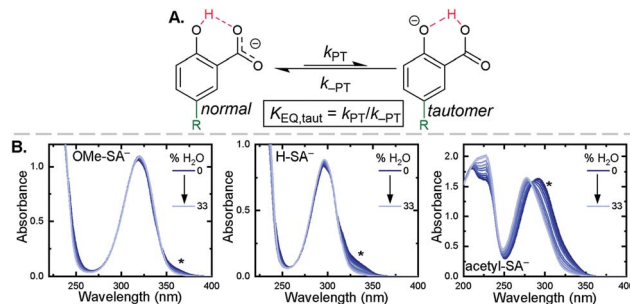


Fig. 2 (A) Salicylate proton-transfer tautomer. The normal and tautomeric forms are labelled. (B) Absorption spectral changes upon the addition of increasing amounts of H_2O to CH_3CN solutions of the indicated salicylate derivatives, $\sim 250 \mu\text{M}$. Asterisks mark the tautomer absorbance.

The absorption spectra of the salicylate derivatives in CH_3CN were fit with the sum of two Gaussian functions to approximate the absorbance of the tautomer. It was assumed that the extinction coefficient of the two species at the maximum absorbance was identical, as was previously done for similar phenol/carboxylate tautomers.⁵¹ Therefore, the ratio between the maxima absorbances gave the equilibrium constant for the tautomerization ($K_{\text{EQ,Taut}}$, Table 1. A measurable increase in the tautomer equilibrium constant coincided with increased electron withdrawing character of the functional group *para*- to the phenolic-OH. This substitution *para*- to the phenolic-OH group is known to decrease the intrinsic $\text{p}K_{\text{a}}$ of the phenolic-OH.¹⁸ This decreases the difference in $\text{p}K_{\text{a}}$ between the carboxylate and phenolic-OH, thus lowering the driving force for proton transfer and increasing the concentration of tautomer in solution.

The apparent reduction potentials of the salicylate derivatives were determined through cyclic voltammetry, $E_{\text{app}}(\text{R-SA}^{\bullet\text{O}}/\text{R-SA}_{\text{OH}}^-)$ where $\text{R-SA}_{\text{OH}}^-$ is the salicylate derivative before oxidation and $\text{R-SA}^{\bullet\text{O}}$ is the oxidized R- SA^- that has undergone an intramolecular proton transfer to the carboxylate functional group. The oxidation of the R- SA^- compounds was completely irreversible as expected for phenolic compounds that undergo an irreversible dimerization after oxidation.⁵² Also, because the proton transfer is coupled with electron transfer in the oxidation, a true one-electron reduction potential $E^\circ(\text{R-SA}_{\text{OH}}^-/\text{R-SA}_{\text{OH}}^{\bullet-})$ for R- SA^- could not be measured. However, the

Table 1 Redox potentials, photophysical properties, $\text{p}K_{\text{a}}$ s, and tautomer equilibrium constants for R- SA^-

R	$E_{\text{app}}(\text{R-SA}_{\text{OH}}^-/\text{R-SA}^{\bullet\text{O}})$ (V vs. NHE)	$\lambda_{\text{max,R-HSA}}$ (nm), $\epsilon_{\text{R-HSA}}$ ($\text{M}^{-1} \text{s}^{-1}$)	$\lambda_{\text{max,R-SA}^-}$ (nm), $\epsilon_{\text{R-SA}^-}$ ($\text{M}^{-1} \text{s}^{-1}$)	$\text{p}K_{\text{a}}$	$K_{\text{EQ,Taut}}$
OH	0.56 ^a	335, 4200	320, 4100	16.9	0.03
OMe	0.79	333, 4200	318, 4100	16.6	0.04
Me	0.97	314, 3800	304, 3600	16.9	0.07
H	1.10	304, 3700	295, 3400	16.7 ^b	0.10
F	1.11	314, 4400	305, 4100	15.8	0.07
Cl	1.16	316, 3500	307, 3200	15.6	0.11
Acetyl	1.30	305, 3100	292, 13 200	15.4	1.2

^a Oxidation of OH-SA^- was quasi-reversible. The value reported is the $E_{1/2}$. ^b Aligns with literature value of 16.7.⁵⁰

apparent reduction potential for the PCET reaction was estimated through a scan rate (ν) dependence, $E_{\text{app}}(\text{R-SA}^{\bullet}\text{O}/\text{R-SA}_{\text{OH}}^-)$.^{52–54} A plot of $\log(\nu)$ versus the oxidative peak potential was linear with slopes of 20–30 mV per decade, within a reasonable deviation from the theoretical 19.7 mV per decade expected for a PCET reaction, Fig. 3.⁵² The y-intercept was corrected for scan rate independent variables and the apparent PCET reduction potentials (E_{app}) are presented in Table 1.

Ion pair formation

Typical ruthenium tris-bipyridyl complexes have a 2+ cationic charge associated with the d^6 ruthenium center. The dicationic 4,4'-bis(trimethylaminomethyl)-2,2'-bipyridine (*tmam*) ligand has been used to increase the charge of the polypyridyl compounds by 2+ per ligand. This has been shown to enhance the formation of ion pairs between cationic ruthenium polypyridyl compounds and anions in acetonitrile (CH_3CN).^{23,44,55–57} The four ruthenium compounds utilized in this study follow the common structure of $[\text{Ru}(\text{tmam})(\text{LL})_2](\text{PF}_6)_4$, where LL was 4,4'-dimethoxy-2,2'-bipyridine (OMe), 4,4'-di-*tert*-butyl-2,2'-bipyridine (Dtb), 2,2'-bipyridine (Bpy), and 2,2'-bipyrazine (Bpz). Herein, the compounds will be denoted **Ru-LL** where LL is the short name of the derivatized bipyridine or bipyrazine ligand (Scheme 2B).

The electrochemical properties of the **Ru-LL** compounds, except **Ru-Bpz** which was previously reported,⁵⁸ were determined through square-wave voltammetry in 0.1 M $\text{TBAClO}_4/\text{CH}_3\text{CN}$ solution. The one-electron oxidation of **Ru-LL** (**Ru-LL**^{5+/4+}) was formally a ruthenium centered oxidation, $\text{Ru}^{\text{III/II}}$. These potentials were measured to be between 1.39 and 2.10 V, with a negative shift of the potential in the order **Ru-OMe** < **Ru-Dtb** < **Ru-Bpy** < **Ru-Bpz**. This trend followed the electron donating ability of the ancillary ligands. The first one-electron reduction of ruthenium polypyridyl compounds, **Ru-LL**^{4+/3+} has been shown to occur at the ligand that is the best π -acceptor.^{59,60} For **Ru-Bpz** this was the Bpz ligand, whereas for **Ru-OMe**, **Ru-Dtb**, and **Ru-Bpy**, the reduction was localized to a *tmam* ligand. The reduction of **Ru-Bpz** was reversible and ~300 mV more positive

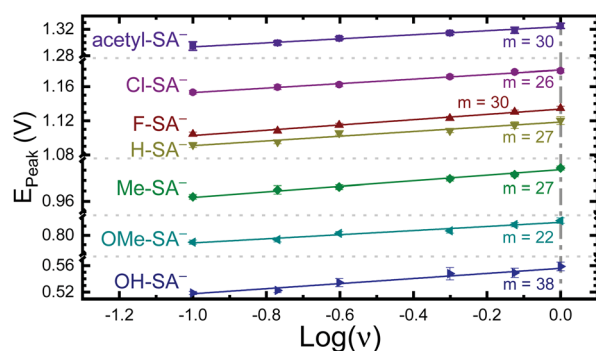


Fig. 3 Peak potential from irreversible cyclic voltammograms of the indicated **R-SA**[−]s vs. the log of the scan rate (ν). The indicated slopes (m) are given in mV per decade. Note the peak potential axis is broken in multiple places for clarity. Error bars are included and in most cases are similar in size to the symbol.

Table 2 Spectroscopic and thermodynamic properties of the **Ru-LL** compounds

Ru-LL	Ru-LL ^{5+/4+} (V)	Ru-LL ^{4+/3+} (V)	ΔG_{ES}^0 ^a (eV)	Ru-LL ^{4+*/3+} (V)	τ_0 (μs)
Ru-Bpz	2.10 ^b	−0.50 ^b	2.09	1.59	1.78
Ru-Bpy	1.57	−0.79	2.02	1.23	0.61
Ru-Dtb	1.51	−0.80	1.97	1.18	0.36
Ru-OMe	1.39	−0.86	1.88	1.02	0.17

^a Data from ref. 45. ^b Data taken from ref. 44.

than the reduction of the other three compounds confirming that Bpz was the ligand reduced. The other three compounds showed irreversible reductions associated with the *tmam* ligand. The irreversible nature of this reduction has been shown for similar compounds and ligands.^{22,61} The reduction potential was therefore estimated from the peak cathodic current of the square-wave voltammogram. The excited-state reduction potential (**Ru-LL**^{4+*/3+}) was estimated through eqn (2), where ΔG_{ES}^0 is the Gibbs free energy change from the ground state to the excited state, which was reported previously through a Franck–Condon line-shape analysis.⁴⁵ All electrochemical values are included in Table 2.

$$E_{\text{o}}(\text{Ru-LL}^{4+*/3+}) = E_{\text{o}}(\text{Ru-LL}^{4+/3+}) + \Delta G_{\text{ES}} \quad (2)$$

All four **Ru-LL** compounds exhibited UV-vis absorption spectra with transitions between 200–650 nm. The low energy absorption bands centered around 460 nm were assigned as metal-to-ligand charge-transfer (MLCT) transitions.⁶² Absorption features in the UV were due to ligand centered transitions. The addition of the **R-SA**[−] derivatives to CH_3CN solutions of **Ru-LL** induced changes in the UV-vis absorption spectra. Fig. 4 shows a representative example of **Cl-SA**[−] with **Ru-Dtb** and **Ru-**

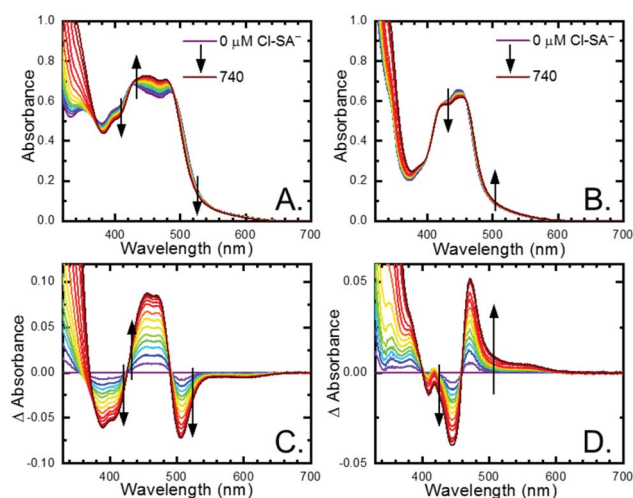


Fig. 4 UV-visible absorption spectra of (A) **Ru-Dtb** and (B) **Ru-Bpz** (~25 μM) upon the addition of 0 to 740 μM **Cl-SA**[−]. Difference spectra calculated by subtracting the absorption spectra at no **Cl-SA**[−] from the spectra with **Cl-SA**[−] present for (C) **Ru-Dtb** and (D) **Ru-Bpz**. Arrows indicate the spectral changes upon **Cl-SA**[−] addition.

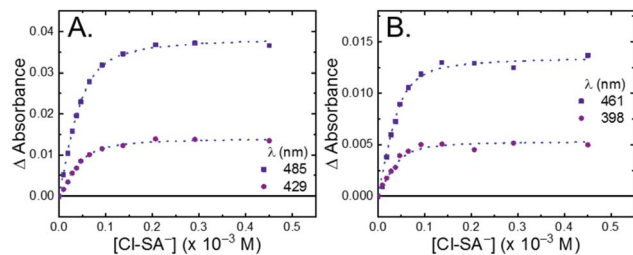


Fig. 5 Change in absorbance at the indicated wavelengths from the titration of (A) Ru-Dtb and (B) Ru-Bpz (25 μM) with Cl-SA $^-$. Wavelengths were chosen as they are the isosbestic points for the formation of the doubly ion-paired species and therefore the changes are only due to the first ion-pair formation. The blue dotted line is a fit to a 1:1 binding model.

Bpz. For Ru-Bpz, a bathochromic shift in the low-energy MLCT and decrease in the MLCT intensity occurred. For Ru-Dtb, Ru-OMe, and Ru-Bpy a hypsochromic shift of the MLCT was accompanied by an increase in MLCT intensity. Isosbestic points were maintained up to 5 equivalents of added R-SA $^-$. These changes were assigned to the formation of the ground-state ion pair, [Ru-LL $^{4+}$, R-SA $^-$] $^{3+}$. These changes could be well described by a 1:1 binding model,⁶³ Fig. 5, which provided the equilibrium constant ($K_{\text{EQ},1}$), for ion-pair formation, [Ru-LL $^{4+}$, R-SA $^-$] $^{3+}$. $K_{\text{EQ},1}$ ranged from 0.5 to 3×10^5 for all 28 combinations of Ru-LL and R-SA $^-$, Table 3. The association constants measured for these exclusively electrostatic ion pairs, formed in the relatively polar CH $_3$ CN, are significantly larger than those generally measured in hydrogen bond systems ($K \leq 10^3$). These electrostatic ion pairs are also orders of magnitude larger than salt-bridged systems in polar organic solvents ($K \approx 10^2$ to 10^3 in DMSO) and on the higher end of systems reported in nonpolar organic solvents ($K \approx 10^4$ to 10^5 in CH $_2$ Cl $_2$).^{31,33–37}

The isosbestic points shifted slightly, by <5 nm, upon R-SA $^-$ additions larger than 5 equivalents for many of the combinations. This shift indicated that a second R-SA $^-$ bound to the ion pair, [Ru-LL $^{4+}$, 2R-SA $^-$] $^{2+}$. A 1:1 binding model was used to fit the changes in the absorbance monitored at the first set of isosbestic points. This allowed the changes in absorbance associated with the second ion pairing to be estimated. This analysis yielded second equilibrium constants ($K_{\text{EQ},2}$) of $1\text{--}5 \times 10^3$; about two orders of magnitude lower than the first ion pair. The doubly ion-paired species will not be discussed further as the

Table 3 First ion-pairing equilibrium constants

Compounds	$K_{\text{EQ},1} (\times 10^5)$			
	Ru-Bpz	Ru-Bpy	Ru-Dtb	Ru-OMe
OH-SA $^-$	2.3 ± 0.8	1.4 ± 0.3	2.0 ± 0.2	1.8 ± 0.5
OMe-SA $^-$	1.5 ± 0.2	0.8 ± 0.1	1.1 ± 0.1	2.5 ± 0.6
Me-SA $^-$	2.4 ± 0.5	1.7 ± 0.2	1.8 ± 0.3	2.4 ± 0.7
H-SA $^-$	1.2 ± 0.2	1.2 ± 0.1	1.7 ± 0.2	3.0 ± 0.6
F-SA $^-$	0.9 ± 0.3	0.50 ± 0.04	1.1 ± 0.1	1.2 ± 0.1
Cl-SA $^-$	1.2 ± 0.2	0.55 ± 0.02	0.9 ± 0.1	1.0 ± 0.1
Acetyl-SA $^-$	1.4 ± 0.2	0.83 ± 0.07	2.0 ± 0.7	1.2 ± 0.1

concentrations utilized to study the ES-PCET reaction were not large enough for an appreciable amount to form.

Selected ion pairs were characterized by ^1H nuclear magnetic resonance (NMR). The addition of 0.5, 1, and 3 eq. of Me-SA $^-$ to CD $_3$ CN solutions of Ru-Bpz and Ru-Dtb led to measurable shifts in the resonances of the protons on the polypyridyl ligands, Fig. 6. Large downfield shifts were observed only in the resonances of the 3,3'- and CH $_2$ -protons on the dicationic tmam ligand, Fig. 6C. No significant shifts were detected for the proton resonances of the Dtb or Bpz ligands. This confirmed that the ion pair was formed through interaction of the R-SA $^-$ with the dicationic tmam ligand. Small upfield shifts were also quantified for the Me-SA $^-$ aromatic protons.

Proton-coupled electron transfer excited-state quenching

Pulsed laser excitation of the Ru-LL compounds in CH $_3$ CN led to room temperature, visible photoluminescence from the $^3\text{MLCT}$ excited state (Ru-LL $^{4+*}$). With no R-SA $^-$ present, exponential decays were fit to a first-order kinetic model. The excited-state lifetimes (τ) decreased with the electron donating ability of the ancillary bipyridine ligands, Ru-OMe < Ru-Dtb < Ru-Bpy < Ru-Bpz, Table 2.

The addition of salicylate anions to solutions of Ru-LL quenched the excited-state photoluminescence in 13 of the 28 [Ru-LL $^{4+}$, R-SA $^-$] combinations. These are listed in Table 4. The other combinations did not show excited-state quenching due to unfavorable driving forces ($\Delta G^0 > 0$) and short-lived excited-state lifetimes. For Ru-Bpz, excited-state quenching was observed with all seven R-SA $^-$ derivatives, Fig. 7. Upon the

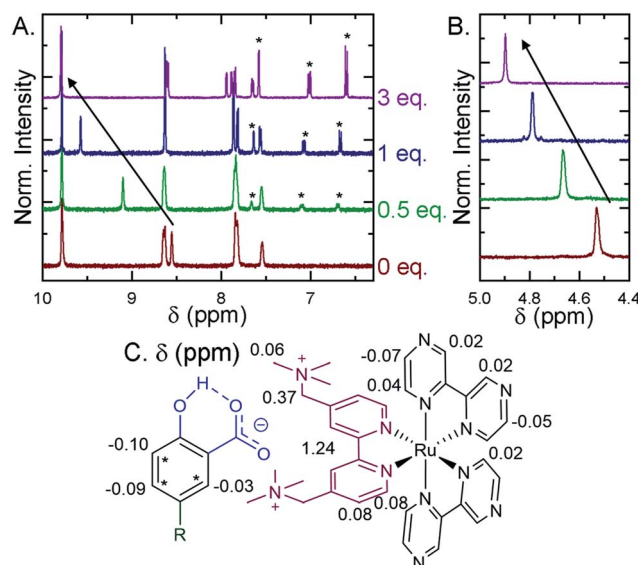


Fig. 6 ^1H Nuclear magnetic resonance spectra recorded in CD $_3$ CN for Ru-Bpz upon the addition of up to 3 eq. of Me-SA $^-$. (A) Shows the aromatic region, asterisks mark the proton resonances of Me-SA $^-$. (B) Shows the methylene resonances on the tmam ligand. Arrows show the downfield shift of the (A) 3,3'-tmam protons and (B) the methylene tmam protons. (C) The total change in chemical shift (Δppm) between 0 and 3 eq. of Me-SA $^-$ shown for all proton resonances on Ru-Bpz and Me-SA $^-$.

Table 4 Excited-state reduction potentials, thermodynamic driving forces, diffusional quenching rate constants, and ion-paired ES-PCET lifetimes and rate constants

Ru-LL	R-SA ⁻	ΔG_{PCET}^0 (eV)	$\Delta G_{\text{ET,b}}^0$ (eV)	$K_{\text{SV}} (\times 10^5 \text{ M}^{-1})$	$k_{\text{q}} (\times 10^{10} \text{ M}^{-1} \text{ s}^{-1})$	$\tau_{\text{PCET}} (\times 10^{-8} \text{ s})$	$k_{\text{PCET}} (\times 10^7 \text{ s}^{-1})$	$k_{\text{ET,b}} (\times 10^7 \text{ s}^{-1})$
Ru-Bpz	OH	-1.03	-1.12	1.15 ± 0.05	6.4	0.8	13	43
	OMe	-0.80	-0.88	1.30 ± 0.04	7.3	0.9	11	28
	Me	-0.62	-0.69	1.06 ± 0.06	5.9	1.4	7.1	10
	H	-0.49	-0.55	1.11 ± 0.05	6.2	2.1	4.8	4.8
	F	-0.48	-0.55	1.11 ± 0.03	6.2	2.4	4.2	6.0
	Cl	-0.44	-0.50	1.20 ± 0.02	6.7	2.6	3.8	3.5
	Acetyl	-0.29	-0.29	1.20 ± 0.04	6.7	4.4	2.3	0.2
Ru-Bpy	OH	-0.67	-0.76	0.07 ± 0.01	1.1	1.9	5.3	18
	OMe	-0.44	-0.52	0.10 ± 0.01	1.6	2.9	3.4	8.6
	Me	-0.26	-0.33	^a	^a	8.6	1.2	1.7
Ru-Dtb	OH	-0.62	-0.71	^a	^a	2.1	4.7	15
	OMe	-0.39	-0.47	^a	^a	4.5	2.2	5.6
Ru-OMe	OH	-0.46	-0.55	^a	^a	1.6	6.3	21

^a Dynamic quenching not observed over the range of concentrations studied.

addition of **R-SA⁻**, the time-resolved photoluminescence decays could not be modelled as first-order decays. Instead, a sum of two exponential decays (biexponential) was needed, eqn (3). The longer lifetime, τ_{d} , decreased with increased **R-SA⁻** concentration. A Stern–Volmer analysis,^{64,65} Fig. 7, of the lifetimes ($\tau_{\text{d}}/\tau_{\text{d,0}}$) was linear vs. the free concentration of **R-SA⁻** and gave a Stern–Volmer constant (K_{SV}) of around $1.1 \times 10^5 \text{ M}^{-1}$ for **Ru-Bpz**, eqn (4). K_{SV} is related to the bimolecular quenching rate constant (k_{q}) by the lifetime ($\tau_{\text{d,0}}$) of the excited state in the absence of quencher ($\tau_{\text{d,0}}$): $K_{\text{SV}} = k_{\text{q}}\tau_{\text{d,0}}$. The quenching rate constant was thus calculated to be $k_{\text{q}} \sim 6.2 \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}$ for all seven **R-SA⁻** derivatives, Table 4. This is close to the diffusion-limited rate constant for electron-transfer quenching of similar ruthenium complexes by iodide, $6.6 \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}$.^{66,67} Therefore, τ_{d} was assigned as the lifetime of the diffusional quenching reaction between non-ion-paired **Ru-Bpz** and **R-SA⁻**.

$$\text{PLI}_t = \text{PLI}_1 e^{-t/\tau_{\text{PCET}}} + \text{PLI}_2 e^{-t/\tau_{\text{d}}} \quad (3)$$

$$\frac{\tau_{\text{d,0}}}{\tau_{\text{d}}} = 1 + K_{\text{SV}}[\text{Q}]_{\text{free}} = 1 + k_{\text{q}}\tau_{\text{d,0}}[\text{Q}]_{\text{free}} \quad (4)$$

The shorter lifetime of the biexponential model was independent of the concentration of **R-SA⁻** and could be fixed throughout the titration. This lifetime varied from 8 ns for **Ru-Bpz** with **OH-SA⁻** to 86 ns for **Ru-Bpy** with **Me-SA⁻**, Table 4. The concentration independence of the rate constants indicated that the photoluminescence quenching occurred within the ion pair, and not from a diffusional reaction. This lifetime was assigned to ES-PCET within the ion pair, (τ_{PCET}), where $1/\tau_{\text{PCET}}$ provided the rate constant for ES-PCET (k_{PCET}).

Due to the shorter intrinsic lifetimes of **Ru-Dtb** and **Ru-OMe**, diffusional quenching was too slow to be detected at the concentrations utilized. A biexponential was still needed to model the excited-state relaxation in the presence of **OH-SA⁻**. The short lifetime still corresponded to the PCET lifetime, and the longer lifetime could be fixed to the lifetime of the complex without quencher.

The ³MLCT excited state of **Ru-Bpz** was produced for nano-second transient absorption spectroscopy through pulsed laser excitation. In ruthenium polypyridyl excited states the electron resides on the most electron withdrawing ligand.⁶⁸ Therefore, for **Ru-Bpz**, the excited state can formally be described as an oxidized Ru^{III} metal center with a reduced Bpz ligand, $[\text{Ru}^{\text{III}}(\text{tmam})(\text{Bpz})(\text{Bpz}^-)]^{4+*}$, **Ru-Bpz^{4+*}**. The appearance of absorption features that correspond to the reduced ligand, a positive delta absorbance at ~380 nm, and the loss of the ground-state MLCT, a negative delta absorbance at 450 nm, were indicative of the excited state, Fig. 8A. These features decayed to the ground-state with an identical lifetime to that of the time-resolved photoluminescence. In the case of **Ru-Bpy**, **Ru-Dtb**, and **Ru-OMe**, the localization of the electron in the MLCT excited state is expected to localize on the quaternary amine ligand (tmam) and thus the excited state can be formally described as $[\text{Ru}^{\text{III}}(\text{tmam}^-)(\text{LL})_2]^{4+*}$, where LL is Bpy, Dtb, or

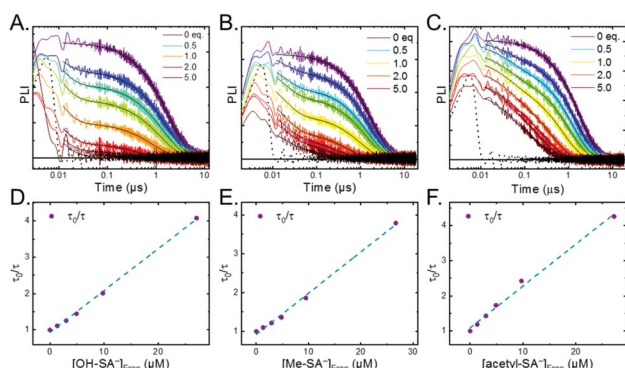


Fig. 7 A–C) Time-resolved photoluminescence decays for **Ru-Bpz** (25 μM) upon the addition of up to 5 equivalents of (A and D) **OH-SA⁻**, (B and E) **Me-SA⁻**, (C and F) **acetyl-SA⁻**. Black dotted lines are the instrument response function based on a scattered laser pulse with no sample present. (D–F) Dynamic Stern–Volmer plots. Blue dashed line is a linear fit to the data.



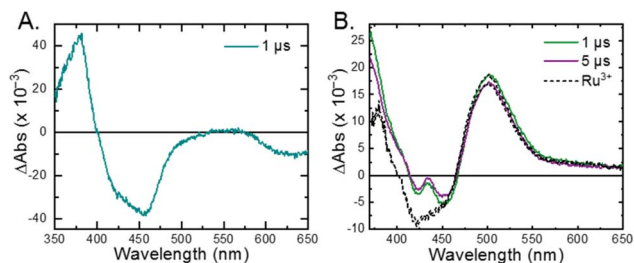
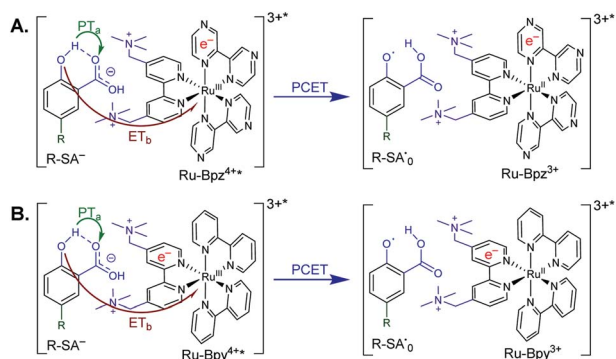


Fig. 8 (A) Excited-state transient absorption spectrum of Ru-Bpz⁴⁺* obtained 1 μs after laser excitation. The negative change in absorbance at wavelengths longer than 570 nm is from uncorrected emission. (B) Transient absorption spectra collected 1 and 5 μs after laser excitation of Ru-Bpz⁴⁺* (25 μM) in the presence of Cl-SA⁻ (75 μM). Overlaid is the normalized Ru-Bpz³⁺ delta absorbance spectra. Deviation from the Ru³⁺ spectra is due to the absorbance of R-SA[•]O.

OMe. This change in excited-state electron localization has been previously proposed in these and similar compounds.^{44,45}

Excitation of a solution of Ru-LL in the presence of R-SA⁻ gave significant differences in the nanosecond transient absorption spectra. Fig. 8B shows representative spectra collected 1 and 5 μs after excitation of the [Ru-Bpz⁴⁺, Cl-SA⁻]³⁺ ion pair. The appearance of an absorption band centered around 510 nm was consistent with the formation of the reduced ruthenium complex ([Ru^{II}(tmam)(bpz)(bpz⁻)]³⁺, Ru-Bpz³⁺). To confirm this, the Ru-Bpz³⁺ delta absorbance spectrum was generated through reductive excited-state quenching by tri-*p*-tolylamine. This spectrum could be normalized to the spectra of the reduced ion pair at lower energy wavelengths (>450 nm). A positive deviation from the reduced complex spectra was present at higher energies, around 430 nm, Fig. 8B. It has been reported that an oxidized phenoxyl radical (PhO[•]) absorbs light in this region, *e.g.* the tyrosine phenoxyl radical has an absorption at 410 nm.^{69–71} Therefore, this absorption was assigned to the oxidized salicylate, in which the proton has transferred to the carboxylate group, R-SA[•]O, Scheme 3. This provided a clear indication that the excited state was quenched by an ES-PCET reaction. Similar spectral features were obtained for all [Ru-LL⁴⁺, R-SA⁻]³⁺ ion pairs that showed photoluminescence quenching.



Scheme 3 Generic mechanism for the ES-PCET reaction within the [Ru-LL⁴⁺, R-SA⁻]³⁺ ion pair, (A) Ru-Bpz and (B) Ru-Bpy. Green arrow shows the proton transfer and red arrow shows the electron transfer.

Single wavelength kinetic analysis at wavelengths near the isosbestic points between the ground and excited state of Ru-Bpz (405 nm and 510 nm) in the presence of R-SA⁻ allowed the formation of Ru-LL³⁺ (510 nm) and R-SA[•]O (405 nm) to be monitored independent of the excited-state decay, Fig. 9. The emission decay was also monitored to directly compare the rates of formation with the excited-state decay. The absorption changes both at 405 nm and 510 nm showed a biexponential signal rise and yielded lifetimes that agreed with the time-resolved photoluminescence titrations, Fig. 9. In fact, the ratio of the pre-exponential factors for the two lifetimes, $A_{\text{PCET}}/A_{\text{d}}$, aligned with those of the excited-state decay. This ratio also aligned with the expected ratio of free ruthenium complex to ion-paired ruthenium, Ru-LL_{free}: [Ru-LL⁴⁺, R-SA⁻]³⁺, based on the equilibrium constant for ion-pair formation, $K_{\text{EQ},1}$. This indicated that quenching through both the diffusional reaction and from the pre-formed ion pair occurred through ES-PCET. Both the oxidized salicylate radical and the reduced ruthenium complex could be identified as primary photoproducts, as both R-SA[•]O and Ru-LL³⁺ had identical formation rate constants that aligned with excited-state decay.

Cage escape quantum yields (ϕ_{CE}) of the photoproducts were estimated at 3 eq. of the respective R-SA⁻ with Ru-Bpz. The change in absorbance at 510 nm, which corresponds solely to the reduced ruthenium complex, was used to estimate the concentration of cage escaped products and the unity internal conversion efficiency of Ru(bpy)₃²⁺ used as an actinometer. Quantum yields of 0.60–0.70 were calculated for all R-SA⁻ derivatives. These cage escape yields are significantly larger than those found in the diffusional excited-state electron-transfer oxidation of iodide by ruthenium excited states (<0.10)^{66,72} and larger than ion-pairs formed between iodide and ruthenium polypyridyl complexes with cage yields of 0.34⁷³ and 0.25–0.50.²¹ For PCET in hydrogen-bonded systems, ϕ_{CE} is often not measured, or irrelevant because the products never dissociate. In cases where measured, ϕ_{CE} has been much lower than here, see *e.g.* ref. 74. The larger cage escape quantum yields measured here can be rationalized based on changes in coulombic forces in tandem with the proton motion coupled

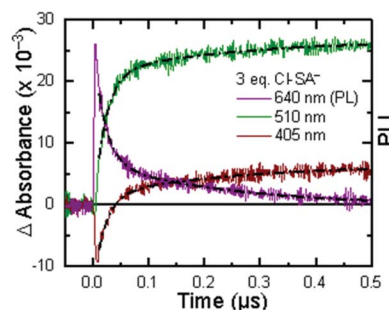


Fig. 9 Single wavelength transient absorption kinetics monitored at 405 nm, red, and 510 nm, green, for [Ru-LL⁴⁺, Cl-SA⁻]³⁺. The purple trace is the excited-state decay. Overlaid in black are biexponential fits with the two lifetimes shared between the three traces. The negative signal is due to small amounts of excited-state bleach still present of the monitored wavelengths.

with electron transfer. Upon oxidation, the salicylate derivatives become neutral, which removes the coulombic attraction of the ion-pair. Secondly, unlike simple single-electron-transfer systems, the oxidation of salicylate involves the movement of a proton from the phenolic-OH to the carboxylate functional group. Thus, recombination must also involve the movement of both the proton and electron. The driving force for this proton-coupled back-electron transfer would necessitate considerations of the driving force for both the electron transfer and proton transfer, which may slow the back-reaction rate constant allowing cage escape to compete with and surpass recombination within the solvent cage.

Mechanistic discussion

The oxidation of salicylate by photooxidized ruthenium complexes is known to occur through a PCET reaction in water.^{19,75} However, to the best of our knowledge, no studies to date have investigated the oxidation mechanism of salicylate in organic solvents. Above, it was shown that the excited state of cationic **Ru-LL** compounds could oxidize salicylate in CH₃CN. Nanosecond transient absorption spectroscopy confirmed that the reaction proceeded through ES-PCET. Time-resolved photoluminescence experiments showed that the ES-PCET reaction occurred through both a diffusional reaction between non-associated pairs and within the photoexcited ion pairs. While the diffusional reaction occurred near the diffusion limit, $6 \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}$, preventing mechanistic analysis, within the pre-formed ion pairs the first-order ES-PCET rate constants were measured directly. Below, we analyse these rate constants along with the above results to probe the ES-PCET mechanism.

As stated in the introduction, this ES-PCET mechanism could proceed through either (1) electron transfer-proton transfer, ET_aPT_b, (2) proton transfer-electron transfer, PT_aET_b, or (3) concerted electron-proton transfer, CEPT. Previously, the concerted mechanism was identified for the oxidation of salicylate in water through flash-quench transient absorption spectroscopy and was initially expected to be the active mechanism in CH₃CN.^{19,75}

The appearance of a ground-state tautomer in the absorption spectra of **R-SA**[−] suggested, however, that the salicylate oxidation mechanism in organic solvents may differ from that in water. In acetonitrile, a stepwise proton-transfer, electron-transfer mechanism could potentially be favored. The ground-state proton transfer would allow electron transfer to occur through the phenolate (**R-O**[−]), which for **H-SA**[−] is known to have a more negative redox potential (0.77 V vs. NHE in water) than the protonated phenol (1.48 V in water).¹⁹ Hence, there is a significantly larger driving force for electron transfer through the phenolate. This also suggests that a stepwise electron-

transfer, proton-transfer reaction is unfavorable, because the driving force for electron transfer from the protonated phenol would be uphill or have a small favorable driving force in the case of **Ru-Bpz**. Therefore, the possibility that the ES-PCET mechanism follows a stepwise proton transfer-electron transfer mechanism was investigated, depicted in Scheme 4.

$$\Delta G_{\text{PCET}}^0 = \mathcal{F}[E_{\text{app}}(\text{R-SA}_0^{\bullet}/\text{R-SA}_{\text{OH}}^{\bullet-}) - E^0(\text{Ru-LL}^{4+/3+})] \quad (5)$$

$$\Delta G_{\text{PT}_a}^0 = -\frac{\ln(K_{\text{EQ,Taut}})}{RT} \quad (6)$$

$$\Delta G_{\text{PCET}}^0 = \Delta G_{\text{PT}_a}^0 + \Delta G_{\text{ET}_b}^0 \quad (7)$$

$$\Delta G_{\text{ET}_b}^0 = \Delta G_{\text{PCET}}^0 - \frac{\ln(K_{\text{EQ,Taut}})}{RT} \quad (8)$$

$$\begin{aligned} \text{Rate} &= \frac{d[\text{Ru-LL}^{3+}, \text{R-SA}_0^{\bullet-}]^{3+}}{dt} = \frac{k_{\text{PT}_a}k_{\text{ET}_b}}{k_{-\text{PT}_a}} [\text{Ru-LL}^{4+}, \text{R-SA}_{\text{OH}}^{\bullet-}] \\ &= K_{\text{EQ,Taut}}k_{\text{ET}_b} [\text{Ru-LL}^{4+}, \text{R-SA}_{\text{OH}}^{\bullet-}] \end{aligned} \quad (9)$$

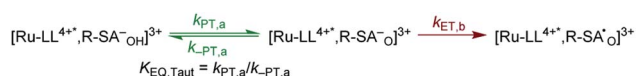
$$k_i = \sqrt{\frac{4\pi^3}{h^2\lambda_i k_B T}} |H_{\text{ab},i}|^2 e^{-\frac{(\lambda_i + \Delta G_i^0)^2}{4\lambda_i k_B T}} \quad (i = \text{ET}_b \text{ or PCET}) \quad (10)$$

$$\Delta G^{\circ} = N_A \{e[E^{\circ}(\text{D}^{+}/\text{D}) - E^{\circ}(\text{A}/\text{A}^{\bullet-})] + \omega(\text{D}^{+}\text{A}^{\bullet-}) - \omega(\text{DA})\} - \Delta G_{\text{ES}} \quad (11)$$

$$\omega = \frac{[z(\text{A}) - z(\text{D}) - 1]e^2}{4\pi\epsilon_0\epsilon_r a} = \frac{[\Delta z - 1]e^2}{4\pi\epsilon_0\epsilon_r a} \quad (12)$$

To start, the Gibbs free energy change for the overall PCET reaction within the ion pair, ΔG_{PCET}^0 , was estimated through eqn (5),¹⁹ E_{app} is the apparent reduction potential for the PCET and encompasses the driving force for proton transfer and the salicylate reduction potential, and is reported in Table 4. This overall driving force can be broken into a sum of the driving force for the initial PT_a between the phenolic-OH and carboxylate, $\Delta G_{\text{PT}_a}^0$, and secondary ET_b from the phenolate to the ruthenium excited state, $\Delta G_{\text{ET}_b}^0$. The estimated tautomerization equilibrium constant, $K_{\text{EQ,Taut}}$, Table 1, was related to $\Delta G_{\text{PT}_a}^0$ through eqn (6). Subtraction of $\Delta G_{\text{PT}_a}^0$ from ΔG_{PCET}^0 provided the driving force for the excited-state electron transfer ($\Delta G_{\text{ET}_b}^0$), eqn (7) and (8). To elucidate the electron-transfer rate constants, the pre-equilibrium approximation was used to develop a rate law based on Scheme 4, eqn (9). Through this assumption, the k_{PCET} measured *via* time-resolved photoluminescence is equal to $K_{\text{EQ,Taut}} \times k_{\text{ET}_b}$, and thus simple division provided k_{ET_b} .

A plot of k_{ET_b} vs. ΔG_{ET} in Fig. 10A showed the increase in rate constant with increasing driving force for electron transfer in the normal to near barrierless region as described by Marcus theory.^{40,76} **Ru-Bpz** provided the most complete data set (7 points) and these data provided a reasonable fit to Marcus theory, eqn (10).^{40,43} The fit, dashed line in Fig. 10A, allowed the two variable parameters of the Marcus equation, λ and H_{ab} , to



Scheme 4 Proposed PTET reaction mechanism within the photoexcited ion pairs.



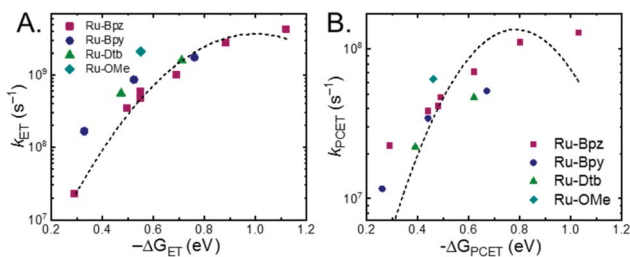


Fig. 10 (A) Relationship between the rate constants for ET_b within the ion pair and the driving force for electron transfer. (B) Relationship between the observed PCET rate constant and driving force for PCET (note the Y-axis spans a much smaller range than in (A)). Dashed lines are best fits to the Marcus equation with a Y^2 weighting restricted to the data points for **Ru-Bpz** (purple, squares). Note, the fit is significantly better for ET over PCET.

float. All constants were fixed to their known values and the temperature was fixed as 298 K (room temperature). Thus, the fit is rather well-defined by the curvature and the approximate rate maximum. From this fit a reorganization energy, λ , of 1.0 eV (8100 cm^{-1}) was approximated. This is close to the expected ~ 1.0 eV for electron transfer with ruthenium polypyridyl compounds.⁷⁷ The electronic coupling, H_{ab} , was estimated as 2.5×10^{-4} eV (2 cm^{-1}), indicative of a non-adiabatic electron transfer. For comparison, a fit to k_{PCET} vs. $-\Delta G_{\text{PCET}}$ using eqn (10) ($i = \text{PCET}$) resulted in a very poor fit (Fig. 10B); a clear indication that a concerted mechanism was not operational under our conditions. No turnover of the kinetics to a Marcus inverted barrier was discernible. Note that the data are first-order rate constants within the ion pairs that are not limited by diffusion. Nevertheless, the minimal curvature in the plots reduces the accuracy of the fit. This may impact the true values of λ and H_{ab} , but their magnitudes should be reasonable approximations.

The other three ruthenium complexes were also quenched through the same mechanism with the most electron donating **R-SA⁻** derivatives, **OH-**, **OMe-**, and **Me-**, Fig. 10A. However, the estimated $k_{\text{ET},b}$ deviated from the Marcus curve defined by **Ru-Bpz**, with most of the rate constants larger than would be expected. The reason for this deviation is unclear, however we believe it comes about due to the difference in the localization of the excited electron in the photoexcited ruthenium complexes. As shown in Scheme 3, the excited state of **Ru-Bpz** localizes the excited electron on the ancillary Bpz ligand away from the direction of charge transfer. Whereas, for **Ru-Bpy**, **Ru-Dtb**, and **Ru-OMe** the electron is localized on the cationic tmam ligand directly between the bound salicylate and the ruthenium. At a first glance this difference in localization should slow the electron transfer as the electron localized on the tmam ligand would repel the anionic salicylate,⁷⁸ lengthening the electron transfer distance. It may also be expected that electron transfer across a reduced bipyridine ligand (by super-exchange or hopping)⁷⁹ would be less favorable, due to the high energy of the (virtual) intermediate with two electrons added to the ligand. However, facile charge transfer across reduced bipyridine

ligands has been observed in the case of iridium³¹ complexes and therefore, we suggest that the orientation of the acceptor Ruthenium d-orbitals relative to the salicylate, which should not depend on the $^3\text{MLCT}$ localization, are more likely to matter than the charge on the bipyridine ligand.

Another factor that electron localization would affect is the electrostatic work terms, ω , for the electron-transfer reaction. These work terms account for the free energy needed to bring the reactants together and to separate the products and should be included in the approximation of the driving force for electron transfer, and through extrapolation a PCET, eqn (11), where D and A are the electron donor and acceptor reactants, N_A is the fundamental charge, and e is the elementary charge.^{56,80} These work terms can be simplified to eqn (12), where z is the charge of the reactant, ϵ_0 is the vacuum permittivity, ϵ_r is the solvent dielectric, and a is the distance between the donor and acceptor. In this study, the difference in charge is $3+$, which means the work term is positive, decreasing the overall driving force for the PCET reaction.

The variation between **Ru-Bpz** and the other complexes may come from the difference in charge distribution in the excited state. For **Ru-Bpy**, **Ru-Dtb**, and **Ru-OMe**, the decrease in positive charge near the salicylate and along the electron-transfer vector due to electron localization on the tmam ligand would lower the magnitude of the work term and increase the overall driving force for the reaction, shifting the data for those three complexes to more negative ΔG values in Fig. 10A. The reverse occurs for **Ru-Bpz**. Therefore, the $\Delta G_{\text{ET},b}^0$ for **Ru-Bpy**, **Ru-Dtb**, and **Ru-OMe** is underestimated relative to **Ru-Bpz** and the rate constants deviate from the fit. Typically, a spherical approximation of the reactants and products could provide a reasonable estimate for the work terms.⁵⁵ However, for the **Ru-LL** compounds studied this would not differentiate the two excited states. Due to the complexity of charge distribution in the ruthenium complexes and salicylate derivatives, developing a reasonably accurate approximation for the electrostatic work terms is difficult.^{46,73,81,82} However, within a single ruthenium complex, such as **Ru-Bpz**, these work terms are not expected to vary significantly, and as such we chose to model solely the **Ru-Bpz** data as it provided the largest uniform series, Fig. 10A.

The weak electronic coupling constant measured for the electron-transfer step of the PCET reaction implies a non-adiabatic electron transfer. This is expected for an outer-sphere electron transfer reaction from salicylate to the excited ruthenium complex. In overall non-adiabatic concerted proton-coupled electron transfer reactions, the CEPT coupling constant (V_{CPET}^2) is approximately equal to the combination of the electronic coupling and overlap between the proton vibrational wavefunctions, $V_{\text{CPET}}^2 \approx V_{\text{ET}}^2 \times S_{\text{PT}}$.¹⁵ The present case is extreme to which the electron transfer can be thought of as gated by internal proton transfer within salicylate. Thus, we were able to deconvolute the electron-transfer rate constants, and in turn estimate the electronic coupling constant for electron transfer (V_{ET} , also referred to as H_{ab}). This non-adiabaticity of the electron transfer does not indicate that the coupling between the electron and proton is weak.^{83–85} This coupling is intrinsic to the coupling of electron



+340 mV in water).^{18–20} In acetonitrile this is not the case as the pK_a s measured for the unsubstituted phenol and benzoate do not account for the free energy change associated with the competitive stabilization of the phenol *vs.* the carboxylate within the conjugated salicylate. However, the ability to measure the tautomer equilibrium constant provides an alternative means to access the driving force for proton transfer which was found to be on average an order of magnitude lower than that in water ($\Delta G_{PT,a} +55$ mV in acetonitrile for **H-SA**[−]). This order of magnitude decrease in the proton-transfer driving force is a key factor in the disparate mechanism of salicylate oxidation in acetonitrile *vs.* water.

The PT_aET_b oxidation of salicylate reported here comes about due to the appearance of the proton transfer tautomer of the strong internal hydrogen bond. Many groups have studied the oxidation of phenols with internal hydrogen bonds.^{84,86} The ability to tune the structure, proton transfer distance, and hydrogen bond strength make these systems valuable for fundamental studies. We used salicylate as the known, strong internal hydrogen bond was hoped to favor H^+ tunneling (due to a large S_{PT}) and potentially a CEPT mechanism, as has been proposed by others under aqueous conditions.^{19,20} However, the present work shows that these strong hydrogen bonds may also favor proton transfer and a PT_aET_b mechanism. The appearance of a low-energy absorption that was assigned to the ground-state proton-transfer tautomer in the salicylate derivatives studied here aligns with a model proposed by Limbach and co-workers who investigated the same phenomena in *intermolecular* hydrogen bonds between phenols and carboxylates. They defined the “localized charge solvation” concept, which states that “an increase in the solvent polarity induces proton transfer in the sense that charge is transferred toward the acceptor less capable of charge delocalization.”⁵¹ Furthermore, they state that aprotic, non-polar solvents are better at stabilizing large, delocalized anions such as a phenolate *vs.* small, localized charges such as a carboxylate. Stated another way, the proton localization along the hydrogen bond between the salicylate phenol and carboxylate oxygens is determined by the difference in stabilization energy between the two anionic groups in the solvent of interest. Therefore, while the equilibrium constants measured in acetonitrile were ≤ 1 , if salicylate were to be dissolved in an even more non-polar solvent such as CH_2Cl_2 , the equilibrium would shift further toward the tautomeric form and the equilibrium constants increase. In the case where no tautomer was observed, water, it was possible to estimate the driving force for *intramolecular* proton transfer through the difference in the pK_a between the unsubstituted phenol and benzoate ($\Delta G_{\text{PT},a} =$

This work highlights the breadth with which ion pairs may facilitate the study of ES-PCET. Without the need to ensure either the ruthenium complexes or salicylates were covalently connected or had directing hydrogen bonding functionality, 28 combinations of ion pairs could be investigated with 13 providing measurable ES-PCET reactivity. The use of ion pairs facilitated direct measurement of the PCET rate constants and evaluation of the PCET mechanism for salicylate oxidation in organic media. Furthermore, the rate constants for electron transfer within the PCET reaction were found to fall within the normal to near-activation-less region of Marcus' parabola. Most systematic studies of PCET mechanisms have reported rate constants within the linear regime of Marcus theory.⁸⁸⁻⁹² Only recently has clear evidence for the Marcus inverted region for concerted PCET been disclosed, in which Mayer and coworkers reported a series of covalently linked donor-acceptor dyads that underwent concerted forward PCET in the normal region of the Marcus parabola and inverted region kinetics for the back reaction.^{9,30} An important feature in these systems was the necessity of covalently linking all three components of the PCET reaction. Ion pairing offers a potential way to remove this limitation while maintaining a first-order reaction. This ability of ion-pairs to reduce the reaction order and overcome diffusion without synthetic difficulty of covalent bonds or need of linked hydrogen bonds has implications not only in fundamental mechanistic studies, but also in applications toward solar fuels^{1,4,93,94} and photosensitized organic synthesis.^{8,95}

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