Nanoscale

PAPER

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Cite this: Nanoscale, 2024, 16, 7011

Received 22nd January 2024, Accepted 5th March 2024 DOI: 10.1039/d4nr00326h

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Introduction

Noble metal nanoclusters (NCs), predominantly composed of gold, silver, and copper, represent a novel and promising class of functional materials exhibiting numerous unique characteristics.¹⁻¹⁹ These atomically precise NCs stand apart from nanoparticles, showcasing significant differences in various properties including electronic and absorption properties,^{1,2,5-8} photoluminescence,^{1,2,6-9} magnetism,^{10,11} catalytic activity,12-14 and biomedical applications.15,16 Recent studies have established that the intrinsic properties of a homometallic NC can be altered through the incorporation of heterometals, leading to improved properties.^{1,2,6,8,17-22} Contemporary research is now largely devoted to designing and tailoring the properties of NCs to suit particular applications in a predictable and manageable manner. Specifically, the stability and catalytic activity of Au NCs can be significantly

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Doping effect on a two-electron silver nanocluster†

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This study investigates the effects of metal addition and doping of a 2-electron silver superatom, $[Ag_{10}\{S_2P(O^{j}Pr)_2\}_8]$ (Ag_{10}). When Ag^+ is added to Ag_{10} in THF solution, $[Ag_{11}\{S_2P(O^{j}Pr)_2\}_8](OTf)]$ (Ag_{11}) is rapidly formed almost quantitatively. When the same method is used with Cu^+ , a mixture of alloys, $[Cu_xAg_{11-x}\{S_2P(O^{j}Pr)_2\}_8]^+$ (x = 1-3, Cu_xAg_{11-x}), is obtained. In contrast, introducing Au^+ to Ag_{10} leads to decomposition. The structural and compositional analysis of Ag_{11} was characterized by single-crystal X-ray diffraction (SCXRD), ESI-MS, NMR spectroscopy, and DFT calculations. While no crystal structure was obtained for Cu_xAg_{11-x} , DFT calculations provide insights into potential sites for copper location. The absorption spectrum exhibits a notable blue shift in the low-energy band after copper doping, contrasting with that of the slight shift observed in 8-electron Cu-doped Ag nanoclusters. Ag_{11} and Cu_xAg_{11-x} are strongly emissive at room temperature, and solvatochromism across different organic solvents is high-lighted. This study underscores the profound influence of metal addition and doping on the structural and optical properties of silver nanoclusters, providing important contributions to understanding the nanoclusters and their photophysical behaviors.

enhanced by doping with Pt, Pd, Ag, Cu, or other foreign metal atoms.^{23–26} Furthermore, Au/Ag bimetallic NCs have been observed to display greater photoluminescence quantum yield (PLQY) when compared to homometallic Ag NCs.^{22,27} It is important to note that the field of NC chemistry still holds a plethora of uncharted potential properties that is yet to be explored and understood.

Copper, as well as the other group 11 metals, possesses multiple coordination versatility and has attracted significant interest due to its outstanding efficacy in CO₂ reduction.^{28,29} However, when incorporated as a dopant into NCs, Cu exhibits distinct influences on the structure and characteristics of the NCs, in contrast to the effects observed with Au or Ag atoms. In the case of Cu alloying, the Cu atoms tend to position themselves in the shell layer of the core-shell NCs.³⁰⁻³⁵ Zhu et al. reported bimetallic NCs, $Au_{36-x}Cu_x(m-MBT)_{24}$ (x = 1-3),³⁰ $Au_{38-x}Cu_x(2,4-DMBT)_{24}$ (x = 0-6, DMBT = 2,4-dimethylbenzenethiolate),³¹ and $[Ag_{62-x}Cu_xS_{12}(SBu^t)_{32}]^{4+}$ (x = 10-21).³² Liu et al. reported bimetallic $Cu_xAg_{20-x}{S_2P(OR)_2}_{12}$ (x = 3, 4; R = ^{*i*}Pr,¹⁹ ^{*n*}Pr³³). These studies consistently demonstrated a marked tendency for Cu atoms to favor surface positions, displaying disordered arrangements across multiple sites instead of situating themselves within the core. This behavior can be attributed to the lower reduction potential of Cu⁺ (0.52 and 0.80 V for Cu^{I}/Cu^{0} and Ag^{I}/Ag^{0} , respectively, vs. SHE, T = 25 °C, 1 atm), which hampers the anti-galvanic reaction (AGR) with the kernel Ag/Au atoms, making it challenging to proceed.³⁶ Consequently, a metal exchange occurring at the surface



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[†]Electronic supplementary information (ESI) available: NMR spectrum, TD-DFT spectrum, XPS, photophysical data. CCDC 2287343. For ESI and crystallographic data in CIF or other electronic format see DOI: https://doi.org/10.1039/ d4nr00326h

M(I) atoms emerges as a more favorable process. To the best of our knowledge, the chemistry of 2-electron Ag/Cu superatomic alloys has never been explored. Kappes et al. reacted the 8-electron cluster $[Ag_{29}(BDT)_{12}(PPh_3)_4]^{3-}$ (BDT = 1,3 benzenedithiolate) with $Cu_{12}S_6(DPPPT)_4$ (DPPPT = bis-(diphenylphosphino) pentane), yielding the alloys $\left[Ag_{29-x}Cu_x(BDT)_{12}\right]^{3-}$ (x = 0-13).³⁷ Although they could not obtain any crystal structures, they performed DFT calculations to determine the Cu locations in the alloy shells and found a preference for the outermost positions. Conversely, in the case of the 46-electron $[Cu_{30}Ag_{61}(SAdm)_{38}S_3]^+$, the Cu atoms are found to be located within the interior of the NCs.³⁸ This compound presents an onion-type structure, conceptualized as [Ag₁₃@Cu₃₀@Ag₄₈(SAdm)₃₈S₃]⁺. Here, the Cu atoms are situated in the second inner shell, sandwiched between the Ag13 icosahedral core and the outermost $Ag_{48}(SAdm)_{38}S_3$ shell. This distinct arrangement highlights the unique distribution of Cu atoms participating in the NC's architecture.

Following our previous research on the doping effect on 8-electron superatoms,^{16,33} we chose the 2-electron Ag NC $[Ag_{10}(dtp)_8] (dtp = S_2P(O^iPr)_2) (Ag_{10} \text{ for short})^{39}$ as the starting template. The subsequent addition of Ag(1) and Cu(1) ions yield $[Ag_{11}(dtp)_8(SO_3CF_3)] (Ag_{11} \text{ for short}) \text{ and } [Cu_xAg_{11-x}(dtp)_8]^+ (x = 0-3) (Cu_xAg_{11-x} \text{ for short}), respectively.$

Results and discussion

In the course of the synthesis (Scheme 1), Ag(OTf) was introduced into a THF solution of Ag₁₀ with a 1:1 molar ratio at ambient temperature. The reaction could be concluded within only one minute, resulting in the production of Ag₁₁ with an impressive yield of up to 97%. This success attests to the viability of adding extra Ag(1) ions to Ag NCs, a strategy corroborated by previous research involving the transformation of $[Ag_7(H)L_6]$ into $[Ag_8(H)L_6]^+$ (L = dithiophosphate or diselenophosphate)⁴⁰ and $[Ag_{20}{S_2P(O^iPr)_2}_{12}]$ into $[Ag_{21}{S_2P(O^iPr)_2}_{12}]^+$.⁴¹ Following the same reaction conditions as those for Ag₁₁, the introduction of one equivalent of [Cu(CH₃CN)₄](OTf) facilitated the formation of Cu_xAg_{11-x} . In contrast, the addition of Au(PPh₃)Cl into Ag10 led to decomposition. Notably, Au-doped NCs, such as $[Au_1Ag_{22}(S-Adm)_{12}]^{3+}$ (8e),⁴² $Au_xAg_{50-x}(Dppm)_6(SR)_{30}$ (20e),⁴³ and $[Ag_{46}Au_{24}(SR)_{32}]^{2+}$ (36e),⁴⁴ exhibit enhanced thermal stability. Au atoms on the surface typically form an S-Au-S linear

Scheme 1 The route of heterometal doping of Ag₁₀.

coordination with ligands.⁴⁵ When the gold atom is located in the kernel, Au is singly bonded to the S atom,⁴⁶ constraining the doping positions available. The challenge becomes pronounced in the Ag_{10} structure, where each Ag atom is bonded to two or three sulfur atoms. Au doping may alter the original coordination environment, resulting in an inability to maintain the structural framework.

The positive-ion ESI-MS of Ag_{11} (Fig. 1) reveals a molecular ion peak at m/z 2892.2122 Da, aligning with its cationic counterpart $[Ag_{11}(dtp)_8]^+$ (calc. 2892.0923 Da). The experimental isotopic distribution showcases a good agreement with the simulated pattern. A fragmentation peak at 2250.5063 Da was observed, which corresponds to $[Ag_9(dtp)_6]^+$ (calc. 2250.2461 Da). A peak at 2178.5623 Da can be assigned to $\left[\operatorname{Ag}_{8}(\operatorname{Cl})(\operatorname{dtp})_{6}\right]^{+}$ (calc. 2178.3088 Da). It is postulated that such species could be formed during sample preparation in CH₂Cl₂ solution, prior to the mass spectrometry measurement. The ESI-MS of Cu_xAg_{11-x} (Fig. 2a) reveals a distribution of $\left[\operatorname{Cu}_{x}\operatorname{Ag}_{11-x}(\operatorname{dtp})_{8}\right]^{+}$ (x = 0-3) with the Cu₁Ag₁₀ (x = 1) species assigned to the most intense band. Such a distribution was also observed in other Cu-doped Ag NCs, such as $[Cu_xAg_{20-x}(dtp)_{12}]$ (R = ^{*i*}Pr, x = 1-7; R = ^{*n*}Pr, x = 0-5),^{19,33} $[Ag_{29-x}Cu_{x}(BDT)_{12}]^{3-}$ (x = 0-13),³⁷ and $[Ag_{62-x}Cu_{x}S_{12}(SBu^{t})_{32}]^{+}$ $(x = 10 - 21).^{32}$

The ³¹P{¹H} NMR spectrum of Ag_{11} (Fig. S2[†]) shows two chemical shifts at 100.8 and 103.2 ppm, with an integration ratio of approximately 6:2, which corresponds to the two coordination environments of the dtp ligands. These resonances are shifted to 104.5 and 105.4 ppm in Ag_{10} . The ³¹P {¹H} NMR spectrum (Fig. S4[†]) of Cu_xAg_{11-x} shows multiple resonances spanning a broad range from 94 to 105 ppm, suggesting that the copper substitution within the Ag_{11} framework may occur at various positions, leading to a diversity of chemical environments.

In the molecular architecture of Ag_{10} , a conspicuous vacant site can be identified within the shell (Fig. 3a), presenting itself as an ideal location for hosting an additional metal anion and creating locally an MS_3 motif. The additional

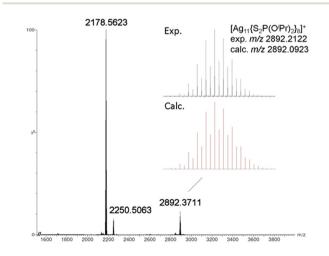


Fig. 1 The positive-ion ESI-MS of Ag11

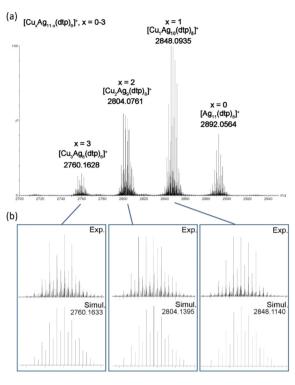


Fig. 2 (a) The positive-ion ESI-MS of Cu_xAg_{11-x} . (b) The experimental and simulated isotopic distribution patterns of Cu_xAg_{11-x} (x = 1-3).

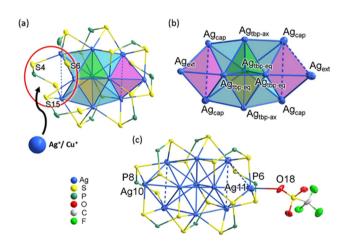


Fig. 3 (a) The structure of Ag_{10} . (b) The metal skeleton in Ag_{11} . (c) The total structure of Ag_{11} . The isopropoxy groups were omitted for clarity.

 Ag^+ ion is indeed attached to this precise location, in the crystal structure of Ag_{11} . The metal skeleton in Ag_{11} is composed of a trigonal bipyramid with four capping (Ag_{cap}) and two external (Ag_{ext}) atoms (Fig. 3b). With the addition of an extra silver atom, the metal framework attains higher symmetry, transitioning from C_1 in Ag_{10} to C_{2v} in Ag_{11} . The average Ag...Ag distances within the yellow (2.8339(7) Å), green (2.8386 (7) Å), and cyan (2.9778(7) Å) tetrahedra in Ag_{11} are similar to those in Ag_{10} (yellow: 2.8562(13) Å; green: 2.8553(13) Å, and cyan: 2.9495(13) Å). However, in Ag_{11} , the two magenta tetrahe-

dra exhibit substantially elongated Ag···Ag distances, averaged at 3.2191(8) Å and 3.3555(9) Å, which are notably longer than the 3.150(1) Å observed in **Ag**₁₀. A weak interaction between the OTf anion and the Ag11 atom (O18–Ag11: 2.61(1) Å) results in a significant elongation between Ag_{ext} and the triangular base it caps (Ag11 – Δ : 2.619 Å, Ag10 – Δ : 2.375 Å, see Fig. 3c). These bond lengths are summarized in Table S1.† The coordination modes of the eight dtp ligands in **Ag**₁₁ vary: P6 and P8 exhibit an η^3 (μ_1 , μ_2) coordination pattern, while the other six ligands demonstrate η^4 (μ_2 , μ_2) pattern. This differentiation is reflected in the ambient-temperature ³¹P{¹H} NMR spectrum, where two chemical shifts are observed. Consequently, these coordination differences contribute to the overall [Ag₁₁(dtp)₈]⁺ molecule possessing approximate C_2 symmetry.

Despite our inability to determine the crystal structures of Cu_xAg_{11-x} directly, ESI-MS data indicate that the species with the highest intensity is Cu_1Ag_{10} . This observation leads us to hypothesize that the doping copper atom predominantly occupies the NC's outermost fringe (*i.e.* Ag_{ext}, see Fig. 3a), and the geometry of Cu_1Ag_{10} closely resembles the core skeleton of $[Cu_xAg_{10-x}(Se)(dsep)_8]$ (x = 3-4, $dsep = Se_2P(O'Pr)_2$, see Fig. S6†).⁴⁷ Moreover, in scenarios where additional copper substitution would occur, it is likely that these copper atoms would prefer being situated at the other Ag_{ext} or Ag_{cap} positions (Fig. S7†). To substantiate this hypothesis, we employed DFT calculations, wherein 1, 2 and 3 Cu atoms were systematically substituted in the original Ag₁₁ structure (*vide infra*).

The absorption spectrum of Ag₁₁ exhibits two prominent bands at 390 and 528 nm (Fig. 4a). The high-energy band at approximately 366 nm is barely seen, embedded in the shoulder. This pattern shows similarity to that of the previously reported Ag10, which exhibits bands at 348, 392, and 520 nm.³⁹ In the case of the copper-doped Cu_xAg_{11-x} NC, a noticeable blue shift of the low-energy band, moving to 506 nm, while the high-energy band at 392 nm remains relatively unchanged. This behavior contrasts with that of the previously investigated 8-electron NCs,^{19,33} where Cu doping did not markedly alter the three primary absorption bands. Notably, the low-energy band around 480 nm remains consistent between the $[Cu_4Ag_{17}(dtp)_{12}]^+$ and $[Ag_{21}(dtp)_{12}]^+$ structures.¹⁹ This could be attributed to the larger core structure of the NC, which likely positions the surface-integrated copper atoms at a greater distance from the core, thereby minimizing their influence on the absorption characteristics.

The time-dependent-DFT (TD-DFT) simulated UV-vis spectra of Ag_{11} and of the most stable isomer of Cu_xAg_{11-x} (x = 1-3) (see the Computational details section) are shown in Fig. S8.[†] They exhibit the same shape as their experimental counterparts (*vide supra*). For all the compounds, the band of lowest energy can be identified as the HOMO (1S) \rightarrow LUMO (1P_x) transition, whereas the high-energy band is of mixed HOMO (1S) \rightarrow LUMO+1 (1P_z) and ligand \rightarrow LUMO (1P_x) character (see the 1S² 1P⁰ electronic structure DFT analysis), with little Cu participation in these two bands in the case of the Cu_xAg_{11-x} alloys. In the case of Ag₁₁, the two simulated peaks appear at 559 and 384 nm, values which compare relatively

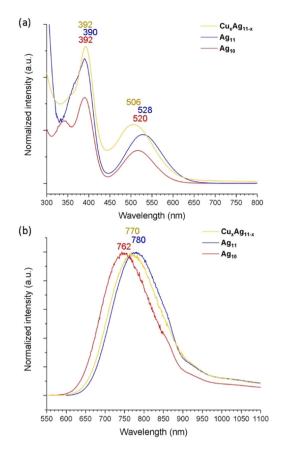


Fig. 4 (a) The ambient-temperature UV-vis absorption spectra and (b) emission spectra of Ag_{10} , Ag_{11} and $Cu_x Ag_{11-x}$.

well with their experimental counterparts, 528 and 390 nm, respectively. In the case of the Cu_xAg_{11-x} alloys, the highenergy band is barely perturbed, while the low-energy band is more blue-shifted as the number of copper atoms increases. This trend matches well with the experimental spectra (*vide supra*). The blue shift of the low-energy band is related to the increase of the $1S \rightarrow 1P_x$ gap (mainly a LUMO destabilization) with the increase in the number of Cu atoms in the NC. On the other hand, the $1S \rightarrow 1P_z$ transition participates to a lesser extent in the high-energy band and, having less Cu character, it is also less sensitive to the Cu content, explaining the inertness of this band to Cu alloying.

The room-temperature emission spectrum (Fig. 4b) exhibits a red shift in peak wavelength from 762 nm for Ag₁₀ to 780 nm for Ag₁₁, marking a transition to the near-infrared (NIR) region upon metal addition, while copper doping to form Cu_xAg_{11-x} results in the emission maxima shifting from 762 nm to 770 nm, a slight blue shift compared with Ag_{11} , which correlates well to the observed blue shift in the low-energy absorption band. This movement toward NIR emission in these NCs is advantageous for bioimaging applications due to deeper penetration capabilities tissue and minimized photodamage.^{15,16,48} On the other hand, Ag₁₁ exhibits an enhanced QY of 14.6%, a significant increase from the 6.0% observed in Ag_{10} . The enhancement can be caused by the

increased nuclearity in Ag11, which effectively stabilizes the surface ligand, thereby reinforcing its rigidity in solution. However, the QY decreases to 7.8% in Cu_xAg_{11-x} compared to Ag_{11} . This reduction may be due to intercluster reactions occurring within the solution, where copper atoms are exchanged between NCs. This phenomenon is akin to the observed intercluster reactions in Ag₇(H)(dtp)₆ and Cu₇(H) $(dtp)_6$, leading to a similar exchange of copper atoms.⁴⁹ Furthermore, the PLQY experiences a notable enhancement in the film state, reaching 37.2% for Ag11 and 16.1% for $Cu_x Ag_{11-x}$ (Table S3[†]). The lifetimes of Ag_{11} and $Cu_x Ag_{11-x}$ are in the nanosecond range, with 5.37 and 6.33 ns at RT (Fig. S11 and S13[†]), and 15.43 and 13.69 ns at 77K (Fig. S12 and S14[†]). The photophysical data are summarized in Table 1. Overall, the result suggests the resilience of superatomic electronic properties, which demonstrates the fine-tuning property following its metal addition or the Cu doping in the ultrasmall NCs.

Structural insights reveal that in Ag₁₁, the OTf anion exhibits a weak interaction with the Agext site, and the AgextS₃ motif is nearly coplanar, suggesting that Agext could readily interact with solvent molecules. The solvatochromic behavior of Ag_{10} , Ag_{11} , and $Cu_x Ag_{11-x}$ shows a more pronounced shift upon solvent polarity for Ag_{11} (76 meV) and Cu_xAg_{11-x} (71 meV) compared to Ag₁₀ (54 meV) (Fig. 5). The emission shifts from negative to positive solvatochromism with increasing solvent polarity. The polarity of solvents is quantified using the polarity parameter $(E_{\rm T})$, which is defined by the molar transition energy (measured in kcal mol⁻¹).⁵⁰ This heightened sensitivity in Ag_{11} , and $Cu_x Ag_{11-x}$ may be attributed to their structural configuration, which features a higher number of accessible Agext/Cuext sites. These sites are proximal to the superatomic core, where the emission is primarily due to the $1P_x \rightarrow 1S$ transition. Thus, solvent molecules have a significant impact on these NCs by affecting the proximity of Agext/Cuext sites to the superatomic core, leading to the substantial solvatochromic effect in observed Ag_{11} , and Cu_xAg_{11-x}.

DFT calculations (see the Computational details section) were first performed on the homometallic Ag_{11} cluster, whose optimized geometry (Table 2) is in good agreement with its SCXRD structure (Table S1†). Whereas the metal framework in this optimized geometry approaches C_{2v} symmetry, that of the whole NC (ligand considered) is very close to C_2 , with the pseudo- C_2 axis collinear with the $(Ag_{tbp-aq})_3$ diagonal that is perpendicular to the plane of Fig. 3b. Accordingly, there are two (slightly) different couples of Ag_{cap} atoms, namely Ag_{cap}^a and Ag_{cap}^b .

The electronic structure of Ag_{11} is strongly related to that of its Ag_{10} precursor. Its Kohn–Sham orbital diagram is shown in Fig. 6. The HOMO can be identified as the 1S superatomic orbital, with some ligand admixture. Its principal metal contribution is from the central trigonal bipyramid, with major participation of the equatorial triangle $(Ag_{tbp-eq})_3$. Both NAO atomic charges and Wiberg bond indices (Table 2) indicate clearly that the Ag_{cap} and Ag_{ext} atoms do not share any significant part of the two *superatomic* electrons and can be conNanoscale

Comp.	State	Absorbance λ_{abs} (nm), ε (M ⁻¹ cm ⁻¹)	Excitation λ_{ex} (nm)	Emission λ_{em} (nm)	Lifetime τ (ns)	Quantum yield, $\Phi_{\rm a}$ (%)	$k_{\mathrm{r}}^{\ a}$ (s^{-1})	k_{nr}^{b} (s^{-1})
Ag ₁₀	2-MeTHF, 298 K	348 (15 500), 392 (21 300), 520 (8100)	338, 395, 517	762	1.95	6.0	3.08×10^{7}	4.82×10^8
	2-MeTHF, 77 K		335, 388, 515	687	15.37			
Ag ₁₁	2-MeTHF, 298 K	390 (20 600), 528 (9000)	391, 530	780	5.37	14.6	2.72×10^7	$1.84 imes 10^8$
	2-MeTHF, 77 K		396, 509	699	15.43			
Cu_xAg_{11-x}	2-MeTHF, 298 K	392(20 100), 506(8800)	391, 527	770	6.33	7.8	$1.23 imes 10^7$	$1.46 imes 10^8$
	2-MeTHF, 77 K		393, 507	687	13.69			

^{*a*} Radiative rate constant, $k_{\rm r} = \Phi_{\rm a} \times (1/\tau)$. ^{*b*} Nonradiative rate constant, $k_{\rm nr} = (1/\tau) - k_{\rm r}$.

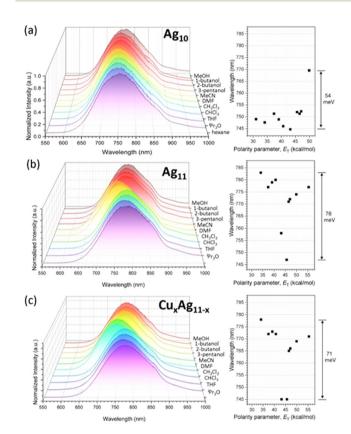


Fig. 5 (a) Solvent-dependent emission spectra of Ag_{10} , (b) Ag_{11} , and (c) $Cu_x Ag_{11-x}$ at RT.

Table 2 Relevant averaged data computed for Ag_{11} . WBI = Wiberg bond index

	Distance (Å)	WBI		NAO charges	
Ag _{tbp-eq} -Ag _{tbp-eq}	2.868	0.185	Ag _{tbp-eq}	0.27	
Ag _{tbp-ax} -Ag _{tbp-eq}	2.923	0.093	Ag _{tbp-ax}	0.61	
Agaab – Agtherax	2.951	0.050	Agacap	0.69	
Agaab – Agtbp-eq	3.006	0.063	Ag ^a _{cap} Ag ^b _{cap}	0.68	
$\begin{array}{l} Ag^{a,b}_{cap} - Ag_{tbp-ax} \\ Ag^{a,b}_{cap} - Ag_{tbp-eq} \\ Ag_{ext} - Ag^{a,b}_{cap} \end{array}$	3.175	0.028	Agest	0.75	
Ag _{ext} -Ag _{tbp-eq}	2.867	0.045	2		

sidered as being in their $+_1$ oxidation state. The two LUMOs can be identified as the superatomic $1P_x$ and $1P_z$ orbitals, with some ligand admixture.

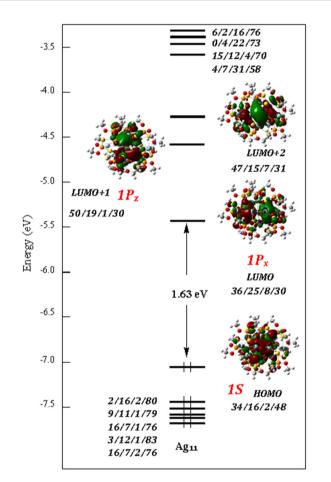


Fig. 6 The Kohn–Sham orbital diagram of Ag₁₁. The orbital contribution (in %) is given in the order: $(Ag_{tbp-})_5/(Ag_{a,b}^{a,b})_4/(Ag_{ext})_2/(dtp)_8$.

In the second step, we investigated the preferred sites of substitution of Ag by Cu in Ag_{11} . The main results are summarized in Fig. 7. Based on the computed relative free energies at 298 K, the doping by one Cu atom is almost equally preferred on the M_{ext} and M_{cap}^{a} positions, which corresponds to those with the more positive charges in Ag_{11} (Table 2). Those with the less positive charges in Ag_{11} , namely the M_{tbp-eq} positions, are the less favored positions. This site preference is maintained when doping with two and three Cu atoms, as shown in the middle and bottom of Fig. 7, which shows the positional isomers of lowest energy for the Cu₂Ag₉ and Cu₃Ag₈ compo-

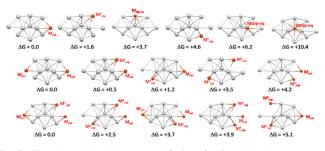


Fig. 7 The six positional isomers of $CuAg_{10}$ (top) and the lowest energy positional isomers of Cu_2Ag_9 (middle) and Cu_3Ag_8 (bottom). The relative free energies at 298 K are given in kcal mol⁻¹.

sitions, respectively. A mixture of three isomers is expected in the case of Cu_2Ag_9 , whereas only one is expected to be obtained (or at least as a largely dominant species) in the case of the trisubstituted Cu_3Ag_8 .

Conclusions

In summary, this investigation delves into the intricacies of metal addition and doping effects, utilizing a 2-electron Ag₁₀ NC as a foundational template. The surficial traits of Ag_{10} uncover an open vacancy that facilitates the attachment of additional Ag⁺/Cu⁺ ions, resulting in the formation of Ag₁₁ and Cu_xAg_{11-x} . Significantly, the solvent-dependent emission spectra shed light on the intricate interplay between the structural characteristics and the potent impact of solvent polarity, offering a window into the solvatochromic behaviors of these materials. Furthermore, the emission of these NCs extends into the near-infrared (NIR) region, which opens up potential applications in areas such as bioimaging and photothermal therapy. These insights underscore the robust nature of superatomic electronic properties, demonstrating their capacity for fine-tuning properties through controlled metal addition and doping, further expanding the functional versatility of metal NCs.

Experimental section

Materials and instrumentation

All chemicals used as received were purchased from commercial sources. Solvents were purified following standard protocols.⁵¹ All reactions were performed in oven-dried Schlenk glassware using standard inert atmosphere techniques. All reactions were carried out under an N₂ atmosphere by using standard Schlenk techniques. Ag₁₀(dtp)₈³⁹ and Au(PPh₃)Cl⁵² were prepared by procedures reported earlier in the literature. The ¹H and ³¹P{¹H} NMR spectra were recorded on a Bruker Avance II 400 MHz NMR spectrometer, operating at 400.13 MHz for ¹H, and 161.98 MHz for ³¹P. The chemical shifts (δ) and coupling constants (J) are reported in ppm and Hz, respectively. The ESI mass spectrum was recorded on a QSTAR® XL (AB SCIEX, Warrington, Cheshire, U.K.). X-ray photoelectron spectroscopy (XPS) spectra were recorded by using a PHI 5000 VersaProbe-Scanning ESCA Microprobe. UV-visible absorption spectra were recorded on an Agilent Cary 60 spectrophotometer using quartz cells with path length of 1 cm. The emission and QY spectra were recorded on a Horiba FluoroMax⁺ fluorescence spectrometer equipped with a 150 W ozone free xenon lamp as the excited source, a detector with R13456P PMT, and NIR InGaAs. The lifetime was recorded on an Edinburgh FLS920 fluorescence spectrometer using the TCPSC technique, using an H₂ pulse lamp as the excited source and R928P as the detector. The QY was determined by a comparative method,^{53,54} and the detailed results of the target complexes are shown in Table S3.[†] Absolute values are calculated by using [Ru(bpy)₃]²⁺ as the standard sample.^{55,56}

Synthesis

[Ag₁₁(dtp)₈(SO₃CF₃)], Ag₁₁. [Ag₁₀(dtp)₈] (0.041 g, 0.0147 mmol) was dissolved in THF (30 mL). Ag(OTf) (0.0038 g, 0.0147 mmol) was added to this mixture and stirred continuously for 1 minute. It was dried under vacuum and the powder was washed with hexane. The precipitate was dried to yield purple powder as [Ag₁₁(dtp)₈(SO₃CF₃)]. Yield: 0.0436 g (97.5%, based on Ag). ³¹P{¹H} NMR (161.98 MHz, CDCl₃, δ , ppm): 103.2, 100.8. ¹H NMR (400.13 MHz, CDCl₃, δ , ppm): 1.40 (m, ³J_{HH} = 4 Hz, 96H, CH₃), 4.82 (br, 16H, CH). UV-vis spectra [λ in nm (ε in M⁻¹ cm⁻¹)]: 391 (20 600), 527 (9900).

[Cu_xAg_{11-x}(dtp)₈](SO₃CF₃), Cu_xAg_{11-x}. [Ag₁₀(dtp)₈] (0.041 g, 0.0147 mmol) was dissolved in THF (30 mL). [Cu(CH₃CN)₄](OTf) (0.0055 g, 0.0147 mmol) was added to this mixture and stirred continuously for 1 minute. It was dried under vacuum and the powder was washed with hexane. The precipitate was dried to yield purple powder as Cu_xAg_{11-x}. ³¹P{¹H} NMR (161.98 MHz, CDCl₃, δ , ppm): 94–105. ¹H NMR (400.13 MHz, CDCl₃, δ , ppm): 1.40 (m, ³J_{HH} = 7 Hz, 96H, CH₃), 4.82 (br, 16H, CH). UV-vis spectra [λ in nm (ε in M⁻¹ cm⁻¹)]: 392 (20 100), 506 (8800).

X-ray crystallography

Single crystals suitable for X-ray diffraction analysis of Ag_{11} were obtained by the slow evaporation of CH_2Cl_2 solution at ambient temperature. Single crystals were mounted on the tip of a glass fiber with Paratone oil. Data were collected on a Bruker APEX II CCD diffractometer using graphite monochromated Mo K α radiation ($\lambda = 0.71073$ A) at 100 K. Absorption corrections for the area detector were performed with SADABS,⁵⁷ and the integration of the raw data frame was performed with SAINT.⁵⁸ The structures were solved by direct methods and refined by least-squares against F^2 using the SHELXL-2018/3 package,⁵⁹ incorporated in SHELXTL/PC V6.14.⁶⁰ All non-hydrogen atoms were refined anisotropically. CCDC no. 2287343† (Ag_{11}) contains the supplementary crystal-lographic data in this article.

Computational details

For the sake of reducing computational effort, all calculations were made on simplified models in which the $S_2P(O^iPr)_2$ ligands were replaced by the smaller $S_2P(OMe)_2$ ones.

Geometry optimizations were performed by density functional theory (DFT) calculations with the Gaussian 16 package,⁶¹ using the BP86 functional⁶² together with Grimme's empirical DFT-D3(BJ) corrections⁶³ and the Def2-TZVP basis set from EMSL Basis Set Exchange Library.⁶⁴ All the optimized geometries were characterized as true minima by vibrational analysis. The NAO charges and Wiberg bond indices were computed with the NBO 6.0 program.⁶⁵ The UV-visible transitions were calculated by means of time-dependent DFT (TD-DFT) calculations, using the B3LYP⁶⁶ functional and the Def2-TZVP basis set. The UV-visible spectra were simulated from the computed TD-DFT transitions and their oscillator strengths by using the SWizard program,⁶⁷ with each transition being associated with a Gaussian function of half-height width equal to 2000 cm^{-1} . The compositions of the molecular orbitals were calculated using the AOMix program.68

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was supported by the National Science and Technology Council of Taiwan (112-2123-M-259-001) and the GENCI Computing Resource (grant A0090807367). The authors gratefully acknowledge the Instrumentation Center of National Taiwan Normal University (NSTC 111-2731-M-003-001) for its support.

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