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# A water-responsive calix[4]resorcinarene system: self-assembly and fluorescence modulation†

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This study explores how water content modulates the self-assembly and fluorescence behavior of two novel calix[4]resorcinarene macrocycles. These macrocycles transition from large, flattened structures in pure THF to large giant vesicles (500–5000 nm) coexisting with small micelles (3.4–3.5 nm) as the water percentage in THF/water mixtures increases up to 53%. At higher water percentages, the assemblies become smaller, forming unimodal micelles with diameters of approximately 140–160 nm. Fluorescence quenching is observed upon water addition, attributed to nonradiative deactivation. These findings highlight water as a key regulator of the assembly and fluorescence of these calix[4]resorcinarene macrocycles, paving the way for further development of water-responsive calixarene systems.

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## 1. Introduction

Surfactants are versatile molecules with hydrophilic and hydrophobic parts, making them indispensable in various industries.<sup>1–7</sup> They serve as essential auxiliary substances in the production of pharmaceuticals and cosmetics, stabilizing mixed systems, enhancing drug solubility, and improving drug permeability.<sup>8–12</sup> Their unique ability to self-assemble results in micelle formation, with the critical micelle concentration (CMC) being a crucial parameter.<sup>13</sup> Below the CMC, surfactants form a layer at interfaces, reducing surface tension across gas, liquid, and solid phases. Conversely, they arrange them into specific spherical and cylindrical micellar structures above the CMC, dispersing as a colloidal solution within the aqueous environment.<sup>9</sup> Resorcinarenes and calixarenes, cyclic compounds composed of aromatic rings derived from resorcinol or phenol linked by carbon bridges, are readily available supramolecular hosts. These compounds, in their original or modified forms, find applications as receptors, cavitands, capsules, and in sensing, storage, reaction nanovessels, and biological fields.<sup>14</sup> Resorcin[4]arenes, comprising four

resorcinol rings connected by four –C(R)–H groups, exhibit inherent amphiphilicity when the substituent R on the carbon bridge is a long aliphatic chain. Resorcinarenes and calixarenes possess conformational flexibility, influencing their ability to form host-guest complexes. Restricting this conformational freedom can enhance host selectivity and affinity. Additionally, the self-assembly behavior of amphiphilic resorcinarenes, including water solubility and aggregate dimensions, is modulated by conformational dynamics. Self-assembly can mitigate conformational mobility by orienting hydrophilic moieties towards the aqueous environment.<sup>15–19</sup> Moreover, controlled manipulation of resorcinarene conformation offers potential for the development of molecular motors and switches.<sup>20</sup>

Resorcinarenes exhibit diverse self-assembly behaviors influenced by their structural features and functionalization.<sup>21</sup> From hexameric and octameric structures to vesicles, micelles, and nanoparticles, these compounds display a wide range of supramolecular architectures.<sup>22–26</sup> Hydrogen bonding and metal coordination are key driving forces in these assemblies.<sup>21</sup> The ability to fine-tune self-assembly through functionalization underscores the potential of resorcinarenes for creating materials with tailored properties.

The aggregation of calixarenes and resorcinarenes significantly impacts their fluorescence. This phenomenon can result in either reduced emission (ACQ) or enhanced emission (AIE) compared to dilute solutions. While AIE is often sought after for applications,<sup>27</sup> ACQ can also be advantageous. For instance, ACQ-based sensors have been developed for pH, ions, and organic pollutants, demonstrating the versatility of these compounds in fluorescence sensing.<sup>28–33</sup>

This study presents a comparative analysis of the influence of water content on the self-assembly and photophysical

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properties of two calix[4]resorcinarene macrocycles, which differ in the chain length by two carbon atoms, by AFM, DLS, and multivariate analysis applied to absorption and fluorescence spectroscopy. Water content significantly influences the self-assembly and fluorescence behavior of calix[4]resorcinarene macrocycles. As water is added to THF/water mixtures, the macrocycles transition from large aggregates to smaller micelles. Fluorescence quenching occurs due to nonradiative deactivation. These findings demonstrate the potential of water-responsive calixarene systems for various applications, especially in sensor design.

## 2. Results and discussion

### 2.1. Synthesis of compounds C1 and C2

The compound C1 was available from previous work.<sup>34</sup> The compound C2 was synthesized in a similar way as C1, *i.e.* through a reaction between compound (a) and resorcinol (see Scheme 1). See ESI† for characterization by ESI-HR-MS, EA, NMR, FTIR, and physical characterization.

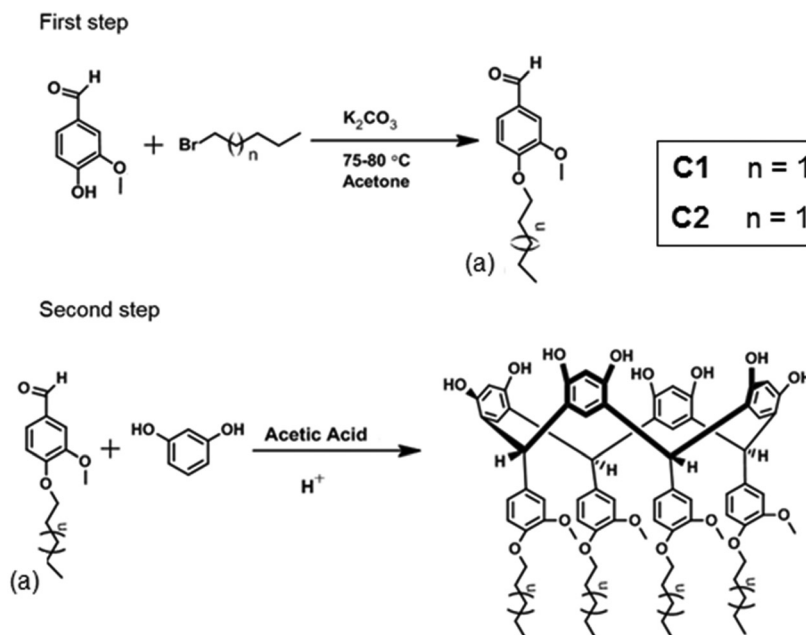
### 2.2. Absorption spectroscopy in THF and THF/H<sub>2</sub>O mixtures and aggregation

The calix[4]resorcinarene macrocycle compound, C2, experiences some minor absorbance changes within 1–24 h after dissolution in THF or THF/H<sub>2</sub>O mixtures, like the absorbances changes experienced by the macrocycle C1 in the same time period,<sup>34</sup> related to aggregation phenomena.<sup>34,35</sup> Consequently, all experiments described in this study were conducted only after the absorbance spectra of the solutions had become stable. To achieve this, C2 solutions were prepared and allowed to stabilize for 24 hours before experimentation.

The macrocycle C2, a calix[4]resorcinarene, exhibits a strong ultraviolet absorption peak when dissolved in THF at 285 nm,

with a less pronounced shoulder at 291 nm. These observations were consistent across a concentration range of  $1\text{--}5 \times 10^{-5}$  M. At higher concentrations (above  $10^{-4}$  M), a broad, weak absorption signal appeared between 330 and 600 nm (Fig. 1 and Fig. S5, ESI†). The overall absorption pattern of C2 closely resembles that of C1<sup>34</sup> and that of a related compound, resorcin[4]arene with  $R = \text{CH}_3\text{--}(\text{CH}_2)_{14}$ , with the same UV-vis spectrum.<sup>36</sup>

Given their structure, resorcin[4]arenes like C1 or C2, lacking chromophoric groups, are expected to be colorless.<sup>36,37</sup> However, it was observed that increasing the concentration of these compounds above a certain point led to a weak absorption in the visible spectrum. This phenomenon, also seen in C1,<sup>34</sup> is attributed to molecular aggregation. To further investigate this, the absorption spectra of C2 were measured in THF solutions with varying water content. The spectra were recorded for solvent mixtures with H<sub>2</sub>O proportions of 0, 20, 33, 47, 53, 60, 67, 75, 82, 88, and 95%. As the water content increased from 0 to 20%, a similar absorption tail in the visible region grew, comparable to the behavior of C1.<sup>34</sup> Additionally, increasing the concentration of C2 in either pure THF or THF with 20% water resulted in a stronger absorption at 355 nm. At water concentrations exceeding 33%, similar to observations with C1,<sup>34</sup> an increase in C2 concentration led to amplified absorbance between 400 and 700 nm, indicating enhanced aggregation. Comparative analysis of C1 and C2 absorption spectra under equivalent experimental conditions (concentrations and solvent mixtures, as depicted in Fig. 1 and Fig. S5, ESI†) revealed subtle disparities suggesting distinct aggregation behaviors. Consequently, we employed principal component analysis (PCA) to comprehensively examine the entire dataset of C1 and C2 absorption spectra. For the PCA, we incorporated additional data for C1 at 75%, 82%, and 88% water content, which were not included in our previous report.<sup>34</sup>



Scheme 1 Synthesis of a functionalized calix[4]resorcinarene macrocycle through a two-step reaction.

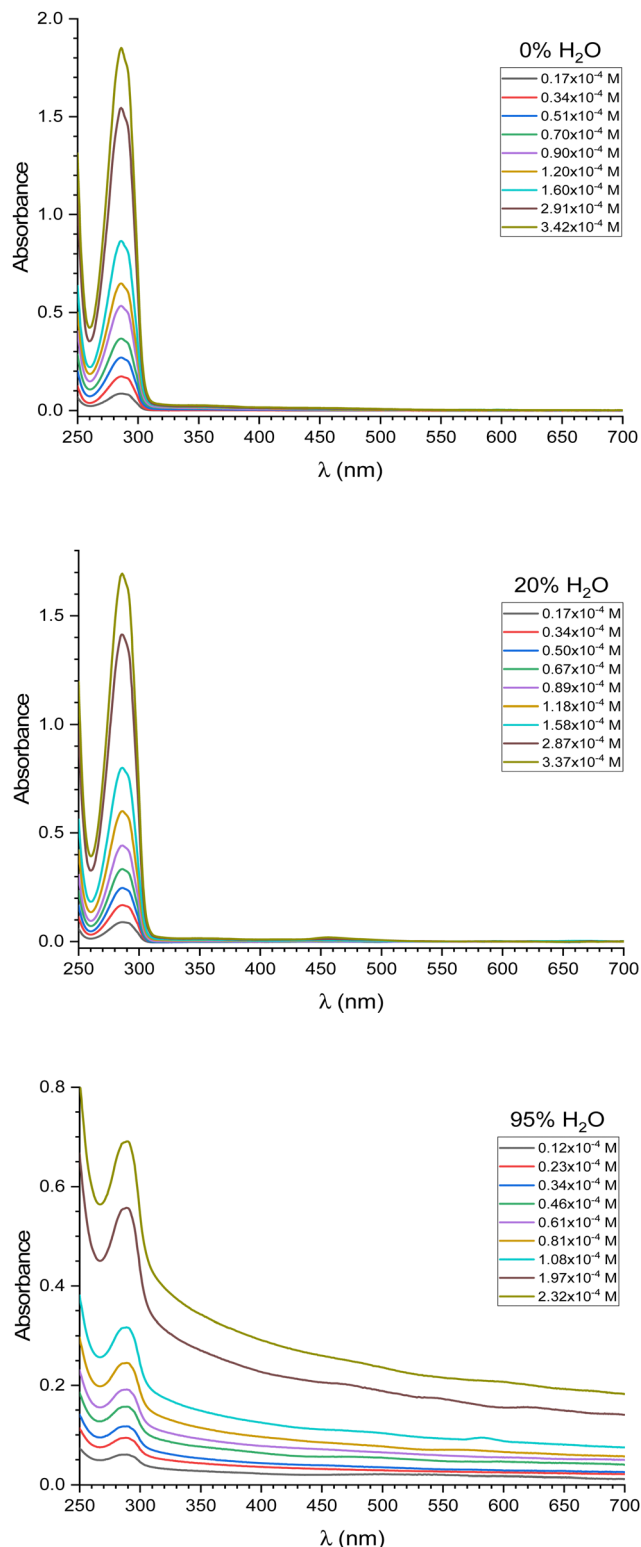


Fig. 1 Absorption spectra of **C2** at several concentrations in THF/H<sub>2</sub>O mixtures.

### 2.3. Principal component analysis

The entire absorption matrix, denoted by  $A$ , encompassing a complete set of around 200 absorption spectra between 250 and

700 nm, was subjected to PCA using two different strategies. This matrix corresponds to the full range of experimental conditions, including solvent compositions varying from 0% to 95% water in THF, and analytical concentrations ( $C_0$ ) ranging from 10 to 300  $\mu\text{M}$  for compounds **C1** and **C2**. The first PCA was performed on the original data, represented by the matrix  $A$ . The second PCA utilized the ratio matrix,  $A/C_0$ , effectively performing PCA on the apparent extinction coefficient matrix,  $E_{\text{app}}$ .

In matrix representation, the model with a given number of components can be expressed by the following equation:  $X = TP^T + E$ , where  $X$  represents the centered data matrix,  $T$  is the scores matrix,  $P^T$  is the transpose of the loadings matrix, and  $E$  is the error matrix. In the present work, we used two alternative definitions of  $X$ . If the original data is denoted by  $A$ , then  $X = A - \bar{A}$ , where  $\bar{A}$  is the mean absorption spectrum matrix. Similarly, if the original data represents apparent extinction coefficients and is denoted by  $E_{\text{app}}$ , then  $X = E_{\text{app}} - \bar{E}_{\text{app}}$ , where  $\bar{E}_{\text{app}}$  is the mean apparent extinction coefficient spectrum.

Within the PCA framework, the scores ( $T$ ) and loadings ( $P$ ) matrices form the crux of the analysis, capturing the most informative, structured variation within the data. These matrices can be likened to a compressed map, highlighting the key features and trends present in the original high-dimensional data space. The residual matrix ( $E$ ) represents the unexplained variation, often referred to as noise, which cannot be effectively modeled by the chosen number of principal components. The core structure of the original data can be reconstructed by multiplying the scores and transposed loadings matrices ( $TP^T$ ). Ideally, if the proper number of principal components is chosen, this reconstruction minimizes the residual matrix, leaving only a minimal amount of random error that is not amenable to meaningful interpretation within the context of the analysis.

Application of PCA revealed that most of the variabilities linked to either the  $A$  or  $E_{\text{app}}$  matrices could be adequately represented by keeping only two principal components. Regarding the  $A$  matrix, PC1 accounted for 85% of the variance, while PC2 captured the rest. Conversely, for the ratio matrix  $E_{\text{app}}$ , the distribution of explained variance shifted slightly, with PC1 explaining 82% and PC2 explaining 17% of the variance.

**2.3.1. Decomposition of the  $A$  and  $E_{\text{app}}$  matrices into principal components.** Fig. 2(a) and (b) depict the PC2 vs. PC1 scores plot for compounds **C1** and **C2**, representing a compressed map of the data across all experimental conditions. ESI,† Fig. S6–S27 visualize the contributions to the core data structure reconstructed by each principal component (PC<sub>i</sub>). These figures present the spectral profiles obtained as  $(TP^T + \bar{A})$  vs.  $\lambda$  as a function of both analytical concentration ( $C_0$ ) and solvent composition. While a direct physical interpretation of the principal components (PCs) can be challenging, PCA enables us to identify underlying patterns and trends. PC1 primarily accounts for the overall increase in spectral intensity with concentration, while PC2 mainly captures the changes in the spectral features induced by variations in environmental conditions. Positive PC2 scores correspond to a

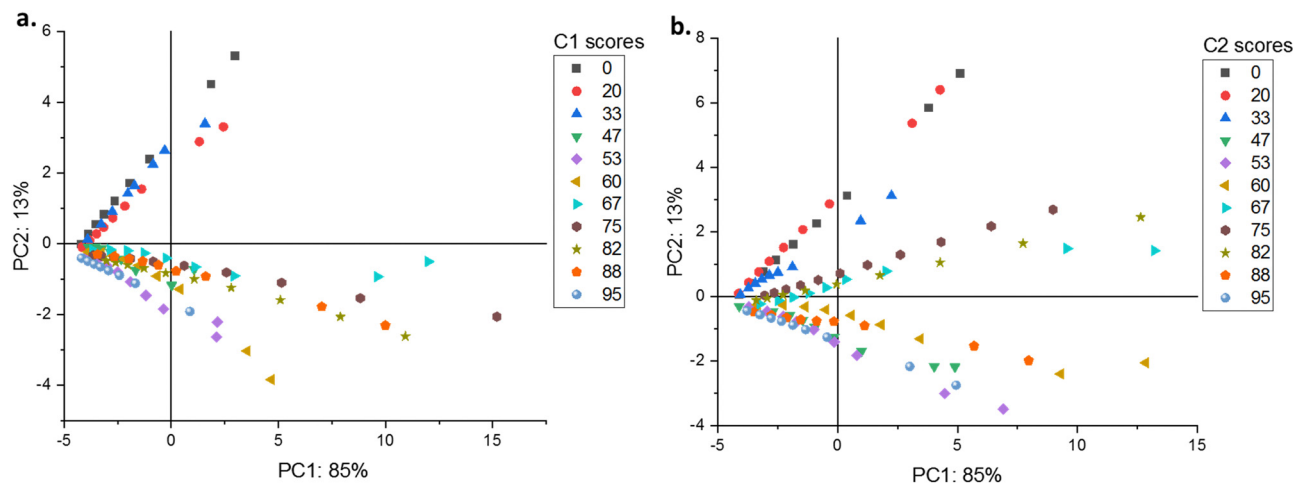


Fig. 2 (a) PC2 vs. PC1 scores for **C1**. (b) PC2 vs. PC1 scores for **C2**.

relative increase in the absorptive contribution, whereas negative scores indicate a greater dispersive contribution to the overall absorption spectrum.

For compound **C1**, the PC2 vs. PC1 scores plots reveal distinct trends based on solvent water content. For water contents between 0% and 33%, the slopes are positive, with the steepest slope observed at 0% and the shallowest at 20%. Interestingly, the slopes become consistently negative for water contents exceeding 47%.

In contrast to compound **C1**, the PC2 vs. PC1 scores plots of compound **C2** exhibit a distinct pattern. The slopes are positive and nearly identical at 0% and 20% water content. The slope becomes slightly less positive at 33% water content. Conversely, water contents of 47%, 60%, 88%, and 95% show negative slopes. However, while the initial slopes for **C2** are positive for water contents ranging from 67% to 82%, they differ from the consistently negative slopes seen for **C1**.

A more detailed analysis of PC behavior is achieved by plotting each score (PC<sub>i</sub>) for a given solvent composition as a function of analytical concentration ( $C_0$ ). Fig. 3(a)–(c) depict these plots for compound **C1**. In Fig. 3(a), PC1 and PC2 scores exhibit a linear increase with increasing  $C_0$  for solvent compositions of 0% and 20% water. At 33% water content, the initially linear trend deviates downward, indicating a decrease in the slope and a downward curvature at higher concentrations exceeding 150  $\mu\text{M}$ . Fig. 3(b) and (c) reveal distinct trends for higher solvent content (47–60% water in THF). While PC1 scores maintain a linear increase with  $C_0$ , PC2 scores exhibit a negative slope. Similarly, for solvent compositions between 67% and 95% water, PC1 scores continue their linear rise with  $C_0$ , whereas PC2 scores display a negative slope.

Compound **C2** exhibits analogous behavior in its first principal component, as evidenced by the PC1 vs.  $C_0$  plots across different water contents (Fig. 3(d)–(f)). In contrast to the **C1** pattern, PC2 scores for **C2** exhibit a more complex profile. The slopes fluctuate, transitioning from positive within the 0–30% water range to negative between 47% and 60% water. Subsequently, they become positive again from 67% to 82%, only to revert to negative between 88% and 95%. Table 1 summarizes

the estimated initial slopes of the plots of PC1 and PC2 scores against the analytical concentrations of **C1** and **C2**.

The core data structures reconstructed by principal component PC<sub>i</sub> for **C1** and **C2** (Fig. S6–S27, ESI†) exhibit shared spectral characteristics and are thus discussed collectively. Across all solvent mixtures (0–95% water in THF), the  $TP^T + \bar{A}$  spectral contributions associated with PC1 comprise an absorption band centered at  $\sim 286$  nm, superimposed on a scattering continuum extending across the entire wavelength range. This scattering component's wavelength dependence aligns with Mie scattering theory.<sup>38</sup>

$$I = C\lambda^{-k} \quad (1)$$

where intensity is inversely proportional to wavelength raised to the power  $k$ . The constant  $C$  in eqn (1) is valid under conditions of uniform particle and medium dispersion. While Rayleigh scattering ( $k = 4$ ) applies to smaller particles, Mie scattering ( $k < 4$ ) is relevant for larger ones. However, in media containing a distribution of particle sizes and irregular shapes, both  $C$  and  $k$  become size-dependent, complicating scattering models. For PC1, both the 286 nm absorption intensity and scattering contribution increase monotonically with  $C_0$ . The contribution associated with PC1 consistently increases more than that of PC2 with increasing analytical concentration, primarily because PC1 accounts for most spectral alterations induced by concentration changes.

Inspection of Fig. S6–S27 (ESI†) reveals that the scattering continuum associated with PC2 is flatter than that of PC1, indicating a lower  $k$  value for PC2 within the solvent composition range analyzed. Comparison of PC1 and PC2 spectral trends suggests that the underlying data structures of **C1** and **C2** can be categorized into four solvent ranges: (a) 0–33%, (b) 47–60%, (c) 67–82%, and (d) 88–95%. Noteworthy, this behavior is in line with the results shown in Fig. 3 and the data presented in Table 1.

**2.3.2. Decomposition of the  $E_{\text{app}}$  matrix into principal components.** Fig. 4 depicts the PC2 vs. PC1 scores plot for compound **C1** and **C2**, representing a compressed map of the data across all experimental conditions. ESI,† Fig. S28–S49

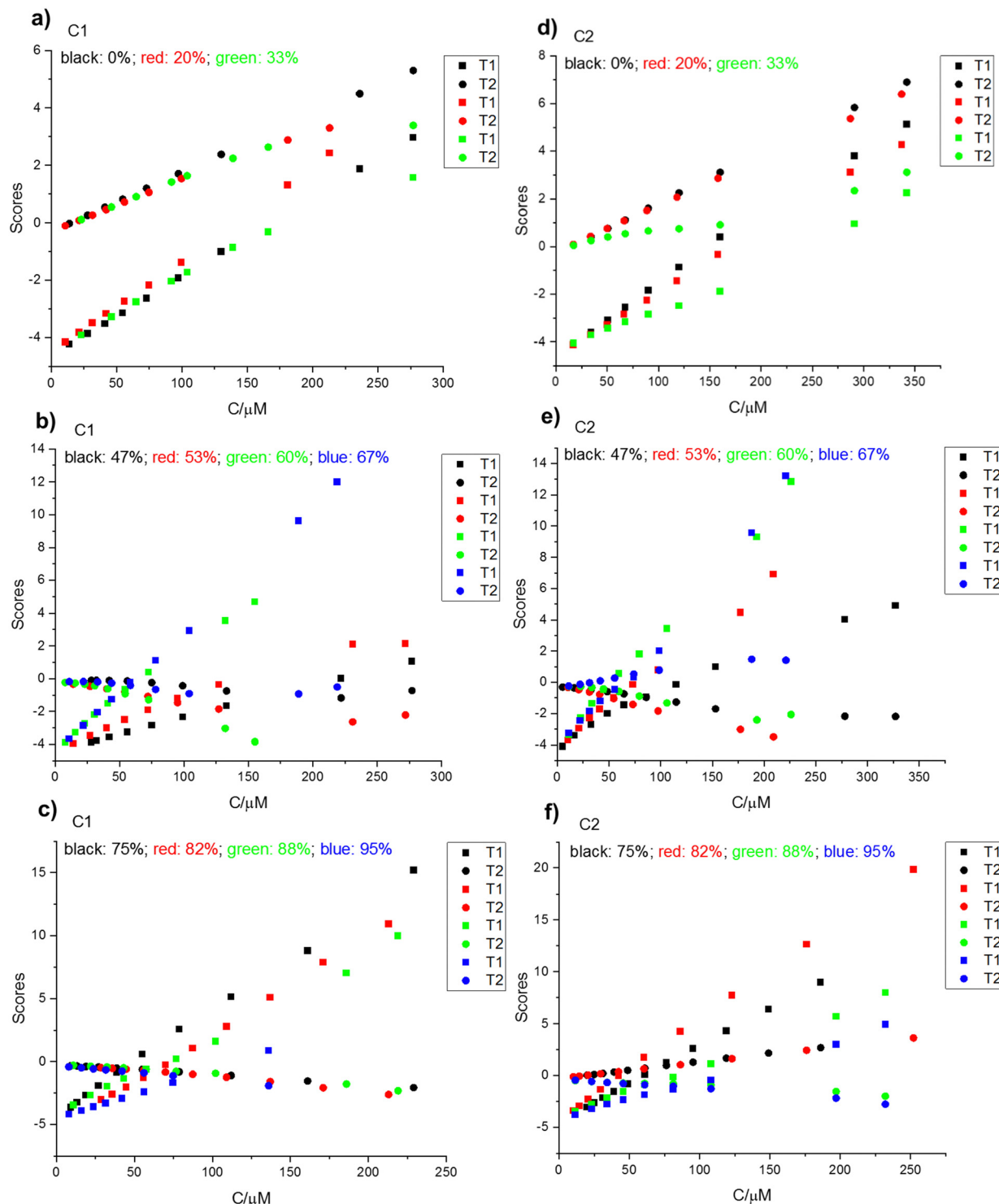


Fig. 3 PC1 vs.  $C_0$  for different H<sub>2</sub>O contents in THF; (a) and (d) for 0, 20, 33%; (b) and (e) for 47, 53, 60, 67%; (c) and (f) for 75%, 82%, 88%, 95%.

visualize the core data structure reconstructed by each principal component (PCi). These figures present the spectral form ( $TP^T + E_{app}$  vs.  $\lambda$ ) as a function of both analytical concentration ( $C_0$ ) and solvent composition.

Fig. 4 differs fundamentally from Fig. 2 due to the transformation of the  $A$  matrix into  $E_{app}$ , which eliminates PC variance

arising from the overall increase of spectral signals associated with higher analytical concentrations.

For compound C1, the PC2 vs. PC1 scores plot (Fig. 4(a)) is divided into four quadrants. The upper left quadrant contains data points representing the 0–33% water in THF range (negative PC1, positive PC2). The lower left quadrant includes data



**Table 1** Estimated<sup>a</sup> initial slopes for PC1 and PC2 scores vs.  $C_0$  plots. The initial slopes were calculated by linear regression analysis taking the appropriate number of points to preserve the linear trend (at least six points were included in this calculation)

% H <sub>2</sub> O	dT1/dC <sub>0</sub> , C1	dT2/dC <sub>0</sub> , C1	dT1/dC <sub>0</sub> , C2	dT2/dC <sub>0</sub> , C2
0	0.0275	0.0202	0.0316	0.0211
20	0.0324	0.0170	0.0265	0.0196
33	0.0254	0.0178	0.0177	0.0097
47	0.0198	−0.0059	0.0392	−0.0081
53	0.0321	−0.0139	0.0502	−0.0162
60	0.0658	−0.0168	0.0682	−0.0125
67	0.0751	−0.0084	0.0581	0.0119
75	0.0840	−0.0081	0.0725	0.0162
82	0.0770	−0.0116	0.0958	0.0155
88	0.0622	−0.0071	0.0483	−0.0053
95	0.0394	−0.0117	0.0356	−0.0091

<sup>a</sup> Relative errors where, in all cases, lower than 6.5%.

for 47%, 53%, and 95% water in THF (negative PC1 and PC2). The upper right quadrant, sparsely populated, contains a small subset of data with positive PC1 and PC2 scores. The lower right quadrant encompasses data for 60%, 67%, 75%, 82%, and 88% water in THF (positive PC1, negative PC2). For compound C2, the PC2 vs. PC1 scores plot (Fig. 4(b)) shows a different distribution. The upper left quadrant contains data points representing the 0–33% water in THF range. The upper right quadrant includes data for 67%, 75%, 82%, and a portion of the 60% water in THF data. The remaining solvent data are distributed across the lower left and right quadrants.

To provide a more detailed examination of PC behavior, the PC2 vs. PC1 scores plots (previously shown collectively in Fig. 4) are presented individually for each solvent composition as a function of analytical concentration ( $C_0$ ) in Fig. S50–S71 (ESI†). In each of these figures, the analytical concentration is labeled directly above the corresponding data point on the PC2 vs. PC1 scores plot.

For C1 and C2, Fig. S50–S71 (ESI†) reveal complex patterns in the relationship between PC1 and PC2 scores as  $C_0$  steadily increases. While some plots show initial steady increases (e.g., Fig. S50, ESI†) or decreases (e.g., Fig. S62, ESI†) in PC2 scores

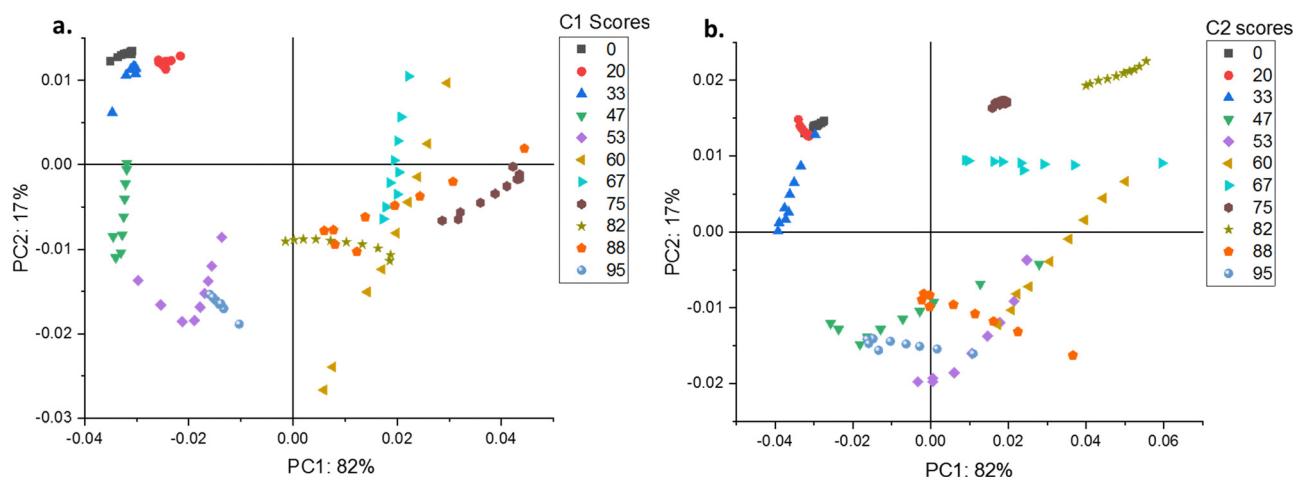
with rising PC1 values, all eventually exhibit more intricate behaviors within specific  $C_0$  ranges. These include abrupt drops, loop-like structures, and even instances of double reversals where the trajectory reverses direction twice. Table 2 provides estimated concentration ranges for these potential turning points.

The core data structures reconstructed by principal component PCi for C1 and C2 (Fig. S28–S49, ESI†) exhibit shared spectral characteristics and are thus discussed collectively. Comparison of the spectral contributions associated with PC1 and PC2 reveals that both exhibit absorption bands centered around 287 nm, overlaid by broad scattering continuums. However, the portion of the signals accounted for PC1 exhibits a less absorptive contribution than that of PC2. In addition, a detailed inspection of the spectral region dominated by scattering effects also shows that the changes in PC1 scores account for those signals associated with higher  $k$  values of eqn (1), whereas the changes in PC2 scores mainly describe scattering signals that correspond to lower values of the exponent.

The PCA analyses of matrices  $A$  and  $E_{app}$  provide complementary insights. While the PCA of matrix  $A$  emphasized the differences between C1 and C2 at low analyte concentrations, the PCA of  $E_{app}$  matrices enabled a comprehensive analysis across all solvent concentrations and revealed the occurrence of turning points. This indicates that phase transitions occur even in pure THF, where scattering is observed at concentrations as low as 70 mM. Although our PCA analysis cannot provide precise CMC values, it suggests multiple phase transitions leading to the final micellar structures (see DLS results below).

## 2.4. Dynamic light scattering

Dynamic light scattering (DLS) confirmed the presence of C2 aggregates, like previous observations for C1,<sup>34</sup> with substantial size variability and large particles detected in pure THF. However, as with C1, data quality precluded reporting precise hydrodynamic diameter ( $d_h$ ) values. Detailed data are presented in Fig. S72–S81 and Tables S1–S10 (ESI†). Table 3 summarizes



**Fig. 4** (a) PC2 vs. PC1 scores for C1. (b) PC2 vs. PC1 scores for C2.

**Table 2** Estimated concentration ranges for potential turning points

% H <sub>2</sub> O (v/v)	C1	C2
	C <sub>0</sub> /μM	C <sub>0</sub> /μM
0	130–236	160–291
20	99.6–181	118–158
33	65–92	160–291
47	133–222	153–278
53	95–127	97.5–177
60	—	106–193
67	104–189	74–98.6
75	18.9–27; 161–229	31.2–39.1
82	55.9; 171–213	—
88	102–186	108–197
95	56.1–74.8	81.2–108

DLS data for **C1** and **C2** at 33, 53, 67, 82, 88, and 95% water in THF.

DLS results for **C2** revealed a trimodal size distribution at 33% water, with  $d_h$  values of  $3.42 \pm 0.07$  nm (78.4%),  $348 \pm 39$  nm (18.7%), and  $5448 \pm 81$  nm (2.9%). Like **C1**, the largest aggregates were of minor importance. However, the relative abundance of the major component (3.42 nm) was lower for **C2** compared to **C1** (3.51 nm, 84%), while the middle size component of **C2** (348 nm) was more prominent than that of **C1** (536 nm). Notably, aggregate sizes for **C1** and **C2** were comparable at this water content. At 53% water, **C2** exhibited a nearly unimodal distribution with  $d_h$  of  $944 \pm 43$  nm (99%) and a minor population at  $5514 \pm 66$  nm (1%), accompanied by a high P.I. of  $0.32 \pm 0.03$ . Regarding overall size, unlike **C1**, **C2** lacked the small aggregates present in **C1** but retained a small population of very large aggregates. A virtually bimodal distribution emerged at 67% water for **C2**, with  $d_h$  values of  $110 \pm 10$  nm (16.6%) and  $470 \pm 26$  nm (82.6%), and a minor population at  $5547 \pm 20$  nm (0.8%). The high P.I. of  $0.53 \pm 0.05$  indicated significant polydispersity. Smaller micellar aggregates ( $\sim 110$  nm), most likely precursors to the

dominant species at higher water content, coexisted with larger, dominant aggregates ( $\sim 470$  nm). Negligible amounts of very large aggregates persisted. This behavior contrasts with the unimodal size distribution of **C1** at 67%. It is worth noting that, although DLS and UV-vis results cannot be straightforwardly correlated, the latter contrast is also observed in Fig. 4, where **C1** and **C2** exhibit a much distinct pattern at 67% water. Unimodal distributions were observed at 82% and 95% water for **C2**, with  $d_h$  values of  $143 \pm 3$  nm (P.I. =  $0.14 \pm 0.01$ ) and  $140 \pm 1$  nm (P.I. =  $0.113 \pm 0.007$ ), respectively. While both **C1** and **C2** showed unimodal distributions at these water contents, **C1** exhibited larger micellar sizes.

The polydispersity index (P.I.) is a measure of the distribution of particle sizes in a sample. A P.I. above 0.5 indicates a broad range of particle sizes.<sup>39</sup> The DLS results show that **C2** in general has a higher P.I. than **C1**, suggesting a broader range of particle sizes in **C2**. In addition, the highest P.I. values are achieved at lower water contents for **C2** (53–67%) compared to **C1** (82%).

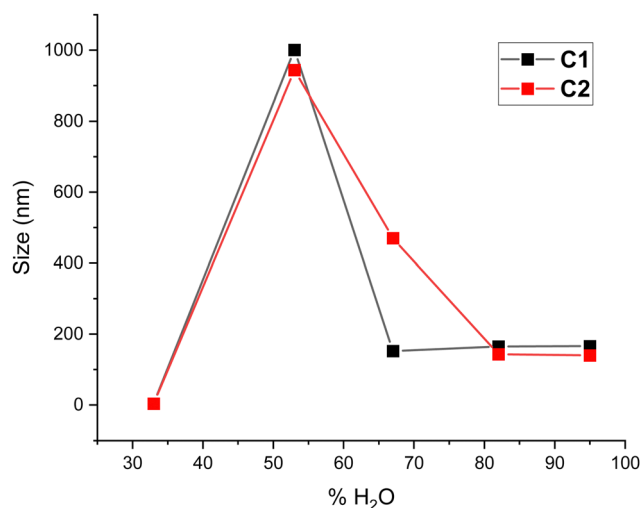
When the DLS sizes of the more important contribution to the scattering (*i.e.* that contributing above 80% to the total scattering) are represented *vs.* % of water content, a plot like that of Fig. 5 is obtained. This figure shows the difference between **C1** and **C2** behaviors of DLS sizes *versus* water content. In line with the results from the PCA of UV-vis data presented in Fig. 4, the major differences between the behavior of **C1** and **C2** were observed at 67% water content.

## 2.5. Fluorescence

Upon excitation at its high-energy absorption maximum ( $\lambda_{\text{ex}} = 285$  nm), **C2** in pure THF exhibited fluorescence emission centered at approximately 315 nm (Fig. 6(a)), with similar spectral characteristics to the fluorescence of **C1**.<sup>34</sup> The quenching of the fluorescence of **C1** and **C2** was studied either by measuring the fluorescence quantum yields by the integrating sphere method and/or by measuring the emission spectra at different water content in THF (for the complete absorption and emission spectra recorded, please refer to Fig. S82–S85, ESI†).

**Table 3** Hydrodynamic radii (DLS) from solutions of **C1** and **C2** prepared from THF or THF/H<sub>2</sub>O mixtures

Water (%)	C1 DLS $d_h$	C2 DLS $d_h$
0	—	—
33	Trimodal distribution $3.51 \pm 0.09$ nm (84.1%) $536 \pm 79$ nm (11.9%) $5204 \pm 203$ nm (3.9%) P.I. = n.d.	Trimodal distribution $3.42 \pm 0.07$ nm (78.4%) $348 \pm 39$ nm (18.7%) $5448 \pm 81$ nm (2.9%) P.I. = n.d.
53	Bimodal distribution $1001 \pm 145$ nm (94.4%) $2.88 \pm 0.04$ nm (5.3%) P.I. = n.d.	Bimodal distribution $944 \pm 43$ nm (99%) $5514 \pm 66$ nm (1%) P.I. = $0.32 \pm 0.03$
67	Unimodal distribution $152 \pm 2$ nm (100%) P.I. = $0.06 \pm 0.02$	Trimodal distribution $110 \pm 10$ nm (16.6%) $470 \pm 26$ nm (82.6%) $5547 \pm 20$ nm (0.8%) P.I. = $0.53 \pm 0.05$
82	Bimodal distribution $165 \pm 6$ nm (96.9%) $4687 \pm 296$ nm (3.1%) P.I. = $0.23 \pm 0.01$	Unimodal distribution $143 \pm 3$ nm (100%) P.I. = $0.14 \pm 0.01$
95	Unimodal distribution $166 \pm 2$ nm P.I. = $0.1 \pm 0.1$	Unimodal distribution $140 \pm 1$ nm (100%) P.I. = $0.113 \pm 0.007$

**Fig. 5** Behavior of major DLS sizes *vs.* % H<sub>2</sub>O for **C1** and **C2**.

Progressive water addition induced fluorescence quenching, resulting in a tenfold intensity reduction at 95% water content compared to the anhydrous solution. Luminescence quantum yields, determined *via* the integrating sphere method, are presented in Table 4 as a function of water percentage. Plots of  $\Phi_0/\Phi$  versus water content for **C1** and **C2** (Fig. 6(b)) deviated upward from linearity even at low water percentages. This observation in addition to the luminescence lifetimes behavior (see below) suggest the contribution of both dynamic and static quenching phenomena. Two key observations emerged from these plots: (i) an initially steeper slope for **C1** compared to **C2**, (ii) two distinct increases in slope, occurring around 60% and above 80% water content, and (iii) a steeper slope for **C2** than **C1** at high water content (*i.e.*, above 80%).

Fluorescence decay of **C2** was best described by a biexponential function. The longer lifetime ( $\tau_{\text{long}}$ ) decreased from 1.5 to 1.0 nanoseconds across the 0–95% water content range, while the shorter lifetime ( $\tau_{\text{short}}$ ) also decreased, varying between 0.5 and 0.2 nanoseconds (Table 4). At 0% water,  $\tau_{\text{long}}$  contributed 44% to the total relative amplitude, with  $\tau_{\text{short}}$  contributing 56%. These relative amplitudes remained relatively constant up to 67–75% water, but thereafter, the contribution of  $\tau_{\text{long}}$  decreased to 23% at 95% water, accompanied by a corresponding increase in the contribution of  $\tau_{\text{short}}$  to 77%.

Amplitude-averaged lifetimes ( $\tau_{\text{ave-amp}}$ ), calculated from  $\tau_{\text{long}}$  and  $\tau_{\text{short}}$ , decreased from 0.75 ns (0–47% water) to 0.2 ns at 95% water. Using  $\tau_{\text{ave-amp}}$ , average radiative ( $k_r$ ) and non-radiative ( $k_{\text{nr}}$ ) deactivation rate constants were determined according to the following equations:

$$k_r = \phi_f / \tau_{\text{ave-amp}}$$

$$k_{\text{nr}} = (1 - \phi_f) / \tau_{\text{ave-amp}}$$

In anhydrous THF,  $k_r$  and  $k_{\text{nr}}$  values for **C2** were determined to be  $1.98 \times 10^8$  and  $1.13 \times 10^9 \text{ s}^{-1}$ , respectively. Upon water addition,  $k_r$  decreased to  $1.1 \times 10^8 \text{ s}^{-1}$ , while  $k_{\text{nr}}$  increased to a

range of  $1.13$  to  $5.2 \times 10^9 \text{ s}^{-1}$  (Table 4). The luminescence decay behavior of **C2** resembled that of **C1**<sup>34</sup> although a more pronounced decrease in lifetimes was observed during quenching, indicating a more significant contribution of the dynamic quenching component to the overall quenching process in **C2** compared to **C1**.

A joint multivariate curve resolution (MCR) analysis of the full absorbance and fluorescence matrices (Fig. S82–S85, ESI†) revealed three primary spectral contributions, hereafter labeled as Sp1, Sp2, and Sp3 (Fig. 7). The first exhibited a relatively pure absorption spectrum and a high quantum yield, while the second showed an almost pure absorption spectrum but a significantly lower quantum yield (approximately 15 times less). The third contribution displayed a predominantly dispersive spectrum with some absorption and a quantum yield somewhat higher than the second. The total fluorescence demonstrated a steep decline between 47% and 60% water content, primarily due to the decrease of the first contribution. The scattering signal peaked at 60% water, but this did not correspond to a fluorescence maximum, suggesting a complex interplay between absorption, emission, and scattering processes. The bottom-center graph suggested a correlation between the high absorptive/low emissive contribution and the state of aggregation as the water percentage increased.

## 2.6. AFM analysis of aggregates

Previously reported AFM analysis of an amphiphilic calixarene revealed two distinct particle populations. Surprisingly, the first consisted of roughly 200 nm diameter spherical objects. More unexpectedly, a second type of structure was observed: a layered system with a thickness of 1.5–2 nm, resembling a supported lipid bilayer.<sup>40</sup> The morphology of **C1** and **C2** aggregates was investigated using AFM. Samples were prepared by depositing solutions onto mica substrates and drying them under vacuum. In pure tetrahydrofuran (THF), AFM imaging of **C1** showed flattened, round bilayer sheets with a trimodal size distribution

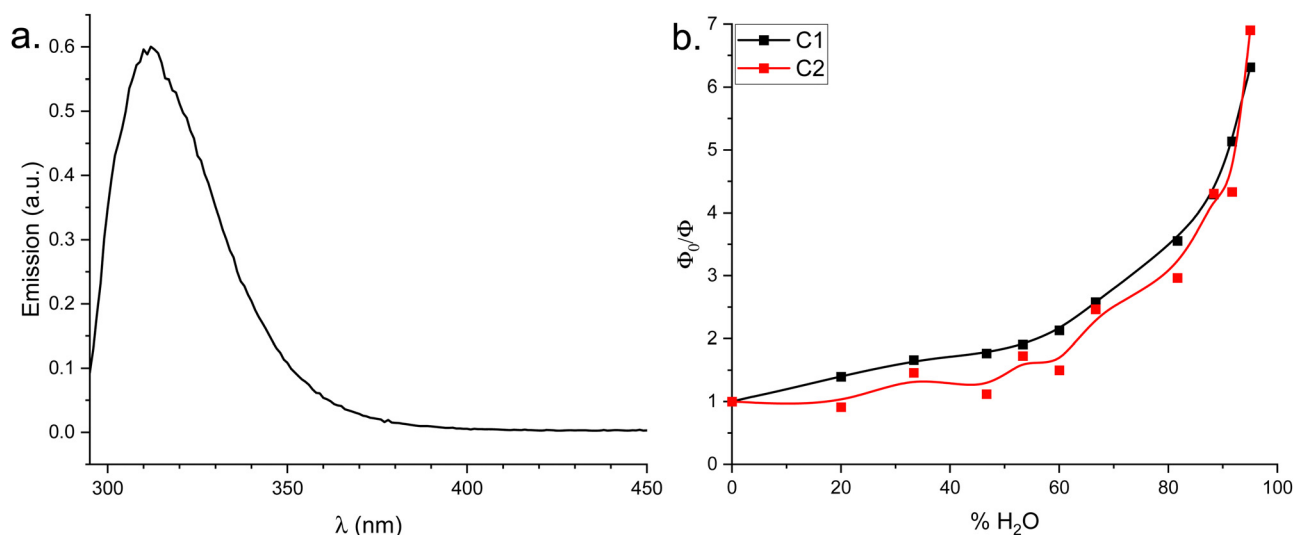
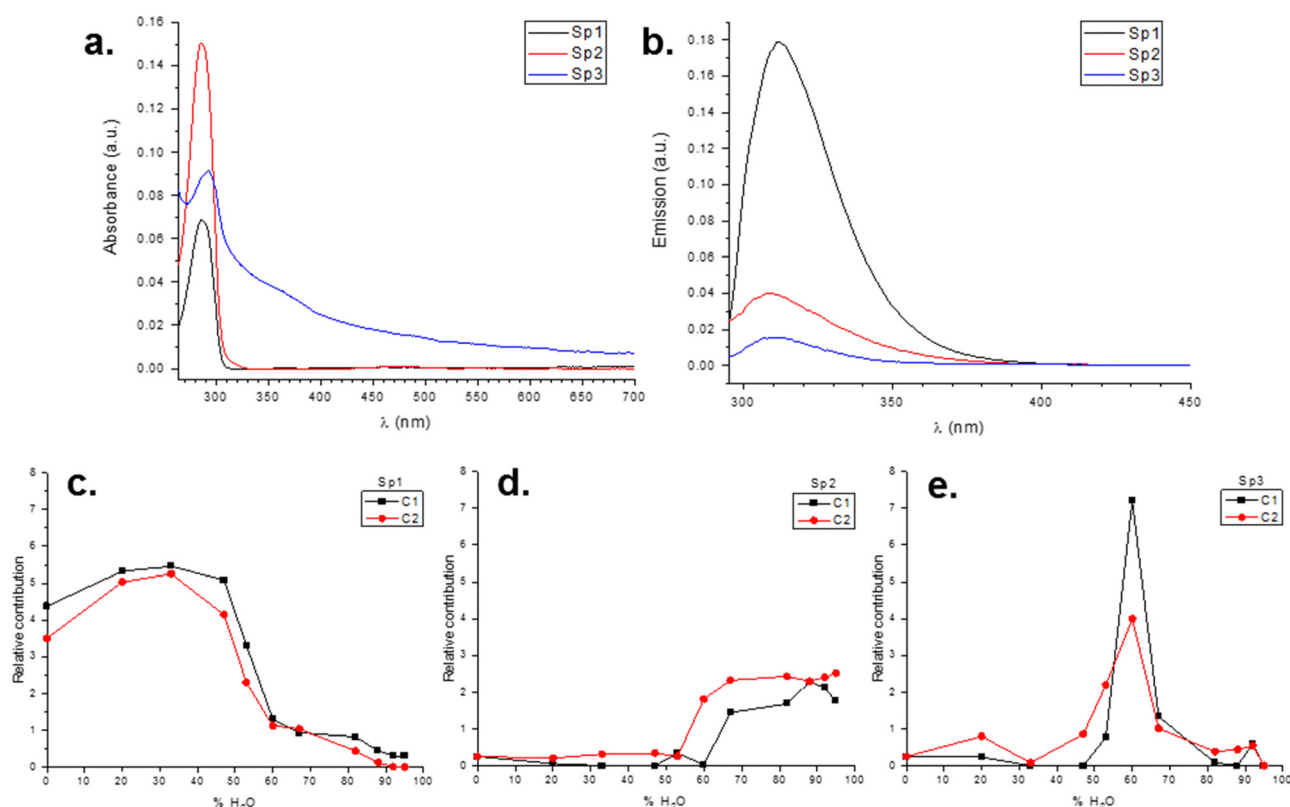


Fig. 6 (a) Luminescence spectra of **C2** in pure THF; (b)  $\Phi_0/\Phi$  values as a function of %  $\text{H}_2\text{O}$  for **C1** and **C2**.



**Table 4** Fluorescence quantum yields ( $\Phi$ ), luminescence lifetimes ( $\tau_{\text{short}}$  and  $\tau_{\text{long}}$ ), amplitude average lifetimes ( $\tau_{\text{ave-amp}}$ ), radiative ( $k_r$ ), and radiationless ( $k_{nr}$ ) deactivation rate constants of **C2** in THF at different % H<sub>2</sub>O

% H <sub>2</sub> O	$\Phi$	$\tau_{\text{short}}/\text{ns}$	$\tau_{\text{long}}/\text{ns}$	$\tau_{\text{ave-amp}}/\text{ns}$	$k_r/10^8 \text{ s}^{-1}$	$k_{nr}/10^9 \text{ s}^{-1}$
0	0.15 $\pm$ 0.02	0.54 $\pm$ 0.02	1.46 $\pm$ 0.04	0.75 $\pm$ 0.03	1.98 $\pm$ 0.09	1.13 $\pm$ 0.05
20	0.16 $\pm$ 0.02	0.54 $\pm$ 0.02	1.50 $\pm$ 0.04	0.75 $\pm$ 0.01	2.06 $\pm$ 0.05	1.13 $\pm$ 0.02
33	0.096 $\pm$ 0.009	0.50 $\pm$ 0.02	1.48 $\pm$ 0.04	0.73 $\pm$ 0.03	1.32 $\pm$ 0.06	1.24 $\pm$ 0.06
47	0.13 $\pm$ 0.01	0.35 $\pm$ 0.07	1.43 $\pm$ 0.09	0.7 $\pm$ 0.1	1.8 $\pm$ 0.3	1.2 $\pm$ 0.2
53	0.082 $\pm$ 0.008	0.28 $\pm$ 0.04	1.34 $\pm$ 0.09	0.55 $\pm$ 0.07	1.5 $\pm$ 0.2	1.7 $\pm$ 0.2
60	0.094 $\pm$ 0.009	0.29 $\pm$ 0.08	1.2 $\pm$ 0.1	0.5 $\pm$ 0.1	2.1 $\pm$ 0.5	2.0 $\pm$ 0.5
67	0.057 $\pm$ 0.006	0.28 $\pm$ 0.02	1.18 $\pm$ 0.04	0.44 $\pm$ 0.04	1.3 $\pm$ 0.1	2.1 $\pm$ 0.2
82	0.048 $\pm$ 0.004	0.26 $\pm$ 0.06	1.0 $\pm$ 0.1	0.37 $\pm$ 0.07	1.3 $\pm$ 0.3	2.6 $\pm$ 0.5
88	0.033 $\pm$ 0.003	0.22 $\pm$ 0.02	1.05 $\pm$ 0.06	0.28 $\pm$ 0.03	1.2 $\pm$ 0.1	4.5 $\pm$ 0.4
92	0.033 $\pm$ 0.003	0.19 $\pm$ 0.02	1.1 $\pm$ 0.1	0.23 $\pm$ 0.02	1.4 $\pm$ 0.3	4.1 $\pm$ 0.4
95	0.020 $\pm$ 0.002	0.15 $\pm$ 0.01	0.96 $\pm$ 0.08	0.19 $\pm$ 0.02	1.1 $\pm$ 0.2	5.2 $\pm$ 0.7

**Fig. 7** MCR analysis of the absorbance and fluorescence: (a) absorption spectra of contributions Sp1, Sp2 and Sp3; (b) emission spectra of contributions Sp1, Sp2 and Sp3; (c) Sp1 relative contribution vs. % H<sub>2</sub>O; (d) Sp2 relative contribution vs. % H<sub>2</sub>O; (e) Sp3 relative contribution vs. % H<sub>2</sub>O.

with diameters varying from 300–600 nm and heights varying from 8–47 nm.<sup>34</sup> Notably, all clusters displayed significant flattening, reminiscent of supported lipid bilayers as described by Shahgaldian *et al.*<sup>40</sup> AFM imaging revealed that **C2** formed a continuous layer on the mica surface, punctuated by depressions resembling a cluster of interconnected doughnuts (Fig. 8(a)).

At 53% water content (Fig. 8(b) and Fig. S86, ESI†), a marked decrease in both the number and size of larger clusters was observed. The distribution became predominantly unimodal with average dimensions of  $d = 150 \pm 42$  nm and  $h = 1.3 \pm 0.7$  nm, suggesting that water addition flattened and reduced the size of these clusters. Increasing the water content to 95%

resulted in the near-complete elimination of large clusters (Fig. 8(c) and Fig. S87, ESI†). Instead, a pronounced unimodal distribution with significantly smaller dimensions emerged ( $d = 73 \pm 18$  nm,  $h = 3.7 \pm 2.1$  nm). Notably, the morphological characteristics of **C2** observed in Fig. 8 and Fig. S86, S87 (ESI†) exhibit similarities to the self-assembly behavior of *para*-carboxy modified amphiphilic calixarene in water.<sup>40</sup>

Table 5 summarizes the sizes of **C2** aggregates measured by AFM, compared to those of **C1**. The trimodal distribution of flattened round bilayer sheets observed in **C1** is not apparent in **C2**, which forms a continuous layer on the mica surface. This suggests increased aggregation in **C2** owing to the longer

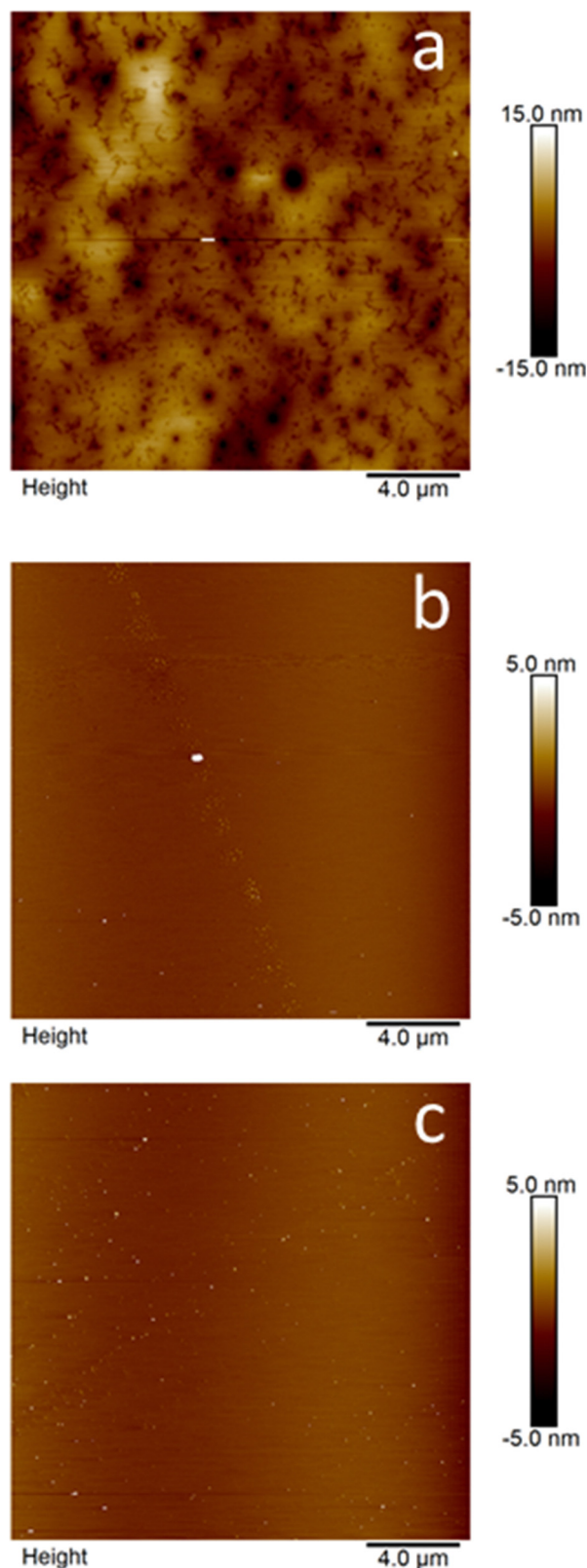


Fig. 8 AFM images of **C2** in: (a) pure THF, (b) THF (47%)/H<sub>2</sub>O (53%) and (c) THF (5%)/H<sub>2</sub>O (95%).

carbon chains. Additionally, at 53% water content, the bilayer sheets in **C2** are significantly smaller than those in **C1**,

indicating a more pronounced aggregation towards micellar-like structures in **C2**. This finding is consistent with the larger values of the hydrodynamic diameter ( $d_h$ ) for the major component of the scattering in DLS experiments for **C2**.

## 2.7. Analysis of solvent-dependent aggregation

PCA analysis of absorption spectra, combined with DLS and AFM experiments, reveals the following insights into the aggregation behavior of **C1** and **C2** molecules in THF/H<sub>2</sub>O mixtures:

1. Even in pure THF, resorcinarene molecules exhibit aggregation. This is evident from scattered light at low concentrations (70 μM) as observed by absorption spectroscopy and from the turning points in the PCA analysis at concentrations above 130 μM, indicating phase transitions.

2. In THF/H<sub>2</sub>O mixtures with 33% water content, DLS experiments suggest the presence of small micelles ( $d_h$  3.4–3.5 nm) as the dominant component (>80%). Additionally, large unilamellar vesicles (LUV) or giant unilamellar vesicles (GUV) with diameters of 536 nm for **C1** and 348 nm for **C2** are observed, along with minor components (3–4%) of very large GUV (5204 nm for **C1** and 5448 nm for **C2**). Based on previous calculations<sup>34</sup> the small micelles likely contain 8–10 molecules. These giant structures are collectively named GV, but the experimental techniques employed do not distinguish between GUV, multilamellar vesicle (MLV), or multivesicular vesicle (MVV) (see Scheme 2 for a schematic representation of micelles, bilayers, and different types of vesicles and their size dimensions).

3. At 53% water content, giant micelles (GM) become the dominant species in the size distribution, with diameters of 1001 and 944 nm for **C1** and **C2**, respectively.

4. AFM images indicate that solvent evaporation induces the coalescence and flattening of solution aggregates. This is evident from the observed flattened structures with diameters significantly larger than their heights.

5. For **C1**, a unimodal distribution of sizes is observed from DLS experiments, varying from 152 nm at 67% water to 165–166 nm at 82–95% water. **C2** shows a slightly different behavior, with a more gradual decrease in size from GUV to micelles as water content increases. At 67% water, the dominant aggregates for **C2** are 470 nm, with smaller aggregates of 110 nm and a minor component of very large structures (5547 nm). This differential behavior of **C1** and **C2** upon water content is also evident in Fig. 5.

## 2.8. Aggregation caused quenching in THF/H<sub>2</sub>O mixtures

Experiments employing absorption spectroscopy, dynamic light scattering, atomic force microscopy, and both steady-state and

Table 5 Diameters and heights of the flattened round layer sheets (AFM) from solutions of **C1** and **C2** prepared from THF or THF/H<sub>2</sub>O mixtures

Water (%)	<b>C1</b> AFM diameter (nm)	<b>C1</b> AFM height (nm)	<b>C2</b> AFM diameter (nm)	<b>C2</b> AFM height (nm)
0	605 ± 317 399 ± 99 336 ± 107	47 ± 34 15 ± 9 8 ± 5	—	—
53	302 ± 89	9 ± 5	151 ± 42	1.3 ± 0.7
95	72 ± 19	2.8 ± 0.5	73 ± 17	3.7 ± 2.1

time-resolved fluorescence revealed the phenomenon of aggregation caused by quenching (ACQ) in **C1** and **C2** molecules due to increasing water content in THF/H<sub>2</sub>O mixtures.

$\Phi_0/\Phi$  versus water content for **C1** and **C2** revealed subtle differences, which may be attributed to the effect of the chain length on the differential aggregation behavior observed by DLS, AFM, and PCA analysis of absorption spectra.

For instance, Fig. 5 shows that DLS sizes corresponding to the shorter-chain macrocycle (**C1**) drop more abruptly after water percent increases above 60% than the DLS of **C2** (the macrocycle with the longer hydrocarbon chain). The distinct behavior is also revealed in the MCR analysis of the Sp3 component (the purely dispersive component) in Fig. 7, which shows a broader shape of Sp3 for **C2** than for **C1**.

The analysis of fluorescence quenching as well as MCR results shows that:

1. The quenching of the fluorescence, evidenced either by the  $\Phi_0/\Phi$  plots (integrating sphere method, Fig. 6(b)) or the  $I_0/I$  or the  $A_0/A$  (intensities at the wavelength maximum or relative areas, Fig. S88 and S89, ESI†) has a relatively small slope which abruptly increases at water concentrations higher than 60% and with a steeper behavior for **C2** than for **C1**.

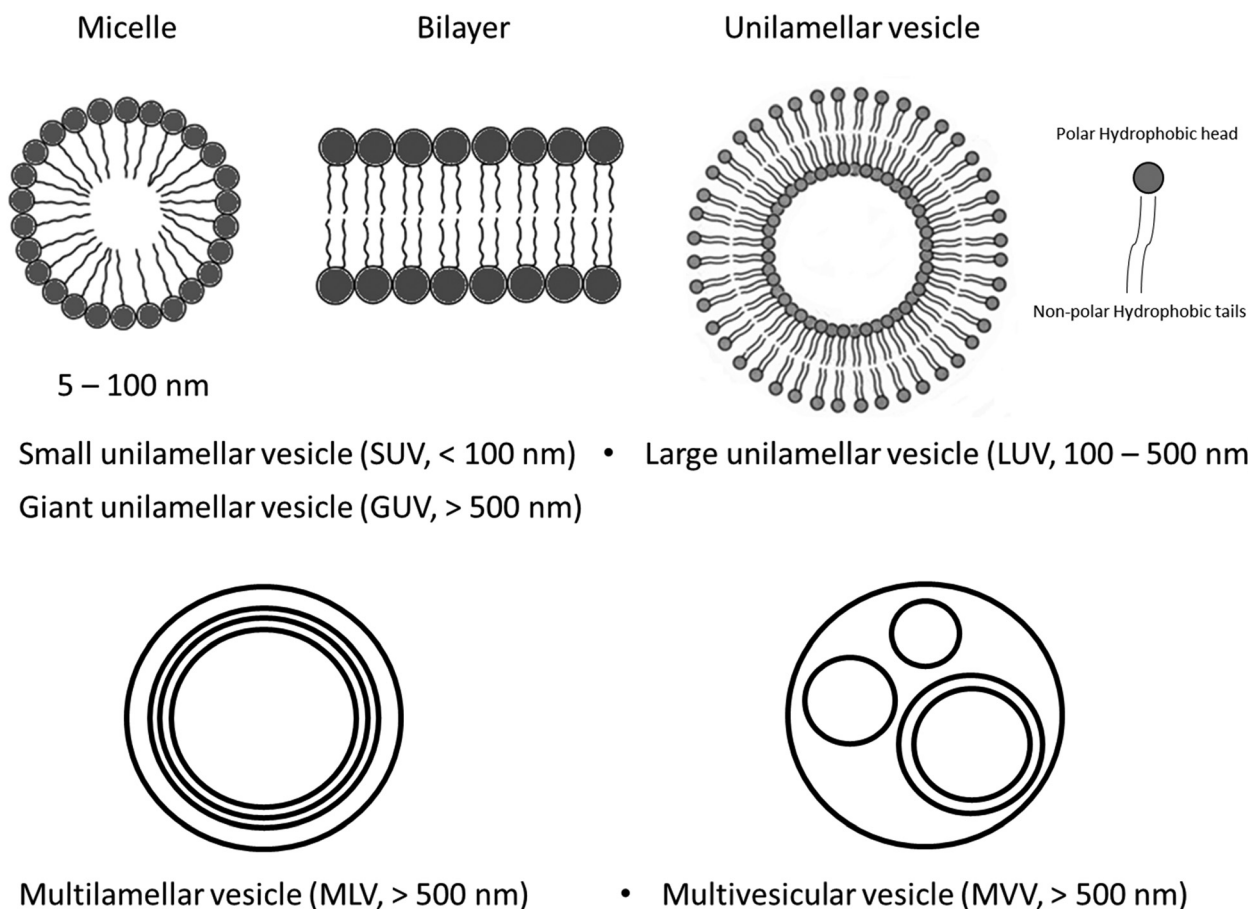
2. Of the three contributions retrieved by MCR, the decline of Sp1 as water% increases accounts for the recorded fluorescence quenching.

3. The poorly emissive contribution Sp2, which has a small scattering component relative to the spectral component, increases abruptly at water concentrations above 53% for **C2** and above 60% for **C1**, indicating differential micelle formation.

4. The Sp3 contribution, with spectral characteristics of a purely dispersive component, shows a water content percentage dependence more abrupt in **C1** than in **C2**.

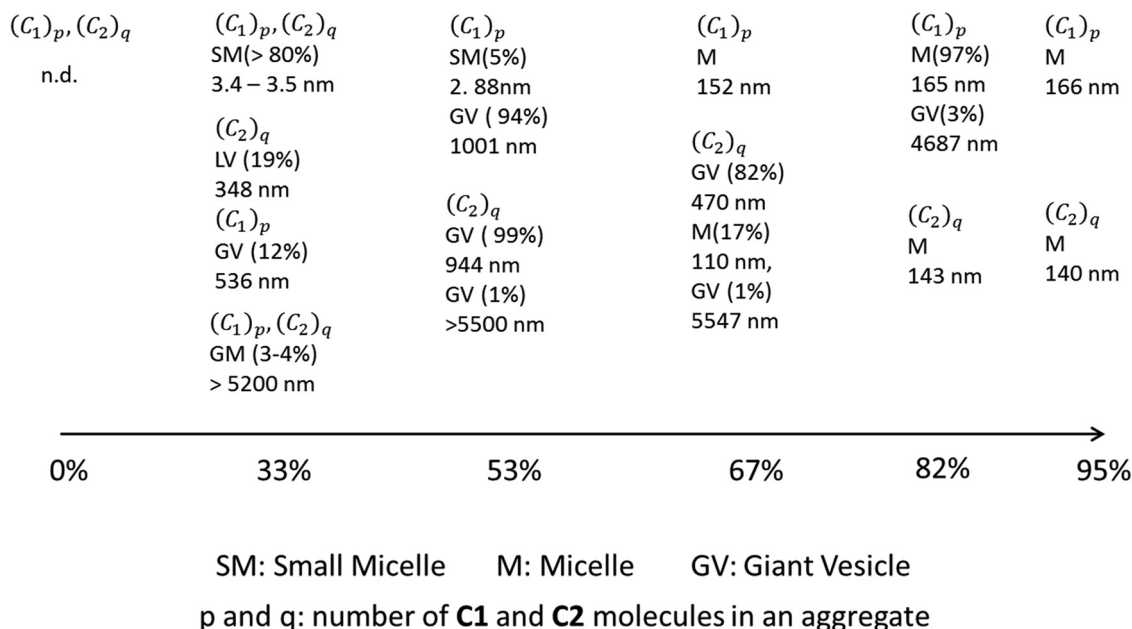
To a certain extent, although AFM and DLS sizes cannot be directly compared due to the flattening process in AFM due to solvent evaporation, this differential aggregation pattern shown by **C1** and **C2** was also observed by AFM, notably showing a trimodal distribution of flattened round bilayer sheets observed in **C1** that is not apparent in **C2**, which forms a continuous layer on the mica surface. This suggests increased aggregation in **C2** owing to the longer carbon chains. Additionally, at 53% water content, the bilayer sheets in **C2** are significantly smaller than those in **C1**, indicating a more pronounced aggregation towards micellar-like structures in **C2**.

Furthermore, PCA shows a differential behavior of PC1 and PC2 versus  $C_0$  in **C1** and **C2**, especially in the different turning points concentration ranges. The deviation from a



Scheme 2 Micelles, bilayers, vesicles and their characteristic dimensions.

## Aggregates types, abundance and size versus water content



Scheme 3 Influence of water content on aggregates types, abundance and sizes.

linear Stern–Volmer plot, even at low percentages of water, indicates the contribution of both static and dynamic quenching. The static quenching observed in the fluorescence of C1/C2 upon the addition of water is likely associated with an enhancement of nonradiative deactivation due to  $\pi$ – $\pi$  stacking of neighboring resorcinol molecules. This stacking is exacerbated by H-bond formation and chromophore crowding as the water content is increased. Within the resorcinarene macrocycle clusters, there may be room for a specific number of chromophores whose fluorescence cannot be quenched by further addition of water, leaving some residual fluorescence at high water (95% or higher) contents.

Scheme 3 summarizes the experimental findings.

### 3. Conclusions

This investigation reveals a significant influence of water content on the self-assembly behavior of two resorcinarenes, C1 and C2. We demonstrate a water-mediated transformation, where resorcinarenes transition from forming large aggregates in THF/H<sub>2</sub>O mixtures with water content below 60% to well-defined, smaller clusters as water concentration reaches 80–95%. Based on the experimental data, the addition of water to THF/H<sub>2</sub>O mixtures leads to the aggregation of C1 and C2 molecules, with the chain length influencing the aggregation behavior. MCR analysis reveals distinct species contributing to fluorescence quenching and aggregation. The static quenching contribution observed is probably due to p–p stacking. C1 and C2 exhibit different aggregation patterns, with C2 showing a more pronounced tendency towards micellar-like structures at higher water concentrations.

This precise control over C1 or C2 self-assembly is accompanied by a phenomenon known as aggregation-induced quenching (ACQ), which effectively modulates the fluorescence intensity. While ACQ may initially appear to be a limitation, it presents a unique opportunity for sensor design. Prior research has established the utility of ACQ in calixarenes and resorcinarenes for selective detection of specific targets, including Cu(II) and Co(II) ions and organic pollutants like 4-nitrotoluene. Our findings suggest that these resorcinarenes, with their water-tunable ACQ, possess significant potential as a platform for the development of next-generation sensors with precise fluorescence control.

### Data availability

The data supporting this article have been included as part of the ESI.†

### Conflicts of interest

There are no conflicts to declare.

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