

Cite this: *Nanoscale*, 2019, **11**, 19225

# Towards fast-charging technologies in Li<sup>+</sup>/Na<sup>+</sup> storage: from the perspectives of pseudocapacitive materials and non-aqueous hybrid capacitors

Haijian Huang  and Markus Niederberger \*

Since the discovery of the pseudocapacitive behavior in RuO<sub>2</sub> by Sergio Trasatti and Giovanni Buzzanca in 1971, materials with pseudocapacitance have been regarded as promising candidates for high-power energy storage. Pseudocapacitance-involving energy storage is predominantly based on faradaic redox reactions, but at the same time the charge storage is not limited by solid-state ion diffusion. Besides the search for pseudocapacitive materials, their implementation into non-aqueous hybrid capacitors stands for the strategy to increase power density by a rational design of the battery structure. Composed of a battery-type anode and a capacitor-type cathode, such devices show great promise to integrate the merits of both batteries and capacitors. Today, the availability of fast-charging technologies is of fundamental importance for establishing electric vehicles on a mass scale. Therefore, from the perspective of materials and battery design, understanding the basics and the recent developments of pseudocapacitive materials and non-aqueous hybrid capacitors is of great importance. With this goal in mind, we introduce here the fundamentals of pseudocapacitance and non-aqueous hybrid capacitors. In addition, we provide an overview of the latest developments in this fast growing research field.

Received 7th July 2019,  
Accepted 12th September 2019  
DOI: 10.1039/c9nr05732c  
rsc.li/nanoscale

## 1. Introduction

In light of the rising environmental issues, developing benign sustainable energy storage technologies has become a global priority. Reflecting its importance, massive efforts from the academic as well as from the industry side have been devoted to the development of Li-ion battery (LIB) technology over the years.<sup>1,2</sup> From microchips to electric vehicles (EVs), LIBs have been tremendously successful as an efficient power source.<sup>3</sup> Particularly, the role of LIBs in enabling the replacement of fossil-fuelled vehicles by EVs was highlighted in order to reduce the consumption of oil resources and the generation of CO<sub>2</sub> gases.<sup>2</sup> However, to achieve a leap-forward development in the field of LIBs several challenges have to be overcome.

Extensive efforts have also been devoted to developing post-Li-ion-battery technologies such as Li-metal batteries,<sup>4,5</sup> Li-S batteries,<sup>6,7</sup> and Li-O<sub>2</sub> batteries<sup>7,8</sup> mainly to boost the energy density of rechargeable batteries. However, also power density remains a technological issue. Recently, Cui *et al.* highlighted the importance of fast-charging technologies for the further promotion of EVs.<sup>9</sup> In fact, when energy density is improved (while keeping the power density constant), also the charging time is

prolonged ( $E = Pt$ , where  $E$ ,  $P$ , and  $t$  correspond to energy, power, and time, respectively). As shown in Fig. 1a, the recharge time of EVs is on the order of tens of hours, almost 200 times longer than the refueling time of a gasoline car. The extremely long charging time of EVs poses a serious barrier for a wider market penetration. The fastest recharge rate achieved by Tesla vehicles is 120 kW, enabling a charging time of 30 min for 200 miles.<sup>9</sup> However, compared to the quick refuelling experience (*ca.* 2 min) of conventional vehicles, this charging capability is still far behind the satisfactory level for the customers. Time is precious in the rush of modern life. High power conversion capability means saving a great amount of time. Flash charging technology will become an attractive feature of EVs and other energy storage devices, significantly strengthening their position in the market. In fact, developing rechargeable batteries with both high energy density and high power output has become one of the main challenges on the way to advanced energy storage devices.

In this review, mainly two topics will be addressed: (i) exploration of high-rate electrode materials, and (ii) rational design of the battery structure.

Regarding the first aspect, electrode materials with pseudocapacitive behaviour turned out to be particularly promising, because they are able to bridge the gap between a capacitor and a conventional battery electrode material. The charge storage mechanism is based either on surface/near-surface faradaic processes or on the fast ion intercalation into tunnels or layers of a redox-active material.<sup>10</sup>

Laboratory for Multifunctional Materials, Department of Materials, ETH Zürich, Vladimir-Prelog-Weg 5, 8093 Zürich, Switzerland.  
E-mail: markus.niederberger@mat.ethz.ch





Fig. 1 (a) Refuelling time and charging time of a gasoline car and an electric vehicle, respectively. (b) Elemental abundance in the Earth's crust. Reproduced with permission from ref. 16. Copyright 2014, American Chemical Society.

Concerning the rational design of battery structures, efforts have been put into the development of hybrid capacitors, in which a battery-type electrode is combined with a capacitor-type electrode to overcome the energy density limitation of conventional electrochemical capacitors.<sup>11,12</sup> When mentioning hybrid capacitors, two systems are usually involved, *i.e.*, aqueous hybrid capacitors and non-aqueous hybrid capacitors. Non-aqueous hybrid capacitors using a non-aqueous electrolyte usually have a higher energy density than aqueous ones due to the high operating voltage.<sup>12-15</sup> Accordingly, this minireview mainly focuses on non-aqueous systems.

In spite of the progress, which has been made towards high-power energy storage devices, the rapid development of EVs and advanced smart phones has massively accelerated the need for fast charging technologies. Thus, exploration of novel pseudocapacitive materials in combination with advanced non-aqueous hybrid capacitors with high power density and high energy density remains in the focus of academic and industrial research.

On the other side, the increasing market share of electronic devices and EVs/hybrid EVs has put Li production under pressure. Not only is the abundance of Li on the earth quite limited (only about 20 ppm in the earth's crust,<sup>16</sup> Fig. 1b), but the export policies of the corresponding countries can vary quickly.<sup>16,17</sup> Na-ion batteries (NIBs), although under investigation since the 1970s and thus as long as LIBs, almost disappeared in the following three decades, because most of the efforts were devoted to the development of LIBs. However, the cost per kW h of energy that Na can provide is much lower than that of Li, moving NIBs back into the centre of interest, especially for large-scale applications.<sup>16-18</sup> Unfortunately, the large ionic radius of Na usually leads to sluggish ion diffusion kinetics, considerably hampering the development of electrode materials with not only high energy density, but also good rate capability.<sup>19</sup> Thus, on the way to high-performance NIBs, suitable electrode materials with favorable ion storage kinetics have to be found.

Our minireview focuses on pseudocapacitive materials and non-aqueous hybrid capacitors for high-power Li<sup>+</sup>/Na<sup>+</sup> storage. The content is divided into two sections. First, the theoretical

background of pseudocapacitance and non-aqueous hybrid capacitors is discussed. Second, based on selected examples rather than an exhaustive literature survey we provide a general overview of the recent advances in pseudocapacitive materials and non-aqueous hybrid capacitors.

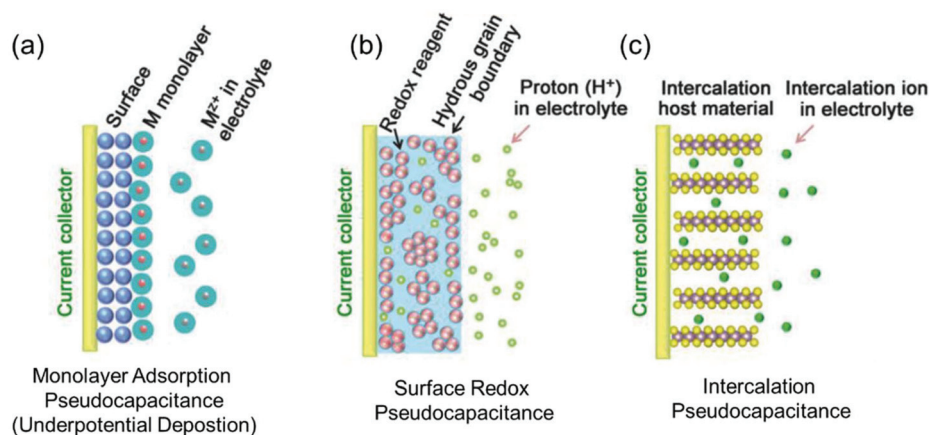
## 2. Pseudocapacitance

### 2.1. Basics of pseudocapacitance

Regarding power density, electric double-layer capacitors (EDLCs) possess a number of desirable properties compared to conventional batteries, including charging within seconds and long-term cycling stability. EDLCs have met with great interest since the first patent was released by General Electric in 1957.<sup>10,20</sup> Due to their extremely high power density, which can be up to 10 times that of rechargeable batteries, this type of devices has been successfully applied in many fields such as computer memory backup power, load lifting equipment, LED flash devices and so on.<sup>21,22</sup> Electrostatic double layer charge storage in EDLCs generally occurs through the adsorption of electrolyte ions onto the surfaces of electrode materials with no faradaic reactions involved. High-surface-area carbon materials are often identified to behave in this manner due to their large amount of surface adsorption sites and extraordinary high electronic conductivity. However, due to the confinement to the electrode surface area, the energy densities of commercial EDLCs are limited to 3–6 W h kg<sup>-1</sup>,<sup>10</sup> which are far below those of batteries (100–200 W h kg<sup>-1</sup> for commercial LIBs). In this context, merging the merits of both EDLCs and rechargeable batteries represents a promising pathway to high-power and high-energy electrode materials.

In the 1990s, Conway and co-workers proposed the concept of pseudocapacitance.<sup>23,24</sup> Energy storage involving pseudocapacitance, which relies predominantly on faradaic redox reactions, while the charge storage is not limited by solid-state ion diffusion, occupies a middle ground between EDLCs and batteries.<sup>25</sup> Generally, pseudocapacitance can be classified into three types:<sup>10</sup> (i) monolayer adsorption pseudocapacitance, (ii)





**Fig. 2** Different types of reversible redox mechanisms that give rise to pseudocapacitance: (a) monolayer adsorption pseudocapacitance (underpotential deposition), (b) surface redox pseudocapacitance, and (c) intercalation pseudocapacitance. Reproduced with permission from ref. 20. Copyright 2018, Wiley-VCH.

surface redox pseudocapacitance and (iii) intercalation pseudocapacitance (Fig. 2).

The origin of monolayer adsorption pseudocapacitance mainly comes from the reversible surface electrochemisorption (Fig. 2a). A typical example is the deposition of Pb on Au.<sup>26</sup> Due to the stronger interaction of Pb–Au compared with that of Pb–Pb in crystalline Pb metal, Pb can be deposited onto Au more easily than onto itself.<sup>26</sup> Such a process is also referred to as underpotential deposition.

Surface redox pseudocapacitance occurs when alkali ions are electrochemically adsorbed onto the surface or near the surface of the electrode materials through charge-transfer processes (Fig. 2b). RuO<sub>2</sub> was the first material found to demonstrate surface redox pseudocapacitive behavior. In 1971, Trasatti and Buzzanca discovered that the cyclic voltammogram (CV) curves of hydrous RuO<sub>2</sub> were rectangular in shape, which is a characteristic feature of a capacitor. However, the charge storage processes (storage of protons in this case) in RuO<sub>2</sub> were found to be of faradaic nature.<sup>27</sup> Subsequent studies revealed that the pseudocapacitive charge storage behavior of RuO<sub>2</sub> was a combination of the following factors: (a) the redox behavior of Ru<sup>4+</sup>, (b) the high electronic conductivity of RuO<sub>2</sub>, (c) short diffusion distances due to the large “outer surface” and (d) the large structural-water-induced “inner surface” in the porous hydrous RuO<sub>2</sub>.<sup>28,29</sup>

Intercalation pseudocapacitance arises, when the intercalation/deintercalation processes are not kinetically limited by the slow solid-state diffusion of the alkali ions in the crystal structure of the electrode materials on the time scale of interest (Fig. 2c). Compared with surface redox pseudocapacitance, which only occurs on or near the electrode surface, intercalation pseudocapacitance exists in the bulk.<sup>30</sup> A clear example is the pseudocapacitance in T-Nb<sub>2</sub>O<sub>5</sub>.<sup>30</sup> Through investigation of the electrochemical features of T-Nb<sub>2</sub>O<sub>5</sub>, Dunn and co-workers proved that the crystal structure of T-Nb<sub>2</sub>O<sub>5</sub> allowed pathways for ion transport with low energy barriers and only negligible changes of the lattice constants during intercalation of Li<sup>+</sup> into

T-Nb<sub>2</sub>O<sub>5</sub> when operating within a voltage window of 1.2–3.0 V.<sup>31,32</sup> The CV measurements showed that the current was linearly proportional to the voltage sweep rate, and the capacity did not vary significantly with charging time even for bulk material. They demonstrated that for charging times as fast as 1 min, there were still no diffusion limitations in bulk T-Nb<sub>2</sub>O<sub>5</sub>.

## 2.2. Identification of pseudocapacitance

The charge storage in pseudocapacitive materials is faradaic in origin as in battery-type materials,<sup>33–35</sup> although  $dQ/dV$  ( $Q$  and  $V$  represent charge and potential, respectively) derived from the charging/discharging curves corresponds to a capacitive behavior.<sup>13</sup> Owing to the non-diffusion controlled ion storage processes, pseudocapacitive materials usually show higher rate capability than battery-type materials. Identification of pseudocapacitance can be carried out through the investigation of electrochemical features characteristic for pseudocapacitive materials using two techniques: (i) galvanostatic charge/discharge curves and (ii) potentiodynamic sweep CV measurements. Additionally, also phase transitions have to be considered. These three aspects will be discussed in the following sections.

**2.2.1. Galvanostatic charge/discharge curves.** The energy storage of EDLCs involves ion adsorption at the surface of the electrode without any redox reactions. The stored charge  $Q_c$  for EDLCs is a linear function of the potential difference  $\Delta U$ :

$$Q_c = C\Delta U \quad (1)$$

where the capacitance  $C$  is an intrinsic property of EDLCs, which can be described as:

$$C = \epsilon_r \epsilon_0 A / d \quad (2)$$

$\epsilon_r$ ,  $\epsilon_0$ ,  $A$ , and  $d$  correspond to the dielectric constant of the electrolyte, the dielectric constant of the vacuum, the charge separation distance, and the electrode surface area, respectively. Accordingly, the profile of the voltage *versus* capacity for EDLCs has a nearly linear shape (Fig. 3a).<sup>36</sup> Pseudocapacitive





**Fig. 3** Galvanostatic charge/discharge curves for different types of electrode materials. (a) Carbon-based EC material. Reproduced with permission from ref. 36. Copyright 2017, Electrochemical Society. (b) Pseudocapacitive  $K_xMnO_2$ . Reproduced with permission from ref. 37. Copyright 2016, American Chemical Society. (c) Battery-type  $LiFePO_4$ . Reproduced with permission from ref. 38. Copyright 2001, Electrochemical Society.

materials exhibit nearly linear charge/discharge profiles without pronounced plateaus, which is analogous to the capacitive behavior (Fig. 3b). Note that the figure demonstrates the relationship between the potential and the time  $t$ . Because the measurement was performed at a constant current  $i$ , and  $Q_c = it$ , the figure also indicates the correlation between the potential and the stored charge.<sup>37</sup> According to eqn (1), the capacitance can be estimated based on the slope of the charge/discharge profiles, *i.e.*,  $C = dQ_c/d\Delta U$ . In contrast, typical battery electrode materials demonstrate distinct charge/discharge potential plateaus related to the phase transformations (Fig. 3c).<sup>38</sup> Therefore, the quasi-linear relationship between capacity and voltage is a distinct electrochemical feature of pseudocapacitive materials when compared with typical battery-type materials.

**2.2.2. Potentiodynamic sweep CV measurements.** In response to potential scanning during CV measurements, the capacitive current  $i_c$  for EDLCs usually follows:

$$i_c = dQ/dt = CdU/dt = Cv \quad (3)$$

$C$  is capacitance and  $v$  is voltage sweep rate. The capacitance hereby can be evaluated according to  $C = i_c/v$ . Due to the potential-independent capacitance, the capacitive current  $i_c$  is also potential-independent, resulting in a rectangular shape of

the CV curves for EDLCs (Fig. 4a).<sup>39</sup> On the other hand, battery-type materials show prominent and clearly separated peaks associated with the redox reactions involved in charge storage (Fig. 4c).<sup>40</sup> Behavior in between these two extremes indicates the presence of pseudocapacitance (Fig. 4b).<sup>41</sup> Here it should be emphasized that materials with CV curves containing intense and clearly separated oxidative and reductive peaks should be classified as battery-type materials instead of pseudocapacitive ones.<sup>25</sup>

Dunn and co-workers demonstrated that CV analysis can be further used to estimate the current contribution from pseudocapacitance through an appropriate experimental design.<sup>10,42</sup> In fact, one can quantitatively separate the contributions from capacitive effects (fast surface/near-surface redox reactions or fast ion insertion) and slow diffusion-controlled processes through analyzing the dependence of the CV curves on the sweep rates. In CV measurements, the current response at a fixed potential ( $i(V)$ ) can be expressed as the sum of the current contributions from the capacitive effects ( $i_c$ ) and the diffusion-controlled insertion reactions ( $i_d$ ):<sup>42</sup>

$$i(V) = i_c + i_d \quad (4)$$

According to eqn (3),  $i_c$  is proportional to the sweep rate  $v$  ( $i_c \propto v$ ). On the other hand, the contribution of the current  $i_d$



**Fig. 4** Galvanostatic charge/discharge curves for different types of electrode materials. (a) Carbon-based EC material. Reproduced with permission from ref. 39. Copyright 2013, Elsevier. (b) Pseudocapacitive  $MnO_2$ . Reproduced with permission from ref. 41. Copyright 1999, Elsevier. (c) Battery-type  $LiFePO_4$ . Reproduced with permission from ref. 40. Copyright 2011, Electrochemical Society.





from the diffusion-limited processes is proportional to the square root of the voltage scan rate  $\nu$  according to the Randles-Sevcik equation,<sup>43</sup> *i.e.*,  $i_d \propto \nu^{1/2}$ . Therefore, eqn (4) can be transformed into:

$$i(V) = k_1\nu + k_2\nu^{1/2} \quad (5)$$

where  $k_1\nu$  and  $k_2\nu^{1/2}$  characterize the capacitive and the diffusion-limited contributions at a given potential. Eqn (5) can be rearranged by dividing both sides by  $\nu^{1/2}$ , resulting in:

$$i(V)/\nu^{1/2} = k_1\nu^{1/2} + k_2 \quad (6)$$

$k_1$  and  $k_2$  can be determined by plotting  $i(V)/\nu^{1/2}$  versus  $\nu^{1/2}$ . This analysis enables the quantitative determination of the relative contributions of the capacitive processes in the total stored charge.

Besides, in CV curves the peak current  $i_p$  also obeys a power law relationship with the sweep rate  $\nu$  described as:

$$i_p = a\nu^b \quad (7)$$

Eqn (7) can be further transformed to:

$$\ln(i_p) = \ln a + b \ln \nu \quad (8)$$

Here,  $a$  and  $b$  are adjustable parameters.<sup>30,42</sup> The parameter  $b$  is determined by the slope of  $\ln(i_p)$ - $\ln(\nu)$  plots. When  $b = 0.5$ , the peak current is proportional to the square root of the sweep rate ( $\nu^{1/2}$ ), indicating a diffusion-controlled faradaic process. At the other end,  $b = 1$  represents a fully capacitive charge storage behavior, *i.e.*, no diffusion limit. A typical example is the study of the pseudocapacitive effects of TiO<sub>2</sub> by Dunn *et al.*<sup>42</sup> Capacitive contributions of 55%, 35%, and 15% of the total lithium-ion storage were calculated for 7, 10, and 30 nm particle sizes, respectively, through the above CV analysis (Fig. 5a and b). In addition,  $b$ -values for the 10 nm TiO<sub>2</sub>

film were in the range of 0.8–1.0 at potentials higher or lower than the peak potential, indicating that the current is predominantly capacitive (Fig. 5c) at these potentials.

Another method based on the relation between capacity  $Q$  and scan rate  $\nu$ , which was developed by Ardzizzone *et al.*,<sup>44</sup> is also suitable to differentiate capacitive and diffusion-controlled processes. The total stored charge  $Q$  can be separated into two parts, *i.e.*, the capacitive contribution  $Q_c$  and the diffusion-controlled contribution  $Q_d$ :

$$Q = Q_c + Q_d \quad (9)$$

According to eqn (1), the capacitive contribution  $Q_c$  is independent of the sweep rate  $\nu$ . The  $Q_c$  value can be obtained at an infinite potential sweep rate ( $\nu \rightarrow \infty$ ), *i.e.*,  $Q_c = Q_{\nu=\infty}$ . On the other hand,  $Q_d$  follows:

$$Q_d = i_d t \quad (10)$$

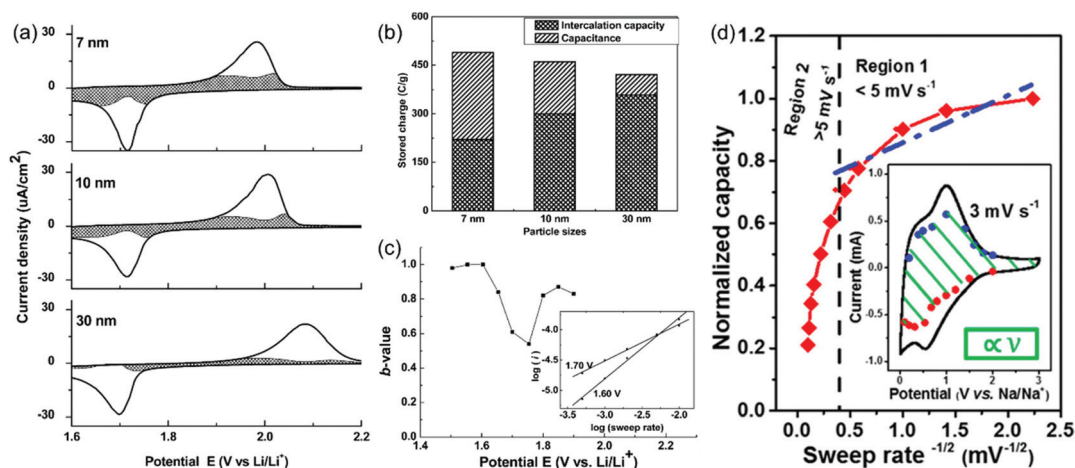
where  $t$  represents the sweep time and follows:

$$t = \Delta U / \nu \quad (11)$$

Therefore,  $t$  is proportional to  $\nu^{-1}$  ( $t \propto \nu^{-1}$ ). As mentioned above,  $i_d$  contributed by the diffusion-limited processes is proportional to the square root of the sweep rate  $\nu$  ( $i_d \propto \nu^{1/2}$ ). Combining eqn (10) with eqn (9), it can be concluded that the stored charge  $Q_d$  is expected to be limited by  $\nu^{-1/2}$ . Accordingly, eqn (9) can be converted into:

$$Q = Q_{\nu=\infty} + k\nu^{-1/2} \quad (12)$$

where  $k$  is a constant. By plotting  $Q$  versus  $\nu^{-1/2}$ ,  $Q_{\nu=\infty}$  can be obtained after extrapolating the linear fit to  $\nu^{-1/2} = 0$ . Hence, the pseudocapacitance-contributed capacity  $Q_{\nu=\infty}$ , which is independent of the scan rate  $\nu$  can be determined. An example is the kinetic analysis of the electrochemical behavior of mesoporous single-crystal-like TiO<sub>2</sub>@graphene nanocomposite



**Fig. 5** (a–c) CV analyses of TiO<sub>2</sub> films. Reproduced with permission from ref. 42. Copyright 2007, American Chemical Society. (a) Voltammetric response (0.5 mV s<sup>-1</sup>) for 7 nm, 10 nm and 30 nm TiO<sub>2</sub> films. (b) Comparison of the capacitive contributions in TiO<sub>2</sub> nanoparticle films (0.5 mV s<sup>-1</sup> sweep rate). (c)  $b$ -Values for the 10 nm film. (d) Plot of capacity vs.  $\nu^{-1/2}$  to separate diffusion-controlled and capacitive-controlled contributions in a mesoporous single-crystal-like TiO<sub>2</sub>@graphene nanocomposite. Reproduced with permission from ref. 45. Copyright 2017, American Chemical Society.



reported by Wu, Lu and co-workers.<sup>45</sup> Utilizing this method, the authors declared that 68% of the total charge achieved by the nanocomposite is contributed from the capacitive process.

**2.2.3. Phase transitions.** Galvanostatic charge/discharge curves and potentiodynamic sweep CV measurements are able to provide helpful information for the identification of pseudo-capacitive materials. However, the definition of pseudocapacitance is still an object of divergent views.<sup>46</sup> Brousse *et al.* insist that the original definition of pseudocapacitive materials should be strictly abided by.<sup>13</sup> A clear example is MnO<sub>2</sub>. Though faradaic in nature, MnO<sub>2</sub> displays a quasi-rectangular CV response without any redox peaks as it is observed for a typical capacitive electrode like carbon. On the other side, Dunn *et al.* presented a less restrictive definition that covers a broader behavior.<sup>10</sup> In this case, the CV features should present a shape that “is rectangular and if peaks are present, they are broad and exhibit a small peak-to-peak voltage separation”.<sup>10</sup> However, independent of the definition of pseudocapacitance, a quasi-rectangular shape of the CV response is always involved. In this context, it is important to note that materials with phase transitions during electrochemical processes should not be considered as intercalation pseudo-capacitive materials. The reason is that such phase transitions most often result in discharge plateaus. According to the definition by Conway, “pseudocapacitance arises when the extent of reaction, *Q*, is some continuous function of potential, *V*, so that the derivative, *dQ/dV*, arises that has the properties of a capacitance”.<sup>47</sup> With discharge plateaus involved, these materials exhibit non-constant capacitance over its operational potential range, which deviates from the definition of pseudocapacitance as mentioned above. At the same time we have to keep in mind that small particles can lead to more favorable ion diffusion kinetics and even first-order phase transitions can be efficiently suppressed in small crystals, which might be attributed to the unfavorable formation of nucleation sites in confined spaces.<sup>48</sup> As an example, bulk MoS<sub>2</sub> is a typical battery-type material with multiple phase transitions during lithiation and delithiation. Tolbert *et al.* found that in nano-sized MoS<sub>2</sub> (average particle size of 22 nm) the phase changes were suppressed, resulting in pseudocapacitive characteristics.<sup>48</sup> Similar observations were made for nano-sized MoO<sub>2</sub>,<sup>49</sup> TiS<sub>2</sub>,<sup>50</sup> and LiCoO<sub>2</sub>.<sup>51</sup> Such materials, which show conventional battery behavior in the bulk, but pseudo-capacitive feature on the nanoscale, have been termed “extrinsic” pseudocapacitive.<sup>51</sup> Table 1 presents an overview of some typical examples of intrinsic and extrinsic pseudocapacitive materials reported in the literature.

Consequently, identification of pseudocapacitance in electrode materials should be based on electrical and electrochemical considerations. Capacitor-like features shown in galvanostatic charge/discharge curves and in potentiodynamic sweep CV measurements are not sufficient proof for pseudocapacitance. Possible phase transitions within the voltage window also have to be taken into account. Only if *dQ/dV* is a constant value within the entire operating voltage window or approaches a constant value, can the material be described as pseudocapacitive.

**Table 1** Summary of some typical intrinsic and extrinsic pseudo-capacitive materials

	Materials	Operating voltage window
Intrinsic pseudoapacitive materials	RuO <sub>2</sub> <sup>28</sup>	0–1 V vs. saturated calomel electrode (SCE)
	MnO <sub>2</sub> <sup>41</sup>	–0.2–1 V vs. SCE
	LiFeTiO <sub>4</sub> <sup>52</sup>	1.5–4.8 V vs. Li/Li <sup>+</sup>
	T-Nb <sub>2</sub> O <sub>5</sub> <sup>30</sup>	1.2–3 V vs. Li/Li <sup>+</sup>
	Ti <sub>3</sub> C <sub>2</sub> T <sub>x</sub> MXene <sup>53</sup>	–0.2–0.4 V vs. Ag/AgCl
Extrinsic pseudocapacitive materials	LiCoO <sub>2</sub> <sup>54</sup>	3–4.2 V vs. Li/Li <sup>+</sup>
	MoS <sub>2</sub> <sup>55</sup>	0.8–3 V vs. Li/Li <sup>+</sup>
	MoO <sub>2</sub> <sup>49</sup>	1.1–3 V vs. Li/Li <sup>+</sup>
	TiS <sub>2</sub> <sup>50</sup>	1.5–3 V vs. Li/Li <sup>+</sup>

### 3. Non-aqueous hybrid capacitors

In comparison to EDLCs, rechargeable batteries offer a high energy density, but a low power density, meaning that it can take hours to recharge after they are empty.<sup>51</sup> This is because their energy storage mechanisms are based on faradaic chemical reactions, which are sluggish. Thus, recharging at high current rates will lead to rapid capacity fading with permanent damage to the battery. On the other hand, EDLCs, which store charges on the surfaces by forming electrical double layers, are distinguished from other types of electrochemical energy storage devices by short charging times and the ability to deliver significantly high power.<sup>56</sup> However, the limitation of EDLCs is their low energy density.<sup>51</sup> The energy limitations of electrical double layer capacitors can be reduced by combining the advantages of batteries and EDLCs in hybrid capacitors.<sup>57</sup>

The concept of hybrid capacitor was introduced more than 10 years ago. In 2001, Pasquier, Zheng and colleagues assembled a non-aqueous hybrid capacitor using activated carbon as the cathode, nanostructured Li<sub>4</sub>Ti<sub>5</sub>O<sub>12</sub> as the anode, and 1 M LiPF<sub>6</sub> in a 2 : 1 volume ratio of ethylene carbonate (EC) and dimethyl carbonate (DMC) as the electrolyte. The asymmetric hybrid cell achieved 90% capacity at 10C charge/discharge rate and 10–15% capacity loss after 5000 cycles. As shown in Fig. 6, a typical non-aqueous hybrid capacitor is composed of a battery-type anode with faradaic energy storage mechanism and a capacitor-type cathode involving. Upon charging, the anode undergoes lithium/sodium ion insertion, whereas the cathode involves electrostatic double layer charge storage processes with anions (*e.g.*, PF<sub>6</sub><sup>–</sup>, ClO<sub>4</sub><sup>–</sup>). During discharging, Li/Na-ion de-insertion occurs in the anode and desorption of anions takes place in the cathode.<sup>58</sup> As a result, the device stores charge through faradaic processes by Li<sup>+</sup>/Na<sup>+</sup> insertion/de-insertion in the anode together with a reversible adsorption/desorption of anions on the surface of the cathode. The combined effects provide an opportunity to enhance the energy density and power capability within one device.<sup>57</sup> When the battery-type electrode is a Li- or Na-intercalating phase, these systems are also referred to as Li- or Na ion hybrid capacitors (LICs/NICs).<sup>59</sup>

However, it is important to remember that the faradaic reactions are much more sluggish than the non-faradaic ones and





Fig. 6 Schematic configurations for a typical battery (left), a typical non-aqueous hybrid capacitor (middle) and a typical capacitor (right).

that the charge storage capability of capacitor-type materials falls behind that of battery-type materials. Therefore, the main challenge in LICs/NICs research is to counteract the imbalance of kinetics and capacity between the faradaic anode and the non-faradaic capacitive cathode.<sup>60</sup> In section 5, we will present several examples, how the advantages of batteries and EDLCs can efficiently be combined by developing suitable electrode materials.

## 4. Recent advancements in pseudocapacitive materials for high-power $\text{Li}^+/\text{Na}^+$ storage

### 4.1 Surface redox pseudocapacitive materials

Pseudocapacitive materials for high-power  $\text{Li}^+/\text{Na}^+$  storage can be divided into the two families of surface redox and intercalation pseudocapacitive materials. For surface redox pseudocapacitive materials, electrochemical reactions are more likely to occur on or near the surface of the electrode materials. Thus, these reactions are not diffusion (into the particles) controlled, leading to a high power density. Moreover, surface redox pseudocapacitive materials can show intrinsic and extrinsic pseudocapacitance. As an example for intrinsic surface redox pseudocapacitance,  $\text{RuO}_2$  exhibits pseudocapacitive features for a wide range of particle sizes and morphologies.<sup>10</sup> In 2007, Dunn *et al.* found that by engineering  $\text{TiO}_2$  (anatase) at the nanoscale, the contribution of pseudocapacitance increased with the reduction of the average size of the nanoparticles.<sup>42</sup> They proved that the diffusion-controlled lithium intercalation processes could be replaced by surface reactions, which are non-diffusion controlled, with decreasing size of the nanoparticles. The study suggested that surface redox pseudocapacitance could be introduced through downsizing of the particles, known as extrinsic pseudocapacitance. Dou *et al.* tested the combination of rGO with  $\text{MnO}$ ,<sup>61</sup> which turned out to exhibit increasing surface redox pseudocapacitive effects with decreasing particle size of  $\text{MnO}$ . Embedding of the  $\text{MnO}$  nanoparticles in between two layers of rGO, the sandwich structure achieved an unprecedented rate

capability ( $331.9 \text{ mA h g}^{-1}$  at  $10 \text{ mA cm}^{-2}$ ,  $379 \text{ mA h g}^{-1}$  after 4000 cycles at  $3.75 \text{ mA cm}^{-2}$ ) with outstanding cycling stability. Other research results on surface redox pseudocapacitive materials harnessing the decreased particle size include ultra-small MoC nanoparticles@nitrogen-doped porous carbon<sup>62</sup> or 5 nm  $\text{MnO}$  nanocrystals-graphene composite.<sup>63</sup>

In addition to particle size, morphological design plays another important role in improving the surface redox pseudocapacitive effects. Optimization of particle shape, in addition to composition and/or 3-dimensional architecture, can efficiently help to increase the surface area, thus providing more surface reaction sites. Sun *et al.* reported a facile and scalable strategy to fabricate hierarchical architectures, where  $\text{TiO}_2$  nanotube clusters were coated with a composite of ultra-fine  $\text{MoO}_2$  nanoparticles embedded in a carbon matrix ( $\text{TiO}_2@\text{MoO}_2\text{-C}$ ).<sup>64</sup> When applied to Na-ion storage, the unique structural design led to high pseudocapacitance contribution with excellent cycling stability up to 10 000 cycles even at a current density of  $7 \text{ mA cm}^{-2}$  and superior rate capacities of 110 and  $76 \text{ mA h g}^{-1}$  at current densities of 7 and  $14 \text{ mA cm}^{-2}$ , respectively. An interesting study by Dunn *et al.* revealed that the surface redox pseudocapacitance greatly benefits from a mesoporous morphology and a specific crystal orientation. Ion intercalation pseudocapacitance was enhanced due to an ideal arrangement of the fast Li ion diffusion channels in the iso-oriented crystalline domains.<sup>56</sup> Another class of porous materials are hollow nanostructures, which also showed enhanced redox pseudocapacitance. Lou *et al.* developed a template-directed strategy to synthesize NiS box-in-box hollow structures with double-shells.<sup>65</sup> The hollow NiS built up by ultrathin nanosheets turned out to exhibit significantly improved specific capacitance compared to dense NiS. Other typical examples falling into this materials category include complex CoS hollow structures,<sup>66</sup>  $\text{NiCo}_2\text{S}_4$  ball-in-ball hollow spheres,<sup>67</sup> and  $\text{Ni}_x\text{Co}_{3-x}\text{S}_4$  hollow nanoprisms,<sup>68</sup> which all demonstrated enhanced pseudocapacitive properties compared with their counterparts with dense structures.

Xia and co-authors used surface engineering to successfully achieve ultrahigh rate capability *via* tunable surface redox pseudocapacitance.<sup>69</sup> In their study, through modifying the surface of  $\text{Fe}_2\text{O}_3$ , they developed oxygen-deficient  $\text{Fe}_2\text{O}_3$







**Fig. 7** Example of a surface redox pseudocapacitive material: oxygen-deficient Fe<sub>2</sub>O<sub>3</sub> quantum dots anchored on graphene nanosheets (V-F@G). (a) Schematic of the fabrication procedure of V-F@G and the adsorption of SO<sub>3</sub><sup>2-</sup> on the V-F@G surface. (b) CV curves of the V-F@G electrode measured in 1 M Na<sub>2</sub>SO<sub>3</sub> electrolyte at high scan rates from 0.2 to 3.2 V s<sup>-1</sup>. (c) Surface-controlled capacity as a function of scan rate. Reproduced with permission from ref. 69. Copyright 2018, Wiley-VCH.

quantum dots anchored on graphene nanosheets (V-F@G) (Fig. 7a). When using a redox active electrolyte such as Na<sub>2</sub>SO<sub>3</sub>, the surface oxygen vacancies enabled greatly enhanced chemical adsorption of SO<sub>3</sub><sup>2-</sup>. At the same time, the negatively charged surface of Fe<sub>2</sub>O<sub>3</sub> with adsorbed SO<sub>3</sub><sup>2-</sup> attracts Na<sup>+</sup> ions, facilitating both the formation of double-layer capacitance and surface redox reactions. The synergistic combination of redox reactions involving both the cation (Na<sup>+</sup>) and the anion (SO<sub>3</sub><sup>2-</sup>) on the Fe<sub>2</sub>O<sub>3</sub> surface enabled rapid charge transport, *i.e.*, high pseudocapacitance (Fig. 7b and c). Similarly, Yang *et al.* prepared oxygen-deficient Li<sub>3</sub>VO<sub>4-δ</sub> with an amorphous surface and also demonstrated the important role of surface engineering in enabling the fast surface reactions.<sup>70</sup>

#### 4.2 Intercalation pseudocapacitive materials

Compared with surface redox pseudocapacitance, which generally needs to be enabled by complicated procedures of nanoscaling or nanostructuring, intercalation pseudocapacitance usually exists in bulk materials, if the ion pathways in the crystalline network are spacious enough to provide fast diffusion. As mentioned above, orthorhombic Nb<sub>2</sub>O<sub>5</sub> (T-Nb<sub>2</sub>O<sub>5</sub>) with an open, layered structure is a typical intercalation pseudocapacitive material. Ganesh *et al.* carefully studied the origin of high-rate intercalation pseudocapacitance in this material.<sup>71</sup> Based on complementary theoretical methods, the authors concluded that the intercalation pseudocapacitive behavior in Nb<sub>2</sub>O<sub>5</sub> was due to the following reasons: (i) the presence of interconnected sheets with multiple adsorption sites of comparable energies; (ii) local charge-transfer at all adsorption sites; (iii) open channels that reduce the diffusion barrier for

lithium to hop between these sites; (iv) the lack of 1<sup>st</sup> order phase transition upon intercalation/deintercalation. In addition to mechanistic studies, efforts were also made to further improve the electrochemical performance of T-Nb<sub>2</sub>O<sub>5</sub>. For example, Long *et al.* reported a simple synthesis method for free-standing orthorhombic Nb<sub>2</sub>O<sub>5</sub> (T-Nb<sub>2</sub>O<sub>5</sub>)@graphene composite papers (Fig. 8a), which showed a superior pseudocapacitor performance (Fig. 8b–d).<sup>72</sup> Using the composite paper, an organic electrolyte-based asymmetric supercapacitor was assembled, delivering a high energy density of 47 W h kg<sup>-1</sup> and a power density of 18 kW kg<sup>-1</sup>. Yin and co-authors prepared three-dimensionally ordered macroporous T-Nb<sub>2</sub>O<sub>5</sub> with mesoporous hierarchical structure, which turned out to offer highly efficient pseudocapacitive Li<sup>+</sup> intercalation behavior.<sup>73</sup> In another study, an in-plane assembled orthorhombic Nb<sub>2</sub>O<sub>5</sub> nanorod film was prepared by Liu *et al.*, which also showed high-rate Li<sup>+</sup> intercalation.<sup>74</sup>

Since the discovery of the intercalation pseudocapacitance in T-Nb<sub>2</sub>O<sub>5</sub>, a lot of efforts were devoted to exploring new materials with similar electrochemical behavior. In addition to T-Nb<sub>2</sub>O<sub>5</sub>, also spinel LiFeTiO<sub>4</sub>,<sup>52</sup> MoO<sub>3-x</sub> with oxygen vacancies,<sup>75</sup> and layered materials such as VOPO<sub>4</sub> nanosheets<sup>76</sup> and Ti<sub>2</sub>CT<sub>x</sub> MXene<sup>76</sup> were proved to demonstrate intercalation pseudocapacitive behaviors. In the case of α-MoO<sub>3-x</sub>, Dunn *et al.* found that the reduced α-MoO<sub>3-x</sub> showed much higher rate capability and cycling stability than fully oxidized α-MoO<sub>3</sub>. The authors proposed three reasons for the enhancement in the electrochemical performance, which all positively affected the intercalation pseudocapacity. X-ray powder diffraction (XRD) and density functional theory (DFT) results revealed an







**Fig. 8** T-Nb<sub>2</sub>O<sub>5</sub>/graphene composite papers as an example of an intercalation pseudocapacitive materials. (a) Schematic of the fabrication process of the T-Nb<sub>2</sub>O<sub>5</sub>/graphene composite papers. (b) CV curves of the T-Nb<sub>2</sub>O<sub>5</sub>/graphene composite papers. (c) Total current (solid line) and capacitive current (shade regions) at 5 mV s<sup>-1</sup> and (d) capacitive contribution at various scan rates for the T-Nb<sub>2</sub>O<sub>5</sub>/graphene composite papers. Reproduced with permission from ref. 72. Copyright 2015, American Chemical Society.

expanded *b*-lattice cell parameter for the reduced  $\alpha$ -MoO<sub>3-*x* relative to the oxidized  $\alpha$ -MoO<sub>3</sub>, which led to a larger interlayer spacing that promoted faster charge storage kinetics. *Ex situ* XRD measurements pointed to an irreversible electrochemically induced phase transition in the oxidized  $\alpha$ -MoO<sub>3</sub> during Li<sup>+</sup> insertion/extraction. In contrast,  $\alpha$ -MoO<sub>3-*x* did not undergo a phase transformation following lithiation to 1.5 V, indicating better structural reversibility. *Ex situ* XPS measurements provided the information that  $\alpha$ -MoO<sub>3-*x*</sub> experienced increased conversion of Mo<sup>6+</sup> to Mo<sup>4+</sup> compared with MoO<sub>3</sub>, which resulted in a capacity difference.</sub></sub>

MXenes, discovered by Gogotsi, Barsoum and colleagues, have been developed as a novel family of 2D materials.<sup>77–80</sup> Due to the large interlayer space, MXenes with intercalation pseudocapacitance are treated as promising candidates for high-power energy storage materials. Based on the pseudocapacitive behavior, Yamada *et al.* successfully applied MXenes in a high-power NIC. The prototype full cell consisting of a Na<sub>2</sub>Fe<sub>2</sub>(SO<sub>4</sub>)<sub>3</sub> positive electrode and a MXene Ti<sub>2</sub>CT<sub>x</sub> negative electrode provided reversible capacities of 90 and 40 mA h g<sup>-1</sup> at 1 and 5 mA cm<sup>-2</sup>, respectively (based on the weight of Ti<sub>2</sub>CT<sub>x</sub>).<sup>81</sup> Regarding the underlying mechanism of the Na-ion intercalation pseudocapacitance, Yamada and colleagues found that the trapped Na<sup>+</sup> after the first sodiation process behaves as a pillar and the penetrated solvent molecules swelled the interlayer space, both of which contributed to keep the interlayer distance constant during the sodiation/desodiation processes.<sup>82</sup> Therefore, no substantial change in the interlayer distance was observed during electrochemical reactions, resulting in high cycling stability in addition to fast Na<sup>+</sup> diffusion in the expanded interlayer space. Zhou *et al.*<sup>83</sup> proved that a surface termination with -F and -OH in MXene blocked electrolyte ion transport and decreased energy storage capacity.

Accordingly, Lin, Sun and co-authors demonstrated that after K<sup>+</sup> intercalation and removal of terminal groups (OH<sup>-</sup>/F<sup>-</sup>), the intercalation pseudocapacitance was three times higher than that of the pristine MXene, which was reflected in a significant enhancement of about 211% of the gravimetric capacitance.<sup>53</sup> In addition to carbide MXenes, also nitride MXenes, *e.g.*, Ti<sub>2</sub>NT<sub>x</sub> MXene,<sup>84</sup> have been proved to show pseudocapacitive charge storage mechanisms.

Recently, various layered metal vanadates, *e.g.*, Ca<sub>0.24</sub>V<sub>2</sub>O<sub>5</sub>·0.83H<sub>2</sub>O,<sup>85</sup> Mg<sub>0.34</sub>V<sub>2</sub>O<sub>5</sub>·0.84H<sub>2</sub>O,<sup>86</sup> Zn<sub>3</sub>V<sub>2</sub>O<sub>7</sub>(OH)<sub>2</sub>·2H<sub>2</sub>O,<sup>87</sup> or NaV<sub>3</sub>O<sub>8</sub>·1.5H<sub>2</sub>O<sup>88</sup> were investigated as intercalation pseudocapacitive materials for high-power energy storage due to their crystal structure characterized by an open framework. In the case of Ca<sub>0.24</sub>V<sub>2</sub>O<sub>5</sub>·0.83H<sub>2</sub>O, Alshareef and co-authors proved that the high diffusion coefficient as a result of the large size of the cavity and the high electronic conductivity positively influenced the intercalation pseudocapacitive behavior of the material, resulting in a high rate capability and cycling stability.

Regarding practical applications of pseudocapacitive materials, the mass loading is a critical parameter. How to achieve high rate capability in thick electrodes has become an important question. Tolbert *et al.* reported a pseudocapacitive behavior in mesoporous MoS<sub>2</sub> thin films, which was related to the ordered porosity and the iso-oriented crystal structure.<sup>55</sup> To further explore the application of such pseudocapacitive materials in thick film electrodes, the authors used precycling steps to convert the 2H phase of MoS<sub>2</sub> into the more conductive 1T phase, which was then processed into a carbon fiber-based highly conductive porous electrode. Synergistic combination of crystal structure engineering with a conductive carbon-based mesh electrode structure successfully preserved the good rate performance in the thicker electrodes.<sup>89</sup>



## 5. Recent advancements in non-aqueous LICs/NICs

Despite the attractive merits of non-aqueous hybrid capacitors mentioned above, the relatively low charge storage capability of the capacitive cathode and the kinetics mismatch between the anode and the cathode remain the main challenges.

### 5.1. Cathode

Activated carbon (AC) is commonly selected as the cathode material in non-aqueous hybrid capacitor research due to its ultrahigh surface area and high electronic conductivity.<sup>60,90–93</sup> However, the energy storage ability of an AC cathode can currently only achieve a capacity of  $\sim 50 \text{ mA h g}^{-1}$ , which is insufficient compared with anode materials. In non-aqueous hybrid capacitors utilizing AC as the cathode, this problem is addressed by choosing the mass loading of the cathode two to four times higher than that of the anode. Therefore, to increase capacity and capacitance, the development of cathode materials with higher charge storage capacity represents a major task.

Carbon sheets with 3D architectures showed great promise for high-performance capacitive cathodes for non-aqueous

hybrid capacitors. Chen *et al.* reported a graphene-based three-dimensional porous carbon material (3DGraphene) with high surface area ( $\sim 3355 \text{ m}^2 \text{ g}^{-1}$ ) as positive electrode material in a non-aqueous hybrid capacitor.<sup>94</sup> Due to the extremely large amount of adsorption sites as well as the high electronic conductivity of this 3DGraphene, the material delivered a high specific capacitance ranging from 148 to  $187 \text{ F g}^{-1}$  at current densities of  $0.096$  to  $9.55 \text{ mA cm}^{-2}$ . Metal–organic frameworks (MOFs) can also be utilized as a platform for the synthesis of porous carbon materials. Yan and Zhang *et al.* reported the synthesis of capacitor-like porous carbon polyhedra and a battery-like  $\text{MoS}_2$ -ZIF composite using the same precursor of MOF (polyhedral ZIF-8).<sup>95</sup> Due to the porous structure of the parent ZIF-8, the obtained porous carbon polyhedra showed a continuous 3D porous network with a high surface area and a well-controlled pore size, resulting in the typical capacitive behavior with a large specific capacitance. Despite the merits that such 3D frameworks of carbon possess, a simple synthesis methodology is important to bridge the gap between laboratory and industry. Yan *et al.* presented a promising approach for the large-scale preparation of highly porous carbon material through direct calcination of sodium citrate.<sup>96</sup> This method works without the use of any additional carbon source, template, or catalyst. Through first-principle calculations, the



**Fig. 9** Example of an advanced non-aqueous hybrid capacitor: A LIC using B and N dual-doped carbon nanofibers as both the cathode and the anode. (a) Schematic of the fabrication process of BNC. (b) Schematic of the charge-storage mechanisms of BNC//BNC LIC. Reproduced with permission from ref. 100. Copyright 2017, Wiley-VCH.



authors verified that the large interlayer spacing and the curved particle shape facilitated the high rate ion storage.

N-Doping is an efficient way to enhance the charge storage capability of carbon materials, because the nitrogen atoms substituting carbon atoms in the graphite matrix are electron donors and promote n-type conductivity.<sup>97,98</sup> Cai and Zhang *et al.* successfully prepared nitrogen-doped activated carbon through a one-step process.<sup>99</sup> The obtained N-doped AC exhibited high surface areas of up to  $2900 \text{ m}^2 \text{ g}^{-1}$ , resulting in a high specific capacity of  $129 \text{ mA h g}^{-1}$  at  $1.6 \text{ mA cm}^{-2}$ . The authors found that the electrochemical performances of the N-doped AC materials were dependent on the N content. AC with 2.97 wt% N (NAC 400) showed better performance than the one with 3.98 wt% N (NAC 600). The authors attributed this observation to the larger concentration of oxidized pyridinic N found in NAC-600 and the relatively lower surface area of NAC-600. The work by Xia, Yu and co-authors demonstrated that the dual doping of B and N in carbon nanofibers exhibited greatly improved electrochemical performance when applied as cathode for LICs (Fig. 9).<sup>100</sup> They proved that the B, N dual doping significantly enhanced the surface area, enlarged the carbon interlayer distance, increased the number of the active sites, and improved the electrode kinetics of the carbon nanofibers. The B,N dual doped carbon nanofibers could not only be used as the cathode, but also exhibited high performance as the anode. A symmetric LIC employing the carbon nanofibers as both binder-free cathode and anode achieved a high energy density of  $104 \text{ W h kg}^{-1}$  at  $22\,500 \text{ W kg}^{-1}$ .

Highly porous carbon materials derived from natural resources are regarded as an eco-friendly way to prepare high-performance cathodes for non-aqueous hybrid capacitors. For instance, Lee and co-workers successfully prepared high-

surface-area activated carbon derived from cinnamon sticks. The bio-inspired AC provided a high cycling stability and an average discharge capacity of  $60 \text{ mA h g}^{-1}$  at a current rate of  $400 \text{ mA g}^{-1}$ .<sup>101</sup> Mitlin *et al.* reported a porous carbon material with sheet-like morphology derived from peanut shells. Based on a tailored synthesis strategy combining a hydrothermal and a chemical activation process, a carbon material with macroscopically open sheet-like morphology was prepared, delivering an extremely high capacity of  $73 \text{ mA h g}^{-1}$  at  $10.24 \text{ mA cm}^{-2}$ .<sup>102</sup>

It has to be noticed that carbon materials with EDLC mechanisms usually have high surface areas, which inevitably result in relatively low volumetric energy densities. In response to this issue, efforts have been made to find intercalation pseudocapacitive cathode materials. For instance, nanoporous  $\text{LiMn}_2\text{O}_4$  with pseudocapacitive characteristics has been proved to show great promise as a fast-charging cathode material.<sup>103</sup>

## 5.2. Anode

The main challenge on the anode side for the fabrication of high-performance non-aqueous hybrid capacitors is to find a material that is able to match the fast kinetics of the cathode. A lot of pseudocapacitive materials have been investigated as the anode in both LICs and NICs. In a recent study, we reported a general method for the synthesis of structurally disordered nanoparticles@reduced graphene oxide (rGO) composite aerogels, which were directly used as high-power pseudocapacitive anodes for LICs.<sup>104</sup> To clarify the effect of the degree of disorder, control samples of crystalline nanoparticles with similar particle size were prepared. A schematic of the preparation route to crystalline and structurally disordered composite aerogels is shown in Fig. 10. Cyclic voltammetry analysis



**Fig. 10** Schematic of the preparation of crystalline and structurally disordered composite aerogels. Reproduced with permission from ref. 104. Copyright 2018, American Chemical Society.







**Fig. 11** (a) Voltammetric response of G-ZVO with separation between total current and pseudocapacitive current at a scan rate of  $1 \text{ mV s}^{-1}$ . (b) Schematic of the G-ZVO||OM-HSTC hybrid device. (c) Long-term cycling stability performance of the G@ZVO||OM-HSTC hybrid device at  $1.5 \text{ mA cm}^{-2}$ . Reproduced with permission from ref. 106. Copyright 2018, Wiley-VCH.

and density functional theory calculations indicated that the structurally disordered samples exhibited a smaller volume expansion during  $\text{Li}^+$  insertion and larger pseudocapacitive effects than the crystalline reference samples. As a result, structurally disordered samples showed much better Li-ion storage performances. Combined with commercial activated

carbon as the cathode, the structurally disordered  $\text{Co}_x\text{Fe}_y\text{O}_z\text{@rGO}$  delivered higher energy densities than crystalline  $\text{CoFe}_2\text{O}_4\text{@rGO}$  when tested at the same power densities. These results indicate that nanomaterials with a low crystallinity and a large number of defects might be interesting for high-power anodes for LICs.

**Table 2** Summary of the performance of recently published LIC and NIC devices

Type of device	Anode/cathode [reference]	Potential range (V)	Energy density ( $\text{Wh kg}^{-1}$ )	Power density ( $\text{W kg}^{-1}$ )	Cycling life/capacity retention	Publication year
LIC	3D inverse opal-structured $\text{FeS-QDs@N-doped carbon//AC}^{107}$	0.5–3.4	72.2	9280	5000 cycles/91%	2019
LIC	$\text{Co}_3\text{ZnC@N-doped carbon//microporous carbon}^{108}$	1.0–4.5	15.2	10 300	1000 cycles/80%	2018
LIC	$\text{P-WO}_{3-x}\text{@N-doped carbon//AC}^{109}$	0.01–4.3	113.84	11 946	6000 cycles/90.7%	2018
LIC	$\text{Li}_3\text{VO}_4\text{@N-doped carbon//AC}^{110}$	1.0–4.0	24.4	11 020	1500/—	2017
LIC	$\text{MoS}_2\text{@RGO//AC}^{111}$	0–4.0	45.3	40 000	10 000 cycles/ $\sim$ 80%	2017
LIC	B and N dual-doped carbon nanofibers (BNC)//BNC <sup>100</sup>	0.02–4.5	104	22 500	5000 cycles/81%	2017
LIC	$\text{Si@C//nitrogen-doped AC}^{99}$	2.0–4.5	141	30 127	8000 cycles/76.3%	2016
LIC	$\text{Si/C//egg white-derived AC}^{112}$	2.0–4.5	147	29 893	15 000 cycles/79.2%	2016
LIC	$\text{TiC//porous nitrogen-doped carbon}^{113}$	0–4.5	23.4	67 500	5000 cycles/ $\sim$ 82%	2016
LIC	$\text{3D VN-RGO//porous carbon nanorods}^{114}$	0–4.0	64	10 000	1000 cycles/83%	2015
LIC	$\text{MnO@graphene//porous N-doped carbon}^{63}$	1.0–4.0	83.25	25 000	3000 cycles/76%	2015
LIC	Anatase $\text{TiO}_2\text{@rGO//AC}^{115}$	1.0–3.0	42	8000	10 000 cycles/80%	2013
NIC	$\text{MoSe}_2\text{@graphene//AC}^{116}$	0.5–3.0	43	6688	5000 cycles/81%	2018
NIC	$\text{Gr-Nb}_2\text{O}_5\text{//AC}^{117}$	1.0–4.3	80.1	5330	—	2018
NIC	3D framework carbon (3DF)//3DFC-derived nanoporous carbon <sup>96</sup>	0–4.0	67	20 000	1000 cycles/95.3%	2018
NIC	$\text{TiO}_2\text{@C//MOFs-derived 3D nanoporous carbon}^{118}$	1.0–4.0	61.8	25 000	10 000 cycles/90%	2018
NIC	$\text{TiO}_2\text{@CNT@C//biomass-derived carbon}^{119}$	1.0–4.0	37.9	12 400	5000 cycles/85.3%	2017
NIC	$\text{NaTi}_2(\text{PO}_4)_3\text{@graphene//2D graphene nanosheets}^{120}$	0–3.0	80	8000	75 000 cycles/90%	2017
NIC	Peanut shell nanosheet carbon/low surface area peanut shell ordered carbon <sup>102</sup>	1.5–3.5	50	16 500	10 000 cycles/72%	2015





For the applications in NICs, Kim *et al.* introduced  $\text{SnS}_2$ @graphene-CNT aerogels with pseudocapacitive behavior, which turned out to result in an energy density of  $26.9 \text{ W h kg}^{-1}$  at power densities of  $6053 \text{ W kg}^{-1}$ , respectively.<sup>105</sup> We also reported a reduced graphene oxide- $\text{Zn}_{0.25}\text{V}_2\text{O}_5 \cdot n\text{H}_2\text{O}$  nanobelt composite (G-ZVO) as a pseudocapacitive anode for NICs.<sup>106</sup> Due to the large interplanar spacings in the crystal structure of  $\text{Zn}_{0.25}\text{V}_2\text{O}_5 \cdot n\text{H}_2\text{O}$  and the high electronic conductivity of the composite, the active material was characterized by high pseudocapacitive contributions of up to 85% at  $1 \text{ mV}^{-1}$  (Fig. 11a) and high rate capability. Pairing it with a hard salt-templated, ordered mesoporous carbon as a high-performance capacitive cathode (OM-HSTC) resulted in a Na-ion hybrid capacitor (Fig. 11b), which delivered a good cycling performance ( $31.7 \text{ W h kg}^{-1}$  (*i.e.*, 87%) retained after 2000 cycles at  $1.5 \text{ mA cm}^{-2}$ , Fig. 11c). The study supports the idea that layer-structured materials with large interlayer spacings might be suitable for NICs involving intercalation mechanisms. To get a general idea of the achievements in the field of non-aqueous hybrid capacitors, some recently published results on LICs and NICs are summarized in Table 2.

## 6. Conclusions and outlook

The goal of this minireview was, on the one hand, to introduce the fundamentals of pseudocapacitance and non-aqueous hybrid capacitors for the non-expert readers and, on the other hand, to provide a general overview of the recent developments in pseudocapacitive materials and Li-ion/Na-ion hybrid capacitors for the broad scientific community interested in novel energy storage devices. Pseudocapacitive materials and non-aqueous hybrid capacitors have shown that they hold great potential to contribute towards fast-charging technologies. However, we are far away from exploiting their full potential. For future advancements, several critical issues have to be considered.

Compared to the progress made in high-rate pseudocapacitive anode materials, the exploration of pseudocapacitive cathode materials with comparable performance is clearly lagging behind. However, from the perspective of the full-cell level, this problem has to be solved. The concept of non-aqueous hybrid capacitors was put forward to push the limit of energy density of electrical double layer capacitors. On the cathode side, the electrostatic double layer charge storage capability of activated carbon materials is dependent on the surface area. Unfortunately, high surface area inevitably results in low volumetric energy densities. To tackle this aspect, the development of advanced intercalation pseudocapacitive cathode materials might be a promising strategy.

In addition, most studies on pseudocapacitive materials are based on very thin electrodes and extremely small particle sizes. Despite that these studies provide instructive understanding of the basic science, they are practically not relevant.<sup>121</sup> The issue here is that increased electrode thickness will result in serious polarization, hence affecting the ion

storage kinetics. Furthermore, with very small particle sizes, unexpectedly poor efficiency and high self-discharge rate will arise. Crystal structure engineering and creation of interconnected porous architectures with continuous electronic and ionic conductive pathways might be possible solutions for these problems.

To bring pseudocapacitive materials closer to the application, we hereby suggest a few guidelines for electrode preparation: The electrodes should be with areal loadings of at least  $2\text{--}4 \text{ mg cm}^{-2}$ ; maximum porosity of 35% is acceptable; the content of inactive components like conductive carbon black and binder should not exceed 8%. Future studies taking these considerations into account will be especially beneficial for the practical application of pseudocapacitive materials.

The research area of pseudocapacitive materials and non-aqueous hybrid capacitors is quickly moving forward and it is not always straightforward to establish clear and universal definitions. To help to identify pseudocapacitive materials and to provide suggestions for the study of non-aqueous hybrid capacitors, we propose three key points: (i) The electrode preparation should be based on the considerations mentioned before; (ii) to be qualified as pseudocapacitive, a material should have a capacitive contribution of more than 70% at  $1 \text{ mV s}^{-1}$  based on CV analysis, intense and separated oxidative and reductive peaks should not be observed in CV curves,  $dQ/dV$  derived from charging–discharging–curves should be approaching a constant value and phase transformation measurements should be carried out and materials with phase transitions within the operating voltage window should not be called intercalation pseudocapacitive materials; and (iii) performances of non-aqueous hybrid capacitors should be tested at power densities  $>5000 \text{ W kg}^{-1}$ .

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

The authors acknowledge ETH Zurich and China Scholarship Council for financial support.

## Notes and references

- 1 J. W. Choi and D. Aurbach, *Nat. Rev. Mater.*, 2016, **1**, 16013–16028.
- 2 S. W. Kim, D. H. Seo, X. Ma, G. Ceder and K. Kang, *Adv. Energy Mater.*, 2012, **2**, 710–721.
- 3 M. Armand and J. M. Tarascon, *Nature*, 2008, **451**, 652–657.
- 4 D. Lin, Y. Liu and Y. Cui, *Nat. Nanotechnol.*, 2017, **12**, 194–206.
- 5 X.-B. Cheng, R. Zhang, C.-Z. Zhao and Q. Zhang, *Chem. Rev.*, 2017, **117**, 10403–10473.



- 6 X. Ji, K. T. Lee and L. F. Nazar, *Nat. Mater.*, 2009, **8**, 500–506.
- 7 P. G. Bruce, S. A. Freunberger, L. J. Hardwick and J.-M. Tarascon, *Nat. Mater.*, 2012, **11**, 19–29.
- 8 Z. Peng, S. A. Freunberger, Y. Chen and P. G. Bruce, *Science*, 2012, **337**, 563–566.
- 9 Y. Liu, Y. Zhu and Y. Cui, *Nat. Energy*, 2019, **4**, 540–550.
- 10 V. Augustyn, P. Simon and B. Dunn, *Energy Environ. Sci.*, 2014, **7**, 1597–1614.
- 11 K. Naoi, W. Naoi, S. Aoyagi, J.-I. Miyamoto and T. Kamino, *Acc. Chem. Res.*, 2012, **46**, 1075–1083.
- 12 K. Naoi, S. Ishimoto, J.-I. Miyamoto and W. Naoi, *Energy Environ. Sci.*, 2012, **5**, 9363–9373.
- 13 T. Brousse, D. Bélanger and J. W. Long, *J. Electrochem. Soc.*, 2015, **162**, A5185–A5189.
- 14 K. Naoi and Y. Nagano, in *Supercapacitors*, 2013, Wiley-VCH Verlag GmbH & Co. KGaA, pp. 239–256.
- 15 K. Naoi, *Fuel Cells*, 2010, **10**, 825–833.
- 16 N. Yabuuchi, K. Kubota, M. Dahbi and S. Komaba, *Chem. Rev.*, 2014, **114**, 11636–11682.
- 17 L. P. Wang, L. Yu, X. Wang, M. Srinivasan and Z. J. Xu, *J. Mater. Chem. A*, 2015, **3**, 9353–9378.
- 18 M. D. Slater, D. Kim, E. Lee and C. S. Johnson, *Adv. Funct. Mater.*, 2013, **23**, 947–958.
- 19 S. Y. Hong, Y. Kim, Y. Park, A. Choi, N.-S. Choi and K. T. Lee, *Energy Environ. Sci.*, 2013, **6**, 2067–2081.
- 20 J. Liu, J. Wang, C. Xu, H. Jiang, C. Li, L. Zhang, J. Lin and Z. X. Shen, *Adv. Sci.*, 2018, **5**, 1700322.
- 21 J. R. Miller and P. Simon, *Science*, 2008, **321**, 651–652.
- 22 P. J. Hall, M. Mirzaei, S. I. Fletcher, F. B. Sillars, A. J. Rennie, G. O. Shitta-Bey, G. Wilson, A. Cruden and R. Carter, *Energy Environ. Sci.*, 2010, **3**, 1238–1251.
- 23 B. E. Conway, *J. Electrochem. Soc.*, 1991, **138**, 1539–1548.
- 24 B. Conway, V. Birss and J. Wojtowicz, *J. Power Sources*, 1997, **66**, 1–14.
- 25 Y. Gogotsi and R. M. Penner, *ACS Nano*, 2018, **12**, 2081–2083.
- 26 B. Conway, *Electrochim. Acta*, 1993, **38**, 1249–1258.
- 27 S. Trasatti and G. Buzzanca, *J. Electroanal. Chem. Interfacial Electrochem.*, 1971, **29**, A1–A5.
- 28 J. Zheng, P. Cygan and T. Jow, *J. Electrochem. Soc.*, 1995, **142**, 2699–2703.
- 29 J. W. Long, K. E. Swider, C. I. Merzbacher and D. R. Rolison, *Langmuir*, 1999, **15**, 780–785.
- 30 V. Augustyn, J. Come, M. A. Lowe, J. W. Kim, P.-L. Taberna, S. H. Tolbert, H. D. Abruña, P. Simon and B. Dunn, *Nat. Mater.*, 2013, **12**, 518–522.
- 31 J. W. Kim, V. Augustyn and B. Dunn, *Adv. Energy Mater.*, 2012, **2**, 141–148.
- 32 C.-P. Liu, F. Zhou and V. Ozolins, in *APS Meeting Abstracts*, 2012.
- 33 A. Djire, P. Pande, A. Deb, J. B. Siegel, O. T. Ajenifujah, L. He, A. E. Sleightholme, P. G. Rasmussen and L. T. Thompson, *Nano Energy*, 2019, **60**, 72–81.
- 34 A. Djire, J. B. Siegel, O. Ajenifujah, L. He and L. T. Thompson, *Nano Energy*, 2018, **51**, 122–127.
- 35 A. Djire, O. Ajenifujah and L. T. Thompson, *Batteries Supercaps*, 2018, **1**, 171–175.
- 36 R. E. Ruther, C.-N. Sun, A. Holliday, S. Cheng, F. M. Delnick, T. A. Zawodzinski and J. Nanda, *J. Electrochem. Soc.*, 2017, **164**, A277–A283.
- 37 N. Jabeen, Q. Xia, S. V. Savilov, S. M. Aldoshin, Y. Yu and H. Xia, *ACS Appl. Mater. Interfaces*, 2016, **8**, 33732–33740.
- 38 A. Yamada, S.-C. Chung and K. Hinokuma, *J. Electrochem. Soc.*, 2001, **148**, A224–A229.
- 39 G. Z. Chen, *Prog. Nat. Sci.*, 2013, **23**, 245–255.
- 40 J. Come, P.-L. Taberna, S. Hamelet, C. Masquelier and P. Simon, *J. Electrochem. Soc.*, 2011, **158**, A1090–A1093.
- 41 H. Y. Lee and J. B. Goodenough, *J. Solid State Chem.*, 1999, **144**, 220–223.
- 42 J. Wang, J. Polleux, J. Lim and B. Dunn, *J. Phys. Chem. C*, 2007, **111**, 14925–14931.
- 43 P. Zanello, *Inorganic Electrochemistry: Theory, Practice and Application*, Royal Society of Chemistry, 2011.
- 44 S. Ardizzzone, G. Fregonara and S. Trasatti, *Electrochim. Acta*, 1990, **35**, 263–267.
- 45 Z. Le, F. Liu, P. Nie, X. Li, X. Liu, Z. Bian, G. Chen, H. B. Wu and Y. Lu, *ACS Nano*, 2017, **11**, 2952–2960.
- 46 C. Costentin, T. R. Porter and J.-M. Savéant, *ACS Appl. Mater. Interfaces*, 2017, **9**, 8649–8658.
- 47 B. E. Conway, *Electrochemical Supercapacitors: Scientific Fundamentals and Technological Applications*, Kluwer-Academic, 1999.
- 48 J. B. Cook, T. C. Lin, H.-S. Kim, A. Siordia, B. S. Dunn and S. H. Tolbert, *ACS Nano*, 2019, **13**, 1223–1231.
- 49 H.-S. Kim, J. B. Cook, S. H. Tolbert and B. Dunn, *J. Electrochem. Soc.*, 2015, **162**, A5083–A5090.
- 50 G. A. Muller, J. B. Cook, H.-S. Kim, S. H. Tolbert and B. Dunn, *Nano Lett.*, 2015, **15**, 1911–1917.
- 51 P. Simon, Y. Gogotsi and B. Dunn, *Science*, 2014, **343**, 1210–1211.
- 52 R. Chen, M. Knapp, M. Yavuz, S. Ren, R. Witte, R. Heinzmann, H. Hahn, H. Ehrenberg and S. Indris, *Phys. Chem. Chem. Phys.*, 2015, **17**, 1482–1488.
- 53 J. Li, X. Yuan, C. Lin, Y. Yang, L. Xu, X. Du, J. Xie, J. Lin and J. Sun, *Adv. Energy Mater.*, 2017, **7**, 1602725.
- 54 M. Okubo, E. Hosono, J. Kim, M. Enomoto, N. Kojima, T. Kudo, H. Zhou and I. Honma, *J. Am. Chem. Soc.*, 2007, **129**, 7444–7452.
- 55 J. B. Cook, H. S. Kim, Y. Yan, J. S. Ko, S. Robbenolt, B. Dunn and S. H. Tolbert, *Adv. Energy Mater.*, 2016, **6**, 1501937.
- 56 T. Brezesinski, J. Wang, S. H. Tolbert and B. Dunn, *Nat. Mater.*, 2010, **9**, 146–151.
- 57 Q. Wang, Z. Wen and J. Li, *Adv. Funct. Mater.*, 2006, **16**, 2141–2146.
- 58 H. Li, L. Peng, Y. Zhu, X. Zhang and G. Yu, *Nano Lett.*, 2016, **16**, 5938–5943.
- 59 D. P. Dubal, O. Ayyad, V. Ruiz and P. Gomez-Romero, *Chem. Soc. Rev.*, 2015, **44**, 1777–1790.
- 60 E. Lim, C. Jo, M. S. Kim, M. H. Kim, J. Chun, H. Kim, J. Park, K. C. Roh, K. Kang and S. Yoon, *Adv. Funct. Mater.*, 2016, **26**, 3711–3719.



- 61 T. Yuan, Y. Jiang, W. Sun, B. Xiang, Y. Li, M. Yan, B. Xu and S. Dou, *Adv. Funct. Mater.*, 2016, **26**, 2198–2206.
- 62 X. Chen, L.-P. Lv, W. Sun, Y. Hu, X. Tao and Y. Wang, *J. Mater. Chem. A*, 2018, **6**, 13705–13716.
- 63 M. Yang, Y. Zhong, J. Ren, X. Zhou, J. Wei and Z. Zhou, *Adv. Energy Mater.*, 2015, **5**, 1500550.
- 64 C. Zhao, C. Yu, M. Zhang, H. Huang, S. Li, X. Han, Z. Liu, J. Yang, W. Xiao, J. Liang, X. Sun and J. Qiu, *Adv. Energy Mater.*, 2017, **7**, 1602880.
- 65 X.-Y. Yu, L. Yu, L. Shen, X. Song, H. Chen and X. W. Lou, *Adv. Funct. Mater.*, 2014, **24**, 7440–7446.
- 66 H. Hu, B. Y. Guan and X. W. D. Lou, *Chem*, 2016, **1**, 102–113.
- 67 L. Shen, L. Yu, H. B. Wu, X.-Y. Yu, X. Zhang and X. W. Lou, *Nat. Commun.*, 2015, **6**, 6694.
- 68 L. Yu, L. Zhang, H. B. Wu and X. W. Lou, *Angew. Chem., Int. Ed.*, 2014, **53**, 3711–3714.
- 69 T. Zhai, S. Sun, X. Liu, C. Liang, G. Wang and H. Xia, *Adv. Mater.*, 2018, **30**, 1706640.
- 70 L. Chen, X. Jiang, N. Wang, J. Yue, Y. Qian and J. Yang, *Adv. Sci.*, 2015, **2**, 1500090.
- 71 A. A. Lubimtsev, P. R. C. Kent, B. G. Sumpter and P. Ganesh, *J. Mater. Chem. A*, 2013, **1**, 14951–14956.
- 72 L. Kong, C. Zhang, J. Wang, W. Qiao, L. Ling and D. Long, *ACS Nano*, 2015, **9**, 11200–11208.
- 73 S. Lou, X. Cheng, L. Wang, J. Gao, Q. Li, Y. Ma, Y. Gao, P. Zuo, C. Du and G. Yin, *J. Power Sources*, 2017, **361**, 80–86.
- 74 B. Deng, T. Lei, W. Zhu, L. Xiao and J. Liu, *Adv. Funct. Mater.*, 2018, **28**, 1704330.
- 75 H.-S. Kim, J. B. Cook, H. Lin, J. S. Ko, S. H. Tolbert, V. Ozolins and B. Dunn, *Nat. Mater.*, 2017, **16**, 454–460.
- 76 Y. Zhu, L. Peng, D. Chen and G. Yu, *Nano Lett.*, 2015, **16**, 742–747.
- 77 M. Naguib, M. Kurtoglu, V. Presser, J. Lu, J. Niu, M. Heon, L. Hultman, Y. Gogotsi and M. W. Barsoum, *Adv. Mater.*, 2011, **23**, 4248–4253.
- 78 M. Naguib, O. Mashtalir, J. Carle, V. Presser, J. Lu, L. Hultman, Y. Gogotsi and M. W. Barsoum, *ACS Nano*, 2012, **6**, 1322–1331.
- 79 O. Mashtalir, M. Naguib, V. N. Mochalin, Y. Dall'agnese, M. Heon, M. W. Barsoum and Y. Gogotsi, *Nat. Commun.*, 2013, **4**, 1716.
- 80 M. R. Lukatskaya, O. Mashtalir, C. E. Ren, Y. Dall'agnese, P. Rozier, P. L. Taberna, M. Naguib, P. Simon, M. W. Barsoum and Y. Gogotsi, *Science*, 2013, **341**, 1502–1505.
- 81 X. Wang, S. Kajiyama, H. Iinuma, E. Hosono, S. Oro, I. Moriguchi, M. Okubo and A. Yamada, *Nat. Commun.*, 2015, **6**, 6544.
- 82 S. Kajiyama, L. Szabova, K. Sodeyama, H. Iinuma, R. Morita, K. Gotoh, Y. Tateyama, M. Okubo and A. Yamada, *ACS Nano*, 2016, **10**, 3334–3341.
- 83 Q. Tang, Z. Zhou and P. Shen, *J. Am. Chem. Soc.*, 2012, **134**, 16909–16916.
- 84 A. Djire, A. Bos, J. Liu, H. Zhang, E. M. Miller and N. R. Neale, *ACS Appl. Nano Mater.*, 2019, **2**, 2785–2795.
- 85 C. Xia, J. Guo, P. Li, X. Zhang and H. N. Alshareef, *Angew. Chem., Int. Ed.*, 2018, **57**, 3943–3948.
- 86 F. Ming, H. Liang, Y. Lei, S. Kandambeth, M. Eddaoudi and H. N. Alshareef, *ACS Energy Lett.*, 2018, **3**, 2602–2609.
- 87 C. Xia, J. Guo, Y. Lei, H. Liang, C. Zhao and H. N. Alshareef, *Adv. Mater.*, 2018, **30**, 1705580.
- 88 F. Wan, L. Zhang, X. Dai, X. Wang, Z. Niu and J. Chen, *Nat. Commun.*, 2018, **9**, 1656.
- 89 J. B. Cook, H. S. Kim, T. C. Lin, C. H. Lai, B. Dunn and S. H. Tolbert, *Adv. Energy Mater.*, 2017, **7**, 1601283.
- 90 E. Lim, C. Jo, H. Kim, M.-H. Kim, Y. Mun, J. Chun, Y. Ye, J. Hwang, K.-S. Ha and K. C. Roh, *ACS Nano*, 2015, **9**, 7497–7505.
- 91 H. Wang, C. Zhu, D. Chao, Q. Yan and H. J. Fan, *Adv. Mater.*, 2017, **29**, 1702093.
- 92 Y. Ma, H. Chang, M. Zhang and Y. Chen, *Adv. Mater.*, 2015, **27**, 5296–5308.
- 93 J. Luo, W. Zhang, H. Yuan, C. Jin, L. Zhang, H. Huang, C. Liang, Y. Xia, J. Zhang and Y. Gan, *ACS Nano*, 2017, **11**, 2459–2469.
- 94 F. Zhang, T. Zhang, X. Yang, L. Zhang, K. Leng, Y. Huang and Y. Chen, *Energy Environ. Sci.*, 2013, **6**, 1623–1632.
- 95 R. Wang, D. Jin, Y. Zhang, S. Wang, J. Lang, X. Yan and L. Zhang, *J. Mater. Chem. A*, 2017, **5**, 292–302.
- 96 B. Yang, J. Chen, S. Lei, R. Guo, H. Li, S. Shi and X. Yan, *Adv. Energy Mater.*, 2018, **8**, 1702409.
- 97 Z. R. Ismagilov, A. E. Shalagina, O. Y. Podyacheva, A. V. Ischenko, L. S. Kibis, A. I. Boronin, Y. A. Chesalov, D. I. Kochubey, A. I. Romanenko and O. B. Anikeeva, *Carbon*, 2009, **47**, 1922–1929.
- 98 M. Terrones, P. Ajayan, F. Banhart, X. Blase, D. Carroll, J.-C. Charlier, R. Czerw, B. Foley, N. Grobert and R. Kamalakaran, *Appl. Phys. A*, 2002, **74**, 355–361.
- 99 B. Li, F. Dai, Q. Xiao, L. Yang, J. Shen, C. Zhang and M. Cai, *Energy Environ. Sci.*, 2016, **9**, 102–106.
- 100 Q. Xia, H. Yang, M. Wang, M. Yang, Q. Guo, L. Wan, H. Xia and Y. Yu, *Adv. Energy Mater.*, 2017, **7**, 1701336.
- 101 R. Thangavel, K. Kaliyappan, K. Kang, X. Sun and Y. S. Lee, *Adv. Energy Mater.*, 2016, **6**, 1502199.
- 102 J. Ding, H. Wang, Z. Li, K. Cui, D. Karpuzov, X. Tan, A. Kohandehghan and D. Mitlin, *Energy Environ. Sci.*, 2015, **8**, 941–955.
- 103 B. K. Lesel, J. B. Cook, Y. Yan, T. C. Lin and S. H. Tolbert, *ACS Energy Lett.*, 2017, **2**, 2293–2298.
- 104 H. Huang, X. Wang, E. Tervoort, G. Zeng, T. Liu, X. Chen, A. Sologubenko and M. Niederberger, *ACS Nano*, 2018, **12**, 2753–2763.
- 105 J. Cui, S. Yao, Z. Lu, J. Q. Huang, W. G. Chong, F. Ciucci and J. K. Kim, *Adv. Energy Mater.*, 2018, **8**, 1702488.
- 106 H. Huang, D. Kundu, R. Yan, E. Tervoort, X. Chen, L. Pan, M. Oschatz, M. Antonietti and M. Niederberger, *Adv. Energy Mater.*, 2018, **8**, 1802800.
- 107 X. Hu, Y. Liu, J. Chen, J. Jia, H. Zhan and Z. Wen, *J. Mater. Chem. A*, 2019, **7**, 1138–1148.



- 108 G. Zhu, T. Chen, L. Wang, L. Ma, Y. Hu, R. Chen, Y. Wang, C. Wang, W. Yan and Z. Tie, *Energy Storage Mater.*, 2018, **14**, 246–252.
- 109 J. Xu, Z. Liao, J. Zhang, B. Gao, P. K. Chu and K. Huo, *J. Mater. Chem. A*, 2018, **6**, 6916–6921.
- 110 L. Shen, H. Lv, S. Chen, P. Kopold, P. A. Van Aken, X. Wu, J. Maier and Y. Yu, *Adv. Mater.*, 2017, **29**, 1700142.
- 111 R. Wang, S. Wang, D. Jin, Y. Zhang, Y. Cai, J. Ma and L. Zhang, *Energy Storage Mater.*, 2017, **9**, 195–205.
- 112 B. Li, F. Dai, Q. Xiao, L. Yang, J. Shen, C. Zhang and M. Cai, *Adv. Energy Mater.*, 2016, **6**, 1600802.
- 113 H. Wang, Y. Zhang, H. Ang, Y. Zhang, H. T. Tan, Y. Zhang, Y. Guo, J. B. Franklin, X. L. Wu and M. Srinivasan, *Adv. Funct. Mater.*, 2016, **26**, 3082–3093.
- 114 R. Wang, J. Lang, P. Zhang, Z. Lin and X. Yan, *Adv. Funct. Mater.*, 2015, **25**, 2270–2278.
- 115 H. Kim, M. Y. Cho, M. H. Kim, K. Y. Park, H. Gwon, Y. Lee, K. C. Roh and K. Kang, *Adv. Energy Mater.*, 2013, **3**, 1500–1506.
- 116 X. Zhao, W. Cai, Y. Yang, X. Song, Z. Neale, H.-E. Wang, J. Sui and G. Cao, *Nano Energy*, 2018, **47**, 224–234.
- 117 X. Wang, Q. Li, L. Zhang, Z. Hu, L. Yu, T. Jiang, C. Lu, C. Yan, J. Sun and Z. Liu, *Adv. Mater.*, 2018, 1800963.
- 118 H. Li, J. Lang, S. Lei, J. Chen, K. Wang, L. Liu, T. Zhang, W. Liu and X. Yan, *Adv. Funct. Mater.*, 2018, **28**, 1800757.
- 119 Y. E. Zhu, L. Yang, J. Sheng, Y. Chen, H. Gu, J. Wei and Z. Zhou, *Adv. Energy Mater.*, 2017, **7**, 1701222.
- 120 R. Thangavel, B. Moorthy, D. K. Kim and Y. S. Lee, *Adv. Energy Mater.*, 2017, **7**, 1602654.
- 121 Y. Gogotsi and P. Simon, *Science*, 2011, **334**, 917–918.

