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Dual supramolecular chirogenesis based on platinum(II) metallotweezers†

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Optically active platinum(II) metallotweezers demonstrate both self-complexation and host-guest complexation capabilities, leading to two distinct supramolecular chirogenic signals in the visible region.

Chirality transfer from the molecular level to the supramolecular level is not only essential in life¹ but is relevant for catalytic, optoelectronic and spintronic applications in materials science.² Supramolecular chirogenesis³ represents an efficient way to express chirality in artificial systems, and involves non-covalent chiral recognition between the host and guest species. Pioneering work in this field has been performed by Inoue and co-workers on the basis of achiral zinc porphyrin tweezers.^{3a,c} Until now, supramolecular chirogenesis has extended to various artificial receptors such as macrocycles,⁴ cages,⁵ and helical foldamers.⁶ Despite the progress achieved, the chirogenic signals appear primarily in the ultra-violet or high-energy visible region due to the following two reasons. One is the lack of large π -conjugated chromophores on host/ guest structures, and the other is the low chirality transfer efficiency because of the remoteness of the chiral center from the host-guest complexation site.⁷ It is intriguing to shift supramolecular chirogenic signals to a low-energy absorption region, which would benefit circular polarized light detection/ emission and chiroptical switch applications.

Platinum(π)-based metallotweezers,⁸ with two cofacial square-planar pincers, represent an ideal candidate to attain this objective. When the cyclometalated Pt(π) pincers are kept

at the distance of 7 Å by a rigid spacer, the metallotweezers are capable of encapsulating a guest molecule into their cavity.⁹ The non-covalent host–guest complexation structure is stabilized by π – π stacking interactions between the Pt(II) pincers and the complementary guest (interplanar distance: ~ 3.5 Å). It can be endowed with fruitful photo-physical properties due to the spin–orbit coupling effect, leading to metal-to-ligand charge transfer (MLCT), ligand-to-ligand charge transfer (LLCT), and metal–metal-to-ligand charge transfer (MMLCT) transitions in the low-energy absorption region.¹⁰ We envisage that supramolecular chirogenic signals could potentially emerge for these electronic transitions, by incorporating a stereogenic center in the receptor of the metallotweezers.

In this study, we have designed the novel Pt(n) metallotweezers **1** (Scheme 1, see Scheme S1 in the ESI† for the synthetic procedure). Unlike previous chiral Pt(n) complexes in which the



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stereogenic center was embedded in the side chains,¹¹ herein, four (1R)-pinene units are fused to the Pt(II) terpyridine pincers in 1 to strengthen the supramolecular chirogenic signals.¹² Interestingly, 1 is prone to associate with other molecules of 1 to form a self-complexed structure (Scheme 1). A stereospecific twist is generated because of the stacking of $Pt(\pi)$ terpyridine [Pt(II)(N^N^N)] pincers, giving rise to the emergent chiroptical signals in the MLCT/LLCT absorption region. With the addition of compound 2 (Scheme 1) as the complementary guest, the self-complexation structure of 1 converts to the sandwich complex $1 \supset 2$. This consequently leads to the chirogenic signal in the MMLCT absorption region, thanks to the participation of Pt(II) - Pt(II) metal-metal interactions for the host-guest entity. Accordingly, dual supramolecular chirogenic signals form in the visible region, by taking advantage of the diverse complexation modes of the Pt(II) metallotweezers.

We first studied self-complexation behavior of the Pt(II) metallotweezers 1. In chloroform, protons H₂ and H₃ displayed downfield shifts upon varying the concentration from 0.20 mM to 20.0 mM (Fig. S12, ESI[†]). Concentration-dependent ¹H NMR measurements provided the self-association constant of $5.34 \times 10^2 \text{ M}^{-1} \ (\pm 34\%)$ for 1 (Fig. S13, ESI†). Generally, two possible self-aggregation modes exist for metallotweezers (Fig. S14, ESI⁺). One is the mutual stacking of pincer units in a quadruple manner.^{8e} The other is sandwiching of the spacer unit into the cavity of the complementary tweezers.^{8f} The latter mode is excluded in the case of metallotweezers 1, since the non-planar diphenylpyridine spacer is unable to be encapsulated into the cavity. The self-complexed structure of 1 was clarified via density functional theory (DFT) computations. Two quadruple stacking structures might form for 1, namely headto-tail and head-to-head binding modes for the neighbouring Pt(II)(N^N^N) pincers (Fig. S15, ESI⁺). For the optimized geometries, the head-to-head binding mode (Fig. 1a) featured a lower Gibbs free energy than that of the head-to-tail mode $(\Delta E = 0.974 \text{ kcal mol}^{-1}, \text{ Fig. S15, ESI}^{\dagger})$. The π - π distances between the Pt(II)(N^N^N) pincers are 3.29 Å, 3.47 Å, and 3.29 Å, respectively. Apparently, the pre-organization effect of the rigid diphenylpyridine spacer, together with the strong stacking tendency of the Pt(II)(N^NN) pincers, guarantees formation of the self-complexation structure for 1. Formation of the headto-head binding structure was further demonstrated via ¹H-¹H ROESY measurements. In particular, strong correlations exist between protons H_4/H_5 and H_4/H_6 (Fig. S16a, ESI[†]), which are absent in the ¹H-¹H COSY spectrum (Fig. S16b, ESI[†]) under the same conditions.

The spectroscopic properties were further examined for **1**. In dilute chloroform (c = 0.10 mM), only 8.9% of **1** existed in the complexed form, denoting the dominance of the monomeric state. The visible light absorbance ranged between 378 and 510 nm ($\varepsilon = 9.07 \times 10^3$ M⁻¹ cm⁻¹ at 395 nm, Fig. 1b), while the emission signal was centred at 560 nm (Fig. S17, ESI†). With reference to previous reports,⁹ these signals were assigned to the admixture of metal-to-ligand and ligand-to-ligand charge-transfer (MLCT/LLCT) transitions of the alkynyl Pt(π)(N^N^N) moiety. Upon switching the solvent from chloroform to



Fig. 1 (a) Optimized structure of the self-complexed dimer $\mathbf{1}_2$. (b) Absorption spectra of $\mathbf{1}$ in acetonitrile (red line) and chloroform (black line) (c = 0.10 mM). (c) Degree of α_{agg} for $\mathbf{1}$ monitored at 500 nm *versus* different temperatures (c = 0.10 mM in acetonitrile). Inset: van't Hoff plot fitting for the self-complexation process of $\mathbf{1}$. (d) Temperature-dependent CD spectra of $\mathbf{1}$ (c = 0.10 mM in acetonitrile). The arrows indicate the spectral change upon decreasing the temperature.

acetonitrile, the MLCT/LLCT emission signal declined for the intensity (Fig. S17, ESI[†]). Moreover, a low-energy shoulder band emerged for **1** ($\varepsilon = 1.70 \times 10^3 \text{ M}^{-1} \text{ cm}^{-1}$ at 510 nm, Fig. 1b). These phenomena suggest a stronger self-complexation capability in acetonitrile. This could be ascribed to the association of the dimer **1**₂ into the oligomeric species, considering that π - π stacking interactions are stronger in acetonitrile than those in chloroform.¹³ The conclusion is manifested by the broadened ¹H NMR peaks (Fig. S18, ESI[†]), together with the larger hydrodynamic diameter from DLS measurements (Fig. S19, ESI[†]).

The intensity of the low-energy band between 510 and 600 nm declined upon increasing the temperature to 353 K, with an isosbestic point at 463 nm (Fig. S20, ESI⁺). The results support the reversible conversion between the monomeric state at high temperature and the complexed state at low temperature. The equal K model¹⁴ was employed to fit the melting curves, acquired by plotting the absorption intensity changes at 500 nm *versus* the temperature (Fig. 1c). The $T_{\rm m}$ values [the temperature at which the degree of aggregation (α_{agg}) is 0.5] increased at higher monomer concentrations (Tm: 288 K at 2.50×10^{-5} M versus 306 K at 2.00×10^{-4} M, Fig. S20, ESI†). According to a modified van't Hoff plot (Fig. 1c, inset), the enthalpy (ΔH) and entropy (ΔS) values were determined to be -84.7 kJ mol⁻¹, and -206 J mol⁻¹ K⁻¹, respectively. Accordingly, this provided the self-complexed binding constant of $1.29 \times 10^4 \text{ M}^{-1}$ at 298 K, which is much higher than that in chloroform [5.34 \times 10² M⁻¹ (±34%)]. We rationalized that the higher self-complexion affinity in acetonitrile involved not only the dimeric stacking but the hierarchical association into oligomeric species.

Since the self-complexed structure adopts a head-to-head binding mode, it provides asymmetry by transferring chirality from the (1R)-pinenes to the Pt(II)(N^N^N) pincers. As can be seen, a weak Cotton effect below 419 nm exists for 1 at 353 K (Fig. 1d), supporting the origin of the molecular chirality from the (1R)-pinene units (Fig. S21, ESI^{\dagger}). Upon decreasing the temperature to 298 K, a bisignate CD signal appeared for 1 in the low-energy MLCT/LLCT absorption region (418-550 nm), with the positive maximum at 489 nm ($\Delta \varepsilon = 2.04 \text{ mol}^{-1} \text{ cm}^{-1}$) and the negative maximum at 437 nm ($\Delta \varepsilon = -1.89 \text{ mol}^{-1} \text{ cm}^{-1}$, Fig. 1d). Accordingly, the self-complexation of 1 prevents carbon-carbon and carbon-platinum bond rotations, exerting a crucial impact on the supramolecular chirogenic behavior. The conclusion is further demonstrated by the weakened Cotton effect in chloroform due to its weakened self-complexation tendency (Fig. S24, ESI[†]).

After elucidating the self-complexation properties of 1, we turned to its host-guest complexation behavior. According to electrospray ionization mass spectrometry, an m/z value of 2501.85 was observed for $1 \supset 2$, corresponding to $[1 + 2 + H]^+$. The color of 1 in acetonitrile solution changed from yellow to orange upon adding the charge-neutral guest 2 in an equivalent ratio (Fig. S26a, ESI[†]). A new absorption band emerged in the low energy region, ranging from 500 to 650 nm (Fig. 2a). This is a characteristic of metal-metal-to-ligand charge-transfer (MMLCT) transitions.¹⁰ Simultaneously, the MLCT/LLCT emission at 578 nm declined in its intensity, with a concomitant increase in the MMLCT emission band at 786 nm (Fig. S26b, ESI[†]). Depending on the molar ratio plot (Fig. S27, ESI[†]), the binding stoichiometry between the metallotweezer receptors 1 and guest 2 was 1:1. By fitting the collected UV/Vis absorbances at 510 nm, 525 nm and 540 nm, the "apparent" constant (K_d) value was determined to be 3.75 \times 10⁵ M⁻¹ $(\pm 27\%)$ in acetonitrile at 298 K (Fig. 2a, inset). Since the selfcomplexation of 1 was involved in the titration process, the "real" binding constant (K_a) value for $1 \supset 2$ was $6.96 \times 10^4 \text{ M}^{-1}$ $(\pm 13\%)$ (eqn (S11), ESI[†]).¹⁵ This value was higher than that in chloroform [K_a : 1.85 × 10³ M⁻¹ (±3.2%), Fig. S29, ESI[†]]. Although the K_a value of $1 \supset 2$ was smaller when observed *via* ¹H NMR titration experiments $[9.32 \times 10^3 \text{ M}^{-1} (\pm 28\%),$

Fig. 2 (a) UV/Vis absorbance changes of 1 at 298 K (c = 0.05 mM in CH $_3$ CN) upon the progressive addition of **2**. Inset: intensity changes in UV/ Vis absorbance at 510 nm, 525 nm, and 540 nm. The solid lines were obtained via a Matlab-based global analysis program. (b) Optimized structure of the host-guest complex $1 \supset 2$ on the basis of DFT calculations.

b

Fig. S30 and S31, ESI[†]], the value was also higher than that in chloroform [K_a : 2.54 × 10³ M⁻¹ (±24%), Fig. S32 and S33, ESI[†]].

The energy-minimized structure of complex $1 \supset 2$ was elucidated via DFT calculations. As expected, 2 is encapsulated into the cavity of the metallotweezers 1 to form a sandwiched complex (Fig. 2b). The inter-planar π -distances between 2 and the two Pt(II)(N^N^N) pincers on 1 are determined to be 3.30 Å and 3.21 Å, validating the presence of two-fold π - π stacking interactions. This conclusion was further validated via ¹H NMR experiments. Upon addition of one equivalent of 2 to 1, the ¹H NMR resonances of protons H₁ and H₂ shifted upfield $(\Delta \delta = -0.36 \text{ and } -0.58 \text{ ppm}, \text{ respectively})$, while protons H₄ varied from 8.43 ppm to 8.71 ppm because of the deshielding effect (Fig. S32, ESI[†]). Meanwhile, the Pt-Pt distances between 1 and 2 are 3.51 and 3.22 Å, respectively. This supports the existence of Pt(II)-Pt(II) interactions in complex $1 \supset 2$, and is highly consistent with the emergence of the MMLCT absorption and emission signals (Fig. S26a and S26b, ESI⁺).

The participation of two-fold $Pt(\pi)-Pt(\pi)$ and $\pi-\pi$ stacking interactions contributes to the high binding affinity for complex $1 \supset 2$. When the control compound 3 (Fig. S21, ESI, † inset) with the mono-nuclear Pt(II)(N^N^N) unit was employed as the host instead of 1, the K_a value for the resulting complex $3 \supset 2$ decreased to be 15.9 M^{-1} M^{-1} (Fig. S34 and S35, ESI⁺), two orders of magnitude lower than that of complex $1 \supset 2$. When the temperature was elevated to 353 K, the K_a value of complex $1 \supset 2$ in acetonitrile was determined to be $3.21 \times 10^4 \text{ M}^{-1} (\pm 5.9\%)$ (Fig. S38, ESI[†]), reaching one half of the value at 298 K. The high binding affinity of $1 \supset 2$ at elevated temperature is ascribed to the weakening of the 1 self-complexation strength upon heating. This buffers the decreased host-guest complexation, and thereby the strong complexation between 1 and 2 persists.

We further investigated the supramolecular chirogenic signal for the resulting host-guest complex. Upon the gradual addition of 2 into an acetonitrile solution of 1, the positive CD signal located at 497 nm became negative ($\Delta \varepsilon$: from 1.79 mol⁻¹ cm⁻¹ to -3.18 mol⁻¹ cm⁻¹, Fig. 3a and Fig. S39, ESI⁺). Meanwhile, the Cotton effect appeared in the MMLCT absorption region $(\Delta \varepsilon = +1.18 \text{ mol}^{-1} \text{ cm}^{-1} \text{ at } 570 \text{ nm})$. In stark contrast, a negligible Cotton effect was observed when employing 3 instead of 1 (Fig. S40a, ESI⁺), because of the weak complexation strength of complex $3 \supset 2$. Accordingly, metallotweezers/guest complexation



Fig. 3 (a) CD spectra of 1 and $1 \supset 2$ at 298 K (c = 0.10 mM for each compound in CH₃CN). (b) Energy-level diagram of $1 \supset 2$ via TD-DFT computations.

а 0.8

0.6

0.2

0.0

Absorband

Equivalence of 2

600

500 Wavelength / nm

with sufficient binding affinity is a prerequisite for supramolecular chirogenesis. The Cotton shape of $1 \supset 2$ was maintained at elevated temperatures because of the robust host-guest complexation, despite the decreased CD intensities (at 497 nm: $\Delta \varepsilon = |1.26| \text{ cm}^{-1} \text{ M}^{-1}$ at 353 K *versus* |3.18| cm⁻¹ M⁻¹ at 298 K, Fig. S41, ESI[†]).

The origin of the low-energy supramolecular chiroptical signals (ranging from 444 nm to 624 nm) was clarified using time-dependent density functional theory (TD-DFT) calculations. As shown in Fig. 3b, the electron density of the LUMO is distributed over the $Pt(\pi)(N^N^N)$ pincers of 1. Meanwhile, the electron density of the HOMO is mainly distributed on the $5d_{z^2}$ orbitals of the Pt(II) atoms in both 1 and 2. Accordingly, the theoretical Cotton effect in the low-energy absorption region is composed of HOMO \rightarrow LUMO transitions (composition: 76.6%) at 544 nm), belonging to metal-metal-to-ligand charge transfer (MMLCT) transitions. For most of previous host-guest systems, supramolecular chirogenic signals arose from the individual or conjoint non-covalent forces of metal-ligand coordination, hydrogen bonding, π - π stacking, and hydrophobic interactions. Complex $1 \supset 2$ represents a rare type of supramolecular chirogenic system with the involvement of Pt(II)-Pt(II) metalmetal interactions.¹⁶

In summary, metallotweezer 1 with optically active $Pt(\pi)(N^{N}N)$ pincers prefers to form a self-complexed structure *via* a head-tohead binding mode, leading to supramolecular chirogenic signals in the MLCT/LLCT absorption region. Furthermore, a metallotweezers/guest complex forms upon adding the complementary guest 2 into 1. This is accompanied by the formation of $Pt(\pi)$ – $Pt(\pi)$ metalmetal interactions, and thereby induces supramolecular chirogenesis in the MMLCT transition region. Therefore, a dual supramolecular chirogenic system in the visible region has been successfully constructed by taking advantage of the diverse complexation modes of $Pt(\pi)$ metallotweezers.

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Conflicts of interest

There are no conflicts to declare.

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