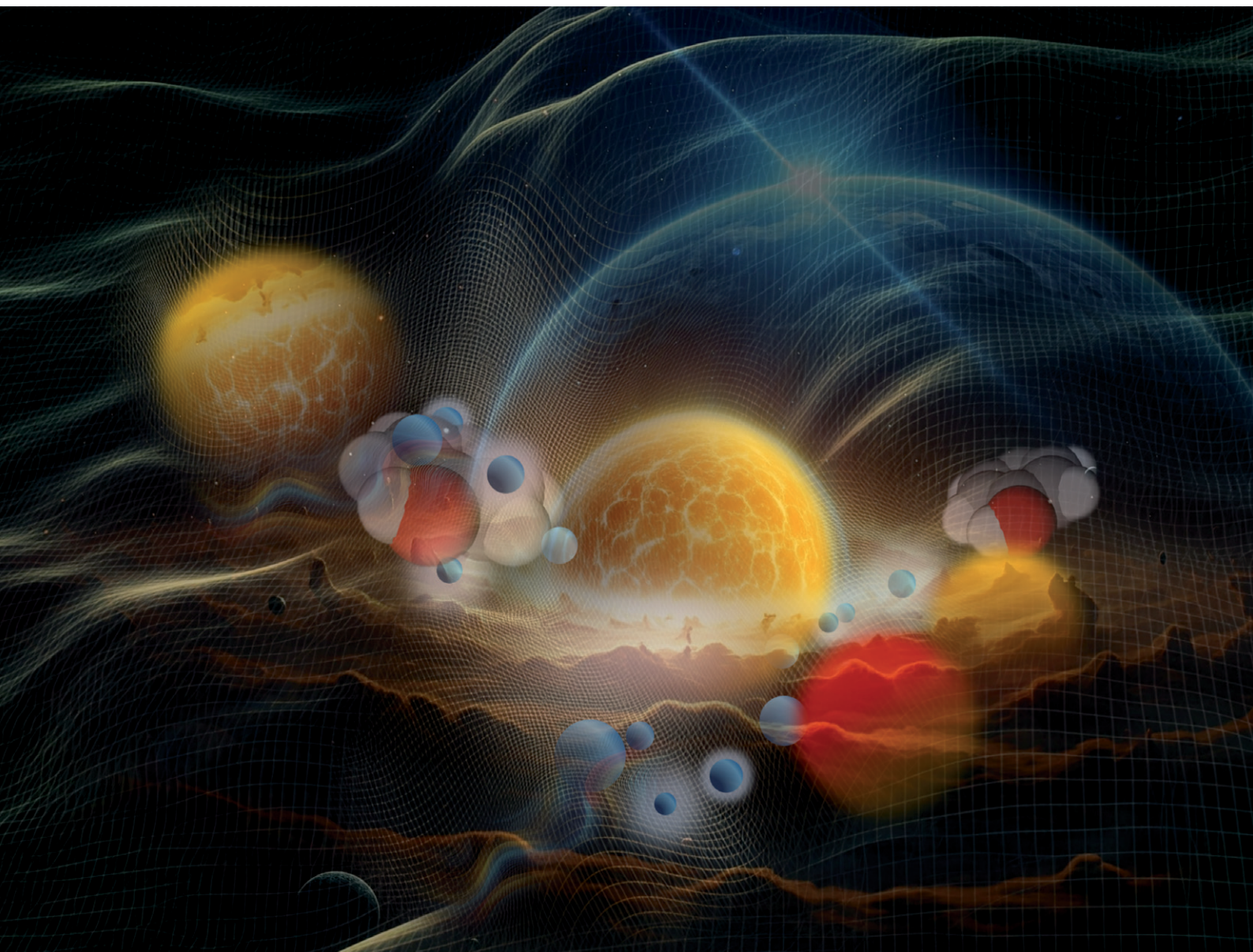


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**PAPER**

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# Environmentally persistent free radicals readily form on TiO<sub>2</sub> but not on ZnO nanoparticles†

Eric P. Vejerano \* and Jeonghyeon Ahn

Environmentally persistent free radicals (EPFRs) are stable organic radicals and pollutants in atmospheric and soil particles. Knowledge of EPFR formation on pristine, unsupported engineered nanoparticles is limited. This study investigates the potential of TiO<sub>2</sub> and ZnO nanoparticles in aiding aromatic precursors in forming EPFRs and is an experimental verification of our earlier prediction on nanoparticles' potential to form EPFRs. TiO<sub>2</sub> formed EPFRs regardless of size, while ZnO exhibited size-dependent behavior. Only smaller ZnO particles readily produced EPFRs. Larger ZnO particles only formed EPFRs when pretreated with ethanol. EPFRs formed on TiO<sub>2</sub> and ZnO near room temperature, challenging the belief that relatively high temperatures are needed for EPFRs to form. The yields of EPFRs on pristine TiO<sub>2</sub> and ZnO were comparable to those found in atmospheric particulate matter. We identified four types of EPFRs: carbon-centered phenoxyl, oxygen-centered phenoxyl, carbon-centered semiquinone, and oxygen-centered semiquinone radicals. Our study suggests the role of band bending in forming EPFRs on TiO<sub>2</sub> and ZnO and the adsorption of aromatic precursors acting as electron acceptors or donors. Our findings suggest that EPFRs may form in unexpected places and under conditions where TiO<sub>2</sub> and ZnO nanoparticles are present.

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## Environmental significance

This study elucidates the interaction of TiO<sub>2</sub> and ZnO nanoparticles with environmental systems, particularly their role in forming environmentally persistent free radicals (EPFRs). By demonstrating that EPFRs can form on these nanoparticles at near-room temperatures and near-ambient pressure, our findings challenge the conventional belief that high temperature and low pressure are necessary for EPFR formation. This work highlights the potential of engineered nanoparticles to contribute to atmospheric and soil pollution through EPFR formation. The size-dependent behavior in EPFR formation enhances our understanding of nanomaterial interactions with natural systems. These insights are crucial for assessing the environmental and human health risks associated with the widespread presence of TiO<sub>2</sub> and ZnO nanoparticles.

## Introduction

Environmentally persistent free radicals (EPFRs) are a new class of important persistent organic pollutants<sup>1</sup> with significant public health impacts.<sup>2</sup> EPFRs are stable and persistent, thus capable of being transported over long distances from their sources. EPFRs have been detected in diverse environmental matrices: atmospheric particulate matter,<sup>3–13</sup> soils/sediments and biochars,<sup>14–25</sup> wood and coal combustion sites,<sup>26,27</sup> microplastics<sup>28</sup> wildfire charcoals,<sup>29</sup> sewage sludge carbonization,<sup>30</sup> to name a few. Recently, persistent free radicals are also found in leaves,<sup>31</sup> dubbed as biogenic PFRs. Many

studies on EPFRs' health effects on murine models demonstrate that EPFRs elicit various health dysfunctions.<sup>2,32–34</sup> Despite their low concentration relative to molecular pollutants in particulate matter (PM), EPFRs exert adverse effects because they can catalytically amplify ROS formation in biological systems.<sup>5</sup>

Many studies on EPFRs have focused on the contribution of transition metal oxide NPs that are supported on surfaces of larger particles. This interest is understandable since the major fraction of ambient PM is from combustion sources that contain clusters of metal oxide nanoparticles (NPs) on surfaces of larger particles. However, pristine NPs, specifically those composed of transition metal oxides, are increasingly used in many commercial products.<sup>35</sup> TiO<sub>2</sub> is among the most widely industrially produced and used NPs.<sup>36</sup> Approximately 4 million metric tons of TiO<sub>2</sub> NPs will be produced by 2025.<sup>37–39</sup> Their eventual disposal warrants investigation if they assist in forming EPFRs from organic components in the waste stream.

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Still, questions exist if pristine engineered NPs can form EPFRs under ambient conditions that typify many environmental systems. NPs, when supported on larger particles (e.g., ambient PM), create Schottky contact due to differences in dielectric properties. Hence, they readily form EPFRs. The Schottky contact is absent in pristine NPs. Only limited information exists on the ability of pristine NPs to stabilize and aid EPFRs.

We posited in our earlier paper that pure NPs can aid in forming EPFRs *via* band bending and that their formation will depend on particle size, chemical composition, and temperature.<sup>40</sup> Here, we investigated the following research questions. First, do NPs with slight differences in their band gap energies (energy difference of the valence and conduction bands) drastically differ in their potential to aid EPFR formation? TiO<sub>2</sub> and ZnO are good candidates since their band gap energies differ only by 0.12 eV. Second, what is the influence of the NP size in forming EPFRs. Third, what is the effect of temperature since thermal perturbation can increase the propensity of electrons to traverse the band gaps? A previous study has demonstrated that TiO<sub>2</sub> can stabilize phenol, forming EPFRs at 220 °C but not at 25 °C.<sup>41</sup> Hence, if large-sized NPs stabilize EPFRs, could their smaller counterpart stabilize and form EPFRs at lower temperatures?

The study described here differs from the earlier studies,<sup>41–43</sup> which were performed under ultralow pressure (high vacuum, <10<sup>−5</sup> Torr). High vacuum can induce a subtle shift in the properties of the material.<sup>44</sup> At sufficiently high temperatures, high vacuum conditions can reduce metal oxides.<sup>45</sup> These conditions do not mimic those occurring in the post-combustion zone during incineration and thermal processing of waste, which most likely occur at high pressure. The study described in this paper was conducted at higher pressures (~14–670 Torr).

## Materials and methods

### NPs

We purchased ZnO (wurtzite, average particle size (APS) 30 nm and 20 nm, 99% purity), TiO<sub>2</sub> (mixed anatase and rutile, APS 20 nm, 99% purity), TiO<sub>2</sub> (rutile, APS 20 nm, 99.8% purity), TiO<sub>2</sub> (anatase, APS 20 nm, 99% purity), and CdO (APS 20–50 nm, 99.8% purity) from Nanostructured & Amorphous Materials, Inc. We used the NPs without any physical or chemical modifications. We used CdO NPs to compare EPFR formation over the ZnO NP surface; both metal oxides are in group 12 (IIB) in the periodic table and have filled d-orbitals. We used pristine TiO<sub>2</sub> (ref. 46) and ZnO since they are among the most used and manufactured NPs and to compare our results to the earlier studies.<sup>41–43</sup> We are also interested in determining the effect of a transition metal with filled d-orbitals (Zn<sup>2+</sup>, [Ar] 3d<sup>10</sup>) and that with multiple unfilled d-orbitals (Ti<sup>2+</sup>, [Ar] 3d<sup>2</sup>).

### Organic precursors for synthesizing EPFRs

We used 2-monochlorophenol (2-MCP, purity ≥99%), phenol (PH, purity ≥99%), and 1,2-dichlorobenzene (1,2-DCB, anhydrous

grade, purity ≥99%) as the organic precursors. These precursors are known precursors of EPFRs on metal oxide particles. We purchased all chemicals from Sigma-Aldrich, which we used without further purification.

### Synthesis of EPFRs

In previous studies, EPFRs were synthesized under a high vacuum of at least 10<sup>−6</sup> Torr. Since EPFRs are detected in atmospheric PM, this suggests that they can form at relatively high pressure. Thus, we synthesized the EPFRs at ~14 Torr and near ambient pressure (670 Torr) to mimic the conditions during combustion and thermal processing. We used a custom-made thermal system (Fig. S1†) for synthesizing EPFRs (details discussed in the ESI†). We placed 100 mg of the metal oxide NPs in each reactor. The reactors were placed inside a ceramic heater and attached to the ports connected to the equilibration chamber. We then pulled a vacuum and slowly opened the valves that connected the reactors to the equilibration chamber to create a final pressure of ~14 or 670 Torr. We set the temperature to 100 °C and heated the NPs for 1 h to reduce moisture. Water removal is critical since a high-water load will compete with the organic adsorbates during the formation of EPFRs. After cleaning, we set the heater to a reaction temperature of 25, 120, 230, 350, or 400 °C. We closed the vacuum line once the reactors had reached the set temperature. We injected 10 μL of the precursor and allowed the vapors to adsorb/react with the NPs for 5 min. We heated the unwrapped glass regions using a heating gun to avoid cold spots during adsorption/reaction. After the reaction, we again pulled a vacuum (14 or 670 Torr) to remove any excess unreacted precursors. We turned off the heater to cool the reactors to room temperature (22 °C). After cooling, we sealed the samples under vacuum, detached the reactors from the exposure chamber, transferred the sample to the side arm, and measured the signal using an electron spin resonance (ESR) spectrometer (details in the ESI†). At 670 Torr, other molecules in the gas phase can compete with the organic precursor.

### Calibration curve using DPPH

We used a three-point calibration curve to bracket the expected concentration using DPPH as the standard. We determined the area counts as the double-integrated area of the ESR signal. We generated the calibration curve by plotting the area counts *v.* the number of molecules of DPPH for each calibration standard. We expressed the [EPFR] in spins per g of NPs. One spin equals one EPFR. We spiked DPPH in silica to obtain a similar matrix with the sample.

### Decay of EPFRs

We measured the [EPFR] at different time intervals after exposure to air. We monitor the sample at different time intervals ranging from days to years. We repeat the cycle of exposing the sample to air, pulling vacuum from the sample, and measuring the EPFR concentration. We plotted decay of the normalized [EPFR] on a semi-log scale as a function of



time. We normalized the concentration to the initial [EPFR]. We determine the pseudo-first order rate constant as the best-fit line. In this case, the  $1/e$ -folding time is the reciprocal of the pseudo-first-order rate constant.

### Material characterization

**X-ray Photo Electron Spectroscopy.** We measured changes in the oxidation state of  $\text{TiO}_2$  and ZnO before and after forming EPFRs. We performed XPS measurement using a Kratos AXIS Ultra DLD XPS system equipped with a hemispherical energy analyzer and a monochromatic Al  $K\alpha$  source. We used the following acquisition parameters: pass energy 20 eV, step size 0.1 eV, dwell time 100 ms. The source was operated at 15 keV and 150 W.

**Transmission electron microscopy (TEM).** We acquired TEM images.  $\text{TiO}_2$  samples before and after forming EPFRs were sent to the Nanoscale Characterization and Fabrication Laboratory at Virginia Tech.

## Results

In our earlier paper, we posited that EPFRs can form on pristine NPs, invoking band bending. Here, we investigated the effect of chemical composition, size, temperature, and pressure on forming EPFRs. Previous studies have generated EPFRs at extremely high vacuum, with pressure exceeding  $10^{-6}$  Torr. Our study investigated the formation of EPFRs under ambient conditions that mimic those during thermal processing.

### Formation of EPFRs on $\text{TiO}_2$ NPs at high pressure

First, we measured the formation of EPFRs on a mixture of anatase- and rutile- $\text{TiO}_2$  NPs at near ambient pressure (670 Torr) and high temperatures.  $\text{TiO}_2$  NPs formed EPFRs when dosed

with 1,2-DCB and 2-MCP (Fig. 1A and B). To test the formation of EPFRs on  $\text{TiO}_2$  NPs and to compare the yield of EPFRs, we performed experiments at a low temperature (25 °C) and high temperature (350 °C). Despite the chlorine substituents destabilizing the aromatic ring in 1,2-DCB, we used it as test precursor molecule since it is symmetric and can form 2-chlorophenoxy and/or *o*-semiquinone radicals.<sup>1</sup> Interestingly, we observed that  $\text{TiO}_2$  NPs generated EPFRs even at 25 °C and 670 Torr for particles >30 nm but less than 100 nm. However, the [EPFR] was 3–4× lower than those dosed at high temperatures. Fig. 1A and B show representative first-derivative ESR spectra of the EPFRs on  $\text{TiO}_2$  NPs and 1,2-DCB at 25 and 350 °C. As shown in Fig. 1A and B, the ESR spectra have a single almost symmetric signal at 350 and 450 °C, whereas at lower temperatures, the signals were highly asymmetric. We attributed this difference to the lower [EPFR] at low temperatures and the types of EPFRs formed. The EPFRs had *g*-values of 2.0035 and 2.0051. Fig. 1B shows the first-derivative ESR spectra of  $\text{TiO}_2$  NPs and 2-MCP at 25, 250, 350, 450 °C. ESR signals generated at room temperature had a *g*-value of 1.9983, which is considerably lower for an organic radical (>2.0023). The *g*-values ranged from ~2.0041 to 2.0049 for all three high temperatures, suggesting the presence of O-centered EPFRs. The [EPFR] increased with increasing temperatures. At 450 °C, we observed the highest [EPFR] of  $1.4 \times 10^{17}$  spins per g of NPs. The *g*-values ranged from 1.9983 to 2.0078, and the  $\Delta H_{\text{p-p}}$  ranged from 4 to 9.6, except those prepared at 25° ( $\Delta H_{\text{p-p}} \sim 2.8$ –26). The ESR spectral characteristics at 350 and 450 °C resemble combustion-generated EPFRs from carbonaceous particles<sup>27</sup> with  $\Delta H_{\text{p-p}}$  of 5.3 to 6.8.

Next, we investigated the formation of EPFRs on surfaces of pure rutile- $\text{TiO}_2$  NPs to determine if a different crystal structure influences EPFR formation using 2-MCP as the precursor. EPFRs formed on rutile- $\text{TiO}_2$  as well. Also, the [EPFR] increased with increasing temperature; the highest

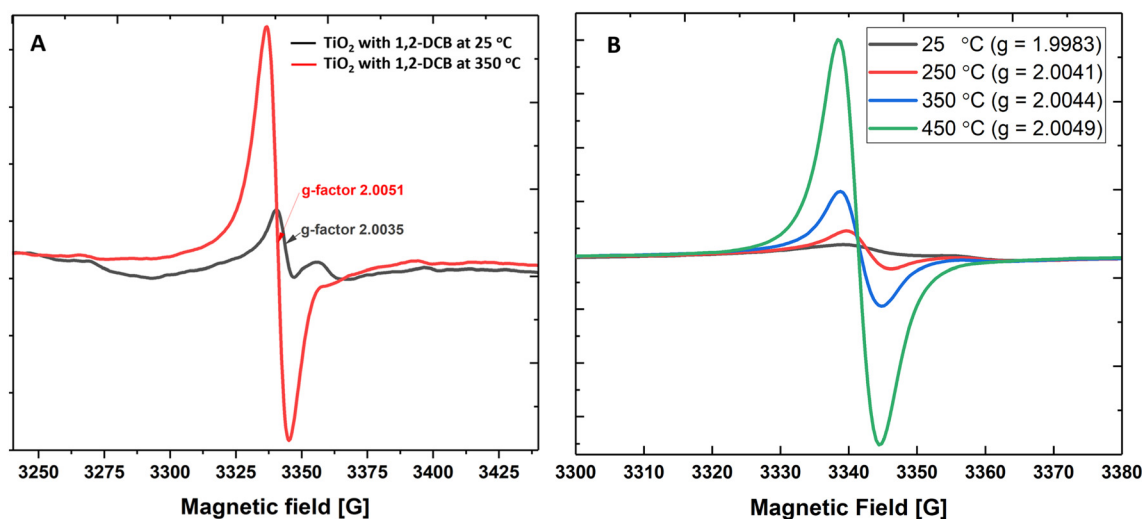
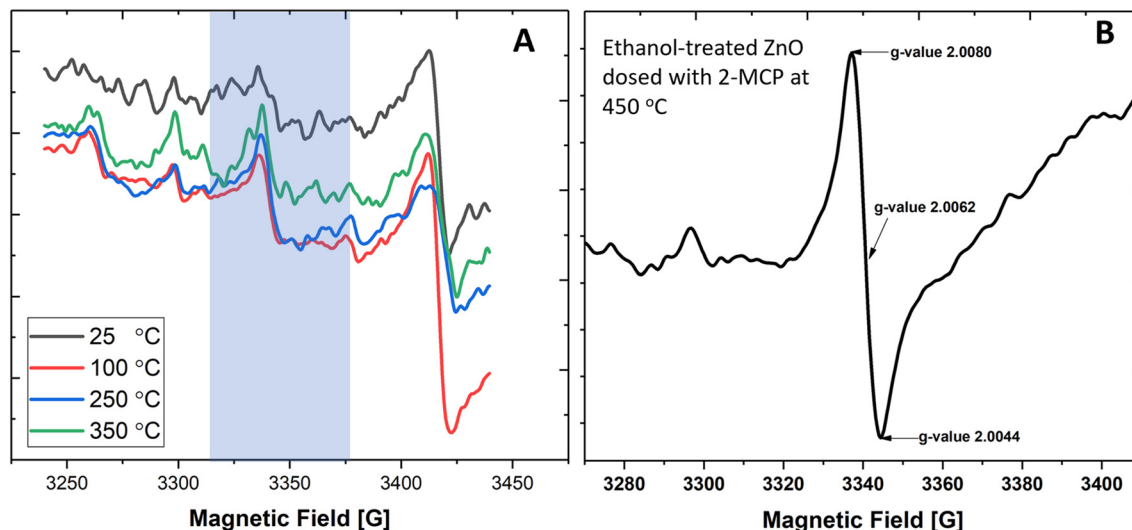


Fig. 1 (A) The ESR intensity of  $\text{TiO}_2$  NPs and 1,2-DCB at low temperature (25 °C) and at high temperature (350 °C) at 670 Torr. (B) The ESR intensity of  $\text{TiO}_2$  NPs and 2-MCP at 25, 250, 350, and 450 °C at 670 Torr.





**Fig. 2** (A) The ESR spectra of ZnO NPs (30 nm) and 2-MCP at 25, 100, 250, and 350 °C at 670 Torr. The shaded region is the magnetic field window for organic radicals (EPFRs). (B) The ESR spectra of ZnO NPs and 2-MCP at 450 °C, pretreated with ethanol at 670 Torr.

[EPFR] of  $1.1 \times 10^{18}$  spins per g of NP was generated at 350 °C. The [EPFR] on rutile-TiO<sub>2</sub> was  $\sim 10\times$  higher than that on anatase-TiO<sub>2</sub> at 250 °C for 2-MCP and 1,2-DCB (Table S1†).

### ZnO forms EPFRs only under unique conditions

While EPFRs readily formed on 30 nm TiO<sub>2</sub> NPs and at 670 Torr, they did not form on ZnO (Fig. 2A). We investigated if modifying the ZnO NP surface can aid in forming EPFRs on 30 nm ZnO NPs at high pressure (670 Torr). We dosed ZnO with methanol since it is among the most abundant simple organic alcohols used in combustion systems.<sup>47</sup> After treating ZnO NPs with methanol, we dosed them with 2-MCP at 670 Torr and 450 °C. However, we did not detect ESR signals. Fig. 2A depicts the representative ESR spectra on ZnO NPs ( $\geq 30$  nm) and 2-MCP dosed at 25, 100, 250, and 350 °C, showing the absence of EPFRs. We then tested a slightly longer chain alcohol (ethanol) to pretreat the ZnO NPs ( $\geq 30$

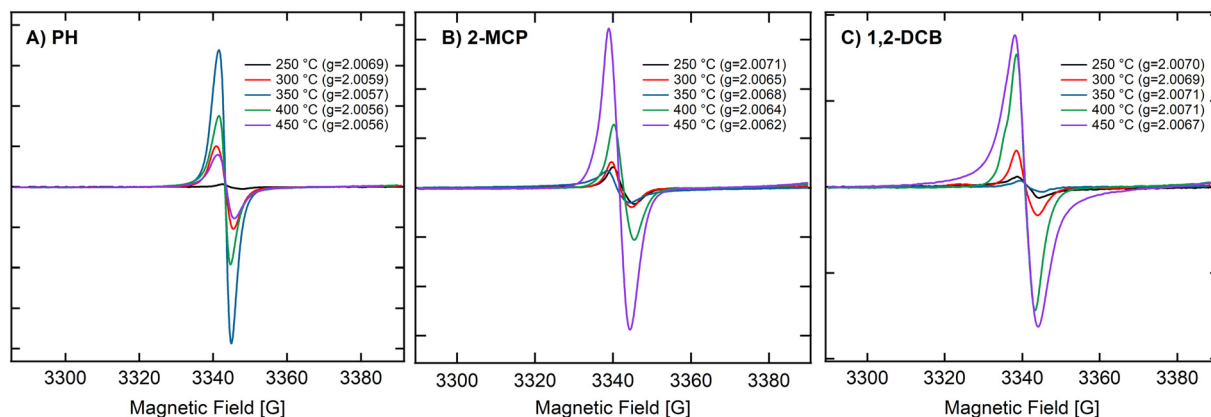
nm) and dosed them with 2-MCP at 450 °C and 670 Torr. EPFRs formed under these conditions (Fig. 2B). Note that the [EPFR] was much lower than those generated at a similar temperature for similar sized TiO<sub>2</sub> NPs. In the absence of 2-MCP, we did not detect any ESR signal for the ZnO NPs pretreated with ethanol. The ESR signal observed at 3400–3450 G (Fig. 2A) corresponds to a *g*-value range that is not associated with organic radicals, such as EPFRs. Overall, larger NPs only produced EPFRs if we pretreat the surface with ethanol.

Thus, we run a series of experiments to see if EPFRs can form on ZnO under specific conditions (*i.e.*, using smaller ZnO NPs and ethanol treatment). We performed these experiments at pressures of 14 and 670 Torr. Results are summarized in Table 1. Representative ESR spectra are depicted in Fig. 2, 3, and S2.† Hereafter, we will use this notation to describe the different experimental conditions, [organic precursor/NP size/pressure/ethanol dosing]. For

**Table 1** Summary of the formation of EPFRs under different pressures on the ZnO surface

Organic precursor	Condition	ZnO NP average particle size (nm)	Pressure (Torr)	Ethanol pretreatment	EPFR formation
2-MCP	A	>30	670	NO	NO
	B	>30	670	YES	YES, weak ESR signals
	C	>30	14	YES	YES, very strong ESR
	D	>30	14	NO	YES, weak ESR
	E	<20	670	NO	YES, strong ESR signal at $\geq 350$ °C
	F	<20	14	NO	YES, strong ESR signal at $\geq 250$ °C
1,2-DCB	G	<20	14	NO	EPFRs formed at $\geq 250$ °C
	H	>30	670	NO	NO
	I	>30	670	YES	YES, very strong ESR signal
	J	<20	670	NO	NO
	K	<20	14	NO	YES, weak ESR signal at $\geq 350$ °C
PH	L	<20	14	NO	EPFRs formed at $\geq 250$ °C
	M	<20	670	NO	EPFRs formed at $\geq 400$ °C
	N	>30	14	NO	NO





**Fig. 3** Representative ESR spectra for the formation of EPFRs for three organic precursors on ZnO NPs. (A) The ESR spectra of ZnO NPs (20 nm) and 2-MCP at 250, 300, 350, 400, and 450 °C at 14 Torr. (B) The ESR spectra of ZnO NPs (20 nm) and PH at 250, 300, 350, 400, and 450 °C at 14 Torr. (C) The ESR spectra of ZnO NPs (20 nm) and 1,2-DCB at 250, 300, 350, 400, and 450 °C at 14 Torr.

example, 2-MCP/30/14T/N means 2-MCP dosed on >30 nm ZnO NPs at 14 Torr without ethanol pretreatment, and Y if the NPs were pretreated with ethanol. Table 1 summarizes the results performed under different conditions. In Table 1, NPs  $\geq 30$  nm did not form EPFRs (conditions A, 2-MCP/30/670T/N and H, 1,2-DCB/30/670T/N) for 2-MCP and 1,2-DCB at high pressure. Maintaining the dosing at 670 Torr, we investigated whether EPFRs will form on smaller ZnO NPs. Under this condition, we observed that smaller ZnO NPs (<20 nm) formed EPFRs (conditions E, 2-MCP/20/670T/N and M, PH/20/670T/N) but only at high temperatures starting at 350 °C for 2-MCP and 400 °C for PH (Fig. 3A–C) without ethanol treatment. However, under similar conditions (J, 1,2-DCB/20/670T/N) but with 1,2-DCB as the organic precursor, ZnO NPs did not generate EPFRs. The corresponding representative ESR spectra of ZnO NPs under different conditions but at a pressure of 670 Torr are shown in Fig. S2†.

We investigated the effect of pressure on EPFR formation. While large ZnO NPs (>30 nm) did not form EPFRs at high pressure and without ethanol pretreatment (condition A, 2-MCP/30/670T/N), at low pressure, they yielded low [EPFR] ESR signal (condition D, 2-MCP/30/14T/N). However, despite multiple dosing attempts with PH, condition N, PH/30/14T/N did not yield EPFR and we expected that it would since PH readily forms EPFRs because of the absence of an electron-withdrawing group (e.g., Cl group), unlike 1,2-DCB or 2-MCP. We did not perform EPFR studies for 20 nm ZnO NPs with ethanol since they readily formed EPFRs at 670 Torr. Thus, they will readily form EPFRs at a lower pressure. Since we only observed a weak ESR signal for 2-MCP, we did not expect 1,2-DCB to form EPFRs under this condition [1,2-DCB/30/14T/N] because of its two electron-withdrawing groups. Thus, we did not perform studies under these conditions.

We observed the temperature dependence of ZnO NPs (<20 nm) at 14 Torr with 2-MCP, PH, and 1,2-DCB. Fig. 3A–C depict the ESR spectra of ZnO NPs (<20 nm) and 2-MCP at 250, 300, 350, 400, and 450 °C. The [EPFRs] ranged from 1.4

to  $7.2 \times 10^{17}$  spins per g of NPs. However, the formation of EPFRs was evident for the smaller size, only at temperatures  $\geq 250$  °C and at 14 Torr for all organic precursors. With increasing temperatures, we observed a shift in the EPFRs'  $g$ -values. The  $g$ -values ranged between 2.0056 and 2.0071, suggesting the presence of O-centered EPFRs. We observed the highest [EPFR] of  $1.8 \times 10^{17}$  spins per g of NPs at 350 °C. In general, the [EPFR] followed the order of 350 °C > 400 °C > 450 °C > 300 °C > 250 °C. We also conducted the same set of experiments at 670 Torr (representative spectra depicted in Fig. S2†). Under this condition, only temperatures exceeding 300 °C led to the formation of EPFRs. PH and 1,2-DCB also produced EPFRs at 450 °C only when dosed at 14 Torr. For ZnO, we observed decreasing  $g$ -values of the EPFRs with increasing temperatures. The [EPFR] ranged from  $1.2 \times 10^{14}$  to  $3.8 \times 10^{18}$  spins per g of NPs for 2-MCP and 1,2-DCB.

We performed a limited series of experiments with 30 nm CdO to confirm if a filled d-orbital would not readily form EPFRs. Indeed, for similar particle size (large >30 nm) and condition (670 Torr), CdO NPs did not form EPFRs at all temperatures ( $\leq 450$  °C) for 1,2-DCB and 2-MCP (data not shown).

### Lifetime and decay of EPFRs

We investigated and quantified the initial concentration of EPFRs from TiO<sub>2</sub> NPs and studied their decay rates. Fig. S3† depicts the [EPFR] decay patterns of TiO<sub>2</sub> NPs (>30 nm) dosed with 2-MCP at 25, 250, and 450 °C. Table S1† shows the ESR characteristics and [EPFR]. The initial [EPFR] values were  $2.33 \times 10^{16}$ ,  $4.63 \times 10^{16}$ , and  $1.43 \times 10^{17}$  spins per g of NPs at 25, 250, and 450 °C, respectively. For EPFRs formed at 250 and 450 °C, the [EPFR] decreased substantially but remained stable after that. Overall, the EPFRs on TiO<sub>2</sub> had a 1/e-folding time of at least 112 days for those formed at 25 °C. Even after more than a year, the [EPFR] remained at  $\sim 10^{15}$  spin per g of NP. Also, EPFRs on ZnO NP (20 nm, dosed with 2-MCP at 350 °C) had a longer lifetime (Fig. S4†)





than those on TiO<sub>2</sub> under the same conditions. The normalized [EPFR] only decayed to 0.167 after 230 days, thus, decaying only by one order of magnitude from its initial [EPFR].

### Effects of heating and EPFR formation on the TiO<sub>2</sub> crystal structure

Since TiO<sub>2</sub> can form EPFRs from low to high temperatures, we investigated if the formation of EPFRs induced the reorientation of the crystal structure. Fig. S5† shows the electron diffraction patterns and fringe lines of the pristine (untreated) and treated TiO<sub>2</sub>. Fig. S5a, c, and e† show the electron diffraction patterns, and Fig. S5b, d, and f† show fringe lines. Fig. S5a, c, and e† depict diffraction ring distances of pure rutile NPs, which is like the heat-treated TiO<sub>2</sub> and those that formed EPFRs. We observed a fringe distance of 0.331–0.332 nm for all three samples, which suggests that thermal treatment does not change the crystallographic properties of TiO<sub>2</sub> under the conditions we experimented with. We did not acquire TEM images for anatase since the transformation of anatase to rutile (phase change) in the presence of air occurs only at temperatures >500 °C,<sup>48</sup> which are above our highest reaction temperature of 450 °C. We did not acquire the electron diffraction and fringe line pattern for ZnO because it is not expected to undergo phase changes at our dosing temperature.<sup>49</sup>

## Discussion

The formation of EPFRs during combustion/thermal processes relies on two components: a surface (*e.g.*, transition metal oxide surface) and an aromatic molecule as a free radical precursor. The single-electron transfer (SET) mechanism is the most invoked conceptual model of EPFR formation.<sup>40</sup> The SET mechanism assumes a one-electron reduction transfer. SET successfully explains the formation of EPFRs on some metal oxides but not on ZnO since it would become Zn<sup>+</sup>, a nonexistent oxidation state.<sup>50</sup>

The idea that electrons can move in both directions (semiconductor-to-organics and *vice versa*) has important consequences. In our earlier paper,<sup>40</sup> we predicted that band bending ( $V_{BB}$ ) is a potential mechanism in forming EPFRs on pristine NPs using the equation developed by Albery,<sup>51</sup>

$$V_{BB} = \frac{eN_d D_0^2}{6\epsilon_0 \epsilon_r}. \quad (1)$$

In this equation,  $V_{BB}$  is the magnitude of band bending,  $e$  is the electronic charge,  $N_d$  is the dopant concentration,<sup>52–55</sup>  $D$  is the depletion layer (absence of charge carriers),  $\epsilon_r$  is the relative dielectric constant,<sup>56,57</sup> and  $\epsilon_0$  is the permittivity of vacuum. For a small particle,  $r_0 \ll \sqrt{3D}$ , where  $r_0$  is the particle radius.<sup>52,58</sup> Such an equation has an important consequence in forming EPFRs due to differences in chemistry ( $\epsilon_r$ ), size ( $r_0$ ), and condition of the physicochemical environment that influences charge transfer (*i.e.*, temperature)<sup>59</sup> during the reaction of an adsorbate on the NP surface. Of the three variables ( $N_d$ ,  $D$ , and

$\epsilon_r$ ),  $\epsilon_r$  is modulated by temperature.  $\epsilon_r$  decreases with increasing temperature.<sup>60,61</sup> Also,  $\epsilon_r$  varies with particle size depending on the nanomaterials.<sup>62–67</sup>

Thus, NPs unsupported on an insulating or a semiconducting layer (*i.e.*, no Schottky contact) can form EPFRs. We posited that EPFRs can form on pure or pristine NPs, depending on their chemical composition and conditions that differ from those in the post-combustion zone below 450 °C.

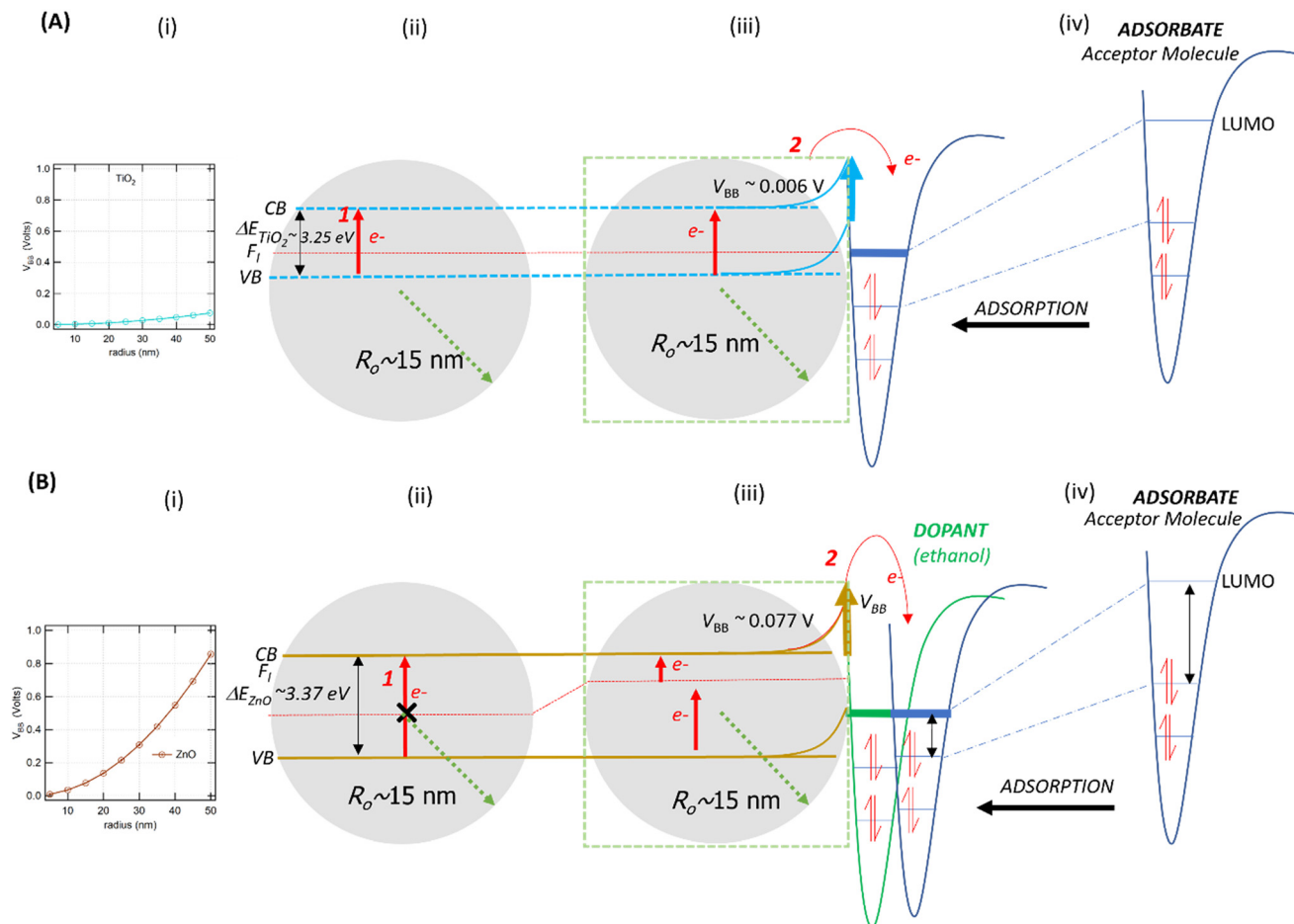
### NP size

From eqn (1), we predicted that small NPs will induce small or negligible  $V_{BB}$  compared to larger NPs since  $V_{BB}$  is proportional to  $r_0$ .<sup>59</sup> For pristine NPs (*i.e.*, NPs with a single chemical composition), this approximation indicates that  $V_{BB}$  for sufficiently small particles depends on  $r_0$  and  $\epsilon_r$ . Using eqn (1), assuming an  $N_d$  of  $10^{24} \text{ m}^{-3}$ , the 30 nm ZnO NPs have a  $V_{BB}$  of 0.077 V (Fig. 4Bi and iii). However, the 30 nm TiO<sub>2</sub> NPs have a  $V_{BB}$  of only 0.007 V (Fig. 4Ai and iii). Such low  $V_{BB}$  values facilitate the easy formation of EPFRs on TiO<sub>2</sub> NPs, consistent with our results. Even the 100 nm TiO<sub>2</sub> NPs have only a  $V_{BB}$  of 0.075 V, which is lower than the  $V_{BB}$  of the 30 nm ZnO NPs (Fig. 4Ai and 4Bi). Smaller ZnO NPs (20 nm), which produced EPFRs, have a  $V_{BB}$  of 0.034 V. This  $V_{BB}$  is lower than the  $V_{BB}$  of the 30 nm ZnO NPs. For this reason, small particles produced EPFRs more easily than their larger counterpart, as we confirmed in our study. This size is equivalent to large TiO<sub>2</sub> NPs (~60 nm), which, by their  $V_{BB}$ , would still form EPFRs. The graph in Fig. 4A and B depict the calculated  $V_{BB}$  as a function of the size of TiO<sub>2</sub> and ZnO NPs, respectively. Small particles have no depletion layer; thus, charge carriers are available on the surface.<sup>59</sup> Thus, in the presence of an organic adsorbate, small NPs easily induce band bending to form EPFRs, as in the case of <20 nm ZnO NPs.

### Effect of the material physicochemical properties

For metal oxide NPs, the interfacial contact between the metal and semiconductor produces an energy barrier known as the Schottky barrier (SB). For this reason, the presence of the SB facilitates the formation of EPFRs on metal oxide NPs supported on large-size silica or alumina particles. Charge carriers in NPs (*e.g.*, electrons) cross the SB to promote electron transfer between the NP and the organic precursor to form EPFRs. However, the SB is absent in pristine NPs. We posited that NPs may form EPFRs whose properties depart from those observed on NPs that are supported on a surface depending on their chemical composition (different  $\epsilon_r$ ). Using eqn (1), the  $V_{BB}$  of ZnO is ~10× higher than that of TiO<sub>2</sub>, assuming all other parameters are equal. We predicted that ZnO will not easily form EPFRs because it has a higher  $V_{BB}$  than TiO<sub>2</sub>.<sup>2</sup> The calculated  $V_{BB}$  for ZnO is higher because of its lower  $\epsilon_r$  (~8.8)<sup>57</sup> than that for TiO<sub>2</sub> ( $\epsilon_r$  ~ 100).<sup>59</sup> Indeed, we have experimentally verified that such is the case for larger (>30 nm) ZnO NPs even at low pressure, even more so under high pressure (670 Torr), where they compete with other gas phase molecules. Thus, no EPFRs were generated under these conditions. Only when we reduced the size that





**Fig. 4** Schematic of the band bending mechanism and electron movement on large (A) TiO<sub>2</sub> NPs and (B) ZnO NPs.  $F_L$  is the Fermi level. CB and VB are the conduction and valence bands, respectively. (i): Plot of variation in band bending ( $V_{BB}$ ) as a function of particle size ( $r$ ) for TiO<sub>2</sub> (A) and ZnO (B). (ii): Energy band diagram of TiO<sub>2</sub> (A) and ZnO (B), showing the CB, VB, and  $F_L$ . (iii) Schematic of the  $V_{BB}$  for both TiO<sub>2</sub> (A) and ZnO (B). In TiO<sub>2</sub>,  $V_{BB}$  is  $\sim 0.006$  V, while in ZnO, it is much higher ( $\sim 0.077$  V). For ZnO, the presence of a dopant (e.g., ethanol) is needed for electron transfer to occur. (iv): Illustration of the electron transfer process from the semiconductor surface to the acceptor molecule.

ZnO NPs formed EPFRs with or without pre-adsorption of a small organic compound, as discussed later.

Although oxygen vacancies (surface defects) are not a prerequisite for generating EPFRs, when present, they enhance EPFR formation.<sup>68,69</sup> The oxygen vacancy can increase the binding energy of both the core level and the valence band electrons.<sup>70</sup> On the ZnO surface, it could enhance band bending in the presence of an organic precursor. Another explanation for why larger 30 nm ZnO NPs do not produce EPFRs is the presence of a few surface-level defects. The ideality factor for spherical shaped ZnO NPs is 1.63,<sup>71</sup> suggesting a low number of surface defects. Consequently, the surface electrons on pristine spherical ZnO NPs are less mobile.<sup>71</sup> This phenomenon may contribute to the difficulty in forming EPFRs on pristine spherical ZnO NPs. We observed that 20 nm ZnO NPs turned yellow, suggesting that oxygen defects formed.<sup>72</sup> Such defects may have contributed to the formation of EPFRs.

Another prediction due to the dielectric constant is the higher relative lifetime of free radicals on ZnO compared to TiO<sub>2</sub> NPs of similar sizes.<sup>56</sup> The lifetimes of EPFRs on TiO<sub>2</sub> and ZnO NPs are comparable to those in other matrices, such

as 5–60 days in soils contaminated with polycyclic aromatic hydrocarbons depending on anoxic conditions or the presence of light irradiation,<sup>21</sup> 0.3–40 days in organic matter of atmospheric particles,<sup>11</sup> 57–59 days in atmospheric haze,<sup>9</sup> 160–401 days in hydrothermal processing of sewage sludge.<sup>30</sup> The [EPFR] on ZnO shifted only by one order of magnitude compared to that on TiO<sub>2</sub> generated under similar conditions. The lifetime for ZnO NPs was much higher than that for ZnO (ref. 50) supported on silica particles. The lower  $\epsilon_r$  for ZnO may explain the longer lifetimes compared to the EPFRs on TiO<sub>2</sub>, consistent with our prediction.<sup>40</sup>

### Effect of temperature

The band gap energy for semiconductors decreases with increasing temperature, as does the  $V_{BB}$ . Increasing temperature enhances lattice vibrations<sup>73</sup> and generates a higher charge-carrier density on the surface.<sup>74</sup> Thermal perturbation can enhance electron transfer, for instance, an electron traversing a band gap and an electron moving between the energy levels of molecular orbitals of a donor/acceptor molecule. ZnO and TiO<sub>2</sub>





are wide-band gap materials; their band gap energies are similar to semiconductors ( $<4$  eV).<sup>75</sup> Thus, the formation of EPFRs may not require the relatively elevated temperature of the post-combustion zone. Such is the case for most EPFRs from various matrices, which often are generated during combustion and thermal processes. Thus, EPFRs can form at room temperature. Indeed, such is the case for  $\text{TiO}_2$  in which even under near ambient pressure conditions and 25 °C, EPFRs easily formed. Note that the yield was lower relative to those formed at higher temperatures. However, the signal was strong, as evidenced by the low-noise ESR spectra (Fig. 1A). An earlier study demonstrated that EPFRs form on ZnO at room temperature.<sup>76</sup> However, this study used a base pressure below  $\sim 10^{-5}$  Torr and heated the ZnO particles at 250 °C before dosing them with the adsorbate at room temperature.<sup>76</sup> As mentioned earlier, at sufficiently high temperatures,<sup>45</sup> high vacuum can modify the surfaces of metal oxides.<sup>44</sup> While such a study is important from a mechanistic standpoint, it does not mimic the conditions at which EPFRs potentially form in the environment.

If there is no Schottky contact, electrons on the NPs are depleted, creating a depletion layer. Under this condition, the energy needed to traverse the band gap becomes larger. This phenomenon is not an issue for  $\text{TiO}_2$  since the band gap is only  $\sim 3.2$  eV (Fig. 4A), easily allowing electrons to flow to the valence band. It is difficult for electrons to flow from the conduction band to the valence band. The thermal energy required for an electron to cross a band gap of 3.2 eV can be estimated using the Boltzmann constant and requires high temperature. Although temperature alone is not a viable mechanism for significant electronic conduction in many materials, increasing temperature increases electronic conduction. For ZnO, in which the band is 0.12 eV higher than that for  $\text{TiO}_2$ , the electrons can only be promoted to the conduction band by thermally activating them at high temperature. Indeed, this hypothesis is consistent with the results that only elevated temperature resulted in higher EPFR yields for the 30 nm ZnO NPs. Thus, doping with an organic precursor, high temperature, and small particle size are needed to generate EPFRs on ZnO NPs.

### Evidence supporting band bending as the likely mechanism for EPFR formation on pristine NPs

Band bending in solid-state materials can be induced by multiple modes.<sup>77</sup> In our study, band bending occurred by two modes: first, the movement of electrons from the valence band to the conduction band of the NPs (Fig. 4A and B) and second, the adsorption of organic molecules acting as acceptors or donors of electrons on surfaces of the NPs, which act as semiconductors (Fig. 4Aiv and Biv). NPs easily sorb volatile organics (e.g., 1,2-DCB, PH, and 2-MCP) to minimize their solid-gas interface's free energy.<sup>78,79</sup> From our XPS analysis, we did not observe  $\text{Ti}^{4+}$  being reduced to  $\text{Ti}^{3+}$  (Fig. S6†). This evidence strengthens the band bending mechanism for  $\text{TiO}_2$  NPs as the potential mechanism for EPFR formation rather than the SET mechanism, which should be evident if  $\text{Ti}^{4+}$  was reduced.

Previous studies conducted with phenol adsorption on  $\text{TiO}_2$  NPs,<sup>41,42</sup> where an organic precursor molecule interacts with an NP, indicated the effect of band bending under ultra-high vacuum conditions. We have shown in this study that such a case also happens under low vacuum (14 Torr) and nearly ambient pressure (670 Torr). We demonstrate that  $\text{TiO}_2$  NPs can generate EPFRs not only at  $\geq 150$  °C but also at 25 °C. The differences in EPFR formation in  $\text{TiO}_2$  NPs at various temperatures are evident from their  $g$ -values, a property of an unpaired electron that reveals chemical information of radicals in any specific environment. Reported  $g$ -values are  $<2.0030$  for C-centered radicals and  $>2.0040$  for O-centered radicals.<sup>80</sup> At high temperatures, thermal excitation and band bending facilitate the easy formation of EPFRs<sup>81</sup> on  $\text{TiO}_2$  from the adsorption of phenol at  $>220$  °C at  $\geq 10^{-3}$  Torr. But band bending could also be a viable mechanism near room temperature and high pressure (670 Torr) for the sorption of 1,2-DCB on  $\text{TiO}_2$ . Laser-induced band bending at room temperature has been observed in some material under ultrahigh vacuum<sup>82</sup> and is reported that small nanoparticles (5–32 nm) have low band bending potential.<sup>77,83–85</sup> Future studies should elucidate the detailed mechanism of EPFR formation on  $\text{TiO}_2$  at room temperature, which is beyond the scope of the current paper.

As illustrated in Fig. 4 and 5, two electron-transfer processes must occur to generate EPFRs: first, electron transfer in the band gap (Fig. 4i and 5i, red arrows) and second, electron transfer from the semiconductor to the aromatic adsorbate (Fig. 4Aiii and Biii). For  $\text{TiO}_2$ , the movement of electrons from the valence band to the conduction band readily happens due to its relatively lower band gap energy, as discussed above. Additionally,  $\text{TiO}_2$  lowers and broadens an organic molecule's lowest unoccupied orbital (LUMO) as it adsorbs on its surface.<sup>59</sup> Adsorption induces electron flow from the  $\text{TiO}_2$  to the adsorbate (upward band bending) (Fig. 4Aiii and iv), forming EPFRs. PH accepts electrons during band bending,<sup>76</sup> and most likely its substituted analogs, 2-MCP and 1,2-DCB, suggesting upward band bending.

ZnO has a wide bandgap energy. A highly pure ZnO NP has an energy gap of 3.31 to 3.49 eV.<sup>86</sup> This band gap does not change beyond 0.1 eV for a size between 0.15 and 80 nm.<sup>87</sup> For ZnO NPs, the promotion of electrons from the valence band to the conduction band does not readily occur because of the 0.12 eV increase in the band gap energy (Fig. 4Bii), especially for larger particles ( $>30$  nm). This slight increase in energy prevents EPFRs from readily forming. However, when doped, materials lower their band gap energies.<sup>88</sup>

In the case of larger ZnO NPs, they produce EPFRs only in the presence of a primary adsorbate (ethanol, Fig. 4Biii). We hypothesize that ethanol unpins the Fermi level of the semiconductor (Fig. 4Biii). This unpinning increases the population of electrons at the Fermi level. Consequently, these electrons easily move to the conduction band of ZnO (red arrows, Fig. 4Biii). However, such unpinning does not immediately result in electron flow from the ZnO to ethanol (i.e., the absence of upward band bending). It is plausible



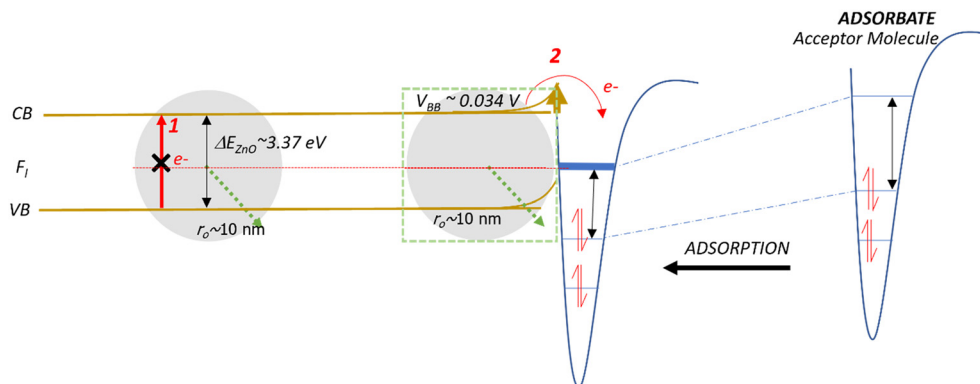


Fig. 5 Schematic of the band bending mechanism and electron movement on a small ZnO NP.  $F_i$  is the Fermi level. CB and VB are the conduction and valence bands, respectively.

that the molecular structure of ethanol could not stabilize the donated electron, but only a larger aromatic molecule can stabilize the donated electron (Fig. 4Biv). The result that ethanol dosed on larger ZnO NPs facilitating the formation of EPFRs is consistent with our earlier studies<sup>89</sup> and that simple acidic/basic organics remove the activation barrier by acting as a bridge between the surface and the aromatic adsorbate to produce EPFRs.<sup>89</sup>

The SET mechanism is also unlikely for several reasons.  $\text{Zn}^{2+}(\text{II})$  is ZnO's only known and confirmed stable valence state. Unlike other metal oxide NPs ( $\text{Fe}_2\text{O}_3$ ,  $\text{NiO}$ ,  $\text{CuO}$ ),<sup>1,90,91</sup> ZnO does not act as an electron acceptor while forming EPFRs. Only a few studies have mentioned that the charge transfers from ZnO to the phenol adsorbate, revealing a mechanism different from SET.<sup>43,92</sup> However, the SET mechanism is not operative in forming EPFRs on pristine ZnO surfaces. The reduction of  $\text{Zn}^{2+}$  to  $\text{Zn}^0$  or its oxidation to  $\text{Zn}^{3+}$  states is still debatable. Our XPS analysis for ZnO (Fig. S6†) supports this explanation. A lone-oxidized  $\text{Zn}^{3+}$  cannot be created thermodynamically;<sup>93</sup> its DFT calculations are heavily contested.<sup>93,94</sup>  $\text{Zn}^+$  exists only under rare conditions when metal-metal bonds dominate,<sup>95,96</sup> which is unlikely in our case. In such a case, only an oxidizing Zn surface is predicted to proceed *via* SET by transferring electrons to the organic precursor, forming EPFRs.<sup>76</sup>

## Conclusions

The insights into the processes of EPFR formation on NP surfaces have important implications. Given that NPs like  $\text{TiO}_2$  (ref. 46) and ZnO (ref. 78) are heavily used in various industrial processes, engineering applications, food processing,<sup>97,98</sup> and consumer products,<sup>99</sup> understanding their ability to generate EPFRs is critical. One important finding of our study is that EPFRs can form on NPs even at room temperature under ambient conditions. Our result challenges existing notions that EPFR formation requires elevated temperatures. Our findings indicate that  $\text{TiO}_2$  NPs readily form EPFRs regardless of size, whereas ZnO nanoparticles exhibit size-dependent behavior. The ease with which  $\text{TiO}_2$  forms EPFRs raises concerns about its widespread use,<sup>46</sup> especially because they are abundant in

air, soil, and water.<sup>100–104</sup> Additionally,  $\text{TiO}_2$  NPs are released into the environment more than any other NPs.<sup>105</sup> Our study highlights the difficulty of predicting the environmental fate and toxicity of NPs. For instance, the fact that large ZnO NPs only form EPFRs when dosed with organic molecules adds a layer of complexity in predicting EPFR formation on NPs. This finding suggests that slight modifications in NP characteristics can yield significantly different behaviors on their ability to form EPFRs. Large ZnO NPs do not easily form EPFRs unless pretreated with small organic molecules. However, organic molecules abound in nature and in consumer products such as sunscreens. The heavy use of ZnO NPs in sunscreen necessitates research on their potential to form EPFRs under ambient condition and *via* photo-driven reactions during their use. Thus, EPFRs may form *via* a light-driven process, which for some reactions is more efficient than thermally activated reactions.<sup>106</sup> Experiments are underway to investigate this hypothesis. The generation and toxicity of EPFRs on transition metal-oxide NPs like  $\text{TiO}_2$  and ZnO and their complexes with organic radicals are yet to be fully characterized. Our findings, however, offer a starting point for identifying which NPs may be suitable for specific applications, considering their environmental and health impacts. Given the widespread use of these materials, further research is urgently needed to inform the safer and more sustainable practices when using nanomaterials on their potential to form EPFRs.

## Data availability

The data for this study are available in the article and its ESI.†

## Conflicts of interest

We declare we have no competing interests.

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