



Cite this: *Mater. Horiz.*, 2024, 11, 5614

Received 9th July 2024,  
Accepted 5th September 2024

DOI: 10.1039/d4mh00882k

rsc.li/materials-horizons

## Defect-engineering of liquid-phase exfoliated 2D semiconductors: stepwise covalent growth of electronic lateral hetero-networks†

Antonio Gaetano Ricciardulli,<sup>a</sup> Christopher E. Petoukhoff,<sup>b</sup> Anna Zhuravlova,<sup>a</sup> Adam G. Kelly,<sup>c</sup> Chun Ma,<sup>a</sup> Frédéric Laquai,<sup>b</sup> Jonathan N. Coleman<sup>c</sup> and Paolo Samori<sup>id</sup> <sup>\*a</sup>

Two-dimensional (2D) in-plane heterostructures display exceptional optical and electrical properties well beyond those of their pristine components. However, they are usually produced by tedious and energy-intensive bottom-up growth approaches, not compatible with scalable solution-processing technologies. Here, we report a new stepwise microfluidic approach, based on defect engineering of liquid-phase exfoliated transition metal dichalcogenides (TMDs), to synthesize 2D hetero-networks. The healing of sulfur vacancies in MoS<sub>2</sub> and WS<sub>2</sub> is exploited to controllably bridge adjacent nanosheets of different chemical nature with dithiolated conjugated molecular linkers, yielding solution-processed nm-scale thick networks with enhanced percolation pathways for charge transport. In-plane growth and molecular-driven assembly synergistically lead to molecularly engineered heterojunctions suppressing the formation of tightly bound interlayer excitons that are typical of conventional TMD blends, promoting faster charge separation. Our strategy offers an unprecedented route to chemically assemble solution-processed heterostructures with functional complexity that can be further enhanced by exploiting stimuli-responsive molecular linkers.

### New concepts

In this work, we have devised a solution-processed strategy based on defect-engineering to control the in-plane growth of thin hetero-networks from inks of different two-dimensional (2D) materials. To enable the formation of lateral 2D hetero-networks,  $\pi$ -conjugated dithiolated systems have been exploited as molecular linkers to yield MoS<sub>2</sub>-WS<sub>2</sub> lateral heterojunctions through the healing of sulfur vacancies, mainly located at the edges of the nanosheets. The selective assembly of adjacent MoS<sub>2</sub> and WS<sub>2</sub> building units is guaranteed by the design of a cyclic and stepwise microfluidic process based on the sequential deposition of MoS<sub>2</sub>,  $\pi$ -conjugated dithiolated molecules and WS<sub>2</sub>. The synthesized covalent lateral 2D hetero-network displays distinct photophysical properties beyond those of pristine MoS<sub>2</sub>-WS<sub>2</sub> networks, such as faster charge separation, as well as superior electrical properties than MoS<sub>2</sub> and WS<sub>2</sub> covalent homo-networks. Hitherto, the lack of component selectivity and control over the edge-to-edge assembly of different 2D materials in conventional deposition techniques hindered the progress and exploration of solution-processed heterostructures. Our modular strategy paves the way toward the assembly of endless low-dimensional complex systems endowed with both tailored and uncharted characteristics.

## Introduction

The emergence of two-dimensional (2D) materials, with their thickness down to the atomic scale, has provided access to radically new physical and chemical features.<sup>1–3</sup> To go beyond the reach of the existing materials and unlock unprecedented

paradigms in physics by pursuing functional diversification, 2D materials have been used as atomically thin building blocks to form heterojunctions endowed with tailored functionalities, either in an out-of-plane or in-plane fashion.<sup>4–6</sup> The latter geometry enables more precise tuning of the heterostructure properties, as the lateral assembly occurs through the edges of the building units, which are more spatially separated.<sup>7</sup> For instance, p–n nodes can be formed at the atomic scale to shrink further optoelectronic devices, beyond the current limits of Moore's law.<sup>8</sup> Noteworthy, interfacing diverse transition metal dichalcogenides (TMDs) to form heterostructures has recently attracted much attention as it could either combine the best characteristics of their components<sup>9</sup> or lead to enticing new physics,<sup>10,11</sup> envisioning the development of innovative multi-functional (opto)electronic devices based on materials by design.<sup>12,13</sup> Research on these heterostructures generally relies on physical vapor transport, chemical vapor deposition (CVD), and epitaxial growth techniques, as they offer atomic control

<sup>a</sup> Université de Strasbourg, CNRS, ISIS UMR 7006, 8 allée Gaspard Monge, 67000 Strasbourg, France. E-mail: samori@unistra.fr

<sup>b</sup> King Abdullah University of Science and Technology (KAUST), Physical Science and Engineering Division (PSE), KAUST Solar Center (KSC), Thuwal, 23955-6900, Kingdom of Saudi Arabia

<sup>c</sup> School of Physics, Centre for Research on Adaptive Nanostructures and Nanodevices (CRANN) and Advanced Materials and Bioengineering Research (AMBER), Trinity College Dublin, Dublin 2, Ireland

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d4mh00882k>



over the assembly. Alternatively, solution processing holds promise for the scalable production of large-area architectures for printed technologies. In this context, films made by distinct 2D materials can perhaps be best defined as hetero-networks because their morphology and orientation significantly differ from conventional 2D heterostructures. However, both terms are used throughout the text.

Liquid-phase exfoliation (LPE) features high-yield, cost-effectiveness and scalability, and it outputs nanosheets dispersed in solvents of choice, being the ideal synthetic platform for large-area (opto)electronic devices compatible with printed technologies.<sup>14–16</sup> Nevertheless, the use of LPE flakes as scaffolds for 2D lateral heterostructures is inhibited by the complete lack of control in terms of component selectivity and preferential edge-to-edge assembly of different species arising from state-of-the-art fabrication strategies, such as spin-coating, spray-coating, drop-casting, *etc.*<sup>17</sup> Indeed, these techniques yield exclusively thick films of overlapped nanosheets. To eliminate this grand challenge, among the different structural defects typically existing in LPE 2D semiconducting nanostructures, chalcogen vacancies located at the edges of TMD nanosheets, which are the most abundant anomalies and conventionally considered as bottlenecks for charge transport,<sup>18–20</sup> can be exploited as reactive sites to covalently bridge adjacent flakes to form ultrathin and continuous films through a novel supramolecular strategy.

In this work, to engineer a uniform and large-area TMD-based hetero-network displaying properties well beyond those of the individual components, dithiolated molecules are employed to covalently connect neighboring molybdenum and tungsten disulfide nanosheets ( $\text{MoS}_2$  and  $\text{WS}_2$ , respectively) *via* microfluidic flow, which regulates the diffusion of molecules in a controlled manner.<sup>21</sup> To ensure selective bridging, cyclic and sequential deposition of  $\text{MoS}_2$ ,  $\pi$ -conjugated dithiolated molecules and  $\text{WS}_2$  is exploited. Moreover, the use of  $\pi$ -conjugated molecular bidentate systems to heal sulfur vacancies in TMDs was demonstrated to be a viable strategy to generate percolation pathways, enabling efficient charge transfer among nanosheets of the same chemical nature, and thus boosting the inter-flake electronic connectivity across the interconnected network.<sup>18</sup> It is worth noting that, compared to previous studies based on typical deposition methods such as drop-casting or spray coating,<sup>18,22</sup> the use of a stepwise approach grants access to a larger number of sulfur vacancies in TMDs as the use of low-concentrated dispersions enables flakes to be more spatially separated, resulting in the maximization of defects healing by thiolated molecules. The synergy of multiscale microscopy and both steady-state and time-resolved spectroscopy analyses confirm that the use of such molecular bridges is critical to both control the synthesis of  $\text{MoS}_2$ - $\text{WS}_2$  lateral heterojunctions by exploiting sulfur vacancies as anchoring sites and enhance the overall electrical performance. Indeed, as revealed by photophysics investigations, faster charge separation is observed in such lateral hetero-networks compared to pristine  $\text{MoS}_2$ - $\text{WS}_2$  blends, suggesting new physical characteristics that hold potential in printed (opto)electronics. Moreover, our microfluidic approach grants access to percolation of charges at reduced thicknesses (single flake-thick). Indeed, field-effect

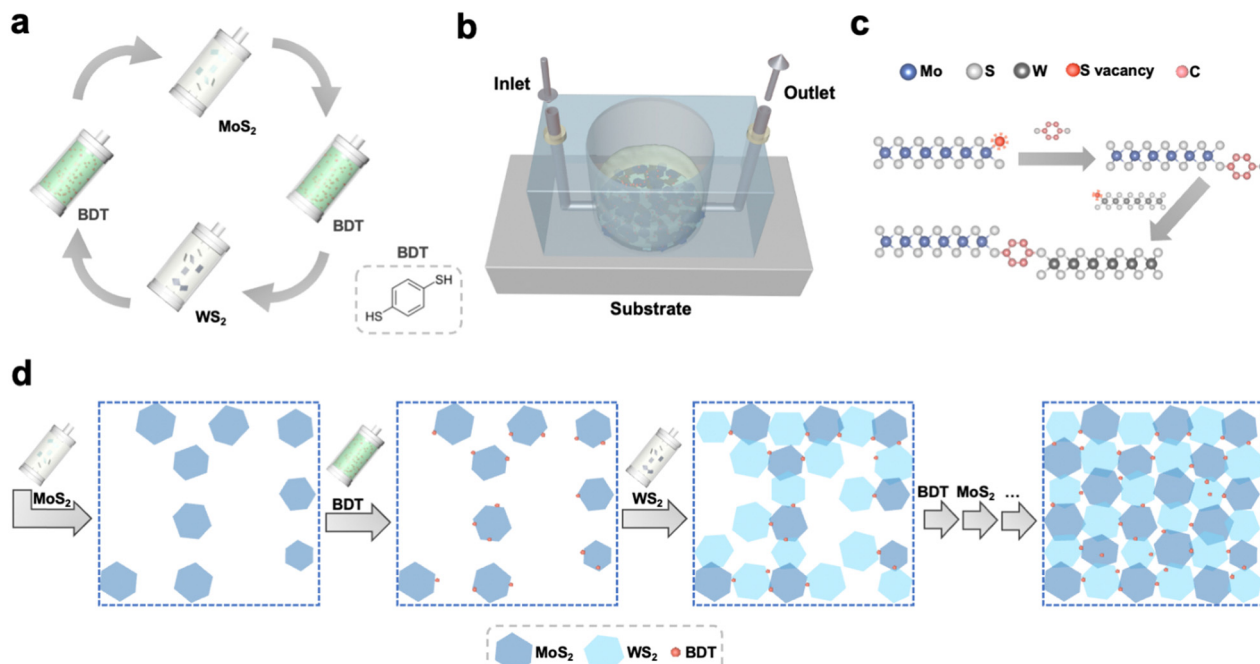
transistors (FETs) based on solution-processed hybrid edge-bridged  $\text{MoS}_2$ - $\text{WS}_2$  heterostructures exhibit charge carrier mobilities matching state-of-the-art films made of LPE TMDs, which are typically measured in films displaying orders of magnitude higher thicknesses. The modular strategy devised in this work can be readily extended to the whole library of solution-processed low-dimensional nanomaterials, including 0D, 1D and 2D structures, and *ad hoc* molecular linkers with tailored multifunctional moieties, and thus serve as a powerful toolkit to assemble endless complex heterostructures with on-demand functionalities through geometrically controlled growth.

## Results and discussion

### Synthesis of solution-processed 2D hetero-networks through defects

The processing steps to grow the 2D networks based on defect-engineered LPE  $\text{MoS}_2$  and  $\text{WS}_2$  are illustrated in Fig. 1. To avoid the randomness of material deposition that occurs with conventional fabrication methods (*i.e.* drop-casting, spin-coating, spray-coating, *etc.* – ESI† Fig. S1) and maximize the amount of  $\text{MoS}_2$ - $\text{WS}_2$  heterojunctions, the synthesis of the heterostructure was carried out under laminar flow assisted by a sequential microfluidic process, using tailored conditions (Fig. S2–S4, ESI†). To demonstrate the validity of the sequential approach to grow heterostructures from solution-processed materials, a basic microfluidic system was set up. First, LPE  $\text{MoS}_2$  ink was flowing on top of a  $\text{Si}/\text{SiO}_2$  substrate, whose wettability was previously tuned *via* the chemisorption of a 3-(aminopropyl)-triethoxysilane (APTES) monolayer to facilitate the adhesion of  $\text{MoS}_2$ . Subsequently, 1,4-benzenedithiol (BDT) was used to control the *in situ* formation of the heterojunctions. The use of a conformationally rigid, bidentate molecule is pivotal to assemble the heterostructure. On the one hand, chalcogen vacancies exposed on the deposited  $\text{MoS}_2$  nanosheets, primarily located at the edges of 2D TMDs, act as anchoring sites for one of the two thiol-terminals of BDT. On the other hand, the second unreacted thiol end-group of the anchored BDT can be used to link covalently further nanosheets in the subsequent step. Thus, LPE  $\text{WS}_2$  ink was introduced in the microfluidic chamber to build the edge-to-edge  $\text{MoS}_2$ - $\text{WS}_2$  heterojunctions *via* BDT bridging. Finally, BDT was further employed to covalently functionalize the exposed sulfur vacancies at the edges of deposited  $\text{WS}_2$ . The alternating steps result in a microfluidic cycle that is repeated multiple times to yield an interconnected, uniform and thin film with a coverage of  $\sim 70\%$  (Fig. S5, ESI†). This feasible and scalable solution-processed growth represents the first example of liquid-phase, structurally controlled, formation of an in-plane heterostructure of 2D materials, exploiting structural defects at the edges of TMDs to drive the preferentially in-plane heterojunction formation by covalently bridging the flakes of the network with dithiolated molecules. While the intrinsic conformational rigidity of the BDT molecules hinders backfolding and linkage of the two thiol end-groups to the same nanosheet, despite the





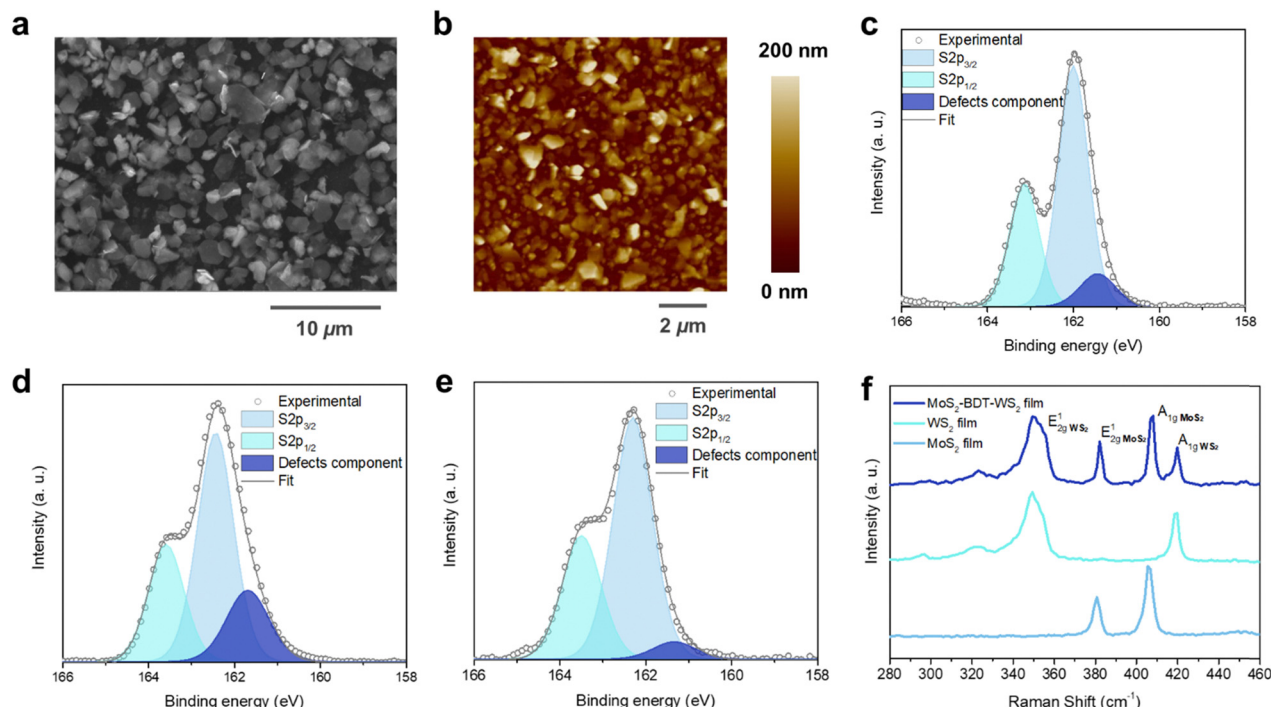
**Fig. 1** Covalent hetero-network formation strategy. (a) Alternating steps of the microfluidic cycle. The process is repeated until the formation of an interconnected in-plane network. To minimize aggregation, IPA is injected between each step. (b) Schematic representation of the microfluidic setup. (c) Illustration of the healing of edge sulfur vacancies, exploited to bridge adjacent TMD flakes by BDT. (d) Schematic top-view of the stepwise growth of the lateral hetero-network, made of MoS<sub>2</sub>–BDT–WS<sub>2</sub> core units. The last arrow illustrates that further alternating steps, following the cycle displayed in Fig. 1a, are repeated to yield the film.

control over the multiple steps of this solution-processed approach, unintentional vertical overlaps of flakes yielding homojunction formations may still occur. The morphology of the hybrid MoS<sub>2</sub>–WS<sub>2</sub> heterostructure was investigated by scanning electron microscopy (SEM) and atomic force microscopy (AFM). To design the optimal procedure for the formation of a smooth solution-processed film, SEM analysis was used to track the nanosheets arrangement after each LPE TMD deposition step (Fig. S6, ESI†). It should be noted that the interconnected film coverage of  $\sim 70\%$  represents the best trade-off in terms of both heterojunction formation and thickness (Fig. 2a) as further addition of either MoS<sub>2</sub> or WS<sub>2</sub> nanosheets would result in a physisorption-driven 3D growth through the overlap of flakes (Fig. S7, ESI†), since the sulfur vacancies of deposited nanosheets become less accessible as a result of the dense packing of the coating. The smoothness of the thin solution-processed lateral hetero-network is further verified by AFM imaging (Fig. 2b).

**Molecular and defect engineered MoS<sub>2</sub>–WS<sub>2</sub> hetero-networks.** The effective *in situ* functionalization of both LPE MoS<sub>2</sub> and WS<sub>2</sub> with BDT was further assessed by X-ray photoelectron spectroscopy (XPS) and Raman spectroscopy measured in various positions over multiple samples. XPS analysis was carried out on BDT-bridged MoS<sub>2</sub>–WS<sub>2</sub> heterostructure as well as on pristine LPE MoS<sub>2</sub> and WS<sub>2</sub> flakes. Notably, high-resolution S 2p spectrum of TMDs containing sulfur as chalcogen unveils the defectivity of the system, thus providing a tool for probing the healing of sulfur vacancies upon selective chemical functionalization.<sup>23,24</sup> Fig. 2c and d

shows the core level S 2p spectra of pristine LPE MoS<sub>2</sub> and WS<sub>2</sub> films, displaying the main two S 2p<sub>3/2</sub> and S 2p<sub>1/2</sub> components of MoS<sub>2</sub> ( $\sim 162.0$  eV and  $\sim 163.2$  eV) and WS<sub>2</sub> ( $\sim 162.4$  eV and  $\sim 163.6$  eV), respectively. A third component for both MoS<sub>2</sub> and WS<sub>2</sub> at  $\sim 161.4$  eV and  $\sim 161.7$  eV, respectively, is further deconvoluted and ascribed to the defects of the crystal structure, such as the vacancy of neighboring sulfur atoms.<sup>25,26</sup> The significant reduction of the relative area of the defects-related component in the BDT-functionalized lateral heterostructure (Fig. 2e), from  $7.2 \pm 1.4\%$  and  $25.3 \pm 2.1\%$  for MoS<sub>2</sub> and WS<sub>2</sub>, respectively, to  $2.8 \pm 1.6\%$ , clearly indicates the decrease of sulfur vacancies, suggesting the efficient healing through BDT covalent functionalization. In contrast, the defective component in MoS<sub>2</sub>–WS<sub>2</sub> bi-component films without BDT treatment, displays a relative area of  $9.0 \pm 0.2\%$  (Fig. S8, ESI†). Further evidence is provided by comparing the substoichiometric Mo 3d and W 4f spectra of BDT-functionalized MoS<sub>2</sub>–WS<sub>2</sub> heterostructure and the pristine LPE MoS<sub>2</sub>–WS<sub>2</sub> blend (Fig. S9 and S10, ESI†). The Raman spectra taken from selected locations of the lateral heterostructures are composed of an overlay of the resonance modes of both MoS<sub>2</sub> and WS<sub>2</sub> domains (Fig. 2f). Noteworthy, the characteristic Raman features of MoS<sub>2</sub> and WS<sub>2</sub> are preserved, suggesting that the interaction with BDT does not induce any apparent change to the intrinsic structural characteristics of the materials. Furthermore, the Raman mapping on the film reveals an equal distribution of both 2D components, indicating an intimate mixing of MoS<sub>2</sub> and WS<sub>2</sub> junctions (Fig. S11 and S12, ESI†), thus the absence of a phase segregation between the different components that would take





**Fig. 2** Morphological and steady-state spectroscopic investigation. (a) SEM image of a film based on MoS<sub>2</sub>-WS<sub>2</sub> heterostructure after the optimized process. (b) AFM topographical images on randomly selected areas of the heterostructure. High-resolution S 2p XPS spectrum of (c) pristine MoS<sub>2</sub>, (d) pristine WS<sub>2</sub> and (e) hybrid MoS<sub>2</sub>-WS<sub>2</sub> lateral hetero-network. (f) Raman spectra comparison of pristine MoS<sub>2</sub> and WS<sub>2</sub> with hybrid MoS<sub>2</sub>-WS<sub>2</sub> lateral heterostructure.

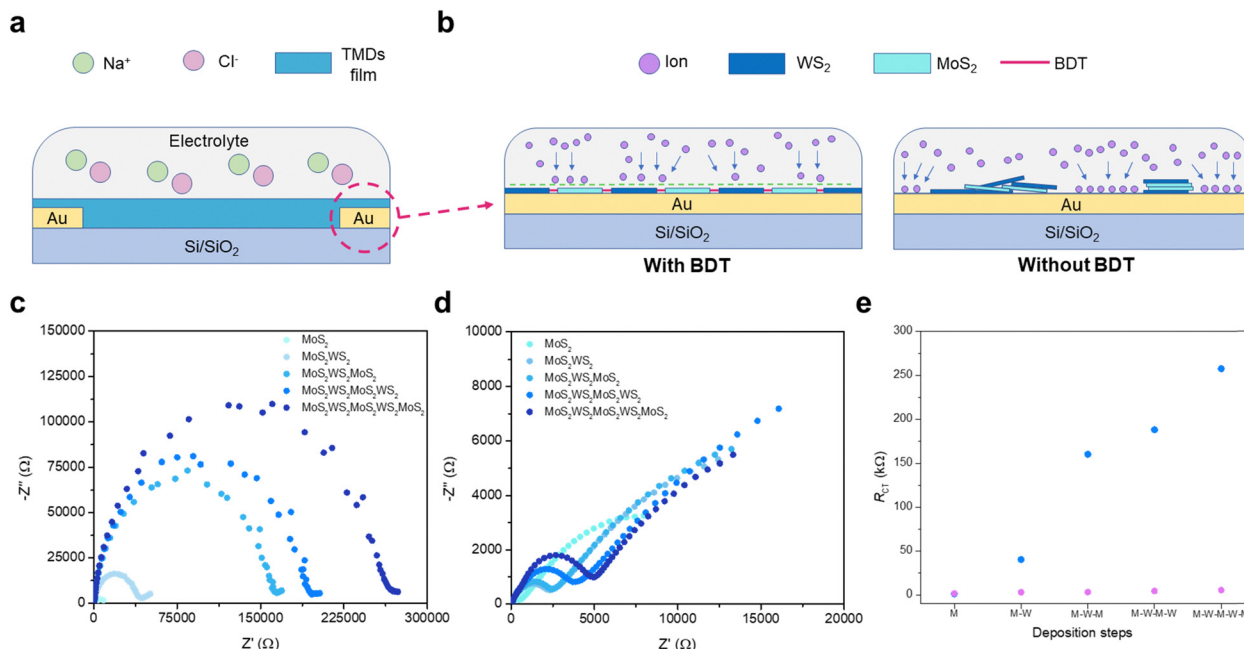
place through a simple mixing. However, the defective presence of homojunctions cannot be fully excluded. In addition to XPS findings, the healing of sulfur vacancies by thiolated molecules is also endorsed by Raman analysis. The full width at half maximum (FWHM) of the E<sub>2g</sub><sup>1</sup> and A<sub>1g</sub> peaks of both MoS<sub>2</sub> and WS<sub>2</sub> witnesses a considerable reduction (as high 32.9% and 20.9% for the E<sub>2g</sub><sup>1</sup> and A<sub>1g</sub> modes of MoS<sub>2</sub>, 6.7% and 10.1% for the E<sub>2g</sub><sup>1</sup> and A<sub>1g</sub> modes of WS<sub>2</sub>, respectively. Table S1, ESI†) compared to their analogues in pristine LPE materials (Fig. 2f).<sup>27</sup> Moreover, the spectra of the heterostructure display a considerable blue shift of the E<sub>2g</sub><sup>1</sup> and A<sub>1g</sub> Raman features (Fig. S13, ESI†), indicating the suppression of defect-induced modes, and thus being consistent with a reduction of sulfur vacancies upon BDT functionalization.<sup>28,29</sup> The shift of the characteristic Raman bands observed in this work exceeds those observed in previous studies on defects passivation in TMDs by thiolated species,<sup>18,22</sup> indicating a considerable increase of sulfur vacancies healing as nanosheets are more accessible by using a simple stepwise microfluidic-assisted strategy compared to typical deposition techniques that yield overlapped and stacked flakes, limiting the access to sulfur vacancies.

**In situ electrochemical impedance spectroscopy.** Electrochemical impedance spectroscopy (EIS) was used to elucidate the synergistic effect of employing dithiolated molecules to propel the heterostructure development through a preferential in-plane growth and promote the creation of percolation pathways, as this technique is extremely sensitive to small variations

within an electrochemical system. As illustrated in Fig. 3a, a tailored 2-electrode electrochemical cell was designed for conducting *in situ* EIS. The growing heterostructure onto patterned interdigitated electrodes on APTES-modified Si/SiO<sub>2</sub> substrates acts as a barrier to the ionic conductivity (Fig. 3b), which is triggered by the application of a fixed sinusoidal voltage across the system (Fig. S14, ESI†). As the steady growth of the 2D assembly hampers the ionic conductivity between electrodes,<sup>30</sup> *in situ* EIS measurements exhibit a significant increase of impedance after each TMD deposition sequence (Fig. 3c), thus acting as internal gauge providing the real-time tracking of stepwise film formation *via* microfluidic approach. The findings indicate larger planar coverage within the interdigitating electrodes is achieved, as a result of the BDT-linking effect within adjacent flakes, propelling the favoured lateral assembly. To validate the critical role of bidentate molecules to assemble the in-plane heterostructure, EIS was performed after each MoS<sub>2</sub> and WS<sub>2</sub> deposition in the absence of BDT as an intermediate process. Transient impedance did not display any significant change upon TMD-loading increase, suggesting partial and uneven coverage of the surface among the interdigitating electrodes (Fig. 3d). A standard Randles equivalent circuit is used to fit the Nyquist plots (model and further discussion in ESI,† Fig. S14) and extract the charge transfer resistance ( $R_{ct}$ ),<sup>31</sup> which is the resistance to the ionic conduction encountered at the interface between electrolyte and electrode. The architecture integrating the MoS<sub>2</sub>-WS<sub>2</sub> heterostructure mediated by BDT witnesses a large contribution of the







**Fig. 3** Tracking the growth of heterostructures via *in situ* EIS. (a) Schematic illustration of the electrochemical cell used for EIS. (b) Representation of the TMDs network acting as a barrier to ionic conductivity with (left) and without BDT (right). (c) Nyquist plot evolution when BDT is inserted in the microfluidic chamber between each MoS<sub>2</sub> and WS<sub>2</sub> step. (d) Nyquist plots after each TMD deposition step for MoS<sub>2</sub>-WS<sub>2</sub> blends in absence of BDT. (e)  $R_{ct}$  comparison as a function of TMD deposition for MoS<sub>2</sub>-WS<sub>2</sub> films with and without BDT, displayed in blue and purple, respectively.

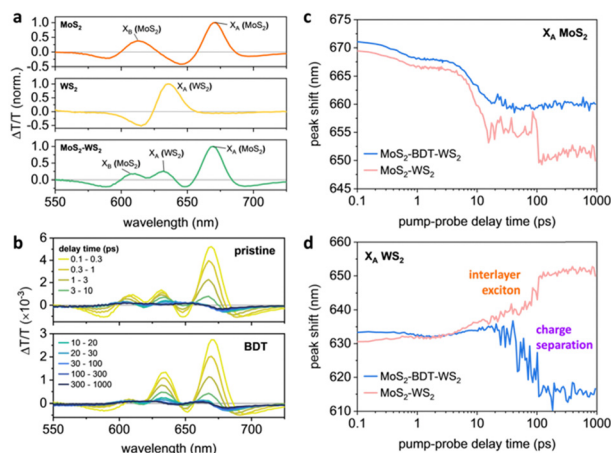
$R_{ct}$  along with the in-plane growth of the network (Fig. 3e). In contrast, the bare deposition of nanosheets without employing any molecular bridge leads to reduced  $R_{ct}$  variation, being consistent with irregular and scarce coating of MoS<sub>2</sub> and WS<sub>2</sub> flakes. Typical SEM images of the film fabricated by alternating only MoS<sub>2</sub> and WS<sub>2</sub> nanosheets provide a distinct footprint of such EIS findings (Fig. S15, ESI†). To exclude the effect of benzene rings in the increase of  $R_{ct}$  and further validate the in-plane growth of the heterostructure mediated by BDT, EIS has been carried out in films where thiophenols have been introduced to heal individual sulfur vacancies in MoS<sub>2</sub> and WS<sub>2</sub> after each deposition step, with the equivalent procedure adopted with BDT. Likewise, for the film formed by alternating MoS<sub>2</sub> and WS<sub>2</sub> only, thiophenol-functionalized coatings exhibit a reduced  $R_{ct}$  comparable to the films without any molecular linker (Fig. S16, ESI†). Indeed, despite the MoS<sub>2</sub> and WS<sub>2</sub> functionalization through the vacancies,<sup>18</sup> thiophenols are not able to bridge adjacent flakes due to the absence of a second thiol like in BDT, inhibiting the formation of a MoS<sub>2</sub>-WS<sub>2</sub> network.

### Photophysics of MoS<sub>2</sub>-WS<sub>2</sub> heterostructures

As BDT enables the sequential linking of adjacent flakes, altering the overall morphology of the microfluidic-assisted grown heterostructure and regulating the distribution of its components, the impact of using such molecular linker on the overall photophysical characteristics of the synthesized MoS<sub>2</sub>-WS<sub>2</sub> hetero-networks was revealed by transient absorption (TA) pump-probe spectroscopy studies. The TA spectra from pristine MoS<sub>2</sub> and WS<sub>2</sub> are dominated by their excitonic

resonances: particularly the A-exciton ( $X_A$ ) and B-exciton ( $X_B$ ) from MoS<sub>2</sub>, and the  $X_A$  from WS<sub>2</sub> within the range of 550 nm to 725 nm (Fig. 4a and Fig. S17, S18, ESI†).<sup>32–34</sup> The spectra have alternating negative and positive differential transmission ( $\Delta T/T$ ) signals, arising from spectral shifts and broadening of the photoexcited states relative to the ground state due to bandgap renormalization.<sup>32–35</sup> The TA spectra of the MoS<sub>2</sub>-WS<sub>2</sub> heterostructures are composed of signals from component materials (Fig. 4a). For the BDT-linked heterostructures, the ratio of  $X_A$  (WS<sub>2</sub>):  $X_A$  (MoS<sub>2</sub>) increased, indicating a higher ratio of WS<sub>2</sub> was present in the linked heterostructures (Fig. 4b). As pump-probe delay time increased, the peaks from each of the excitons shifted spectrally, with each of the excitons in the pristine materials blue-shifting as their bandgaps returned to their ground-state values (Fig. S18–S20, ESI†). This was similar for MoS<sub>2</sub>'s  $X_A$  signals in the heterostructures (Fig. 4c); however, MoS<sub>2</sub>'s  $X_B$  and WS<sub>2</sub>'s  $X_A$  exhibited different trends (Fig. S19b and Fig. 4d, respectively, ESI†). For the unlinked heterostructure, the WS<sub>2</sub>'s  $X_A$  red-shifted with increasing pump-probe delay times – indicating further narrowing of WS<sub>2</sub>'s bandgap upon formation of a long-lived interlayer exciton between MoS<sub>2</sub> and WS<sub>2</sub>, which is typical for vertical heterostructures.<sup>36</sup> This was not the case for the BDT-linked heterostructure: in this case, WS<sub>2</sub>'s  $X_A$  blue-shifted back to its ground state value in  $\sim 100$  ps, further suggesting the different architecture through the preferential in-plane growth mediated by BDT. This rapid blue shift of WS<sub>2</sub>'s  $X_A$  indicates the lack of formation of tightly bound interlayer excitons, indicating that the BDT-linked heterostructures promote faster charge separation. This was further evidenced by the faster relaxation dynamics of WS<sub>2</sub>'s  $X_A$



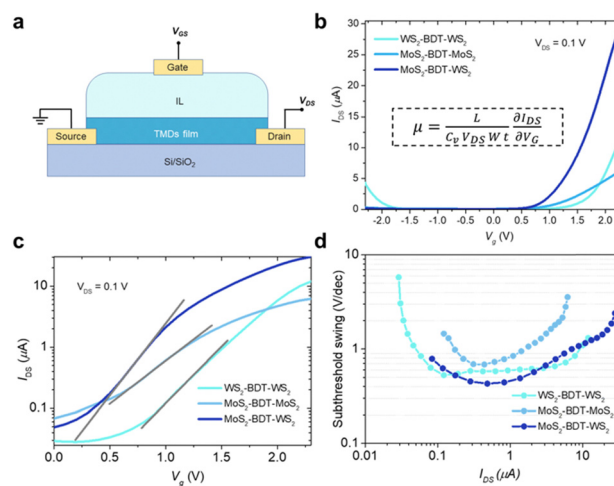


**Fig. 4** Transient absorption (TA) spectra and peak shifts. (a) TA spectra of pristine MoS<sub>2</sub>, WS<sub>2</sub>, and MoS<sub>2</sub>-WS<sub>2</sub> heterostructure averaged over pump-probe delay times of 0.1–0.3 ps. The dominant species observed within this spectral range are indicated. (b) TA spectra of MoS<sub>2</sub>-WS<sub>2</sub> heterostructures, with or without BDT. The spectra are averaged over the pump-probe delay times listed in the legend. (c) and (d) Peak shifts for the reported species in the MoS<sub>2</sub>-WS<sub>2</sub> heterostructures, with or without BDT, as determined from Gaussian fits to the individual peaks: (c) X<sub>A</sub> of MoS<sub>2</sub>; (d) X<sub>A</sub> of WS<sub>2</sub>.

in the BDT-linked heterostructures compared to the unlinked heterostructures (Fig. S21 and Table S2, ESI†), suggesting potential applications in light harvesting systems.

### Electrical characteristics

To unveil the electrical characteristics, we have fabricated proof-of-concept bottom-contact, top-gated FETs based on BDT-connected MoS<sub>2</sub>-WS<sub>2</sub> lateral heterostructures (Fig. 5a and Fig. S24, ESI†). In this device configuration, electrolyte gating induces large carrier densities and enables strong drain-current modulation at low voltage operation.<sup>37,38</sup> It is worth noting that the choice of an aromatic dithiolated linker over an aliphatic one to yield the covalent network of flakes is critical to grant enhanced inter-flake electronic connectivity and formation of additional percolation pathways, which go beyond the electrical features of films based on defective pristine solution-processed 2D materials.<sup>18,22</sup> Direct comparison of the electrical performances with films exclusively made of LPE TMDs with similar thickness is not possible as their fabrication in absence of molecular linkers leads to disconnected and rough coatings, inhibiting charge transport across the device (Fig. S15, ESI†), thus exhibiting negligible transistor response. Fig. 5b displays the transfer characteristics of FET devices based on BDT-connected MoS<sub>2</sub>-WS<sub>2</sub> lateral heterostructure as well as thin films comprising a single 2D material being either MoS<sub>2</sub> or WS<sub>2</sub>. MoS<sub>2</sub> and WS<sub>2</sub> networks, also assembled *via* microfluidic approach, exhibit characteristic n-type and ambipolar behavior, respectively.<sup>12,38,39</sup> As expected, the lateral heterostructure, composed of a network of alternated MoS<sub>2</sub> and WS<sub>2</sub> flakes connected in series by BDT, generates n-type dominated conduction, because only when both MoS<sub>2</sub> and WS<sub>2</sub> components are biased to be at ON state, the overall channel is



**Fig. 5** Electrical characteristics of BDT-linked TMD films. (a) Schematic illustration of the FETs. 1-Ethyl-3-methylimidazolium bis(trifluoromethanesulfonyl)imide was readily cast on top of the thin film as ionic liquid (IL). (b) Transfer curves for in-plane BDT-linked TMD films with V<sub>g</sub> sweeping from -2.3 to +2.3 V at V<sub>ds</sub> = 0.1 V. The inset shows the equation for calculating the field-effect mobility,  $\mu$ . L, channel length; C<sub>v</sub>, volumetric capacitance; W, channel width; t, film thickness (c) subthreshold slopes obtained from logarithmic-scale current characteristics of BDT-linked TMD films. (d) Subthreshold swing as a function of I<sub>DS</sub> in films based on lateral BDT-linked TMD architectures.

turned ON. Further, the linear and symmetrical output characteristics (Fig. S25, ESI†) suggest that ohmic contact is achieved. Interestingly, the device including the hybrid MoS<sub>2</sub>-WS<sub>2</sub> lateral heterostructure, exhibiting lower threshold voltage, shows a higher ON current, which gives rise to a higher ON/OFF current ratio. As the formation of MoS<sub>2</sub>-WS<sub>2</sub> heterojunctions facilitates the electron conduction at the ON state, the fabricated heterostructure yields enhanced field-effect mobility (10<sup>-1</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>), which outperform of one order of magnitude the mobility values reached by the MoS<sub>2</sub> and WS<sub>2</sub> networks (Table S3, ESI†). Importantly, field-effect mobilities obtained in approximately single-flake-thick films by our methodology match state-of-the-art micrometer-thick coatings based on liquid-phase exfoliated TMDs (Table S4, ESI†).<sup>16,18,22</sup> Therefore, in-plane conductivity in the film is maximized at reduced thickness by molecularly bridging flakes through edge defects, which leads to a continuous thin network without compromising the overall thickness.

To estimate the degree of trap density into semiconducting films, the slope of transfer curves was analyzed (Fig. 5c). The steeper slope of the I-V characteristics generated by the BDT-linked MoS<sub>2</sub>-WS<sub>2</sub> heterostructure indicates a lower amount of trap states, which we ascribe to the formation of an intimate network of heterojunctions. To quantitatively assess the trap density (D<sub>it</sub>) of the devices, the parameter is extracted from the subthreshold swing (SS, Fig. 5d) by using the following formula:<sup>40,41</sup>

$$SS = \frac{kT}{q} \ln(10) \left( 1 + \frac{C_{dep} + C_{it}}{C_{ox}} \right) \quad (1)$$



$$D_{it} = \frac{C_{it}}{q} \quad (2)$$

where  $k$  is the Boltzmann's constant,  $T$  is the temperature of the measurement,  $q$  is the electron charge,  $C_{ox}$  is the oxide capacitance (corresponding to the specific capacitance in IL-gated systems),  $C_{dep}$  and  $C_{it}$  are the depletion capacitance and the interface trap capacitance in the weak inversion region, respectively. As the effective substrate bias approaches infinity,  $C_{dep}$  approaches zero. Thus, the SS can be estimated with the following approximate equation:

$$SS = \frac{kT}{q} \ln(10) \left( 1 + \frac{C_{it}}{C_{ox}} \right) \rightarrow \frac{kT}{q} \ln 10 \quad (3)$$

$$= 60 \text{ mV decade}^{-1} \quad \text{at } T = 300 \text{ K}$$

The trap density  $D_{it}$  in the film based on the lateral heterostructure was calculated to be  $3.4 \times 10^{14} \text{ V}^{-1} \text{ cm}^{-2}$ , being lower than the mono-component BDT-networks with  $\text{MoS}_2$  and  $\text{WS}_2$  which amount to  $8.0 \times 10^{14} \text{ V}^{-1} \text{ cm}^{-2}$  and  $4.4 \times 10^{14} \text{ V}^{-1} \text{ cm}^{-2}$ , respectively. These results are consistent with the decrease of inter-flake trap density induced by the heterojunction formation in the synthesized  $\text{MoS}_2$ - $\text{WS}_2$  lateral heterostructure.

## Conclusions

In summary, we have demonstrated the planar growth of solution-processed hetero-networks based on 2D TMDs by exploiting intrinsic defects *via* a simple sequential microfluidic approach. Sulfur vacancies in  $\text{MoS}_2$  and  $\text{WS}_2$ , mainly located at the edges, have been healed by rigid, bidentate  $\pi$ -conjugated molecules to synergistically anchor adjacent nanosheets and promote the preferential in-plane heterostructure formation, providing percolation across the system at flake density far below state-of-the-art solution-processed films. The fabricated BDT-linked  $\text{MoS}_2$ - $\text{WS}_2$  heterostructures exhibited photophysical features that diverge from conventional blends in the absence of linkers, promoting faster charge separation. Moreover, they are endowed with superior field-effect mobility than mono-component networks due to reduced trap density resulting from the formation of heterojunctions. Our strategy to assemble heterostructures from solution-processed nanosheets can be easily extended to other classes of low-dimensional materials by making use of suitably functionalized bidentate linkers, opening up new avenues in the exploration of unprecedented optoelectronic features that go beyond the current 0D, 1D and 2D materials portfolio. Moreover, the selection of anchoring groups which selectively bind a given nanostructure and cores that can impart novel functionality to the hybrid heterostructure, thereby offering endless combinations for further integrating unprecedented solution-processed low-dimensional assemblies in miniaturized next-generation digital (opto)electronic devices.

## Author contributions

A. G. R. and P. S. conceived the experiments and designed the study. A. G. K. and J. N. C. produced the raw materials.

A. G. R. designed and performed the multiscale characterizations on the final functionalized materials. A. G. R. and A. Z. carried out Raman spectroscopy, XPS and impedance analysis. A. G. R. and C. M. performed electrical measurements and interpreted the data. C. E. P. and F. L. designed and performed transient absorption spectroscopy measurements and studies. All authors discussed the results and contributed to the interpretation of data. A. G. R. and P. S. co-wrote the paper with input from all co-authors. All authors have given approval to the final version of the manuscript.

## Data availability

The authors confirm that the data supporting the findings of this study are available within the article and its ESI.†

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

We thank Prof. Thomas Hermans for kindly providing the microfluidic cell. We acknowledge funding from the European Commission through the ERC projects SUPRA2DMAT (GA-833707) and FUTURE-PRINT (GA-694101) and the HORIZON-CL4-2023-DIGITAL-EMERGING-01-CNECT project 2D-PRINTABLE (GA-101135196) as well as the Agence Nationale de la Recherche through the Interdisciplinary Thematic Institute SysChem *via* the IdEx Unistra (ANR-10-IDEX-0002) within the program Investissement d'Avenir, the Foundation Jean-Marie Lehn and the Institut Universitaire de France (IUF). J. N. C. thanks Science Foundation Ireland (SFI) for support. C. E. P. acknowledges support from King Abdullah University of Science and Technology (KAUST) Global Fellowship Program under Award No. ORA-2022-5002. The authors thank Hao Mei (Shanghai Jiao Tong University) for contributing Fig. 1.

## References

- Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, *Nat. Nanotechnol.*, 2012, 7, 699–712.
- A. Rodin, M. Trushin, A. Carvalho and A. H. Castro Neto, *Nat. Rev. Phys.*, 2020, 2, 524–537.
- K. S. Novoselov, A. Mishchenko, A. Carvalho and A. H. Castro Neto, *Science*, 2016, 353, aac9439.
- A. G. Ricciardulli, S. Yang, J. H. Smet and M. Saliba, *Nat. Mater.*, 2021, 20, 1325–1336.
- E. Pomerantseva and Y. Gogotsi, *Nat. Energy*, 2017, 2, 1–6.
- J. R. Schaibley, P. Rivera, H. Yu, K. L. Seyler, J. Yan, D. G. Mandrus, T. Taniguchi, K. Watanabe, W. Yao and X. Xu, *Nat. Commun.*, 2016, 7, 13747.
- M.-Y. Li, Y. Shi, C.-C. Cheng, L.-S. Lu, Y.-C. Lin, H.-L. Tang, M.-L. Tsai, C.-W. Chu, K.-H. Wei, J.-H. He, W.-H. Chang, K. Suenaga and L.-J. Li, *Science*, 2015, 349, 524–528.



- 8 X. Duan, C. Wang, J. C. Shaw, R. Cheng, Y. Chen, H. Li, X. Wu, Y. Tang, Q. Zhang, A. Pan, J. Jiang, R. Yu, Y. Huang and X. Duan, *Nat. Nanotechnol.*, 2014, **9**, 1024–1030.
- 9 C. Huang, S. Wu, A. M. Sanchez, J. J. P. Peters, R. Beanland, J. S. Ross, P. Rivera, W. Yao, D. H. Cobden and X. Xu, *Nat. Mater.*, 2014, **13**, 1096–1101.
- 10 A. Pospischil, M. M. Furchi and T. Mueller, *Nat. Nanotechnol.*, 2014, **9**, 257–261.
- 11 J. S. Ross, P. Klement, A. M. Jones, N. J. Ghimire, J. Yan, D. G. Mandrus, T. Taniguchi, K. Watanabe, K. Kitamura, W. Yao, D. H. Cobden and X. Xu, *Nat. Nanotechnol.*, 2014, **9**, 268–272.
- 12 Y. Gong, J. Lin, X. Wang, G. Shi, S. Lei, Z. Lin, X. Zou, G. Ye, R. Vajtai, B. I. Yakobson, H. Terrones, M. Terrones, B. K. Tay, J. Lou, S. T. Pantelides, Z. Liu, W. Zhou and P. M. Ajayan, *Nat. Mater.*, 2014, **13**, 1135–1142.
- 13 Z. Zhang, Z. Huang, J. Li, D. Wang, Y. Lin, X. Yang, H. Liu, S. Liu, Y. Wang, B. Li, X. Duan and X. Duan, *Nat. Nanotechnol.*, 2022, **17**, 493–499.
- 14 G. Fiori, F. Bonaccorso, G. Iannaccone, T. Palacios, D. Neumaier, A. Seabaugh, S. K. Banerjee and L. Colombo, *Nat. Nanotechnol.*, 2014, **9**, 768–779.
- 15 R. Raccichini, A. Varzi, S. Passerini and B. Scrosati, *Nat. Mater.*, 2015, **14**, 271–279.
- 16 A. G. Kelly, T. Hallam, C. Backes, A. Harvey, A. S. Esmaily, I. Godwin, J. Coelho, V. Nicolosi, J. Lauth, A. Kulkarni, S. Kinge, L. D. A. Siebbeles, G. S. Duesberg and J. N. Coleman, *Science*, 2017, **356**, 69–73.
- 17 F. Bonaccorso, A. Bartolotta, J. N. Coleman and C. Backes, *Adv. Mater.*, 2016, **28**, 6136–6166.
- 18 S. Ippolito, A. G. Kelly, R. Furlan de Oliveira, M.-A. Stoeckel, D. Iglesias, A. Roy, C. Downing, Z. Bian, L. Lombardi, Y. A. Samad, V. Nicolosi, A. C. Ferrari, J. N. Coleman and P. Samori, *Nat. Nanotechnol.*, 2021, **16**, 592–598.
- 19 H.-P. Komsa, J. Kotakoski, S. Kurasch, O. Lehtinen, U. Kaiser and A. V. Krasheninnikov, *Phys. Rev. Lett.*, 2012, **109**, 035503.
- 20 S. McDonnell, R. Addou, C. Buie, R. M. Wallace and C. L. Hinkle, *ACS Nano*, 2014, **8**, 2880–2888.
- 21 Y. Xia, S. Sevim, J. P. Vale, J. Seibel, D. Rodríguez-San-Miguel, D. Kim, S. Pané, T. S. Mayor, S. De Feyter and J. Puigmartí-Luis, *Nat. Commun.*, 2022, **13**, 7006.
- 22 S. Ippolito, F. Urban, W. Zheng, O. Mazzarisi, C. Valentini, A. G. Kelly, S. M. Gali, M. Bonn, D. Beljonne, F. Corberi, J. N. Coleman, H. I. Wang and P. Samori, *Adv. Mater.*, 2023, **35**, 2211157.
- 23 Z. Lin, B. R. Carvalho, E. Kahn, R. Lv, R. Rao, H. Terrones, M. A. Pimenta and M. Terrones, *2D Mater.*, 2016, **3**, 022002.
- 24 X. Zheng, A. Calò, T. Cao, X. Liu, Z. Huang, P. M. Das, M. Drndic, E. Albisetti, F. Lavini, T.-D. Li, V. Narang, W. P. King, J. W. Harrold, M. Vittadello, C. Aruta, D. Shahrjerdi and E. Riedo, *Nat. Commun.*, 2020, **11**, 3463.
- 25 N. S. McIntyre, P. A. Spevack, G. Beamson and D. Briggs, *Surf. Sci.*, 1990, **237**, L390–L397.
- 26 M. Donarelli, F. Bisti, F. Perrozzi and L. Ottaviano, *Chem. Phys. Lett.*, 2013, **588**, 198–202.
- 27 S. Mignuzzi, A. J. Pollard, N. Bonini, B. Brennan, I. S. Gilmore, M. A. Pimenta, D. Richards and D. Roy, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2015, **91**, 195411.
- 28 S. Bae, N. Sugiyama, T. Matsuo, H. Raebiger, K. Shudo and K. Ohno, *Phys. Rev. Appl.*, 2017, **7**, 024001.
- 29 B. Chakraborty, A. Bera, D. V. S. Muthu, S. Bhowmick, U. V. Waghmare and A. K. Sood, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, **85**, 161403.
- 30 S. Wang, J. Zhang, O. Gharbi, V. Vivier, M. Gao and M. E. Orazem, *Nat. Rev. Methods Primers*, 2021, **1**, 1–21.
- 31 E. P. Randviir and C. E. Banks, *Anal. Methods*, 2013, **5**, 1098–1115.
- 32 P. D. Cunningham, K. M. McCreary, A. T. Hanbicki, M. Currie, B. T. Jonker and L. M. Hayden, *J. Phys. Chem. C*, 2016, **120**, 5819–5826.
- 33 S. Sim, J. Park, J.-G. Song, C. In, Y.-S. Lee, H. Kim and H. Choi, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **88**, 075434.
- 34 P. D. Cunningham, A. T. Hanbicki, K. M. McCreary and B. T. Jonker, *ACS Nano*, 2017, **11**, 12601–12608.
- 35 E. A. A. Pogna, M. Marsili, D. De Fazio, S. Dal Conte, C. Manzoni, D. Sangalli, D. Yoon, A. Lombardo, A. C. Ferrari, A. Marini, G. Cerullo and D. Prezzi, *ACS Nano*, 2016, **10**, 1182–1188.
- 36 P. Rivera, J. R. Schaibley, A. M. Jones, J. S. Ross, S. Wu, G. Aivazian, P. Klement, K. Seyler, G. Clark, N. J. Ghimire, J. Yan, D. G. Mandrus, W. Yao and X. Xu, *Nat. Commun.*, 2015, **6**, 6242.
- 37 S. H. Kim, K. Hong, W. Xie, K. H. Lee, S. Zhang, T. P. Lodge and C. D. Frisbie, *Adv. Mater.*, 2013, **25**, 1822–1846.
- 38 F. Wang, P. Stepanov, M. Gray, C. N. Lau, M. E. Itkis and R. C. Haddon, *Nano Lett.*, 2015, **15**, 5284–5288.
- 39 D. Braga, I. Gutiérrez Lezama, H. Berger and A. F. Morpurgo, *Nano Lett.*, 2012, **12**, 5218–5223.
- 40 I. Ferain, C. A. Colinge and J.-P. Colinge, *Nature*, 2011, **479**, 310–316.
- 41 S. M. Sze and K. K. Ng, 2006, *Physics of Semiconductor Devices*, Wiley VCH-Verlag GmbH, 3rd edn.

