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## A two-dimensional semiconducting Cu(I)-MOF for binder and conductive additive-free supercapattery†

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Two-dimensional (2D) semiconductive metal–organic frameworks (MOFs) have emerged as potential electrode materials for electrochemical energy devices, combining both the merits of a battery and supercapacitor owing to their porous structure containing redox active metal centers. These electrode materials often suffer from poor conductivity and low cyclic stability and often require conductive additives. Herein, we report the synthesis of a CuCN-based semiconducting MOF (CuCN-MOF) formed from the reaction of CuCN with a nitrogen-rich multidentate 2-(1*H*-1,2,4-triazol-1-yl)pyridine (2TzPy) linker. The ligand with multiple binding sites along with  $\mu_2$ -C:N and  $\mu_3$ -C:C:N coordination modes of cyanide forms a 2D structure, which further generates a 3D architecture due to  $\pi$ - $\pi$  stacking with cyanide and triazole rings of different layers. Due to the presence of redox-active Cu(I) centers, porous structure, and conducting nature, the applicability of materials has been investigated as a binder and conductive additive-free electrode material. The investigation confirms the supercapattery behavior of the material with a specific capacity of 508 C g<sup>-1</sup> at 1 A g<sup>-1</sup>. Further, the symmetric device formed by CuCN-MOF delivers outstanding energy and power densities ( $^{max}E_d = 68.175 \text{ W h kg}^{-1}$ ,  $^{max}P_d = 5.54 \text{ kW kg}^{-1}$ ) with exceptionally high cyclic stability (96.5% up to 10 000 cycles).

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## Introduction

Energy consumption is rapidly increasing with the growing population. It is now necessary to move towards renewable energy sources because of depleting fossil fuels and the alarming rise in CO<sub>2</sub> levels. The transition from fossil fuel to renewable sources needs effective energy storage systems to support the uninterrupted energy supply.<sup>1,2</sup> Therefore, there is much interest in developing high-performance energy storage devices, such as batteries and supercapacitors. Batteries are energy storage devices having high energy density due to the involvement of charge transfer across the electrode–electrolyte interface; however, since these processes are relatively slow, batteries exhibit poor power density.<sup>3,4</sup> On the other hand, charge storage in supercapacitors occurs due to the simple adsorption of the electrolyte ions on the electrode surface or due to surface redox reactions. Though the amount of energy

stored by simple ion adsorption/desorption on the electrode surface is much less than through charge transfer processes, these processes are much faster than the charge transfer process. Therefore, supercapacitors feature low energy density but very high power density.<sup>5–8</sup> Attempts have been made to improve the energy density of supercapacitors by invoking charge transfer-based energy storage, which has led to a wide range of supercapacitor types.

Based on the charge–discharge mechanisms, supercapacitors are classified as electrical double-layer capacitors (EDLCs) and pseudocapacitors. Pseudocapacitors store charge as a result of surface redox reactions, whereas EDLCs store charge due to the accumulation of ions at the electrode/electrolyte interface, where no electron transfer is involved.<sup>9,10</sup> Supercapacitors also offer the advantages of fast discharging, long cyclic stability, being lightweight, and simple to use.<sup>11,12</sup> Recently, efforts have been made to combine the properties of both supercapacitors and batteries in a single electrode material to develop a new category called “supercapattery” that combines the benefits of high energy as well as high power density.<sup>13–17</sup> This idea has led to the exploration of new types of electrode materials, such as metal oxides/sulfides, MXenes and their composites with different morphologies.<sup>18–23</sup> In addition, efforts have also been made to develop various nanostructured porous electrode materials containing redox active metal centers to attain a greater specific surface area.<sup>24–26</sup> Such

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materials offer charge storage by both double-layer formation as well as faradaic processes and, therefore, render improved energy storage performance.<sup>27</sup> In this regard, metal–organic frameworks (MOFs) have been widely investigated as positive electrode materials for lithium and sodium ion batteries and supercapacitors due to their controllable structure, high specific surface area, and high porosity.<sup>28–32</sup> The channeled structures of MOFs function as conduits for seamless ion transport within the electrode, facilitating high charge storage and high power density.<sup>33</sup> It has been observed that three-dimensional (3D) MOFs formed by a 2D layer through non-covalent interaction offer high surface area and thus provide a greater electrode–electrolyte interface, resulting in improved charge storage performance.<sup>34</sup> Moreover, the extended  $\pi$ -conjugation in MOFs enables them to exhibit better conductivity than other cathode materials, making them promising electrode materials.<sup>29,35–40</sup>

Among different MOFs, Cu-based MOF have been extensively explored as electrode materials because of the high natural abundance and low cost of Cu precursors. For example, Chen *et al.* reported a Cu(II) based 3D MOF, featuring a gravimetric capacitance of  $479 \text{ F g}^{-1}$  at  $0.2 \text{ A g}^{-1}$  with 80% capacitance retention after 2000 cycles.<sup>41</sup> Liu *et al.* reported a Cu(II) based electrode material  $[\text{Cu}(\text{hmt})(\text{tfbdc})(\text{H}_2\text{O})]$  (hmt = hexamethylenetetramine, tfbdc = 2,3,5,6-tetrafluoroterephthalate), showing a very high specific capacitance of  $1274 \text{ F g}^{-1}$  at  $1 \text{ A g}^{-1}$  with cyclic stability of 88% after 2000 cycles.<sup>42</sup> Our group also fabricated an rGO composite of Cu(II)-MOF, assembled using a piperazine-based tetracarboxylic linker, which shows a specific capacity of  $867.09 \text{ F g}^{-1}$  at  $1 \text{ A g}^{-1}$  with ultrahigh cyclic stability (135% capacitance after 10 000 cycles). The material's exceptional stability feature is attributed to the presence of the nitrogen-rich piperazine units, which increases the wettability of the electrode material.<sup>43</sup> We observed that the similar rGO/Cu(II) MOF, formed by a tetracarboxylic linker without any nitrogen, exhibited only 83.3% capacity retention after the same number of cycles.<sup>44</sup>

Though the applications of Cu(II)-based MOFs are ubiquitous as electrode materials for energy storage, Cu(I)-based MOFs have rarely been explored despite the fact that electron-rich Cu(I)-complexes have a greater tendency to form clusters, which can tolerate structural changes during charging and discharging, and are therefore expected to impart improved cyclic stability.<sup>45,46</sup> Wu *et al.* have reported polyoxometalate based copper(I) MOFs that showed a gravimetric capacitance of  $828 \text{ F g}^{-1}$  at  $1 \text{ A g}^{-1}$  current density.<sup>47</sup> Recently, Zhu *et al.* have reported a CuCN-based MOF as an electrode material for a supercapattery with a specific capacity of  $677.5 \text{ C g}^{-1}$  at  $1 \text{ A g}^{-1}$  current density.<sup>48</sup>

In the continuation of research involving the design and synthesis of these interesting MOFs, herein, we report on the synthesis of a new 2D copper(I) cyanide-based semiconductive MOF using a 1*H*-1,2,4-triazole substituted pyridine ligand as a linker. The superior conductivity and the presence of Cu(I) redox-active centers in its 2D structure, which eventually transform into an interesting 3D-framework with suitable channel structure due to  $\pi$ - $\pi$  stacking, encouraged us to investigate it as

a binder and conducting additive-free electrode material in energy storage devices. Moreover, the polymer also consists of a nitrogen-rich linker, a desirable feature to derive improved charge storage performance as established previously.<sup>43</sup> The electrochemically inactive binder and conducting additive do not play a part in the charge storage process, thus diminishing the gravimetric capacity of the device.<sup>49–53</sup> Therefore, materials that can be cast on the current collector in neat form without compromising device cyclic stability are of much interest. The present MOF shows a specific capacity of  $508 \text{ C g}^{-1}$  at  $1 \text{ A g}^{-1}$  with excellent capacity retention of 92% after 10 000 cycles. Additionally, the symmetrical device fabricated using the metal–organic framework (MOF//MOF) features a maximum energy density of  $68.175 \text{ W h kg}^{-1}$  and power density of  $5.54 \text{ kW kg}^{-1}$  with exceptional device stability up to 96.5% after 10 000 cycles.

## Experimental section

### Synthesis of CuCN-MOF

In an oven-dried Schlenk tube, 2-(1*H*-1,2,4-triazol-1-yl)pyridine (2TzPy) (0.34 mmol, 50 mg) was dissolved in a mixture of solvent ACN : DCM (1 : 1, 6 mL) and stirred at  $65^\circ\text{C}$  for 5 min, then CuCN (1.36 mmol, 123 mg) was dispersed in 10 mL ethanol solution. The reaction mixture was allowed to stir for a further 72 hours at  $65^\circ\text{C}$  which gives a yellow precipitate. The precipitate was filtered and washed with DCM ( $5 \text{ mL} \times 2$ ) and ACN ( $5 \text{ mL} \times 2$ ) and dried under vacuum for 6 hours. Yield = 92%.

CuCN-MOF was dissolved in acetonitrile by heating and yellow-coloured crystals, suitable for single-crystal X-ray analysis, were obtained by keeping the solution at room temperature for several days.

## Characterisation

### Physicochemical characterisation

The phase purity of the coordination polymer was confirmed using an Empyrean Malvern Panalytical X-ray diffractometer with monochromatic Cu  $K\alpha$  (0.1540 nm) radiation in the  $2\theta$  range of  $5$ – $50^\circ$ . The Brunauer–Emmett–Teller (BET) surface area and Barrett–Joyner–Halenda (BJH) pore size distribution were calculated using  $\text{N}_2$  gas adsorption–desorption on an Autosorb iQ (Quantachrome Instruments) and the data was collected using AsiQwin software. Attenuated total reflectance Fourier transform infrared spectroscopy (ATR FT-IR) was performed on a Bruker Alpha II spectrophotometer of the powdered sample in the range of  $4000$ – $400 \text{ cm}^{-1}$ . Morphological characterization through field emission scanning electron microscopy (FE-SEM) was done by using a JEOL JSM-7400F.

### Electrode fabrication for electrochemical analysis

Firstly, slurries were made by sonicating 10 mg of electrode materials in 1 mL of ethanol for 2 hours. Simultaneously,  $1 \times 1 \text{ cm}^2$  of Ni-foam stripes were cut and cleaned properly before further use. Thereafter, the as-prepared slurries with a concentration of 10 mg of electrode material in 1 mL of ethanol were uniformly coated on the nickel foam with a net electrode

material loading of 1 mg. The coated electrodes were dried in a vacuum for 4 hours. 1 M KOH electrolyte solution was used for all the electrochemical tests.

### Electrochemical studies

**Charge storage study.** All electrochemical tests were done at room temperature on a NOVA software-controlled Autolab PGSTAT 204N. Ag/AgCl was used as the reference electrode, platinum as the counter electrode, and Ni-foam coated with the electrode material was used as the working electrode. Cyclic voltammetry (CV) at different scan rates and galvanostatic charge/discharge (GCD) at different current densities were conducted using an electrochemical workstation. Also, electrochemical impedance spectroscopy (EIS) was conducted at 10 mV AC amplitude in the frequency range of 10 mHz to 100 kHz. Using the GCD discharge curve, the following equation was used to determine the specific capacities.

$$C = \frac{I\Delta t}{m} \quad (1)$$

where  $C$  is specific capacity ( $\text{C g}^{-1}$ ),  $I$  is current,  $\Delta t$  is the discharge time, and  $m$  is the mass of the coated active material (g).

**Fabricating a symmetrical solid-state device.** A Swagelok cell with a diameter of 3 cm was used to fabricate the solid-state symmetrical charge storage device. An electrode slurry was prepared by sonicating 10 mg of electrode materials in 1 mL ethanol solution for 2 h. The resulting slurry was coated on a circular Ni-foam cutting and then placed in an oven at 60 °C for 8 h. The electrolyte solution was prepared by adding 1 g of PVA in 10 mL of water; subsequently, the resultant solution was heated at 90 °C for 3 h with continuous stirring, followed by the addition 10 mL of 1 M KOH solution. Thereafter, the as-prepared electrodes were dipped in PVA-KOH gel electrolyte and dried partially. Afterward, two such electrodes were assembled in a Swagelok cell using cellulose paper as a separator to complete the device's fabrication.

From the GCD curves, the following equations were used to derive the specific capacity, energy density ( $E$ ), and power density ( $P$ ) of the fabricated device, where,  $m$  is coated.

$$C = \frac{2I\Delta t}{m} \quad (2)$$

$$E\left(\frac{\text{W h}}{\text{kg}}\right) = \frac{0.5 \times C\Delta V^2}{3.6} \quad (3)$$

$$P\left(\frac{\text{W}}{\text{kg}}\right) = \frac{3600 \times E}{\Delta t} \quad (4)$$

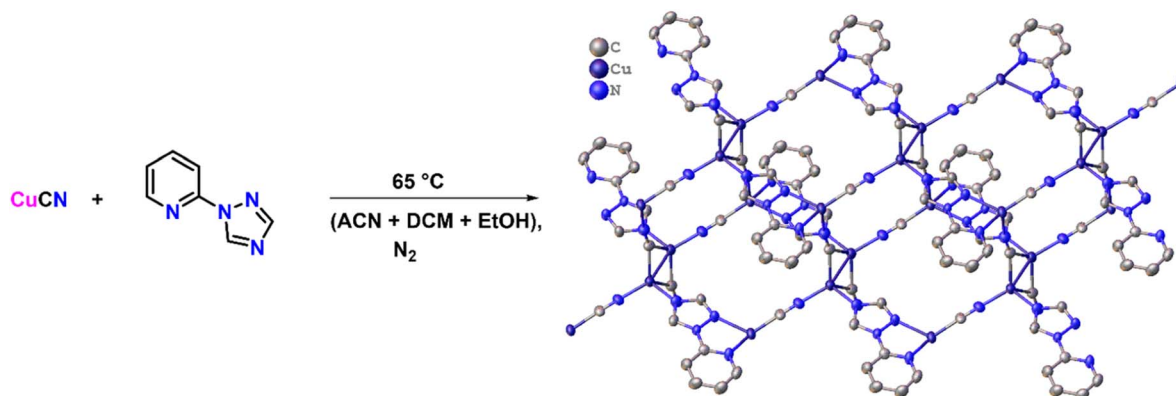
active material mass (g),  $\Delta t$  is discharge time (s),  $I$  is current (A),  $\Delta V$  is potential range (V), and  $C$  is specific capacity ( $\text{C g}^{-1}$ ).

## Results and discussion

To synthesize the **CuCN-MOF**, 2TzPy ligand was selected, which has three coordination sites (Scheme 1). CuCN was reacted with the ligand in a 4 : 1 ratio to get the MOF in 92%. Reaction in different metal:ligand ratios and in different sets of solvents yielded the same product, with some impurities. Crystals suitable for single crystal X-ray analysis were grown by slow evaporation of acetonitrile solution. The crystallographic parameters, selected bond lengths, and bond angles are presented in Tables S1–S3.†

### Structural analysis

The 2D **CuCN-MOF** crystallizes in a triclinic system with  $P\bar{1}$  space group consisting of two crystallographically independent copper(I) ions, two cyanides, and a ligand in the asymmetric unit (Fig. 1a). The crystal structure analysis reveals the occurrence of a dimer [ $\{\text{Cu}(\mu_3\text{-CN})_2\text{Cu}\}$ ] and monomer  $\text{Cu}(\mu_2\text{-CN})$  secondary building units (SBUs) in the polymeric structure. Recently, we have reported one-dimensional Cu(I) CP with the same ligand, where only one nitrogen of the triazole ring is coordinated with the metal center.<sup>54</sup> Interestingly, all the nitrogens are coordinated with the Cu(I) center and the polymeric structure shows the presence of two different bonding modes of cyanides;  $\mu_2\text{-C:N}$  and  $\mu_3\text{-C:C:N}$ . Though the occurrence of different coordination modes of cyanide in a single MOF is unique,  $\mu_2\text{-C:N}$  and  $\mu_3\text{-C:C:N}$  have been previously reported in CuCN-CPs by the Zhou, Xing Li, and Giordana groups.<sup>55–57</sup> The Cu1 adopts tetrahedral geometry, coordinated by triazolyl N4,



Scheme 1 Reaction scheme for the synthesis of the 2D MOF.

two  $\mu_3$ -cyanide C8 and two  $\mu_2$ -cyanide N5 atoms. The Cu2 also has a tetrahedral geometry, coordinated by N1 and N3 of ligand, a  $\mu_3$ -cyanide N6, and a  $\mu_2$ -cyanide C9 atom (Fig. 1b). The bridged ligand and cyanides lead to a 2D polymeric structure (Fig. S1†).

The Cu1–Cu1<sup>1</sup> bond distance is 2.447(17) Å, shorter than the sum of van der Waals radii, and is close to the previously reported for Cu2– $\mu_2$ C:N units.<sup>58,59</sup> The bond lengths of Cu1–N4 and Cu1–N5 are 2.070(5) and 1.930(5), and those of Cu2–N1<sup>2</sup>, Cu2–N3<sup>2</sup>, and Cu2–N6<sup>3</sup> are 2.134(5), 2.153(5), and 2.004(6) Å, respectively along the symmetry operation (2) + x, 1 + y, –1 + z, (3) + x, 1 + y, + z. The bond angles between Cu and  $\mu_3$ -C:N: Cu1–C8–N6 and Cu1<sup>1</sup>–C8–N6 are 151.4(7) and 138.9(6), respectively. Further, the bond angles between Cu and  $\mu_2$ -C:N: Cu2–C9–N5 and Cu1–N5–C9 are 174.5(6) and 177.4(6)° suggesting a linear arrangement. The C–N bond lengths of 1.153(9) and 1.161(8) Å for C8–N6 and C9–N5 are close to the reported C–N bond length in **CuCN-MOFs**.<sup>60</sup> The bond length of Cu1<sup>1</sup>–C8 is 2.281(8) whereas the Cu2–C9 bond length is 1.874(6) Å. The Cu–Cu distances between each rectangular face are shown in Fig. S2.†

Furthermore, the building blocks extend to form a 16-membered metallacycle with a zig-zag 1D chain consisting of six copper, two CN ions, and six atoms of the organic ligand, with another symmetrical zig-zag chain cross-over to form a rectangular layer with a size calculated to be 11.09 × 9.05 Å<sup>2</sup> (Fig. S3†). These rectangular layers are separated by a distance of 4.952 and 5.072 Å (two different distances of the  $\mu_3$ -C:C:N unit) with a CN bridge along the *a*-axis. The largest metallacycle present in **CuCN-MOF** is a 20-membered hexagon consisting of six copper, four CN ions and six atoms of the organic ligand (Fig. S4†). The third type of cycle involves 12 atoms from four copper and four cyanides. These layers combine to form a 2D polymeric chain, where each unit forms a cuboid with one missing edge (Fig. S2a†) and three rectangular faces with a diagonal length range from 8.289 to 9.835 Å (Fig. S2b†). These cuboids stack with each other, having opposite faces connected alternatively with each side of the edge to form the 2D **CuCN-MOF** (Fig. 2a). This **CuCN-MOF** arranges in a way such that the 2TzPy unit occupies the peripheral position where the centroid

distance between the pyridinic ring is 8.209 Å and the Cu(I) metal center is at the inner side of the MOF (Fig. 2b). The centroid distances between two pyridyl/triazolyl rings are the same, confirming that both the rings are in the same plane. The monolayer of the 2D MOF has multiple voids present with different sizes and the crystalline material has a density of 1.974 g cm<sup>–3</sup>. The packing arrangement of 2D **CuCN-MOF** is shown in Fig. S5,† where different layers are stacked by the  $\pi$ – $\pi$  interaction between the CN anions and triazolyl rings of the 2TzPy unit with a distance of 3.215 Å (Fig. S6†). The crystal structures along the *b* and *c* axes confirm the 2D structure of the **CuCN-MOF** (Fig. S7†).

### PXRD, IR, NMR and TGA analysis

The formation of 2D **CuCN-MOF** was confirmed by single-crystal X-ray analysis. Further, we performed powder X-ray diffraction (PXRD) analysis to check the bulk property of the **CuCN-MOF** within the 2 $\theta$  range of 5–50°. The experimentally observed peaks from PXRD data correspond well with the simulated pattern obtained from SCXRD analysis (Fig. 3a). This confirms the phase purity of the bulk sample, and sharp peaks suggest high crystallinity. The crystallite size calculated from PXRD data is found to be 8.9 Å. The FT-IR of the MOF shows characteristic stretching frequencies of different cyanide ions at 2167, 2124 and 2079 cm<sup>–1</sup>. The bands at 1442–1602 cm<sup>–1</sup> correspond to the stretching frequencies of C–C and C–N of the pyridyl ring. The weak peaks at 1372, 1321, and 1292 cm<sup>–1</sup> correspond to the  $\nu$ (N–N) and  $\nu$ (C–N) of the triazole ring, respectively (Fig. 3b and S9†). Further, solid-state NMR of **CuCN-MOF** also confirms the presence of cyanides (broad signal between 191 and 230 ppm) and ligands (112–149 ppm) (Fig. S10†). Thermal stability is an important criterion for the aptness of a complex for a particular application. TGA was performed on the **CuCN-MOF** at a temperature range of 25–800 °C at a flow rate of 10 °C min<sup>–1</sup> (Fig. 3c). The TGA graph suggests the MOF's stability up to 168 °C, and a subsequent rise in temperature causes a stepwise degradation of the organic ligands up to 420 °C.

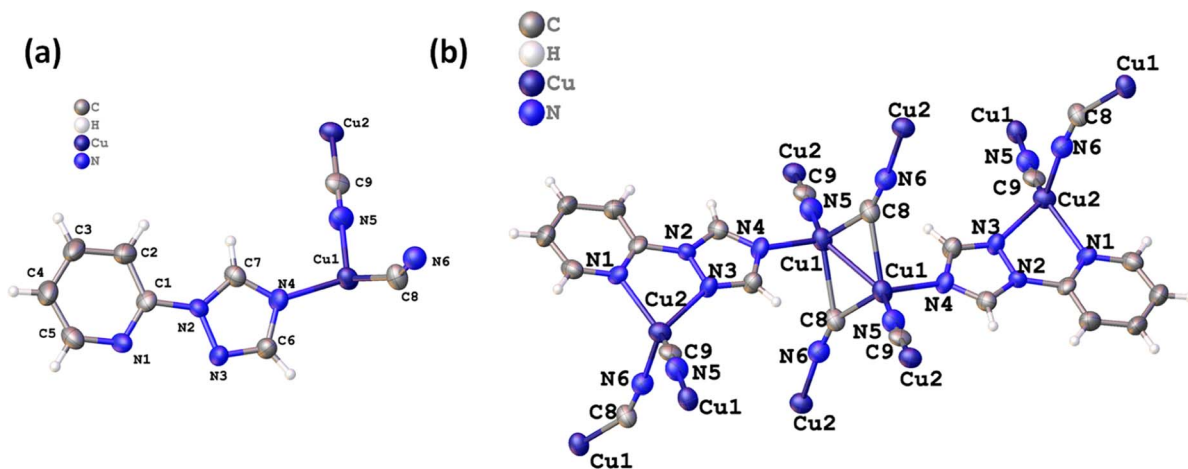


Fig. 1 (a) Asymmetric unit of the 2D MOF, and (b) different coordination modes of CN<sup>–</sup> ions in chain propagation.

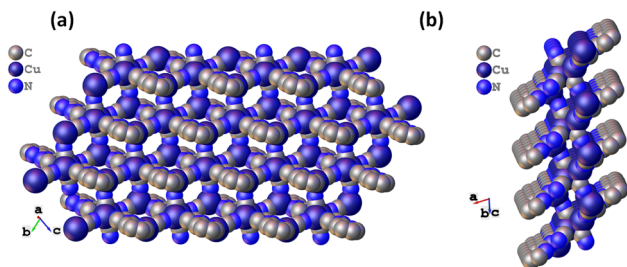


Fig. 2 (a) The space filling model of CuCN-MOF along the *a* axis, and (b) space filling model along the *b* axis.

### BET, SEM and electrical conductivity analysis

To confirm the specific surface area and pore size distribution of the **CuCN-MOF**, a Brunauer–Emmett–Teller (BET) isotherm was recorded using the N<sub>2</sub> adsorption/desorption method. Less surface area is expected for the MOF as the layers are densely packed in the crystal system by interlocking themselves between the spaces in the peripheral gaps forming a 3D arrangement. These interlocked layers are arranged in a staggered conformation in an ABAB stack sequence and minimize the effective surface area (Fig. S8†).<sup>57</sup> The BET isotherm represents the type IV isotherm suggesting a mesoporous structure with a surface area of 8.036 m<sup>2</sup> g<sup>−1</sup> (Fig. S11†). The pore size distribution was obtained by the Barrett–Joyner–Halenda (BJH) method, revealing a unimodal pore size distribution of 3.42 nm, confirming the mesoporous structure (Fig. S12†).

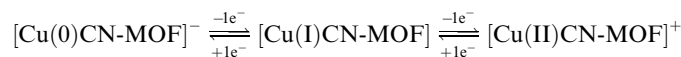
The microstructure and morphology of the 2D **CuCN-MOF** were characterized by field emission scanning electron microscopy (FE-SEM) (Fig. 4a–c) at various magnifications, which shows an irregular arrangement of various sheets with rough edges. Interestingly, the surface of 2D sheets of the MOF reveals a trough and crest pattern (Fig. 4c), which originates from its unique stacked crystalline structure (Fig. S8†). Such folding in the 2D sheets is advantageous from the charge storage point of view, as it facilitates the seamless diffusion of electrolytes throughout the bulk of the electrode material. Elemental mapping confirms the presence of Cu, C, and N elements in the material (Fig. 4d–f) with an appropriate ratio, which was confirmed by EDS analysis (Fig. 4g). The electrical

conductivity of **CuCN-MOF** was measured using a two-probe direct current method, at 300 K, on pressed pellets, in the voltage range of −10 V to +10 V. The linear (*I/V*) graph for the MOF validates its ohmic behaviour with a conductivity value of  $9.61 \times 10^{-7}$  S cm<sup>−1</sup> (Fig. S13†).

### Electrochemical study

To evaluate the supercapacitive efficacy of materials, cyclic voltammetry (CV), galvanostatic charge–discharge (GCD), and electrochemical impedance spectroscopy (EIS) were performed using an electrochemical workstation. In order to optimize the performance of the **CuCN-MOF** electrode, the working potential window was investigated methodically. To ascertain the optimal potential window (Fig. S14a–d†), cyclic voltammetry profiles of the electrode were recorded at a scan rate of 10 mV s<sup>−1</sup> in both the positive (0.0 to 0.5 V) and the negative potential windows (−0.5 to 0.0 V). The CV and GCD investigations reveal that the electrode material exhibits similar behaviour in terms of current density and discharge time; however, in the positive potential window, well-discernible redox peaks in CV and a smaller IR drop in GCD are observed (Fig. S14†). Based on these results, all subsequent electrochemical studies were conducted within the potential window of −0.5 to 0.5 V.

Fig. 5a depicts the CV profiles of **CuCN-MOF** recorded in the potential range −0.5 to 0.5 V at varying scan rates. The CV profile reveals the presence of two cathodic and anodic peaks originating from Cu redox centers suggesting a battery-type behaviour of the electrode material and also giving a dual electron charge transfer mechanism within the **CuCN-MOF**. Based on the redox steps involving four electron transfer and the formula weight of **CuCN-MOF** (empirical formula C<sub>9</sub>H<sub>6</sub>Cu<sub>2</sub>N<sub>6</sub>, formula weight 325.84), its theoretical capacity is calculated to be 1184.45 C g<sup>−1</sup>. The two cathodic and anodic peaks may be attributed to the redox steps given below:



Further, as expected, the scan rate increase leads to the current rise. Interestingly, variation in scan rate does not lead to any considerable shift in the peak position, highlighting the

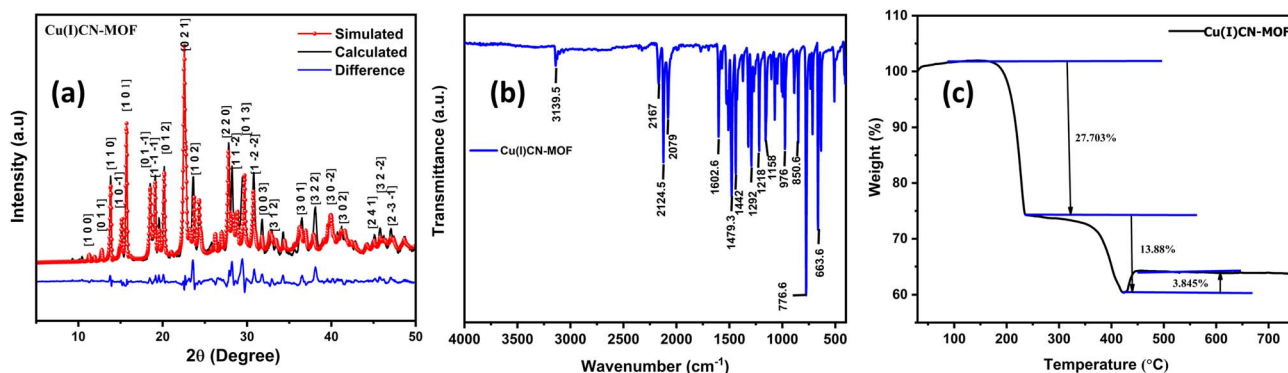


Fig. 3 (a) PXRD pattern and (b) IR spectrum of the MOF, and (c) TGA graph.

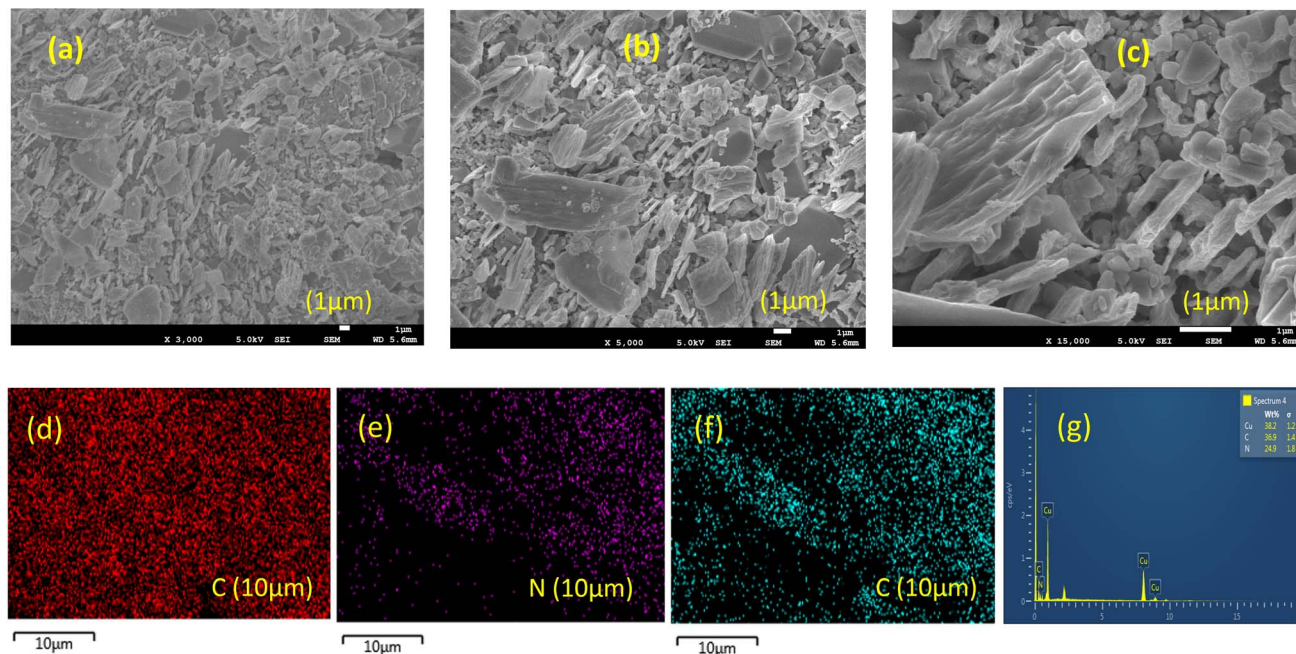


Fig. 4 SEM images at various magnifications (a)  $\times 3000$ , (b)  $\times 5000$  (c)  $\times 15\,000$  magnification, (d)–(f) EDS elemental mapping images, (g) EDS spectrum.

material's low inherent resistance with insignificant structural alteration during charging and discharging, which are desirable features for cyclic stability and better rate capability.

Further, to comprehend the charge storage mechanism, the power law method was used. As per this law, the peak current response in cyclic voltammetry (CV) is related to the scan rate ( $\nu$ ), as described by eqn (5).

$$i = a\nu^b \quad (5)$$

where  $a$  and  $b$  are constants for a system.

The slope of  $\log i$  vs.  $\log \nu$  provides the value of  $b$ , which describes the charge storage mechanism of an electrode. If  $b \approx 0.5$ , charge storage is considered diffusion-controlled as in battery-type materials, whereas if  $b \approx 1$ , charge storage is surface-controlled as in EDLC-type materials. The linear plot for the electrode material is shown in Fig. 5b, and the value of  $b$  was found to be 0.74. This suggests that the charge storage occurs due to the combined effect of faradaic and ion adsorption processes, suggesting a supercapattery behaviour of the material.

To evaluate the quantitative partial contributions of faradaic and surface-controlled processes to the overall charge storage capacity of the material, Dunn's method has been used. Dunn's equation expresses the current at a given potential as a function of the scan rate ( $\nu$ ) as a combination of surface-controlled and diffusion-controlled processes, as shown in eqn (6).

$$i(V) = k_1\nu + k_2\nu^{1/2} \quad (6)$$

where  $i(V)$  represents the current at a given potential and  $\nu$  represents the scan rate, while  $k_1$  and  $k_2$  represent the constants

for the surface- and diffusion-controlled processes and the values are listed in Table S4,<sup>†</sup> respectively.

Fig. 5c depicts the bar graph displaying the percentage capacity contributions of the surface and diffusion-controlled process to the overall capacity of the electrode at various scan rates. At lower scan rates, the overall charge storage is dominated by diffusion, *i.e.*, involving a faradaic process. However, as the scan rate increases, the diffusion-based charge storage contribution gradually decreases, whereas the charge storage due to the surface-controlled process gradually increases. At 50  $\text{mV s}^{-1}$  scan rate, both charge storage mechanisms almost equally contribute to the overall capacitance of the material. The decrease in diffusion capacity with increasing scan rate can be explained by the fact that, at higher scan rates, the diffusion rate is outpaced by the scan rate, causing surface capacity to predominate. For comparison purposes, the proportion of the diffusion process to the overall CV profile of the material at 10  $\text{mV s}^{-1}$  is shown in Fig. S15.<sup>†</sup>

The specific charge storage capacities of the electrode material were evaluated by galvanostatic charge–discharge (GCD) experiments at current densities ranging from 1 to 10  $\text{A g}^{-1}$  within a potential window of  $-0.5$  to  $0.5$  V. Fig. 5d shows the GCD profiles of the electrode at different current densities. It is observed that increasing the current density leads to a decrease in the discharging time (Fig. 5d). At 1  $\text{A g}^{-1}$  current density, the calculated specific capacity of the electrode was found to be 508  $\text{C g}^{-1}$ . The variation of the specific capacity of the material at different current densities is depicted in Fig. 5e, showing a rapid reduction in specific capacity up to 3  $\text{A g}^{-1}$ , followed by a gradual decay. This is mainly due to the fact that, at low current density, the charge storage is mainly dominated by

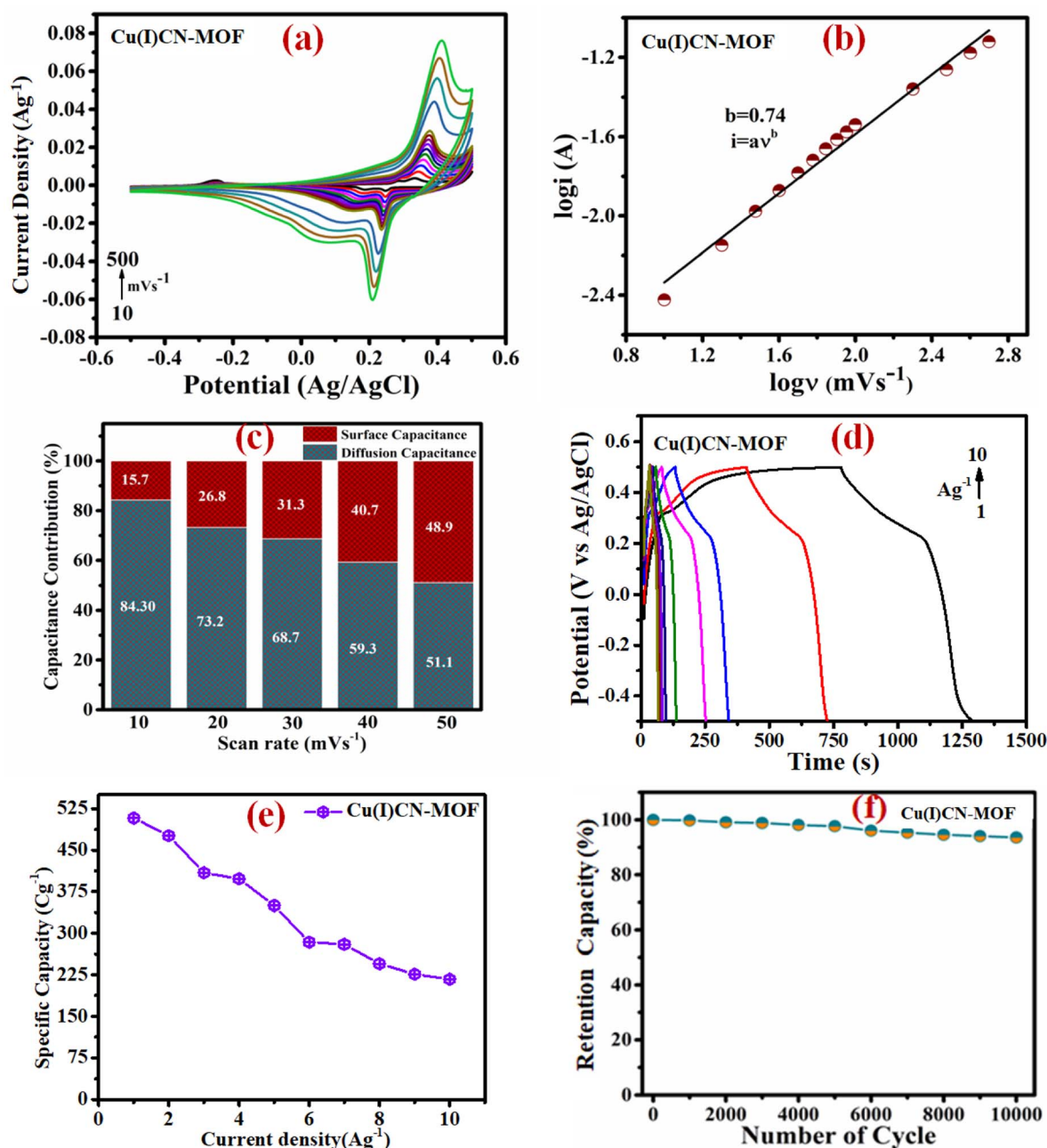


Fig. 5 (a) CV of CuCN-MOF at different scan rates. (b) Charge storage kinetics: log *i* vs. log *v* plots of CuCN-MOF. The slope of each plot represents the value of *b*. (c) Percentage surface and diffusion-controlled capacity contributions at different scan rates to the overall capacity of CuCN-MOF. (d) GCD at different current densities of CuCN-MOF. (e) Specific capacity value of CuCN-MOF at different current densities (1–10 A g<sup>-1</sup>). (f) Cyclic stability at 10 A g<sup>-1</sup> for 10 000 cycles.

a diffusion-controlled process, which wanes rapidly due to the limitation imposed by the ion diffusion. After 3 A g<sup>-1</sup> current density, the surface-controlled process has a significant contribution to the overall capacity, which is not much affected by the current density, resulting in a gradual reduction in the specific capacity.

The improved performances of CuCN-MOF can be explained on the basis of unique structural features. The 2D molecular structure is transformed into the supramolecular structure *via*  $\pi$ - $\pi$  interaction between the layers that help in facilitating the

ion transport and easing the charge transfer. Also, due to the densely packed structure and due to small bridging cyanide ligands, the redox-active centers  $[\{\text{Cu}(\mu_3\text{-CN})_2\text{Cu}\}]$  and  $\{\text{Cu}(\mu_2\text{-CN})\text{Cu}\}$  are closer, which facilitates the electron transfer process. This structural feature overall enhances the supercapattery performances. This idea is also supported by Jin *et al.*, who described that the high concentration of metal centers helps in the supercapacitor performances even for a nonporous coordination polymer CuAg<sub>4</sub>BHT.<sup>61</sup>

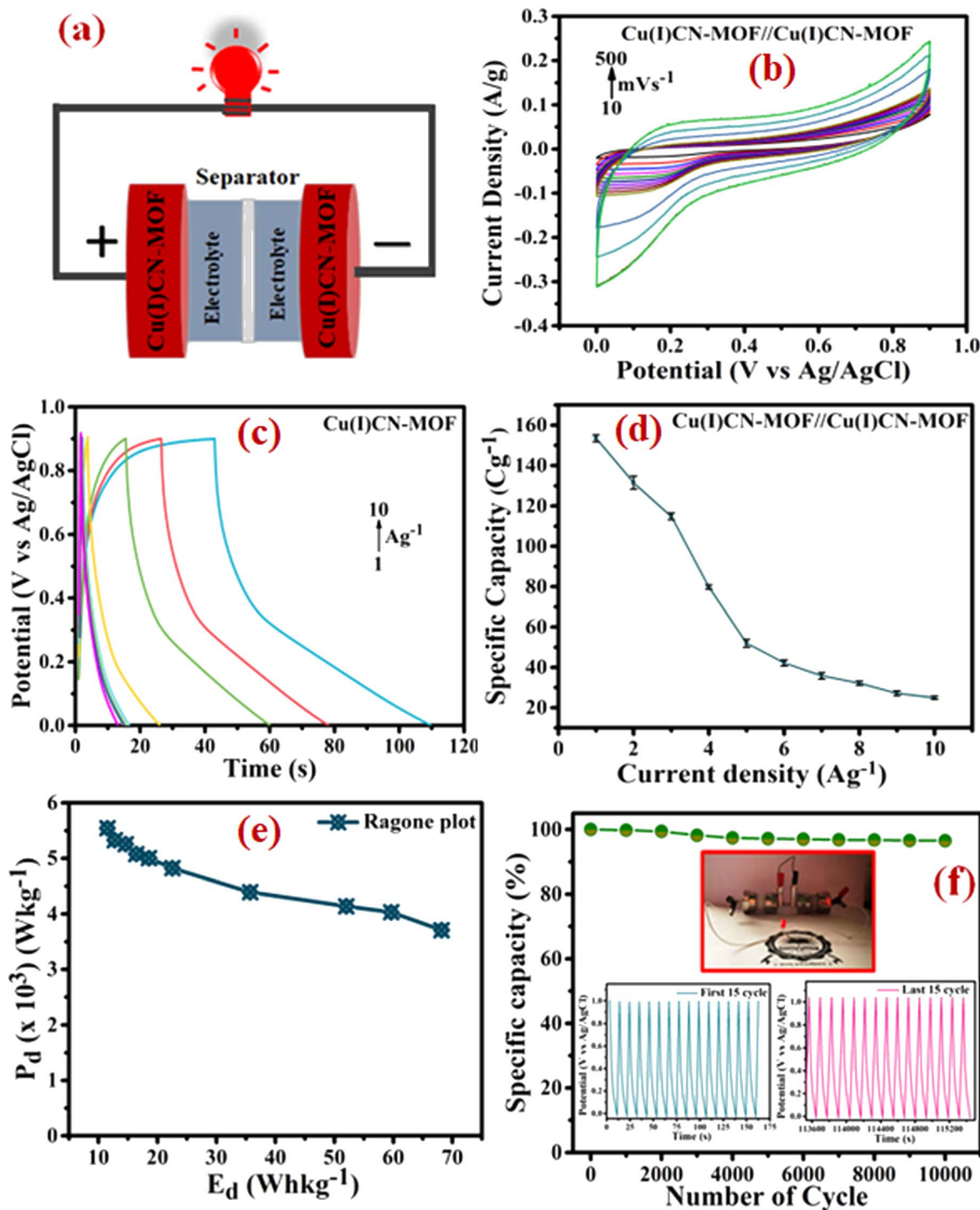


Fig. 6 (a) Schematic of the symmetric device, (b) CV profiles of the symmetric device at different scan rates (10–500  $\text{mV s}^{-1}$ ), (c) GCD profiles of the symmetric device at different current densities (1–10  $\text{A g}^{-1}$ ), (d) specific capacity variations of the device with changing current densities (1–10  $\text{A g}^{-1}$ ), (e) Ragone plot of energy density vs. power density of the device, (f) cycling performance of the device at 10  $\text{A g}^{-1}$  for 10 000 consecutive cycles (insets show glowing LED using the fabricated device).

The cyclic stability of the electrode material was assessed for 10 000 galvanostatic charge–discharge (GCD) cycles at 10  $\text{A g}^{-1}$ , as shown in Fig. 5f. After 10 000 cycles, the material shows excellent capacity retention of 92%. Even though the material was used as a binder-free electrode, no trace of electrode material stripping was observed after the long cyclability

experiment. This observation highlights the potential of the present coordination polymer as a binder-free electrode material for practical applications.

To assess the practical applicability of the electrode, a symmetric device with a cellulose membrane as a separator was assembled, as shown in Fig. 6a. Further, a Ni-foam coated

**CuCN-MOF** electrode device was characterized in a two-electrode configuration with a 0.9 V optimal cell voltage. The CV curves of the device at various scan rates (10 to 500 mV s<sup>-1</sup>) exhibited a non-rectangular shape that retained their characteristics even at higher scan rates (Fig. 6b), indicating the device's high-rate capability. GCD experiments were conducted at various current densities (1.0 to 10 A g<sup>-1</sup>) to evaluate the specific capacity of the device (Fig. 6c). The asymmetry in the GCD curves indicates that the overall charge storage is a combination of EDLC and faradaic processes. Similar to the three-electrode configuration, the calculated specific capacity is observed to decrease upon increasing the current density. At 1 A g<sup>-1</sup> current density, the device exhibits a specific capacity of 136.48 C g<sup>-1</sup>, which reduces to 23.0 C g<sup>-1</sup> at 10 A g<sup>-1</sup> (Fig. 6d). To ensure the charge storage capacity of the device is mainly because of the **CuCN-MOF**, we also evaluated the specific capacity of the device with bare Ni-foam as an electrode, which showed a charge storage capacity of 32 C g<sup>-1</sup> (Fig. S16†). This confirms that **CuCN-MOF** is the main contributor in the charge storage device.

To analyse the energy density ( $E_d$ ) and power density ( $P_d$ ) of the symmetric device, a Ragone plot was constructed (Fig. 6e). The plot shows a maximum energy density of 68.175 W h kg<sup>-1</sup> at a power density of 3.7 kW kg<sup>-1</sup> and a maximum power density of 5.54 kW kg<sup>-1</sup> at an energy density of 11.52 W h kg<sup>-1</sup>. Further, the cyclic stability of the device was determined by executing 10 000 GCD cycles at a current density of 10 A g<sup>-1</sup> within a potential window of 0–1.0 V. After 10 000 cycles, the device retained 96.5% of its original capacity, demonstrating exceptional cyclic stability (Fig. 6f). To demonstrate the device's practical application, two similar symmetric supercapacitor devices were connected in series to power a red light-emitting diode (LED) (inset of Fig. 6f). Overall, the energy storage device fabricated using **CuCN-MOF** reported in this paper exhibits outstanding electrochemical performance, including high-rate capability, cyclic stability, and energy and power densities, making it suitable for practical applications. A comparison of the performance of the present electrode material with some previously reported similar systems is summarised in Table S5.† The table clearly suggests that the binder-free electrode fabricated using **CuCN-MOF** presented in this work exhibits superior performance compared to most of the earlier reported MOFs in terms of cyclic stability, specific capacity, and energy and power densities.

## Conclusion

In summary, we have reported a facile one-pot synthesis of a 2D MOF using a nitrogen rich 2TzPy ligand with three different coordinating sites. The single crystal X-ray analysis shows that all the nitrogen sites are coordinated to the copper(i) metal center along with  $\mu_2$ -C:N and  $\mu_3$ -C:C:N coordination modes of cyanides, resulting in a 2D polymeric structure. Due to  $\pi$ - $\pi$  stacking of these 2D polymer structures, it further extends to a 3D architecture, offering suitable channels to facilitate the ion transport to the bulk of the material, making it a promising electrode material. Moreover, the extended  $\pi$ -conjugation provided by

CuCN leads to conductivity up to  $9.61 \times 10^{-7}$  S cm<sup>-1</sup>, suggesting the semiconductor behavior of the MOF. The electrochemical investigations suggest that the **CuCN-MOF** features an excellent charge storage capacity of 508 C g<sup>-1</sup> at 1 A g<sup>-1</sup> as a binder and conductive additive-free electrode with a cyclic stability of 92% after 10 000 cycles. The charge storage mechanism analysis reveals that the material exhibits a supercapattery type charge storage, combining the features of battery and EDLC type electrodes. The fabricated symmetric device shows an outstanding cyclic stability of 96.5% after 10 000 cycles, with high maximum energy and power densities of 68.175 W h kg<sup>-1</sup> and 5.54 kW kg<sup>-1</sup>, respectively. This work highlights the use of a Cu(i) MOF in developing cost-effective, highly stable, and binder and conductive additive-free electrode materials for high energy as well as high power density energy storage devices.

## Conflicts of interest

The authors declare no conflicts.

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## References

- 1 B. Chen, L. Xu, Z. Xie and W.-Y. Wong, *EcoMat*, 2021, **3**, e12106.
- 2 A. González, E. Goikolea, J. A. Barrena and R. Mysyk, *Renewable Sustainable Energy Rev.*, 2016, **58**, 1189–1206.
- 3 Poonam, K. Sharma, A. Arora and S. K. Tripathi, *J. Energy Storage*, 2019, **21**, 801–825.
- 4 K. V. G. Raghavendra, R. Vinoth, K. Zeb, C. V. V. Muralee Gopi, S. Sambasivam, M. R. Kummara, I. M. Obaidat and H. J. Kim, *J. Energy Storage*, 2020, **31**, 101652.
- 5 Y. Jiao, J. Pei, C. Yan, D. Chen, Y. Hu and G. Chen, *J. Mater. Chem. A*, 2016, **4**, 13344–13351.
- 6 B. Dunn, H. Kamath and J.-M. Tarascon, *Science*, 2011, **334**, 928–935.
- 7 P. Simon, Y. Gogotsi and B. Dunn, *Science*, 2014, **343**, 1210–1211.
- 8 P. Simon and Y. Gogotsi, *Nat. Mater.*, 2020, **19**, 1151–1163.
- 9 Z. Yu, L. Tetard, L. Zhai and J. Thomas, *Energy Environ. Sci.*, 2015, **8**, 702–730.
- 10 J. Kim, J. H. Ryu, M. Jang, S. Park, M. Kim, K. H. Lee, S. Choi, Y. Yoon, H.-K. Jung, S. S. Lee and K.-S. An, *Small Methods*, 2023, **7**, 2201539.
- 11 W. Li, J. Liu and D. Zhao, *Nat. Rev. Mater.*, 2016, **1**, 1–17.
- 12 T. K. Ghosh and G. R. Rao, *Dalton Trans.*, 2023, **52**, 5943–5955.

- 13 G. Zou, Z. Tian, V. S. Kale, W. Wang, S. Kandembeth, Z. Cao, J. Guo, J. Czaban-Jóźwiak, L. Cavallo, O. Shekhah, M. Eddaoudi and H. N. Alshareef, *Adv. Energy Mater.*, 2023, **13**, 2203193.
- 14 A. Khan, F. Shaheen, M. Roman, R. Ahmad, K. Mehboob and M. H. Aziz, *J. Energy Storage*, 2023, **58**, 106317.
- 15 X. Ma, X. Wei, Z. Wei, Y. Guo, W. Wang, C. Zhang and Z.-Y. Jiang, *Electrochim. Acta*, 2021, **385**, 138434.
- 16 K. Wang, Q. Li, Z. Ren, C. Li, Y. Chu, Z. Wang, M. Zhang, H. Wu and Q. Zhang, *Small*, 2020, **16**, 2001987.
- 17 K. Karuppasamy, D. Vikraman, S. Hussain, P. Santhoshkumar, R. Bose, P. Sivakumar, A. Alfantazi, J. Jung and H.-S. Kim, *Small*, 2022, **18**, 2107284.
- 18 H. Liu, X. Liu, S. Wang, H.-K. Liu and L. Li, *Energy Storage Mater.*, 2020, **28**, 122–145.
- 19 T. Brezesinski, J. Wang, S. H. Tolbert and B. Dunn, *Nat. Mater.*, 2010, **9**, 146–151.
- 20 G. Zhang and X. W. (David) Lou, *Sci. Rep.*, 2013, **3**, 1470.
- 21 M. Hu, H. Zhang, T. Hu, B. Fan, X. Wang and Z. Li, *Chem. Soc. Rev.*, 2020, **49**, 6666–6693.
- 22 N. M. Deyab, M. M. Taha and N. K. Allam, *Nanoscale Adv.*, 2022, **4**, 1387–1393.
- 23 B. L. Ellis, P. Knauth and T. Djenizian, *Adv. Mater.*, 2014, **26**, 3368–3397.
- 24 R. Liu, J. Duay and S. B. Lee, *Chem. Commun.*, 2011, **47**, 1384–1404.
- 25 Y. Fang, X.-Y. Yu and X. W. (David) Lou, *Matter*, 2019, **1**, 90–114.
- 26 S. Bi, H. Banda, M. Chen, L. Niu, M. Chen, T. Wu, J. Wang, R. Wang, J. Feng, T. Chen, M. Dincă, A. A. Kornyshev and G. Feng, *Nat. Mater.*, 2020, **19**, 552–558.
- 27 C. Largeot, C. Portet, J. Chmiola, P.-L. Taberna, Y. Gogotsi and P. Simon, *J. Am. Chem. Soc.*, 2008, **130**, 2730–2731.
- 28 B. Y. Guan, X. Y. Yu, H. B. Wu and X. W. (David) Lou, *Adv. Mater.*, 2017, **29**, 1703614.
- 29 D. Shen, Y. Sha, C. Chen, X. Chen, Q. Jiang, H. Liu, W. Liu and Q. Liu, *Dalton Trans.*, 2023, **52**, 7079–7087.
- 30 C. Dong and L. Xu, *ACS Appl. Mater. Interfaces*, 2017, **9**, 7160–7168.
- 31 J. Liu, M. Zheng, S. Wu and L. Zhang, *Coord. Chem. Rev.*, 2023, **483**, 215084.
- 32 Y. Wu, Y. Chen, M. Tang, S. Zhu, C. Jiang, S. Zhuo and C. Wang, *Chem. Commun.*, 2019, **55**, 10856–10859.
- 33 S. Bi, H. Banda, M. Chen, L. Niu, M. Chen, T. Wu, J. Wang, R. Wang, J. Feng, T. Chen, M. Dincă, A. A. Kornyshev and G. Feng, *Nat. Mater.*, 2020, **19**, 552–558.
- 34 K. Wang, S. Wang, J. Liu, Y. Guo, F. Mao, H. Wu and Q. Zhang, *ACS Appl. Mater. Interfaces*, 2021, **13**, 15315–15323.
- 35 T. Chen, J.-H. Dou, L. Yang, C. Sun, N. J. Libretto, G. Skorupskii, J. T. Miller and M. Dincă, *J. Am. Chem. Soc.*, 2020, **142**, 12367–12373.
- 36 X. Huang, S. Zhang, L. Liu, L. Yu, G. Chen, W. Xu and D. Zhu, *Angew. Chem.*, 2018, **130**, 152–156.
- 37 Y. Chen, M. Tang, Y. Wu, X. Su, X. Li, S. Xu, S. Zhuo, J. Ma, D. Yuan, C. Wang and W. Hu, *Angew. Chem., Int. Ed.*, 2019, **58**, 14731–14739.
- 38 Y. Chen, Q. Zhu, K. Fan, Y. Gu, M. Sun, Z. Li, C. Zhang, Y. Wu, Q. Wang, S. Xu, J. Ma, C. Wang and W. Hu, *Angew. Chem.*, 2021, **133**, 18917–18924.
- 39 K. Fan, C. Fu, Y. Chen, C. Zhang, G. Zhang, L. Guan, M. Mao, J. Ma, W. Hu and C. Wang, *Advanced Science*, 2023, **10**, 2205760.
- 40 K. Fan, J. Li, Y. Xu, C. Fu, Y. Chen, C. Zhang, G. Zhang, J. Ma, T. Zhai and C. Wang, *J. Am. Chem. Soc.*, 2023, **145**, 12682–12690.
- 41 J. Liu, Y. Zhou, Z. Xie, Y. Li, Y. Liu, J. Sun, Y. Ma, O. Terasaki and L. Chen, *Angew. Chem., Int. Ed.*, 2020, **59**, 1081–1086.
- 42 Q. Liu, X. Liu, C. Shi, Y. Zhang, X. Feng, M.-L. Cheng, S. Su and J. Gu, *Dalton Trans.*, 2015, **44**, 19175–19184.
- 43 S. Krishnan, A. K. Gupta, M. K. Singh, N. Guha and D. K. Rai, *Chem. Eng. J.*, 2022, **435**, 135042.
- 44 M. K. Singh, A. K. Gupta, S. Krishnan, N. Guha, S. Marimuthu and D. K. Rai, *J. Energy Storage*, 2021, **43**, 103301.
- 45 A. Schlachter and P. D. Harvey, *J. Mater. Chem. C*, 2021, **9**, 6648–6685.
- 46 J. Troyano, F. Zamora and S. Delgado, *Chem. Soc. Rev.*, 2021, **50**, 4606–4628.
- 47 K. Wang, Z. Wang, S. Wang, Y. Chu, R. Xi, X. Zhang and H. Wu, *Chem. Eng. J.*, 2019, **367**, 239–248.
- 48 Z.-R. Zhang, Z.-H. Ren, C.-Y. Luo, L.-J. Ma, J. Dai and Q.-Y. Zhu, *Inorg. Chem.*, 2023, **62**, 4672–4679.
- 49 J. Kang, S. Zhang and Z. Zhang, *Adv. Mater.*, 2017, **29**, 1700515.
- 50 T. Jin, Q. Han and L. Jiao, *Adv. Mater.*, 2020, **32**, 1806304.
- 51 K. Shen, S. Zhai, S. Wang, Q. Ru, X. Hou, K. San Hui, K. Nam Hui and F. Chen, *Batteries Supercaps*, 2021, **4**, 860–880.
- 52 J. Ni, X. Zhu, Y. Yuan, Z. Wang, Y. Li, L. Ma, A. Dai, M. Li, T. Wu, R. Shahbazian-Yassar, J. Lu and L. Li, *Nat. Commun.*, 2020, **11**, 1212.
- 53 K. Shen, Z. Zhang, S. Wang, Q. Ru, L. Zhao, L. Sun, X. Hou and F. Chen, *Energy Fuels*, 2020, **34**, 8987–8992.
- 54 S. Mishra, D. Pandey, K. Mishra, L. Viau and A. Raghuvanshi, *New J. Chem.*, 2023, **47**, 19751–19759.
- 55 S.-W. Liang, M.-X. Li, M. Shao and Z.-X. Miao, *Inorg. Chem. Commun.*, 2006, **9**, 1312–1314.
- 56 A. Giordana, E. Priola, G. Gariglio, E. Bonometti, L. Operti and E. Diana, *Polyhedron*, 2021, **198**, 115059.
- 57 H. Xu, B.-Y. Zhou, K. Yu, Z.-H. Su, Z.-F. Zhao, W.-L. Sun and B.-B. Zhou, *New J. Chem.*, 2015, **39**, 1301–1307.
- 58 X. Liu, G.-C. Guo, A.-Q. Wu, L.-Z. Cai and J.-S. Huang, *Inorg. Chem.*, 2005, **44**, 4282–4286.
- 59 G.-L. Wang, M.-L. Wang, W. Zhong, M. Afzal, A. Alarifi, S. Singh, A. Kumar and J. Jin, *J. Mol. Struct.*, 2022, **1264**, 133332.
- 60 Y.-L. Qin, H. Sun, Y. Jing, X.-P. Jiang, G.-F. Wang and J.-F. Qin, *Acta Crystallogr., Sect. C: Struct. Chem.*, 2019, **75**, 1517–1523.
- 61 Y. Jin, S. Wu, Y. Sun, Z. Chang, Z. Li, Y. Sun and W. Xu, *Mater. Horiz.*, 2023, **10**, 3821–3829.