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A novel negative thermal expansion material of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$

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A novel negative thermal expansion (NTE) material of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ was synthesized. Its microstructure, crystal structure and thermal expansion properties were studied in detail. Structure analysis shows that $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ possesses a cubic crystal structure with the space group $Pa\bar{3}$ (no. 205) at room temperature (RT). The experimental results show that $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ exhibits a stable isotropic negative thermal expansion property and maintains a cubic structure from 103 to 773 K without any phase transition. The average linear coefficients of thermal expansion (CTEs) are calculated to be -3.75×10^{-6} K⁻¹ (163–673 K) and -4.50×10^{-6} K⁻¹ (RT-773 K). The intrinsic linear CTEs are calculated to be -3.84×10^{-6} K⁻¹ (103–473 K) and -4.22×10^{-6} K⁻¹ (RT-673 K). And the phase transition temperature of ZrV_2O_7 has been reduced below RT, this may be attributed to the incorporation of the cubic $ZrMo_2O_8$ structure and partially occupied Zr atoms (4a). The NTE property of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ can be attributed to the quasi-rigid unit modes (QRUMs).

1 Introduction

Since Sleight et al.1 reported that the cubic ZrW2O8 exhibits isotropic negative thermal expansion (NTE) from 0.3 to 1050 K in 1995, the NTE materials have received extensive attention. NTE materials have potential application in the field of precise optics, aerospace, biomedical, integrated circuits and sensors due to their unique thermal expansion properties.²⁻⁴ However, the existing NTE materials still have many deficiencies from the perspective of application requirements. For instance, the typical NTE material ZrW2O8 is built of an open network of rigid ZrO6 octahedra and WO₄ tetrahedra. It supports lower energy (~ 0.5 meV) phonon modes and has a metastable phase at room temperature. However, ZrW2O8 undergoes an order-disorder transition at 430 K ($\alpha \rightarrow \beta$, CTEs from -8.8×10^{-6} K⁻¹ to $-4.9 \times$ 10⁻⁶ K⁻¹) and it is easily to be decomposed above 770 K.⁵⁻¹⁰ In addition, its phase transition pressure is relatively low and will undergo a phase transition from α to γ at 0.21 GPa (CTEs from $-8.8 \times 10^{-6} \text{ K}^{-1}$ to $-1.2 \times 10^{-6} \text{ K}^{-1}$)¹¹ Similarly, the cubic ZrMo₂O₈ which is isostructural to ZrW₂O₈, has an average expansion coefficients of $-6.9 \times 10^{-6} \text{ K}^{-1}$ (2–200 K) and $-5.0 \times$ $10^{-6} \text{ K}^{-1} (250-502 \text{ K}).^{12} \text{ A}_2 \text{M}_3 \text{O}_{12}$ (A is a rare earth or a transition metal, M is Mo or W) series NTE materials have the advantages of big CTEs, wide-temperature range and large chemical flexibility. However, its anisotropy and hydroscopicity are serious obstacles for achieving excellent NTE properties. 4,13,14 ZrV2O7 has a strong isotropic NTE characteristic from 375 to 1075 K (CTE $-7.1 \times$

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 $10^{-6}~{\rm K}^{-1}$) and is stable at room temperature which is different from that of ZrW₂O₈. Nevertheless, it displays two first phase transitions at 350 K and 375 K. It belongs to a 3 × 3 × 3 superstructure below 350 K and transforms to 1 × 1 × 1 cubic phase above 375 K, meanwhile, it transforms from positive thermal expansion to NTE. In addition, ZrV₂O₇ undergoes a reversible phase transition at 1.38–1.58 GPa from α (cubic) to β (pseudotetragonal). ^{15–19} It should be highlighted that, this material will have an important application prospects if the phase transition temperature can be reduced to below room temperature.

From 1995 to 2001, Sleight *et al.* has done deep researches on $ZrV_{2-x}P_xO_7$ series materials. ^{15,17,21} In addition, Yamamura *et al.* has studied the phase transition temperature of $Zr_{1-x}Hf_xV_2O_7$. ²⁰ However, it's regrettable that the phase transition temperature hasn't been reduced to below room temperature successfully.

In this paper, we report a novel NTE material $Zr_{0.70}V_{1.33}$ - $Mo_{0.67}O_{6.73}$ for the first time, which exhibits a stable isotropic NTE property in a wide temperature range from 103 to 773 K. Its crystal structure, microstructure and thermal expansion properties are studied in detail. The mechanism of the NTE property was also discussed. Its excellent NTE property integrated with the facile preparation method suggests potential applications of this novel material in many fields.

2 Experimental procedures

Analytica-grade ZrO_2 , V_2O_5 and MoO_3 powders were used as starting materials and mixed according to the molar ratio of 0.70:1.40:0.70 as eqn (1).

 $0.70\text{ZrO}_2 + 1.33\text{V}_2\text{O}_5 + 0.67\text{MoO}_3 = \text{Zr}_{0.70}\text{V}_{1.33}\text{Mo}_{0.67}\text{O}_{6.73}$ (1)

The excess V_2O_5 and MoO_3 is to compensate for their evaporation during heating. Raw materials were thoroughly mixed and ground for 2 h with absolute ethyl alcohol in an agate mortar. The homogenized raw materials were pressed into cylindrical pellets with diameter of 5 mm and thickness of 3–5 mm by 769YP-15A powder tableting machine in 8 MPa and dwelled 5 min. The pellets were put into a tubular furnace preheated to sintering temperature at 953 K and maintained for 4 h before quenched, then dried at 473 K for 1 h. To get more uniform sample, the process of ground, press and sinter was performed for a secondary round according to the process described above. Final sample was obtained after annealed at 673 K.

A D8 ADVANCE X-ray diffractometer (Germany Bruker) was used for room, high and low temperature XRD measurements; a Quanta250 scanning electron micrograph (SEM, USA FEI) and a APOLLO XP energy dispersive spectroscopy (EDS, USA EDAX) were used to analysis microstructure and chemical of sample; a ESCALAB 250Xi X-ray photoelectron spectroscopy (XPS, Germany Thermo Fisher) was used to analyze the composition of the samples and states of elements; a STA 449F3 simultaneous thermal analyzer (Germany Netzsch, in the temperature range of 300-873 K with heating and cooling rates of 10 K min⁻¹) was used to differential scanning calorimetry (DSC) and thermogravimetric measurements (TG); a LabRAM HR Evolution Raman spectrometer (France HORIBA Jobin Yvon S.A.S.) was used to measure room temperature and temperature-dependent Raman spectral; a LINSEIS DIL L75 dilatometer (Germany LINSEIS, in the temperature range of RT-773 K with heating rate of 5 K min⁻¹) and a LINSEIS DIL L76 dilatometer (Germany LINSEIS, in the temperature range of 163-673 K with heating rate of 5 K min⁻¹) were used to measure the relative length changes depended temperature.

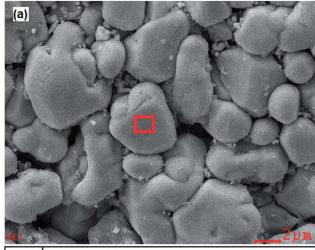
3 Results and discussion

3.1 Microstructure

Fig. 1a shows the SEM image of the sample. It can be seen that the ceramic sample is composed of irregular particles with different sizes and smooth surface. The average particle sizes are approximately 1–6 μ m. Fig. 1b shows the EDS spectrum corresponding to Fig. 1a. The EDS analysis of the grains presents primary elements of Zr, V, Mo and O.

To further confirm the composition of sample, we also performed XPS analyses. Fig. 2a shows the survey scan XPS spectrum of sample, and Zr, V, Mo and O signals can be clearly observed. The Zr 3d, V 2p, Mo 3d and O 1s XPS spectra are shown in Fig. 2b–e. The Zr 3d, V 2p and Mo 3d peak values are in agreement with the existing reports, $^{22-24}$ confirming the state of Zr $^{4+}$, V $^{5+}$ and Mo $^{6+}$ in Zr $_{0.70}$ V $_{1.33}$ Mo $_{0.67}$ O $_{6.73}$. We also calculated the atomic ratio of Zr, V, Mo and O by the relative sensitivity factor (RSF) method based on eqn (2) and (3).

$$C_{\rm x} = \frac{n_{\rm x}}{\sum_{i} n_{i}} = \frac{\frac{I_{\rm x}}{S_{\rm x}}}{\sum_{i} \frac{I_{i}}{S_{i}}}$$
(2)



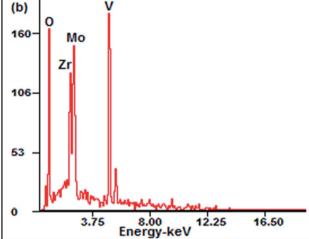


Fig. 1 SEM image and EDS spectrum. (a) SEM image of sample; (b) EDS spectrum corresponding to (a).

$$\frac{n_i}{n_j} = \frac{\frac{I_i}{S_i}}{\frac{I_j}{S_i}} \tag{3}$$

where C_x , I_x , S_x and n_x denote the relative content, characteristic spectra integral area, RSF and content of x element, respectively. The results were show in Table 1. It can be clearly seen that the atomic ratio of Zr, V, Mo and O are about 0.7:1.33:0.67:6.73, which in good agreement with the molar ratio of starting materials.

3.2 Crystal structure

The crystal structure is refined by Rietveld analysis with software of Fullprof. Fig. 3 exhibit the Rietveld analysis for the RT XRD pattern. It must be highlighted that the refinement has reference the result of Sahoo et~al. about ${\rm ZrV}_{2-x}{\rm Mo}_x{\rm O}_{7+\delta}$. The analysis shows that the sample adopts a cubic crystal structure with space group of $Pa\bar{3}$ (no. 205). The cell lattice parameters and volume are calculated to be a=b=c=8.93063 Å, V=712.27316 Å 3 , respectively, with the acceptable values of, $R_{\rm p}=17.6\%$, $R_{\rm wp}=14.4\%$ and $R_{\rm exp}=11.17\%$.

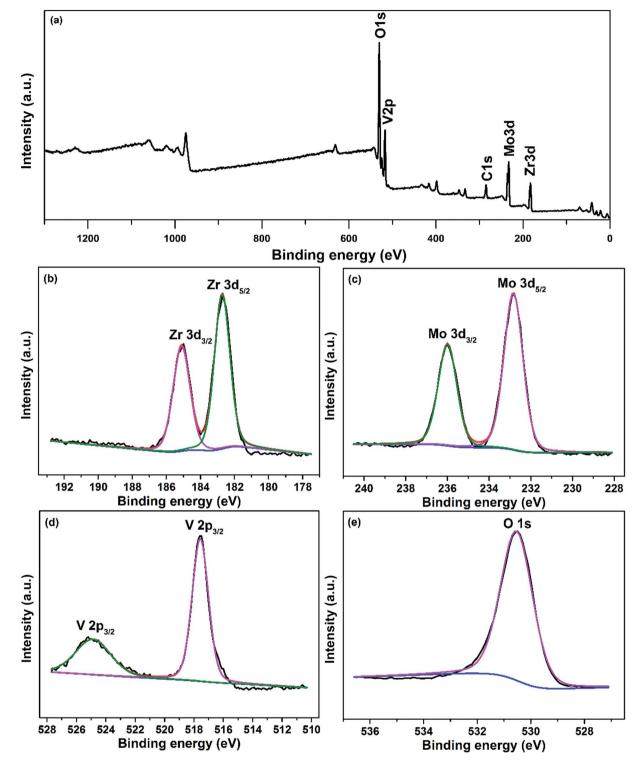


Fig. 2 The XPS spectra of sample: (a) the survey scan XPS spectrum; (b) Zr 3d; (c) V 2p; (d) Mo 3d; (e) O 1s.

Fig. 4a presents the schematic diagram of $Zr_{0.70}V_{1.33}Mo_{0.67}$ - $O_{6.73}$ building block depending on the XRD refinement result. The crystal structures are characterized by linkages of corner shared VO_4 or MoO_4 tetrahedra and ZrO_6 octahedra. The ZrO_6 octahedra share all the oxygen atoms with the adjacent VO_4 or MoO_4 tetrahedra. On the other hand, the two VO_4 tetrahedra

share an oxygen atom and make a V_2O_7 group. Each MoO_4 tetrahedron, however, share only three of its four oxygens with the neighboring ZrO_6 octahedra, and consequently, has one unshared vertex. The ZrV_2O_7 crystallizes in the centrosymmetric $Pa\bar{3}$ (no. 205),¹⁵ however, it undergoes a phase transformation from the space group $P2_13$ to $Pa\bar{3}$ at 483 K. Whereas cubic

Table 1 The integration and atomic ratio results of Zr 3d, V 2p, Mo 3d and O 1s

Element	RSF	Area	Area/RSF	Atomic ratio
Zr 3d	9.032	20 893,893	2303.354	0.700
V 2p	9.280	40 464.230	4360.370	1.325
Mo 3d	11.008	24 716.019	2249.366	0.684
O 1s	2.881	64 939.260	22 540.528	6.850

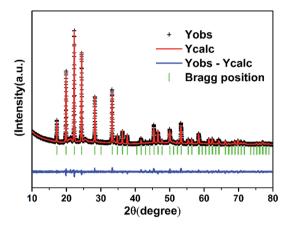


Fig. 3 XRD pattern of the sample at RT and the result of Rietveld analysis.

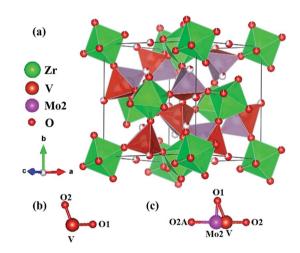


Fig. 4 Schematic diagram of Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}. (a) Polyhedral representation of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$; (b) asymmetric unit of ZrV_2O_7 ; (c) asymmetric unit of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$.

 ZrW_2O_8 crystallize in the noncentrosymmetric space group $P2_13$ (no. 198).8 As cubic ZrW₂O₈ is isostructural to cubic ZrMo₂O₈, we will compare the $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ with HT-ZrW₂O₈. Table 2 provides the atomic coordinates of the current compound of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$, ZrV_2O_7 and HT- ZrW_2O_8 . A comparison of the asymmetric unit of Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73} with ZrV₂O₇ is shown in Fig. 4b and c. It is clear that the formula derived from the XRD refinement result is $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ and the structure of Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73} is a hybrid of ZrV₂O₇

and cubic ZrMo₂O₈. The incorporation of cubic ZrMo₂O₈ structure and partially occupied Zr atoms (4a) maybe the main reasons of the phase transition temperature decrease. And the NTE in $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ can be mainly attributed to the QRUMs, which will be discussed later.

Thermal expansion properties

Fig. 5 shows the relative length (%) change of Zr_{0.70}V_{1.33}-Mo_{0.67}O_{6.73} ceramic with temperature measured by hightemperature dilatometer (RT-773 K) and low-temperature dilatometer (163–673 K), respectively. It shows that $Zr_{0.70}V_{1.33}$ $Mo_{0.67}O_{6.73}$ ceramic exhibits the NTE from 163 to 773 K, and the average linear CTEs are calculated to be $-3.75 \times 10^{-6} \text{ K}^{-1}$ $(163-673 \text{ K}) \text{ and } -4.50 \times 10^{-6} \text{ K}^{-1} \text{ (RT-773 K)}.$

Fig. 6a and c show the low and high temperature XRD patterns of Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73} ceramic. The XRD patterns remain unchanged from 103 to 473 K and from RT to 673 K except successful shifts of the diffraction peaks to higher angles, suggesting that $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ keeps the cubic structure from 103 to 673 K without phase transitions. The successive shifts of the diffraction peaks to higher angles reveals a contraction in lattice constants with increasing temperature in complete agreement with the result of Fig. 5. Fig. 6b and d exhibit the changes of the lattice constants along a, b, c axes (a = b = c) and lattice volume V with temperature which is calculated by PowderX. The intrinsic linear CTEs of *a*-, *b*-, and *c*-axes are calculated to be $\alpha_a = \alpha_b = \alpha_c = -3.84 \times 10^{-3}$ $10^{-6} \,\mathrm{K}^{-1} \,(103\text{-}463 \,\mathrm{K})$ and $\alpha_a = \alpha_b = \alpha_c = -4.22 \times 10^{-6} \,\mathrm{K}^{-1} \,(\mathrm{RT}$

Table 2 Atomic coordinates of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$, ZrV_2O_7 and HT-ZrW₂O₈

	$Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$	ZrV ₂ O ₇ (ref. 15)	HT-ZrW ₂ O ₈ (ref. 8)
Atom	Zr (4a)	Zr (4a)	Zr (4a)
x = y = z	0	0	0
Occupancy	0.7	1	1
Atom	V (8c)	V (8c)	W1 (8c)
x = y = z	0.40072	0.38174	0.33945
Occupancy	0.66667	1	0.5
Atom	Mo2 (8c)		W2 (8c)
x = y = z	0.34932		0.60355
Occupancy	0.33333		0.5
Atom	O2 (4b)	O2 (4b)	O3 (8c)
x = y = z	0.5	0.5	0.50550
Occupancy	0.66667	1	0.5
Atom	O2A (8c)		O4 (8c)
x = y = z	0.23221		0.23224
Occupancy	0.16667		0.5
Atom	O1 (24d)	O1 (24d)	O1 (24d)
x	0.45206	0.43634	0.05494
γ	0.18844	0.20582	-0.20893
z	0.41926	0.40730	-0.06713
Occupancy	0.95556	1	1

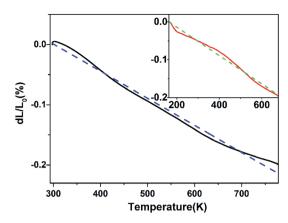


Fig. 5 Relative length (%) changes of the $\rm Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ ceramic with temperature from RT to 773 K and from 163 to 673 K (inset). The dotted lines are linear fits to the experimental data.

673 K). The volume CTEs are calculated to be $\alpha_V = -11.5 \times 10^{-6}$ K⁻¹ (103–463 K) and $\alpha_V = -12.7 \times 10^{-6}$ K⁻¹ (RT–673 K), which agree well with the results measured by dilatometry (see Fig. 5).

Fig. 7 shows the DSC and thermogravimetric plots of the $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$. The mass increase and dip in DSC before 335 K are an instrumental artifact and thermal drift. Neither evident endothermic/exothermic peaks nor obvious weight loss

from RT to 873 K appear, confirming that $\rm Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ possesses perfect thermal stability and does not exhibit phase transition from RT to 873 K.

3.4 Thermal shrink mechanism discussed

Fig. 8a shows the RT Raman spectrum of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$. The Raman modes in the ranges 310-400 cm⁻¹, 480-550 cm⁻¹ 700-850 cm⁻¹ and 850-1050 cm⁻¹ are identified as symmetric bending (ν_2) , asymmetric bending (ν_4) , asymmetric stretching (ν_3) and symmetric stretching (ν_1) modes of the MoO₄ and VO₄ tetrahedra, respectively. And those peaks below 300 cm⁻¹ are vibration of Zr atomic and translational and librational modes of MoO₄ and VO₄ tetrahedra. 14,19,26,27 Fig. 8b shows the temperaturedependent Raman spectra of ZrMoV₂O₁₀ from 4 to 464 K. It reveals that there are no distinct appear or disappear of Raman peaks. The Raman shifts occur successively redshift from 4 to 24 K and blue shift from 24 to 54 K, so it can be speculated that the temperature range of 24-54 K is the transition of superstructure. 15,16 Meanwhile, the Raman shifts occur redshift from 54 to 464 K without changing the amount of Raman peaks, suggesting that $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ keeps the cubic structure from 54 to 464 K without phase transitions. It can be also confirmed by the result of DSC (Fig. 7). The rotations of MoO₄ and VO₄ tetrahedra with increasing temperature result in the change of symmetric

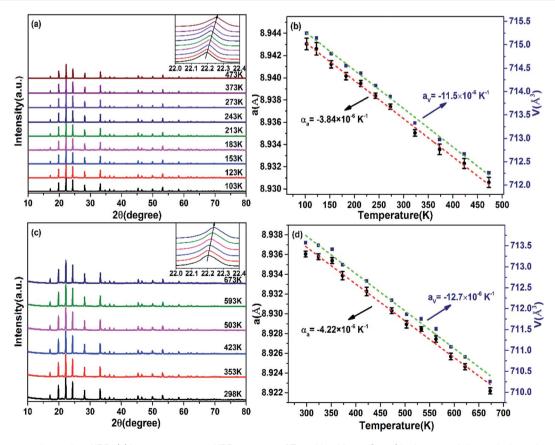


Fig. 6 Temperature dependent XRD. (a) Low temperature XRD patterns of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ (the inset on right top is the enlarged view of the main peak between 22.0° and 22.4°); (b) change in lattice constant and cell volume of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ of low temperature XRD; (c) high temperature XRD patterns of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ (the inset on right top is the enlarged view of the main peak between 22.0° and 22.4°); (d) change in lattice constant and cell volume of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ of high temperature XRD.

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Fig. 7 Typical DSC and thermogravimetric plots for $\rm Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}.$

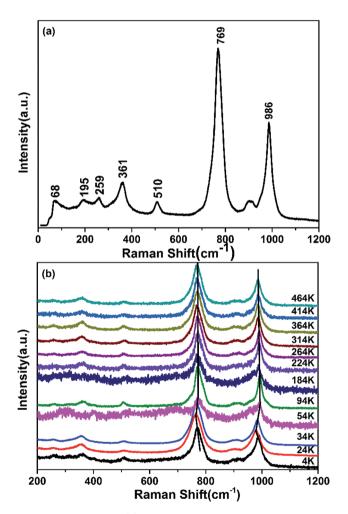


Fig. 8 Raman spectra. (a) Room temperature Raman spectrum of $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$; (b) low temperature Raman spectra of $Zr_{0.70}-V_{1.33}Mo_{0.67}O_{6.73}$ from 4 to 464 K.

bending (ν_2) , asymmetric bending (ν_4) , asymmetric stretching (ν_3) and symmetric stretching (ν_1) vibrational modes, which is the main reason for the redshift of Raman bands.

From the results of XRD refinement and temperature dependent Raman spectra, we deduced that Zr_{0.70}V_{1.33}Mo_{0.67}- $O_{6,73}$ belongs to a 3 \times 3 \times 3 superstructure below 24 K then starts to transition from $3 \times 3 \times 3$ superstructure to $1 \times 1 \times 1$ cubic phase and completely transition to $1 \times 1 \times 1$ cubic phase above 54 K. It exhibits stable NTE property above 54 K. Considered the similarities of structure and properties between $Zr_{0.7}V_{1.4}Mo_{0.7}O_{7}$ and $ZrV_{2}O_{7}$, 15,17,27 we regard MoO_{4} , VO₄ tetrahedra and ZrO₆ octahedra in Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73} crystal as quasi-rigid, when the temperature is higher than 54 K, the polyhedra vibrate as the temperature increasing. The tilts and rotations of MoO4 and VO4 tetrahedra, will necessarily drag the ZrO₆ octahedra along with it, and pull the rest of the structure in, giving a volume reduction. Since this tilts and rotations, and hence volume reduction, will be larger at higher temperature, there will be a NTE property. It is of interest to note that in ZrV₂O₇, the negative thermal expansion begins at 375 K. Nevertheless, in $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$, it begins at 54 K. This indicates that the phase transition temperature has been reduced to 54 K. According to the preceding crystal structure analysis in 3.3 and the NTE property of cubic ZrMo₂O₈, 8,15 we deduced that it may be attributed to the incorporation of cubic ZrMo2O8 structure and partially occupied Zr atoms (4a).

4 Conclusion

In summary, we developed a novel NTE material Zr_{0.70}V_{1.33}-Mo_{0.67}O_{6.73}, which possesses excellent stable isotropic NTE property and keeps cubic crystal structure over a wide temperature range from 103 to 773 K. Structure analysis shows that $Zr_{0.70}V_{1.33}Mo_{0.67}O_{6.73}$ has a cubic structure with the space group $Pa\bar{3}$ (no. 205). The average linear CTEs are calculated to be $-3.75 \times 10^{-6} \text{ K}^{-1}$ (163–673 K) and $-4.50 \times 10^{-6} \text{ K}^{-1}$ (RT–773 K). The intrinsic linear CTEs are calculated to be -3.84×10^{-6} ${
m K}^{-1}$ (103–473 K) and $-4.22 \times 10^{-6}~{
m K}^{-1}$ (RT–673 K). ${
m Zr}_{0.70}{
m V}_{1.33}$ - $Mo_{0.67}O_{6.73}$ belongs to a 3 \times 3 \times 3 superstructure below 24 K and $1 \times 1 \times 1$ cubic phase above 54 K while it exhibits the stable isotropic NTE property. The NTE property can be attributed to the QRUMs. And the phase transition temperature of ZrV₂O₇ has been reduced below RT, this may be attributed to the incorporation of ZrMo₂O₈ structure and partially occupied Zr atoms (4a). The excellent NTE property integrated with the facile preparation method implies potential application of Zr_{0.70}-V_{1,33}Mo_{0,67}O_{6,73}. This work could open up new opportunities for developing NTE materials.

Acknowledgements

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