


 CrossMark  
click for updates
Cite this: *RSC Adv.*, 2017, 7, 1776Received 2nd December 2016  
Accepted 16th December 2016

DOI: 10.1039/c6ra27673c

www.rsc.org/advances

# Sensitivity enhancement of graphene Hall sensors modified by single-molecule magnets at room temperature†

Yuanhui Zheng,<sup>‡,ab</sup> Le Huang,<sup>‡,c</sup> Zhiyong Zhang,<sup>\*,c</sup> Jianzhuang Jiang,<sup>\*,d</sup>  
Kaiyou Wang,<sup>\*,e</sup> Lian-Mao Peng<sup>\*,c</sup> and Gui Yu<sup>\*,ab</sup>

Single-molecule magnets (SMMs) possess many unique magnetic properties and thus attract a wide range of attention. However, the applications of SMMs always need a strict atmosphere, *i.e.* low temperature, high vacuum and strong magnetic field. In this work, we report the preparation and characterization of sensitivity enhanced graphene Hall elements (GHEs) decorated with Tb-core SMMs. By comparing the magnetic sensing and electronic tests of the GHEs before and after the SMMs modifications, the sensitivity of the GHEs increases by 44.9% in voltage mode and 59.0% in current mode compared with pristine GHEs. The increase of sensitivity may result from the magnetic center introduction of SMMs at room temperature. Moreover, the magnetic molecules may affect the graphene field environment leading to a Hall signal change. In addition, the SMMs modified GHEs present excellent linearity, offset voltage, repeatability and stability in magnetic sensing. This study paves the way to apply SMMs into practical use at room temperature and atmospheric pressure without strong magnetic field excitations.

Furthermore, SMMs are considered as a promising material for building quantum computers,<sup>6,7</sup> spintronic devices<sup>8,9</sup> and magnetic memories.<sup>10,11</sup> Generally, SMMs demonstrate ferromagnetic characteristics under low temperature.<sup>12</sup> That is why the vast majority of investigations and applications surrounding SMMs were carried out under low temperature.<sup>13–15</sup> However, the strict testing environment limits the practical use of SMMs-based devices and circuits. When the temperature exceeds the Curie temperature of SMMs, the SMMs would transfer to paramagnetic characteristics due to the increasing thermal motions excited by the increased temperature.<sup>12</sup> Unlike the ferromagnetic materials, the magnetization of paramagnetic materials would have a linear relationship with the external magnetic field.<sup>16,17</sup>

Actually, graphene is a terrific platform for investigating the proximity effect and magnetic properties of SMMs.<sup>9,13</sup> This is originated from the ultra-thin body and excellent electronic properties of graphene. On one hand, the extremely high surface-volume ratio of graphene helps to maximize the impact from the surface SMMs-induced effects on graphene.<sup>18</sup> On the other hand, the high mobility and conductivity of graphene enable the transfer of SMMs magnetic information into the electronic signals of graphene effectively.<sup>13</sup> These two advantages of graphene are essential to confirm the existence and develop the potential applications of SMMs.

At the same time, graphene is regarded as a fabulous material for constructing high-performance Hall sensors.<sup>19</sup> Benefitted from the high mobility, thin body, unique band structure, low noise, thermal stability and flexibility of graphene, graphene Hall elements (GHEs) exhibit high sensitivity, resolution, linearity and stability in magnetic detections.<sup>19–23</sup> Moreover, the realization of integrating GHEs with silicon Complementary Metal Oxide Semiconductor (CMOS) integrated circuits (ICs) makes GHEs more promising in practical use.<sup>21,24</sup>

In this work, we report the fabrication and characterization of sensitivity enhanced GHEs decorated with TbPc<sub>2</sub> SMMs. Based on this method, the sensitivity of the GHEs increases by 44.9% in voltage mode and 59.0% in current mode compared with pristine GHEs. This may be attributed to the introduction

## Introduction

In recent years, single-molecule magnets (SMMs) have been intensively investigated due to their unique magnetic hysteresis and quantum behavior phenomenon under low temperature.<sup>1–5</sup>

<sup>a</sup>Beijing National Laboratory for Molecular Sciences, Institute of Chemistry, Chinese Academy of Sciences, Beijing 100190, P. R. China. E-mail: yugui@iccas.ac.cn

<sup>b</sup>University of Chinese Academy of Sciences, Beijing 100049, P. R. China

<sup>c</sup>Key Laboratory for the Physics and Chemistry of Nanodevices, Department of Electronics, Peking University, Beijing 100871, P. R. China. E-mail: zyzhang@pku.edu.cn; lmpeng@pku.edu.cn

<sup>d</sup>Beijing Key Laboratory for Science and Application of Functional Molecular and Crystalline Materials, Department of Chemistry, University of Science and Technology Beijing, Beijing 100083, P. R. China. E-mail: jianzhuang@ustb.edu.cn

<sup>e</sup>SKLSM, Institute of Semiconductors Chinese Academy of Sciences, Beijing 100083, P. R. China. E-mail: kynwang@semi.ac.cn

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c6ra27673c

‡ These authors contributed equally to the paper.



of magnetic center of SMMs at room temperature. In addition, the magnetic molecules can influence the graphene field environment leading to Hall signal change.<sup>13</sup> Additionally, the SMM modified GHEs present excellent linearity, offset voltage, repeatability and stability in magnetic sensing.

## Experimental section

### Fabrication methods of GHEs

Graphene sample was grown by CVD method based on Pt foil substrate.<sup>20</sup> Afterwards, the graphene was transferred onto a SiO<sub>2</sub> (285 nm)/Si substrate with a bubbling method.<sup>20</sup> Electron beam lithography (EBL) based on PMMA was then carried out to pattern the symmetric cross-shaped graphene channel, followed by a reactive ion etching (RIE) process to remove the needless region of graphene. The second EBL process was used to pattern the metal contacts, followed by an electron beam evaporator (EBE) deposition of Ti/Au (5/30 nm) and a standard lift-off process. The electronic and magnetic response data were collected by a Keithly 4200 SCS semiconductor parameter analyser.

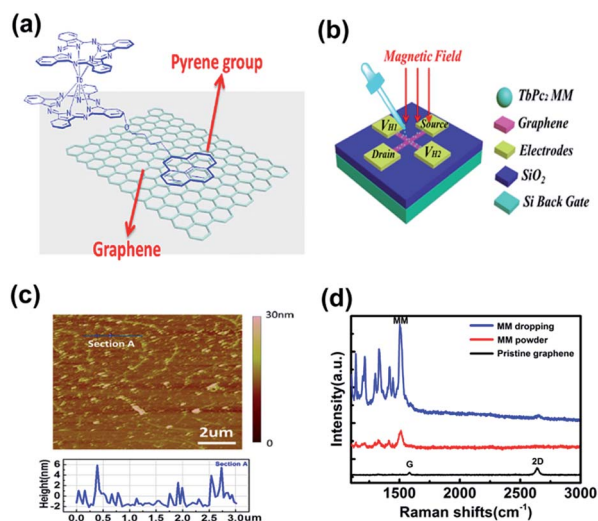
### Preparations and modifications of SMMs

The SMMs deposited on graphene are pyrenyl-substituted heteroleptical bis(phthalocyaninato) terbium(III) complexes, which is referred as TbPc<sub>2</sub> in this work (Fig. 1b). This molecule

consists of a single magnetic ion (Tb<sup>3+</sup>) coordinated to two phthalocyanine ligands (see scheme of Fig. 1a). In order to improve the adhesion to graphene, one of the two phthalocyanine macrocycles was substituted by a pyrene group. The pyrene group exhibits an attractive interaction with sp<sup>2</sup> carbon materials, maximizing the intermolecular interaction.<sup>25</sup> The molecule was deposited by drop casting of a TbPc<sub>2</sub> solution in dichloromethane (DCM) with a molecule concentration ranging from 10<sup>−7</sup> to 10<sup>−3</sup> mol L<sup>−1</sup>. After 10 seconds, the sample was rinsed in DCM for 1 minute to remove the DCM. Afterwards, residual DCM was removed by a second rinse in isopropanol. Finally, the sample was dried by nitrogen flow.<sup>26</sup>

## Results and discussion

As is shown in Fig. 1a, the SMM used in this work is pyrene-substituted terbium(III) bis(phthalocyanine). This molecule was synthesized as reported before.<sup>25</sup> The SMMs consist of a single magnetic Tb<sup>3+</sup> ion coordinated to two phthalocyanine ligands. In order to improve its adhesion to graphene, one of the two phthalocyanine macrocycles was substituted by a pyrene group. While covalent bonding might severely alter the performance of graphene and SMMs, non-covalent  $\pi$ -stacking should preserve the intrinsic features of them.<sup>26</sup> The graphene was grown by CVD approach and transferred onto the substrate by a bubbling method.<sup>22</sup> After the transfer of graphene, GHEs were constructed through micro-fabrications based on electron beam lithography (EBL) (Fig. 1S†). The channel is patterned as crossed shape so as to work as Hall magnetic sensing tunnels.<sup>21</sup> The channel length and width are 50 and 15  $\mu$ m, respectively. After the fabrications of GHEs, SMMs diluted with DCM are dropped onto the sensors to modify graphene (Fig. 1b). It is mentionable that the method of SMMs modifications avoids introducing extra complicated processes. Therefore, the whole process of GHEs fabrications and modifications is easy and compatible with silicon technology.<sup>23</sup> To investigate the surface topography of SMMs-modified graphene, we conducted an AFM test and the result is shown in Fig. 1c. It is clear that the grafting of SMMs on graphene is homogeneous with a surface roughness of 3.24 nm. By contrast, SMMs is hardly adsorbed on bare SiO<sub>2</sub> substrate (Fig. S2†). The comparison points to a fact that SMMs is adhered to graphene with selectivity because of the  $\pi$ - $\pi$  interactions between pyrene group and graphene.<sup>27</sup> In addition, we carried out Raman spectrum test for pristine graphene, SMMs powder and SMMs-graphene hybrid system to verify the modification of SMMs onto graphene (Fig. 1d). Based on the Raman results of pristine graphene sample, the 2D/G ratio of pristine graphene is 2.0. This value indicates that our monolayer graphene is of high quality, and the results in SAED also indicate the sample single-crystalline feature (Fig. S3†),<sup>28</sup> which is necessary for constructing high performance magnetic sensors. Since all Raman modes of the TbPc<sub>2</sub> molecule follow the same behavior in our experiments, we use the most intense band (MM) for the Raman spectra. The Raman signal of the graphene-TbPc<sub>2</sub> hybrid system is a simple superposition of the graphene and SMMs signals without any shift or disappearance, indicating that both graphene and SMMs remain chemically



**Fig. 1** Graphene-Single Molecular Magnets (SMMs) hybrid structure and characterizations. (a) Schematic view of TbPc<sub>2</sub> molecule attached to graphene. The magnetic core is Tb<sup>3+</sup>. One of the terminal ligands is substituted by a pyrene group. (b) Schematic of graphene Hall elements (GHEs) decorated with SMMs. The channel length and width are 50 and 15  $\mu$ m, respectively. The concentration of SMMs is 10<sup>−7</sup> mol L<sup>−1</sup>. (c) Atomic force microscope (AFM) image of SMM-modified graphene. The scale bar is 2  $\mu$ m. The concentration of SMMs is 10<sup>−7</sup> mol L<sup>−1</sup>. (d) Raman spectrum of the pristine graphene (bottom, black), TbPc<sub>2</sub> powder (middle, red), and the hybrid system of graphene and TbPc<sub>2</sub> SMMs (top, blue). The G and 2D modes of graphene and the MM band of SMMs are indicated. The wavelength of the excitation laser is 633 nm. The concentration of SMMs is 10<sup>−3</sup> mol L<sup>−1</sup>.



unchanging after the grafting process.<sup>26</sup> The Raman intensity of SMMs is much stronger and well-distinguished due to the easy charge transfer between graphene and molecules, and the similarity of the chemical structure between graphene and pyrene substituted SMMs, the vibrational coupling between them may be contribute to the Raman enhancement.<sup>27,29,30</sup> The G band of graphene almost did not shift and without any broadening or splitting. These results indicate that the interaction between graphene and TbPc<sub>2</sub> SMMs is weak.<sup>26</sup>

The back-gate transfer characteristics of the pristine and SMMs-modified GHEs are shown in Fig. 2a. Two opposite electrodes (source and drain) were biased by a constant voltage ( $V_{ds} = 50$  mV). A back gate sweeping voltage from  $-40$  to  $40$  V was applied at silicon substrate during the measurements. The carrier mobility of pristine graphene was extracted to be  $3155 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  based on the fitting method reported in ref. 31. The high carrier mobility further indicates that the as-fabricated graphene devices are of high-quality. Afterwards, GHEs were decorated with  $10^{-7} \text{ mol L}^{-1}$  SMMs and tested again. The carrier mobility decreases but still can reach  $2000 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ . This result is originated from the scattering introduced by the

SMMs. Additionally, the Dirac point voltage of graphene-based device shifts positively from  $16$  to  $21$  V, indicating a p-type charge transfer from SMMs to graphene. The charge transfer is caused by the residual DCM of SMMs during the modification process.<sup>26</sup>

The magnetic responses of the GHEs were then measured on the Hall probe station. All the measurements were carried out at room temperature and ambient atmosphere, which is a huge advantage over traditional applications of SMMs. During the measurement, a constant voltage/current bias ( $V_C/I_C$ ) was applied between the electrodes named drain and source. At the same time, an external magnetic field ( $B_0$ ) normal to the plane of sensors would generate a constant Hall voltage ( $V_{\text{Hall}}$ ) between the electrodes  $V_{H1}$  and  $V_{H2}$  due to the Hall effect.<sup>24</sup>

Sensitivity is a key parameter to benchmark the performance of Hall sensors. The Hall sensors with high sensitivity can detect magnetic field with high signal-noise ratio and thus reduce the cost of amplifications.<sup>23,32</sup> The absolute sensitivity ( $S_A$ ) of GHEs is defined as the ratio of Hall voltage and external magnetic field,<sup>19,24</sup>

$$S_A = \frac{V_{\text{Hall}}}{B_0} \quad (1)$$

in which  $V_{\text{Hall}}$  is the Hall voltage and  $B_0$  is the external magnetic field.

We first consider Hall sensors powered by a constant voltage. The absolute sensitivity of pristine GHEs working at voltage mode is determined as follows,<sup>19,24</sup>

$$S_A = \frac{WV_C}{L} \mu \quad (2)$$

$W$  and  $L$  represent the width and length of graphene channel.  $V_C$  is the applied voltage and  $\mu$  is the mobility. Based on eqn (2), degraded mobility caused by SMMs modification would lead to a decrease of sensitivity in Hall sensing. However, the magnetic response test demonstrates a contradictory result.  $V_{\text{Hall}}-B$  curves of GHEs before and after SMMs modification are shown in Fig. 2b. The power supply is set as  $1$  V. The pristine GHE exhibits an excellent linear  $V_{\text{Hall}}-B$  relation as expected with a slope of  $0.06 \text{ V T}^{-1}$ . After the SMMs modification ( $10^{-7} \text{ mol L}^{-1}$ ), the sensitivity increases up to  $0.088 \text{ V T}^{-1}$ . The sensitivity increase is originated from the magnetic field enhancement effect of SMMs. To be detailed, molecular magnets show paramagnetic features at room temperature.<sup>1-3</sup> That is to say, the magnetic field ( $B_{\text{MM}}$ ) of SMMs magnetization may be added in the hybrid system with the external magnetic field ( $B_0$ ),<sup>17</sup>

$$B = B_{\text{MM}} + B_0 \quad (3)$$

$B$  is the total magnetic field applied at Hall sensors which should be the sum of  $B_0$  and  $B_{\text{MM}}$ . Fig. 2c depicts the schematic of the magnetic field enhancement effect of SMMs. Thus  $V_{\text{Hall}}$  of the graphene-SMMs hybrid system can be represented as follows:

$$V_{\text{Hall}} = \frac{WV_C}{L} \mu' B \quad (4)$$

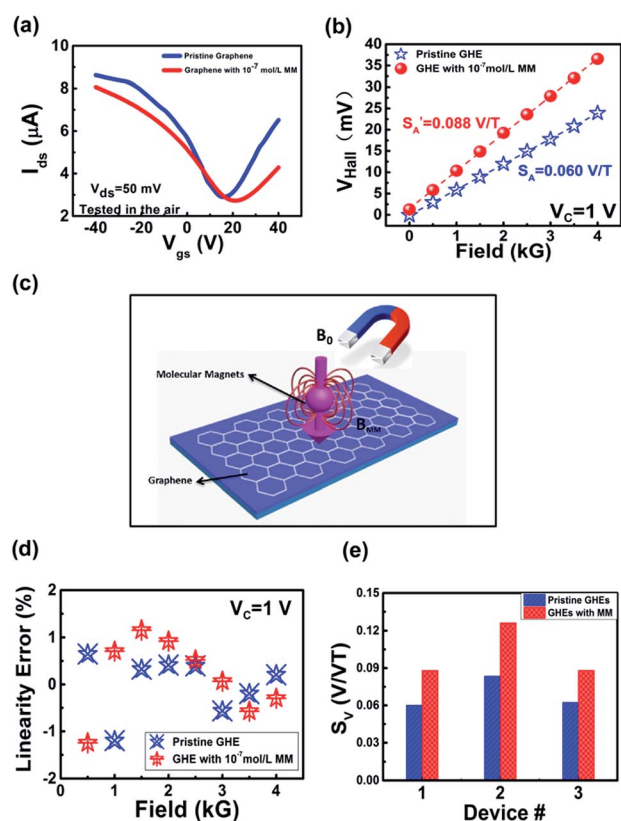


Fig. 2 Comparison of electronic properties and magnetic responses between pristine and SMM-modified GHEs in voltage mode. (a) The back-gate transfer characteristics of the pristine and SMM-modified GHEs. (b)  $V_{\text{H}}-B$  curves of pristine GHEs (blue line) and SMM modified GHEs (red line) under voltage bias of  $1$  V. The dashed lines denote the linear fitting results of the  $V_{\text{H}}-B$  curves. (c) Schematic of mechanism of GHEs sensitivity enhancement. (d) The linearity errors of the pristine and SMM modified GHEs. (e) Sensitivity statistics of GHEs before and after the modification of SMMs.





$\mu'$  is the graphene mobility after SMMs modifications. Based on eqn (1), the absolute  $S_A$  of the SMMs modified GHEs is derived,

$$S'_A = \frac{WV_C}{L} \mu' \frac{B}{B_0} \quad (5)$$

It is evident that the sensitivity of graphene-SMMs hybrid Hall sensors is the compromise between the mobility decrease and magnetic field influence. The  $\mu'$  is decreased, but the  $S'_A$  is increased. From the eqn (5), we can conclude that the  $S'_A$  enhancement can be ascribed to the  $B/B_0$ , and it may be because the magnetic center influence of SMMs affects the total magnetic field applied at Hall sensors (Fig. 2c). Fig. 2d demonstrates the linearity errors of GHEs before and after decorating molecular magnets. Both the pristine and SMMs modified GHEs exhibit excellent linearity with linearity error less than 1.5%. SMMs modifications avoid decreasing the sensing accuracy of Hall sensors. This is attributed to the linear relationship between  $B_{MM}$  and  $B_0$  as shown in eqn (3). The voltage-related sensitivity, defined as  $S_V = S_A/V$ , is widely used as an important metric to benchmark the magnetic sensitivity of different Hall sensors excluding the impact of driving voltage. The voltage sensitivity has an average increase of 44.9% (from 0.069/T to 0.100/T) after SMMs modifications (Fig. 2e). It is worth mentioning that the  $S_V$  of the SMMs-graphene hybrid Hall element outperforms silicon based Hall elements ( $S_V = 0.07/T$ ).<sup>21</sup>

We then conducted the magnetic detection experiment in current mode. The Hall sensors were powered by a constant current of 0.1 mA. Based on theoretical calculations, the absolute sensitivity of the Hall sensor in current mode is inversely proportional to the carrier concentration,<sup>19,24</sup>

$$S_A = \frac{I_C}{ne} \quad (6)$$

in which  $I_C$  is the applied current,  $n$  is the carrier density of channel and  $e$  is electron charge. Since the Dirac point voltage of graphene shifted positively, the carrier concentration of the device would increase. Based on eqn (6), the sensitivity of GHEs ought to decrease. By contrast, the sensitivity increases from 0.0536 V T<sup>-1</sup> to 0.0959 V T<sup>-1</sup> as shown in Fig. 3a after the SMMs modifications. The Hall voltage of graphene-SMMs hybrid sensors can be described as follows,

$$V_{Hall} = \frac{I_C}{e} \frac{1}{n'} B \quad (7)$$

in which  $n'$  is the carrier density of graphene after the SMMs modifications.

Therefore, the absolute sensitivity can be described this way,

$$S'_A = \frac{I_C}{e} \frac{1}{n'} \frac{B}{B_0} \quad (8)$$

As for typical back-gate transistors, the carrier density of graphene is proportional to its Dirac point voltage,

$$n = \frac{C_{ox} V_{Dirac}}{e} \quad (9)$$

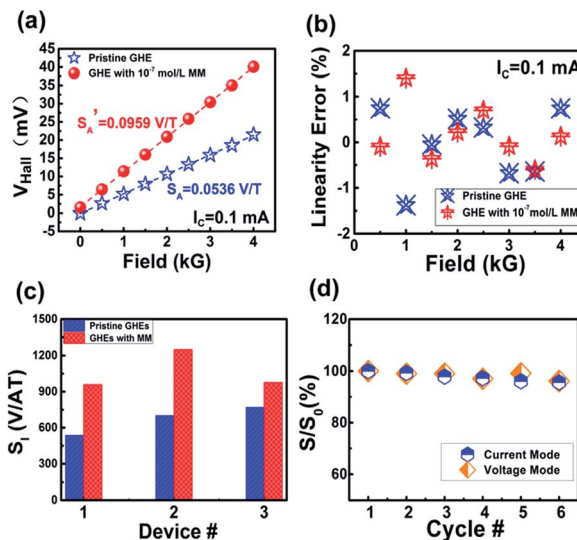


Fig. 3 Comparison of magnetic responses between pristine and SMM-modified GHEs in current mode. (a)  $V_H$ - $B$  curves of pristine GHEs (blue line) and SMM modified GHEs (red line) under current of 0.1 mA. The dashed lines denote the linear fitting results of the  $V_H$ - $B$  curves. (b) The linearity errors of the pristine and SMM modified GHEs. (c) Sensitivity statistics of GHEs before and after the modification of SMMs. (d) Cycle test of the magnetic responses within six magnetic detection loop after modification of 10<sup>-7</sup> mol L<sup>-1</sup> molecular magnets in voltage (denoted by orange diamonds) and current (denoted by blue hexagons) modes.  $S_0$  represents the magnetic sensitivity of the GHEs during the first magnetic sensing cycle.

$C_{ox}$  is the oxide capacitance of SiO<sub>2</sub> with a thickness of 285 nm.

Based on eqn (8) and (9), we can get the relation between  $S'_A$  and  $V'_{Dirac}$  in current mode,

$$S'_A = \frac{I_C}{C_{ox} V'_{Dirac}} \frac{B}{B_0} \quad (10)$$

As discussed in voltage mode, after modifying SMMs, the  $V'_{Dirac}$  shifted positively, but  $S'_A$  increase to 0.0959 V T<sup>-1</sup>. The enhancement may result from the magnetic center effect of SMMs and the magnetic molecules can affect graphene field environment because of the very close range interaction between graphene and SMMs at room temperature.<sup>13,33,34</sup>

After the modifications, the Hall sensors still shows tiny linearity error of less than 1.5% (Fig. 3b). The current sensitivity has an average increase of 59.0% (from 668 to 1062 V A<sup>-1</sup> T<sup>-1</sup>) after SMMs modifications (Fig. 3c). To make a comparison, the  $S_I$  of the SMMs-graphene hybrid Hall element outperforms silicon based Hall elements ( $S_I = 100$  V A<sup>-1</sup> T<sup>-1</sup>) and commercialized III-V compounds Hall sensors ( $S_I = 200$  V A<sup>-1</sup> T<sup>-1</sup>).<sup>24</sup>

Fig. 3d demonstrates the cycle test of the SMMs-modified GHEs within six magnetic detection loop in current (denoted by orange diamonds) and voltage (denoted by blue hexagons) modes. The floating range is less than 5%, indicating that the repeatability of our SMMs-modified GHEs is excellent.

Finally, we investigated the SMMs-concentration dependent magnetic detection performance of the GHEs. Fig. 4a shows the



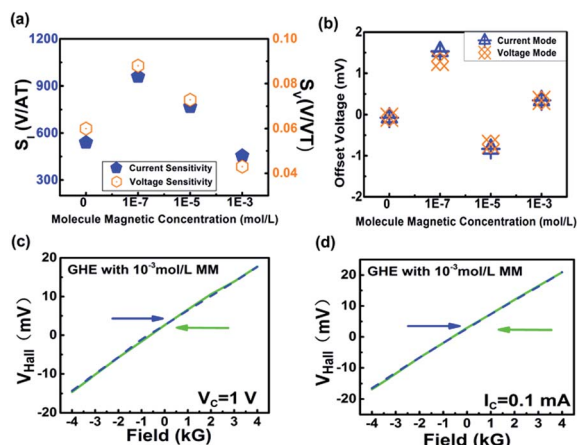


Fig. 4 Concentration dependent performance and magnetic sweeping direction test of SMM-modified GHEs. (a) SMM concentration dependent magnetic sensitivity of GHEs in current and voltage modes. (b) SMM concentration dependent offset voltages of GHEs in current and voltage modes. (c)  $V_H$ - $B$  curves of SMM-modified GHEs working at voltage mode and (d) current mode with opposite magnetic field sweeping direction. The concentration of SMM is  $10^{-3} \text{ mol L}^{-1}$ . The blue dashed lines show the magnetic response with sweeping direction from south to north and the green solid lines show the magnetic response with sweeping direction from north to south.

SMMs concentration dependent magnetic sensitivity of GHEs in current and voltage modes. It is clear that the sensitivity of GHEs increases when SMM concentration reaches  $10^{-7} \text{ mol L}^{-1}$  but declines when concentration exceeds  $10^{-7} \text{ mol L}^{-1}$ . The sensitivities are even lower than those of the pristine GHEs at the concentration of  $10^{-3} \text{ mol L}^{-1}$ . Based on the electronic measurements, graphene completely changes to p-type transport behavior when the concentration reaches  $10^{-3} \text{ mol L}^{-1}$  (Fig. S4†). The strong p-type carrier charge transfer is originated from the increasing residue dichloromethane with the additional modifications processes.<sup>26</sup> At the same time, the mobility significantly declines to as low as  $420 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ . The additional scattering is due to the scatterings of clusters and crystallizations formed at high SMMs concentrations.<sup>26</sup> Some clusters and crystallizations become visible under optical microscopes (see the ESI, Fig. S5†). As for Hall sensors, the offset voltage refers to the output Hall voltage when no magnetic field stimulation is applied. Generally, the offset voltage should be tuned as small as possible so as to reduce the cost of back-end calibrations. Fig. 4b shows the offset voltage as a function of the SMMs concentrations. In both the current and voltage modes, the offset voltages are within  $\pm 2 \text{ mV}$ . These values are competitive compared with the silicon, III-V compounds and graphene Hall sensors.<sup>19,24</sup> In addition, a sweeping direction test was carried out to further benchmark the reliability of our SMMs-modified GHEs. We selected a device decorated with  $10^{-3} \text{ mol L}^{-1}$  molecular magnets to conduct the experiment. Fig. 4c and d demonstrate the magnetic responses of GHEs with different sweep directions in voltage and current mode, respectively. These almost identical curves prove that our SMMs-graphene hybrid Hall sensors are independent with the

external magnetic field sweeping directions. To sum up the magnetic response tests, graphene-SMMs hybrid structures enhance the sensitivity while maintain excellent linearity, offset voltage, repeatability and stability in magnetic sensing.

## Conclusions

We report herein an easy way to apply the magnetic properties of SMMs at room temperature and enhance the sensitivity of GHEs by constructing a graphene-SMMs hybrid structure. The SMMs modification technique is directly applicable *in situ* onto GHEs without introducing complicated process. During the magnetic test, the sensitivity of the GHEs has an average increase of 44.9% in voltage mode and 59.0% in current mode. The sensitivity enhancement may result from the magnetic center introduction of SMMs at room temperature. The magnetic molecules may affect the graphene field environment leading to Hall signal change. Meanwhile, the Hall devices maintain other outstanding properties in magnetic detection, *i.e.* linearity (within  $\pm 1.5\%$ ), offset voltage (within  $\pm 2 \text{ mV}$ ), repeatability (sensitivity floating percentage within 5% during cycle tests) and stability (independent with magnetic field sweeping directions). Based on the graphene-SMMs hybrid structure, we are capable of utilizing the magnetic properties of SMMs at room temperature to enhance the sensitivity of magnetic detections.

## Acknowledgements

We are grateful for the financial support from the National Natural Science Foundation of China (61390502), the National Major State Basic Research Development Program (2013CB933403), and the Strategic Priority Research Program of the Chinese Academy of Sciences (XDB 12030100).

## Notes and references

- 1 N. W. Daniel, E. P. W. Richard and A. L. Richard, *Chem. Rev.*, 2013, **113**, 5110.
- 2 N. Domingo, E. Bellido and D. Ruiz-Molina, *Chem. Soc. Rev.*, 2012, **41**, 258.
- 3 X. Y. Wang, A. Carolina and R. D. Kim, *Chem. Soc. Rev.*, 2011, **40**, 3213.
- 4 L. Bogani and W. Wernsdorfer, *Nat. Mater.*, 2008, **7**, 179.
- 5 P. Dawid, H. I. Southerland, C. Avendaño, A. Prosvirin, C. Sanders, W. Wernsdorfer, K. S. Pedersen, J. Dreiser, R. Clérac, J. Nehr Korn, G. G. Simeoni, A. Schnegg, K. Holldack and K. R. Dunbar, *J. Am. Chem. Soc.*, 2015, **137**, 14406.
- 6 M. N. Leuenberger and D. Loss, *Nature*, 2001, **410**, 789.
- 7 J. Lehmann, A. Gaita-Arino, E. Coronado and D. Loss, *J. Mater. Chem.*, 2009, **19**, 1672.
- 8 J.-P. Cleuziou, W. Wernsdorfer, V. Bouchiat, T. Ondarcuhu and M. Monthieux, *Nat. Nanotechnol.*, 2006, **1**, 53.
- 9 C. Cervetti, A. Rettori, M. G. Pini, A. Cornia, A. Repollés, F. Luis, M. Dressel, S. Rauschenbach, K. Kern, M. Burghard and L. Bogani, *Nat. Mater.*, 2016, **15**, 164.



- 10 M. Affronte, *J. Mater. Chem.*, 2009, **19**, 1731.
- 11 K. Hong and W. Y. Kim, *Angew. Chem., Int. Ed.*, 2013, **52**, 3389.
- 12 K. Katoha, H. Isshiki, T. Komedab and M. Yamashita, *Coord. Chem. Rev.*, 2011, **255**, 2124.
- 13 A. Candini, S. Klyatskaya, M. Ruben, W. Wernsdorfer and M. Affronte, *Nano Lett.*, 2011, **11**, 2634.
- 14 A. Giusti, G. Charron, S. Mazerat, J.-D. Compain, P. Mialane, A. Dolbecq, E. Riviere, W. Wernsdorfer, R. N. Biboum, B. Keita, L. Nadjo, A. Filoramo, J.-P. Bourgoin and T. Mallah, *Angew. Chem., Int. Ed.*, 2009, **48**, 4949.
- 15 L. Bogani, C. Danieli, E. Biavardi, N. Bendiab, A.-L. Barra, E. Dalcanele, W. Wernsdorfer and A. Cornia, *Angew. Chem., Int. Ed.*, 2009, **48**, 746.
- 16 J. F. Schenck, *Med. Phys.*, 1996, **23**, 815.
- 17 J. S. Miller, *Chem. Soc. Rev.*, 2011, **40**, 3266.
- 18 S. Sasha, A. D. Dmitriy, H. B. D. Geoffrey, M. K. Kevin, J. Z. Eric, A. S. Eric, D. P. Richard, T. N. SonBinh and S. R. Rodney, *Nature*, 2006, **442**, 282.
- 19 H. Xu, L. Huang, Z. Zhang, B. Chen, H. Zhong and L.-M. Peng, *Appl. Phys. Lett.*, 2013, **103**, 112405.
- 20 R. Shi, H. Xu, B. Chen, Z. Zhang and L.-M. Peng, *Appl. Phys. Lett.*, 2013, **102**, 113102.
- 21 L. Huang, H. Xu, Z. Zhang, C. Chen, J. Jiang, X. Ma, B. Chen, Z. Li, H. Zhong and L.-M. Peng, *Sci. Rep.*, 2014, **4**, 5548.
- 22 H. Xu, Z. Zhang, R. Shi, H. Liu, Z. Wang, S. Wang and L.-M. Peng, *Sci. Rep.*, 2013, **3**, 1207.
- 23 L. Huang, Z. Zhang, B. Chen, X. Ma, H. Zhong and L.-M. Peng, *Appl. Phys. Lett.*, 2014, **104**, 183106.
- 24 R. S. Popovic, *Hall Effect Devices*, Institute of Physics Publishing, London, 2004.
- 25 S. Kyatskaya, J.-R. Galan-Mascaros, L. Bogani, F. Hennrich, M. Kappes, W. Wernsdorfer and M. Ruben, *J. Am. Chem. Soc.*, 2009, **131**, 15143.
- 26 M. Lopes, A. Candini, M. Urdampilleta, A. Reserbat-Plantey, V. Bellini, S. Klyatskaya, L. Marty, M. Ruben, M. Affronte, W. Wernsdorfer and N. Bendiab, *ACS Nano*, 2010, **4**, 7531.
- 27 X. Ling, L. Xie, Y. Fang, H. Xu, H. Zhang, J. Kong, M. S. Dresselhaus, J. Zhang and Z. Liu, *Nano Lett.*, 2010, **10**, 553.
- 28 D. Geng, H. Wang, Y. Wan, Z. Xu, B. Luo, J. Xu and G. Yu, *Adv. Mater.*, 2015, **27**, 4195.
- 29 N. S. Sariciftci, A. J. Heeger, V. Krasevec, P. Venturini, D. Mihailovic, Y. Cao, J. Libert and J. L. Bredas, *Synth. Met.*, 1994, **62**, 107.
- 30 X. Dong, Y. Shi, Y. Zhao, D. Chen, J. Ye, Y. Yao, F. Gao, Z. Ni, T. Yu and Z. Shen, *Phys. Rev. Lett.*, 2009, **102**, 135501.
- 31 S. Kim, J. Nah, I. Jo, D. Shahrjerdi, L. Colombo, Z. Yao, E. Tutuc and S. K. Banerjee, *Appl. Phys. Lett.*, 2009, **94**, 062107.
- 32 B. Chen, L. Huang, X. Ma, L. Dong, Z. Zhang and L.-M. Peng, *Carbon*, 2015, **94**, 585.
- 33 G. M. Gusev, C. A. Duarte, A. A. Quivy, T. E. Lamas and J. R. Leite, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2005, **71**, 165311.
- 34 K. Tauber, A. Honemann, M. Gradhand and I. Mertig, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2015, **91**, 220404.

