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First-principles study of decomposition mechanisms of $Mg(BH_4)_2$ 2NH₃ and $LiMq(BH_4)$ ₃ \cdot 2NH₃

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The decomposition mechanisms of Mg(BH₄)₂. 2NH₃ and LiMg(BH₄)₃. 2NH₃ were studied by using density functional theory calculations. Compared to that of $Mg(BH_4)_2$: $2NH_3$, the incorporation of LiBH₄ with the formation of LiMg(BH₄)₃ \cdot 2NH₃ slightly increased Bader charges of B atoms, meanwhile it decreased Bader charges of N atoms. Mg(BH₄)₂ 2NH₃ shows a low ammonia vacancy diffusion barrier, but relatively high ammonia vacancy formation energy, which lead to a low concentration of $NH₃$ vacancies and limit NH₃ transportation. In contrast to that of Mg(BH₄)₂.2NH₃, LiMg(BH₄)₃.2NH₃ has a relatively high ammonia vacancy formation energy and diffusion barrier, which suppresses ammonia release. The incorporation of LiBH₄ and Mg(BH₄)₂.2NH₃ does not decrease but increases the hydrogen formation barrier of LiMg(BH₄)₃. 2NH₃, resulting in a slight increase in the dehydrogenation peak temperature, consistent with experimental results. PAPER

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Introduction

Recently, many efforts have been devoted to B–N based chemical hydrides as potential hydrogen storage materials because of their high theoretical hydrogen capacity. $1-3$ For instance, ammonia borane (AB), with a high H-capacity of 19.6 wt%, is a typical B–N based hydride for chemical hydrogen storage.¹ However, upon decomposition of AB, accompanied volatile compounds of ammonia, diborane, and borazine are evolved, which lead to a reduction of dehydrogenation capacity and are fatal for fuel cell applications.^{1,4} Many different approaches have been adopted to facilitate hydrogen release from AB during the last decade.⁵⁻¹⁰ Recent studies show that the substitution of H atoms in the $NH₃$ unit of AB by alkali metals with the formation of single or double metal amidoborane (MAB) is an effective way to improve the dehydrogenation properties of AB in terms of the reduced H_2 release temperatures, accelerated H_2 release kinetics, and minimized borazine evolved.5,6,11–¹⁷

Ammine metal borohydrides (AMBs), which show favourable hydrogen storage properties competitive with ammonia borane, have been developed recently as promising materials for hydrogen storage.¹⁸–³¹ However, many of these composites suffer from the release of undesirable gas of ammonia during

dehydrogenation. Further experimental results show that the purity of gas released and dehydrogenation temperature of AMBs can be improved by using double-cation substitutions approach and tuning $BH₄/NH₃$ ration.^{15,24,32,33} The experimental and theoretic studies indicate that ammonia is weakly bound to the metal cations with low electronegativity $($ 4.2) in AMBs, therefore tend to release ammonia at low temperature.^{32,34} Although these studies have provided valuable insight for understanding the decomposition processes of single metal cation AMBs, the results may not be applicable to double cations AMBs. For instance, $Mg(BH_4)_2$ 2NH₃ (with electronegativity of 1.31 for Mg cation) mainly release hydrogen along with a small amount of ammonia.²³ The incorporation of ${\rm LiBH_4}$ (with low electronegativity of 0.98 for Li cation) and $Mg(BH_4)_2$ 2NH₃ with the formation of double cations ammine borohydride, $LiMg(BH₄)₃$ 2NH₃ results in improving the purity of gas released compared to $Mg(BH_4)_2 \cdot 2NH_3$.³⁵ Further improved dehydrogenation of ammine magnesium borohydride by tuning the NH₃/BH₄ ratios and combining Mg(BH₄)₂ 2NH₃ with MgH₂ and NaAlH₄ were reported.³⁶⁻³⁹

The mixed-cation strategy offers a promising route toward tuneable dehydrogenation of ammine metal borohydrides, however, a detail study of the dehydrogenation mechanism is still needed for further improving their dehydrogenation performance. Herein, we presented a comparison study of the electronic structure and dehydrogenation mechanisms of $Mg(BH_4)_2$ 2NH₃ and LiMg(BH₄)₃ 2NH₃ by density functional theory (DFT) calculation.

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Computational method

 $Mg(BH_4)_2$ 2NH₃ crystallizes in the orthorhombic structure with space group of *Pcab* and lattice parameters of $a = 17.4872(4)$ Å, $b = 9.4132(2)$ Å, $c = 8.7304(2)$ Å.²³ LiMg(BH₄)₃.2NH₃ has a hexagonal structure with space group $P6₃$ and lattice constants of $a = b = 8.0002(1)$ Å and $c = 8.3944$ Å.³⁵ The geometric structures were optimized by DFT calculation as implemented in MedeA@VASP code.⁴⁰ To describe the weak van der Waals H⁺···⁻⁻H dihydrogen bonds, the optB86b-vdW functional⁴¹⁻⁴³ was adopted for geometric optimization. Plane waves with kinetic energy cutoff of 500 eV were used. The generalized gradient approximation (GGA) of Perdew–Burke–Ernzerhof (PBE) was adapted to treat the exchange and correlation of electronics.44,45 The projector-augmented wave (PAW) approach was used to describe the electron-ion interactions⁴⁶ with 1s2s2p of Li, s2p1 of B, s2p3 of N, s2p0 of Mg as the explicit valence electrons. The Brillouin zones were sampled by Monkhorst– Pack k -point meshes⁴⁷ with meshes points spacing less than 0.05 per A for both $Mg(BH_4)_2$ 2NH₃ and LiMg(BH₄)₃ 2NH₃. Structural relaxations of atomic positions were carried out until the total energies and residual forces were less than 1.0 \times 10^{-5} eV and 0.02 eV $\rm \AA^{-1},$ respectively. For the calculation of NH $_3$ vacancy formation energies and H₂ formation energies, $1 \times 2 \times$ 2 supercells of Mg(BH₄)₂ \cdot 2NH₃ and 2 \times 2 \times 2 supercells of $Limg(BH₄)₃$ 2NH₃ were used. Our tests showed that the used of $1 \times 2 \times 2$ supercells of Mg(BH₄)₂ \cdot 2NH₃ and 2 \times 2 \times 2 supercells of LiMg $(BH_4)_3$ 2NH₃ with k-point mesh spacing less than 0.05 per Å yield energies that converged within 0.01 eV $(\mathrm{f.u.})^{-1}.$ The NH₃ diffusion barriers and H₂ formation barriers were estimated by using climbing image nudged elastic band (CI-NEB) method.^{48,49} **EXAMELIBY ARTIFIC COMPUTATE (The CALL COMPUTATE CO**

The $NH₃$ vacancy formation energy was estimated using the following equation:

$$
E_{\rm c}=E_{\rm total}-E(\rm AMBs-NH_3)-E(NH_3)
$$

where E_{total} is the total energy of the AMBs supercells; $E(NH_3)$ represents the energy of isolate $NH₃$ molecule; $E(AMBs-NH₃)$ is

the total energy of the AMBs supercells after $NH₃$ molecules are removed. The positive energy of E_c indicates that the creation of $NH₃$ vacancy is an endothermic process; while the negative energy of E_c indicates that the creation of NH_3 vacancy is an exothermic process.

The concentration of ammonia vacancy in $Mg(BH_4)_2 \cdot 2NH_3$ and $LiMg(BH₄)₃·2NH₃$ could be estimated by the following equation⁵⁰

$$
c = N_{\text{sites}} N_{\text{config}} \exp(E_{\text{c}}/k)
$$

where E_c is the formation energy of NH₃ vacancy; N_{sites} represents the number of sites that the defect can be incorporated; N_{config} is the number of configurations per site in which the vacancy can be formed; k and T represent Boltzmann constant and temperature, respectively.

Results and discussion

Electronic structure

The electron localization function (ELF) and charge transfer between the H, N, B atoms and metal cations (Li and Mg) were analysed to understand the bonding characters of $Mg(BH_4)_{2}$ - $2NH_3$ and LiMg(BH₄)₃ $2NH_3$. The H atoms bond to N atom and B atom are represent as (N) H and (B) H, respectively. As shown in Fig. 1, the calculated ELF shows the covalent bonding of N–H and B–H. Although the Mg–H bonds are mainly ionic, the distorted ELF isosurfaces around (B)H, (N)H and Mg indicate partial covalent bond feature of Mg–H. The low ELF value around Li indicates the essentially ionic bonding character of Li–H. Table 1 shows the Bader charges of $(B)H$, $N(H)$, N, B and Mg for $Mg(BH_4)_2 \cdot 2NH_3$ are $-0.58/-0.64$, 0.44, -1.30 , 1.59 and 1.65, respectively. The Bader charge of Li is 0.90 for $Limg(BH₄)₃$ · 2NH₃, indicates a strong ionization of the Li cation. Hence, Li cation transfers most of its 2s electron to neighbouring $BH₄$ unit, similar to that of LiBH₄. Compared to $Mg(BH_4)_2$ 2NH₃, the incorporation of LiBH₄ with the formation of LiMg $(BH_4)_3$ 2NH₃ barely affects the charge distribution of H

Fig. 1 The calculated electron localization function (ELF) for (a) Mg(BH₄)₂ 2NH₃ and (c) LiMg(BH₄)₃ 2NH₃ plotted as yellow-colored transparent isosurfaces at a level of 0.6; (b) and (d) present zoomed-in view showing more details for Mg(BH₄₎₂ $2NH_3$ and LiMg(BH₄)₃. 2NH₃, respectively. The green, pink, orange, purple and blue colors represent B, H, Mg, Li and N atoms, respectively.

Table 1 Bader charge of $Mg(BH_4)_2 \cdot 2NH_3$ and LiMg(BH₄)₃ $\cdot 2NH_3$

	Bader charge		
Atom	$Mg(BH_4)_2 \cdot 2NH_3$	$LiMg(BH4)3 \cdot 2NH3$	
Li		0.90	
Mg	1.65	1.64	
B	1.59	1.67	
N	-1.30	$-1.32/-1.34$	
(B)H	$-0.58/-0.64$	-0.63	
(N) H	0.44	0.45	

and Mg. The Bader charge of B is slightly increased and Bader charge of N is slightly decreased in $Limg(BH_4)_3 \cdot 2NH_3$.

Ammonia vacancy formation energies and diffusion barriers

As demonstrated by previous report, $Mg(BH_4)_2$ 2NH₃ started to release hydrogen at temperature around 120 °C, with a maximum hydrogen release rate at $205 °C.^23$ A small amount of $NH₃$ was released along with hydrogen evolution from $Mg(BH_4)_2$ 2NH₃. The LiMg(BH₄)₃ 2NH₃ shows dehydrogenation performance comparable to that of $Mg(BH_4)_2$ 2NH₃, with dehydrogenation peak located at 221 °C.³⁵ In addition, incorporation of LiBH₄ with Mg(BH₄)₂ 2NH₃ suppresses ammonia release.

The formation and transport properties of $NH₃$ vacancy are crucial to the thermodynamics and kinetics of ammonia release from AMBs. To understand the microscopic mechanisms behind the release of ammonia, the formation and diffusivity of NH₃ were studied. The NH₃ vacancy was created by directly removed a NH₃ unit from $Mg(BH_4)_{2} \cdot 2NH_3$ and $LiMg(BH₄)₃·2NH₃.$

As shown in Table 2, the calculated $NH₃$ removal energies are 1.81 and 1.97 eV for $Mg(BH_4)_2 \cdot 2NH_3$ and $LiMg(BH_4)_3 \cdot 2NH_3$, respectively. The relatively high NH₃ removal energies indicates that the formation of $NH₃$ vacancies are thermodynamics unfavourable at low temperature, resulting in low concentration of ammonia vacancy for those two composites. The high formation energy of $NH₃$ vacancies can attribute to the coordination bond of Mg-N and $H^+ \cdots$ ⁻H dihydrogen network.

In addition to the formation energies of $NH₃$ vacancies, the diffusive of $NH₃$ vacancies is also importance for ammonia release. The diffusion paths were calculated by moving a NH₃ unit from a nearby lattice site into the vacancy. The diffusion barrier is defined as the energy difference between the saddle point and the ground state. The activation energy (O) for selfdiffusion of ammonia can be obtained by combining the

Table 2 Calculated NH₃ vacancy formation energies (E_c), diffusion barriers (E_b) and activation energies ($Q = E_b + E_c$) for Mg(BH₄)₂ 2NH₃ and $LiMg(BH₄)₃$ 2NH₃

	E_c (eV)	$E_{\rm b}$ (eV)	Q (eV)
$Mg(BH_4)_2 \cdot 2NH_3$	1.81	0.26	2.07
$LiMg(BH4)3 \cdot 2NH3$	1.97	1.31	3.28

calculated vacancy formation energy with the diffusion barrier. As summarized in Fig. 2 and Table 2, for $Mg(BH_4)_2 \cdot 2NH_3$, the calculated energy barrier and activation energy of ammonia diffusion are 0.26 and 2.07 eV, respectively. It should be noted that the $NH₃$ diffusion barrier is relatively low, the formation energy of $NH₃$ vacancy is the dominate term in the activation energy for ammonia diffusion. The relatively high formation energy would result in low concentration of $NH₃$ vacancy, which limit its transport in $Mg(BH_4)_2 \cdot 2NH_3$. This is in agreement with previous report that only a small amount of $NH₃$ was released during decomposition of $Mg(BH₄)₂ \cdot 2NH₃.²³$

The calculated ammonia vacancy diffusion barrier and activation energy of $Limg(BH_4)_3 \cdot 2NH_3$ are 1.31 and 3.28 eV, respectively. Compared to that of $Mg(BH_4)_2 \cdot 2NH_3$, the relatively high ammonia diffusion barrier and activation energy indicate that low concentration and mobility of ammonia vacancy in $LIMg(BH₄)₃$ 2NH₃, in inconsistent with experimental results that the dehydrogenation purity of $Mg(BH_4)_2$ 2NH₃ can be improved by introducing LiBH4 with the formation of $LiMg(BH₄)₃ \cdot 2NH₃$.³⁵

Hydrogen formation energies and barriers

Our previous studies suggest that the initial dehydrogenation of AMBs is achieved by combination of H atoms from $NH₃$ and H atoms from BH_4 groups.^{51,52} Therefore, H_2 formation energies were calculated by moving one (N)H and one (B)H atom away from host N or B atom to form a hydrogen molecule with H–H distance of 0.74 Å in the supercell of AMBs. The geometry optimization was first performed by fixed the $H₂$ positions and relaxed the rest of the atoms, following by full relaxed all of the atoms in the supercell. In agreement with our previous studies, the formation of H_2 molecules lead to significant rearrangement of the surrounding lattice, which may result in overestimated the hydrogen formation energies. In addition, both $Mg(BH_4)_2 \cdot 2NH_3$ and $LiMg(BH_4)_3 \cdot 2NH_3$ started to release hydrogen at temperature higher than their melting point. In other word, the crystal structure of $Mg(BH_4)_2$ 2NH₃ and $Limg(BH₄)₃$ · $2NH₃$ disappeared before hydrogen evolved. Therefore, we further calculated the hydrogen formation energies by using the molecule model in which two formula units of $Mg(BH_4)_2$ 2NH₃ and LiMg(BH₄)₃ 2NH₃ were placed in a cubic cell with lattice parameter of 20 A. Paper

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> As shown in Table 3, the calculated hydrogen formation energies by using supercell of AMBs are 0.84 and 1.22 eV for $Mg(BH_4)_2$ 2NH₃ and LiMg(BH₄)₃ 2NH₃, respectively. In consistent with our previous theoretical study, the dissociation of H_2 results in dramatic movement of around atoms.⁵¹ The hydrogen formation energies calculated by molecule model are -0.11 and 0.08 eV for $Mg(BH_4)_2 \cdot 2NH_3$ and LiMg(BH₄)₃ \cdot 2NH₃, respectively. The combination of the (N)H and (B)H results in rearrangement of the surrounding atoms, similar with our previous report.⁵¹ The NH₂ and BH₃ units reoriented and BH₃ units moved toward NH₂ to form NH₂-BH₃ complexes. The N–B distances reduce to 1.58 \AA , indicating the formation of N–B bond during dehydrogenation, in agreement with experimental observation.^{14,35} However, the lengths of Li-

Fig. 2 The calculated energetic profiles, initial (IS), transition (TS) and final (FS) geometric structure of NH₃ diffusion for (a) Mg(BH₄)₂·2NH₃ and (b) LiMg(BH₄)₃ · 2NH₃. E_b represents the calculated energy barrier. The green, pink, orange, purple and blue colors represent B, H, Mg, Li and N atoms, respectively.

Table 3 Formation energies of H_2 release via (N)H and (B)H combination by using crystal model (E_{H_2-C}) and molecule model (E_{H_2-M})

	$E_{\rm H, -C}$ (eV)	$E_{\text{H}_2\text{-M}}$ (eV)
$Mg(BH_4)_2 \cdot 2NH_3$	0.84	-0.11
$LiMg(BH4)3 \cdot 2NH3$	1.22	0.08

N, Mg-N, B-H and N-H bonds keep almost the same after structural rearrangements.

The low H₂ formation energies of $Mg(BH_4)_2$ ²NH₃ and $L iMg(BH₄)₃$ 2NH₃ suggest that the formation of hydrogen is thermodynamic favourable at low temperature. Therefore, the energy barrier of $H₂$ formation is the key of those two composites release hydrogen at temperature above 100 °C.

We further calculated the energy barriers of $H₂$ formation and the results were summarized in Fig. 3. The calculated energy barrier of H_2 formation from $Mg(BH_4)_2 \cdot 2NH_3$ is 2.20 eV. The transition geometric structure of hydrogen release from $Mg(BH_4)_2$ 2NH₃ (Fig. 3(a)) shows the broken of B–H and N–H bonds with the formation of H_2 molecule. The H_2 molecule is located between NH_2 and BH_3 units with H_2-NH_2 and H_2-BH_3 distances of 2.47 and 2.31 Å, respectively. Meanwhile, the $NH₂$ unit move toward Mg cation and lead to slightly reduce the Mg– N distance from 2.16 to 1.95 Å. The calculated hydrogen formation energy barrier of $LiMg(BH₄)₃$ · $2NH₃$ is 2.55 eV, which is 0.35 eV higher than that of $Mg(BH_4)_2 \cdot 2NH_3$. Previous experimental results show dehydrogenation peak of 205 \degree C and

221 °C for $Mg(BH_4)_2$ 2NH₃ and LiMg(BH₄)₃ 2NH₃, respectively.23,35 The relatively high dehydrogenation peak of $LiMg(BH₄)₃$ and be attributed to the high hydrogen formation barrier. The transition geometric structure of hydrogen release from $LIMg(BH₄)₃$ 2NH₃ is similar to that of $Mg(BH_4)_2 \cdot 2NH_3$. The H_2 molecule is located between NH₂ and $BH₃$ unit. The H₂-NH₂ distance in transition structure of LiMg $(BH_4)_3$ \cdot 2NH₃ is 2.24 Å, which is 0.23 Å shorter than that in $Mg(BH_4)_2$ 2NH₃. And the H₂–BH₃ distance in transition structure of LiMg $(BH_4)_3$ 2NH₃ is 2.11 Å, which is 0.20 Å shorter than that in $Mg(BH_4)_2 \cdot 2NH_3$. In addition, the NH₂ unit shortening its distance to the Mg cation from 2.19 to 1.95 \AA , similar to that of $Mg(BH_4)_2 \cdot 2NH_3.$

Although the above calculations show a low $NH₃$ diffusion barrier for $Mg(BH_4)_2$ 2NH₃, the formation energy of NH₃ vacancy is relatively high, which results in low concentration of NH₃ vacancy in Mg(BH₄)₂ · 2NH₃. Therefore, the Mg(BH₄)₂ · 2NH₃ mainly releases hydrogen accompany with a small amount of ammonia during decomposition. In contrast to that of $Mg(BH_4)_2 \cdot 2NH_3$, Li $Mg(BH_4)_3 \cdot 2NH_3$ shows relatively high formation energy and diffusion barrier of $NH₃$ vacancy, which limit both the concentration and transport of ammonia, therefore improve the dehydrogenation purity. The calculated hydrogen formation barrier of $LiMg(BH₄)₃$ 2NH₃ is slightly higher than that of $Mg(BH_4)_2 \cdot 2NH_3$, therefore incorporation of LiBH₄ with Mg(BH₄)₂ \cdot 2NH₃ may not decrease the dehydrogenation temperature.

Fig. 3 The calculated energetic profiles, initial (IS), transition (TS) and final (FS) geometric structure of H₂ formation for (a) Mg(BH₄)₂. 2NH₃ and (b) LiMg(BH₄)₃. 2NH₃. E_b represents the calculated energy barrier. The green, pink, orange, purple and blue colours represent B, H, Mg, Li and N atoms, respectively.

Conclusions

First-principles calculations based on density functional theory were carried out to investigate the decomposition mechanisms of $Mg(BH_4)_2$ 2NH₃ and LiMg(BH₄)₃ 2NH₃. The electronic structure analysis indicates that Mg–H interaction in those two composites are mainly ionic with partial covalent bond feature. The incorporation of LiBH₄ and Mg(BH₄)₂.2NH₃ with the formation of $LiMg(BH_4)_3 \cdot 2NH_3$ barely affects the charge distribution of H and Mg. The Bader charge of B is slightly increased and Bader charge of N decreased due to the incorporation of LiBH₄. Although the NH₃ diffusion barrier for Mg(BH₄)₂ · 2NH₃ is low, the relatively high formation energy of $NH₃$ vacancy lead to low concentration of NH₃ vacancy and limit its transportation, in agreement with experimental results that $Mg(BH_4)_2 \cdot 2NH_3$ mainly releases hydrogen along with a small amount of ammonia. The $LiMg(BH₄)₃$ 2NH₃ shows relatively high ammonia vacancy formation energy and diffusion barrier, which suppress ammonia release compared to $Mg(BH₄)₂ \cdot 2NH₃$. The incorporation of LiBH₄ and Mg(BH₄)₂ \cdot 2NH₃ does not decrease the hydrogen formation barriers, instead slightly increase the hydrogen formation barriers of $Limg(BH_4)_3 \cdot 2NH_3$, in agreement with experimental results that $\text{Limg}(BH_4)_3 \cdot 2NH_3$ shows a dehydrogenation peak slightly higher than that of $Mg(BH_4)_2 \cdot 2NH_3$.

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