RSC Advances



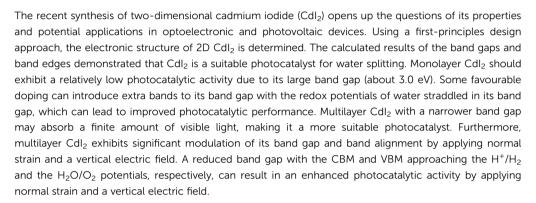
PAPER View Article Online View Journal | View Issue



Cite this: RSC Adv., 2017, 7, 53653

Computational design of enhanced photocatalytic activity of two-dimensional cadmium iodide†

Lin Tao and Le Huang **



Received 31st August 2017 Accepted 1st November 2017

DOI: 10.1039/c7ra09687a

rsc.li/rsc-advances

Introduction

The conversion of solar light into hydrogen by photocatalytic water splitting is an attractive technology for producing an alternative energy resource to fossil fuels without environmental pollution.1-4 To achieve high efficiency, the photosatisfy several requirements.^{5,6} photocatalyst should have suitable band edges to straddle the redox potentials of water. Specifically, its conduction band minimum (CBM) should be higher than the hydrogen reduction potential (H⁺/H₂) and the valence band maximum (VBM) should be lower than the water oxidation potential (H_2O/O_2) . The standard H+/H2 and H2O/O2 potentials with respect to the vacuum level are -4.44 eV and -5.67 eV, respectively.7 In addition to this, the photocatalyst should possess a band gap between 1.5 and 3.0 eV in order to utilize UV light and visible light.8,9

Recently, various two-dimensional (2D) materials have been predicted to have remarkable photocatalytic activities. ¹⁰⁻¹² First-principles calculations have predicted that single-layer group-III monochalcogenides are suitable for photocatalytic water splitting. ¹³ Some previous works demonstrated that single-layer MoS₂, WS₂, PtS₂ and PtSe₂ are potential photocatalysts due to their suitable band edges with respect to the redox potentials of water. ¹⁴⁻¹⁶ 2D SnS and GeSe materials are demonstrated to be

School of Materials and Energy, Guangdong University of Technology, Guangzhou, Guangdong 510006, China. E-mail: huangle@semi.ac.cn

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c7ra09687a

promising alternative materials for optoelectronic applications. 17-20 Furthermore, heterobilayer GeSe/SnS could have promising applications in photovoltaic devices by modifying its electronic structure. 11 These 2D materials maximize the surface area available for water splitting and minimize the distance of generated electron and hole migration, which enhances their catalytic performance. 22

Beyond 2D chalcogenide materials, metal halide materials with layered crystal structures (*e.g.*, PbI₂ and CdI₂) have also attracted significant attention due to their importance as precursors for the synthesis of halide perovskites.^{23–27} Our previous works have demonstrated that layered PbI₂ possesses great potential in photovoltaic devices.^{28,29} 2D CdI₂ nanoplates have also been synthesized by using a vapor transport and deposition approach.^{25,30,31} However, the physical properties of 2D CdI₂ are far from clear. In particular, we focus on layered CdI₂ as a possible photocatalyst and propose some effective approaches to improve its photocatalytic performance.

Computational methods

All calculations are implemented in the VASP code^{32,33} within the projector-augmented plane-wave method.³⁴ The generalized gradient approximation (GGA) of the Perdew, Burke and Ernzerhof (PBE) functional³⁵ is adopted for the electron exchange and correlation. The plane-wave cutoff energy is set as 450 eV and the HSE06 hybrid functional^{36–38} is used to calculate the band structure of monolayer CdI₂. A vacuum larger than 15 Å is used to simulate the isolated sheet. The pair-wise force field in the DFT-D2 method of Grimme³⁹ is employed in the

RSC Advances

calculations of multilayer CdI2 to correctly describe the vdW interactions between layers. For the relaxation of CdI₂ of various thicknesses, a (15 × 15 × 1) Monkhorst-Pack grid⁴⁰ is used, whereas for the relaxation of a doped CdI2 monolayer, the first Brillouin zone is sampled with a $(5 \times 5 \times 1)$ Monkhorst-Pack grid. All structures are fully relaxed with a force tolerance of 0.02 eV Å^{-1} .

Results and discussion

Compared with layered transition metal dichalcogenides, CdI₂ has a layered structure with a hexagonal unit cell and it crystallizes in the space group P63mc. This indicates that the 1T phase of CdI₂ is the most stable, rather than the 2H phase. As shown in Fig. 1a, each sheet of Cd atoms is sandwiched between two sheets of I atoms through strong ionic bonding, 41,42 and the atomic layers combined together via weak van der Waals (vdW) interactions. The optimized lattice constant of monolayer CdI₂ is 4.33 Å. Its band structure calculated on the PBE level indicates an indirect band gap of 2.53 eV in Fig. 1b. According to the calculated density of states (Fig. 1c), the VBM is mainly contributed to by Cd_p and I_p orbitals and the CBM is composed of I_p and Cd_s orbitals. In order to correct the PBE band gaps, a hybrid Heyd-Scuseria-Ernzerhof functional is adopted to describe the electron exchange and correlation (Fig. S1†). The general features of the band structure are well preserved except for a larger band gap than that of the PBE functional. Furthermore, the spin-orbit coupling (SOC) effect is taken into consideration in Fig. S1.† An obvious spin-orbit splitting is observed at the VBM of the band structure. As a result, the initial quadruple degenerate VBM becomes doubly degenerate. Meanwhile, the SOC exerts little influence on the

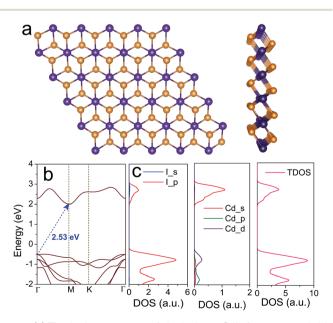


Fig. 1 (a) The lattice structure of single-layer Cdl₂ from top and side views. The purple and brown balls denote the Cd and I atoms, respectively. (b) The band structure of single-layer Cdl₂ and (c) the projected density of states (PDOS) and total density of states (TDOS).

CBM of the CdI₂ bands. This is attributed to the strong coupling between the Cd d and I p orbitals at the VBM and the absence of the Cd_d orbital at the CBM in the band structure of CdI₂. Although an indirect band gap is not favourable for water splitting, layered CdI₂ has suitable band edges to straddle the redox potentials of water. The photocatalytic activity can be enhanced when combining CdI2 with other photocatalysts, such as C₃N₄, to construct a direct Z-scheme type photocatalysis system for water splitting.43,44 What is more, the imaginary part of the dielectric function and light absorption (Fig. S2†) indicate good absorption in solar light. Our results are in good agreement with previous works except for an energy shift that is due to the underestimation of the GGA-PBE functional. CdI₂ thin films are reported by experimental works to have a direct optical energy gap of ~3.5 eV which seems too large for photocatalysis.30,31 It is indicated by our results and previous experimental works that CdI2 exhibits good absorption in ultraviolet light and has great potential in photocatalysis. Efficient ways to modulate its band edges are desired to improve the optoelectronic properties and the photocatalytic activity.

It seems that the band gap of monolayer CdI2 is too large to efficiently utilize the visible light in photocatalysis. It is well known that the electronic properties including the band gaps of most layered transition metal dichalcogenides are strongly dependent on the thickness of the materials. 45,46 Thus we checked the thickness-dependent electronic structure of CdI₂ as shown in Fig. 2a and b. It is indicated that both the CBM and the VBM approach the Fermi level, resulting in a decrease of the band gap of CdI2 with thickness. The vacuum level is taken as a reference in Fig. 2b. It is indicated in Fig. 2b that the energies of both the CBM and the VBM increase with the number of layers. Furthermore, the band gap decreases gradually with the number of layers. The evolution of the band edges and band gap with the thickness is not very significant because of the rather weak interlayer coupling. What is more, the redox potentials of water are straddled between the band edges of CdI2 with various thicknesses. The reduced band gap can make CdI2 a good photocatalyst with it absorbing solar light over a wider wavelength range. Consequently, the photocatalytic performance can be improved to some extent by increasing the thickness of CdI₂ to a suitable value.

According to Fig. 2a and b, the band edges and band gap of CdI₂ exhibit weak dependence on the thickness with respect to most 2D transition metal dichalcogenides. An efficient approach to improve its photocatalytic performance is desired. It has been demonstrated that 2D materials exhibit significant modulation of their electronic structures by applying a normal strain. The evolution of the band structures of bilayer Cd₂ with an applied normal strain is shown in Fig. S3.† The normal strain is calculated as $\varepsilon = (d_0 - d)/d_0$. Our results indicate that the band gap of bilayer CdI2 decreases with the applied normal strain (Fig. 2c). Clear insight can be obtained from the variation of the band edges of bilayer CdI₂ with the applied normal strain (Fig. 2d). The VBM increases gradually with the applied strain while the CBM does not show a monotonous trend because of the CBM's shift from the M point to the Γ point. Furthermore, the redox potentials of water are always straddled between the

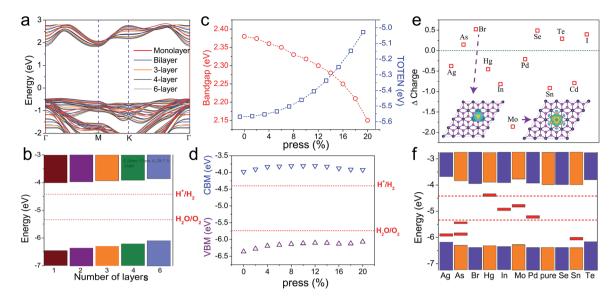


Fig. 2 (a) Band structures and (b) band alignments of Cdl_2 of different thickness. Evolution of (c) band gap (red circles), total energy (blue squares) and (d) band edges of the Cdl_2 bilayer as functions of the applied normal strain. (e) Transfer charge results from Bader charge analysis in the doped Cdl_2 monolayer. The insets are the charge density difference in the Mo-doped and the Br-doped Cdl_2 monolayer. (f) Band alignments of the doped Cdl_2 monolayer. The red stripes in the band gap are the dopant bands introduced by substitution. The vacuum level is set as zero and is taken as a reference in (b), (d), and (f). The redox potentials of water are plotted as red dashed lines.

VBM and the CBM of the pressed CdI₂ bilayer. The CdI₂ bilayer with a reduced band gap from the application of normal strain is a good photocatalyst with improved performance due to the widened wavelength range for light absorption.

It should be mentioned that the interlayer binding in the CdI_2 bilayer is relatively weaker than that in most of the transition metal dichalcogenide bilayers such as those of MoS_2 . Fig. 2c also shows the evolution of the total energy of the CdI_2 bilayer. The charge density difference (Fig. S4†) shows slight charge transfer between layers in the CdI_2 bilayer upon binding together. The weak binding gives rise to the relatively weak response of the electronic structure and strain energy to the applied strain. Similar evidence comes from the integrated charge density of the CdI_2 bilayer under different normal strain (Fig. S5†). It is found that applying a normal strain will bring about extra charge exchanging between layers. It can be concluded that applying a normal strain is an efficient way to improve the photocatalytic performance.

Besides applying normal strain, substituted doping has also been demonstrated to be an effective method to modulate the band alignment of a semiconductor. In Fig. 2e and f, several non-magnetic atoms are employed to make substituting dopants in a (4 \times 4 \times 1) CdI $_2$ monolayer supercell. Dopant bands are induced either in the band gap or near the band edges (Fig. S6†), which should exert a significant influence on the optoelectronic and photovoltaic properties of monolayer CdI $_2$. In Fig. 2e, Bader charge analysis demonstrates that charge transfer occurs between dopants and it defected the CdI $_2$ monolayer.

The band alignments of these doped CdI_2 monolayers are plotted in Fig. 2f. It is found that In, Mo and Pd will introduce dopant bands in the redox potential of water, making the

corresponding doped CdI_2 not suitable for photocatalysis, while other dopants employed here will not introduce extra bands in the redox potential of water. Ag, As and Sn dopants will induce extra bands between the VBM and the H_2O/O_2 potential, and Hg can induce extra bands between the H^+/H_2 potential and the CBM, making the corresponding doped CdI_2 suitable for photocatalysis in water splitting.

In many optoelectronic and photovoltaic devices, the materials are always subjected to a vertical electric field (E_{\perp}) . Multilayer 2D materials usually exhibit significant modulation of their electronic structures by applying an E_{\perp} . As shown in Fig. 3a, the band structures of the CdI₂ bilayer under different E_{\perp} are calculated. Our results indicate that its band gap is significantly reduced by the applied E_{\perp} . Further insight can be obtained from the variation of the band edges with the E_{\perp} in Fig. 3b. It is clear that the CBM exhibits a linear descending trend with the E_{\perp} while the VBM is hardly influenced by the E_{\perp} . As a photocatalyst, when subjected to a vertical electric field, the CBM of CdI₂ will approach the H⁺/H₂ level. This will improve the absorption of solar light, further leading to enhanced photocatalytic activity. However, a strong E_{\perp} will reduce the CBM of CdI₂ to the position below the H⁺/H₂ level, making CdI₂ unsuitable as a photocatalyst.

To make the distinct responses of the VBM and the CBM to the applied E_{\perp} clear, the band decomposed charge density of the CBM and the VBM at the four high-symmetry k-points (Γ , M, K, A) is shown in Fig. 3c and d. A is the k-point between K and Γ , which is given in Fig. 3a. As can be seen, the lowest-energy carriers (holes) are almost completely localized around the I atoms, forming π bonds. The absence of a polarized distribution between the I and Cd atoms results in the weak response of the VBM to the applied E_{\perp} , while at the CBM, the lowest-energy

RSC Advances Paper

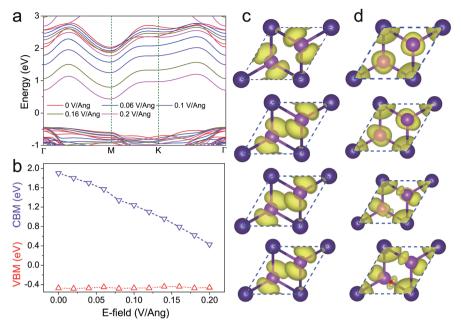


Fig. 3 Evolution of the (a) band structure and (b) band edges of the Cdl_2 bilayer as functions of the applied vertical electric field. (c) and (d) plot the band decomposed charge density of the VBM and CBM at four high-symmetry k-points (Γ , M, K, A) of the Cdl_2 bilayer, respectively. The A point is the k-point between the K and the Γ points.

carriers (electrons) are neither localized on the I atoms nor the Cd atoms. Charge exchange exists between I and Cd atoms. As a result, the applied E_{\perp} exerts a significant influence on the CBM.

Conclusions

In summary, the electronic structure of 2D CdI₂ is explored by density functional theory calculations. The calculated results of the band gaps and band edges demonstrated that CdI₂ is a suitable photocatalyst for water splitting. Monolayer CdI₂ should exhibit a low photocatalytic activity due to its large band gap (about 3.0 eV). Multilayer CdI₂ with a narrower band gap may absorb a finite amount of visible light, making it a more suitable photocatalyst. What is more, its photocatalytic activity can be enhanced by applying normal strain and a vertical electric field with two distinct mechanisms. Both the CBM and VBM of bilayer CdI₂ can be tuned by applying a normal strain, while a vertical E_{\perp} is only able to modulate the CBM of bilayer CdI2. Some favourable substituting dopants can also make the band alignment of CdI₂ more suitable for photocatalysis. These results provide valuable guidance for enhancing the photocatalytic activity of 2D CdI2.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was financially supported by the National Natural Science Foundation of China (Grant no. 11674310, 61622406,

61571415, 51502283, and 51602065), the National Key Research and Development Program of China (Grant No. 2016YFB0700700), "Hundred Talents Program" of Chinese Academy of Sciences (CAS), and the CAS/SAFEA International Partnership Program for Creative Research Teams.

References

- T. Simon, N. Bouchonville, M. J. Berr, A. Vaneski, A. Adrovic,
 D. Volbers, R. Wyrwich, M. Döblinger, A. S. Susha,
 A. L. Rogach, et al., Nat. Mater., 2014, 13, 1013.
- 2 X. Chen, L. Liu, Y. Y. Peter and S. S. Mao, *Science*, 2011, 331, 746.
- 3 R. Asahi, T. Morikawa, H. Irie and T. Ohwaki, *Chem. Rev.*, 2014, **114**, 9824.
- 4 J. Ran, J. Zhang, J. Yu, M. Jaroniec and S. Z. Qiao, *Chem. Soc. Rev.*, 2014, 43, 7787.
- 5 Y. Gai, J. Li, S.-S. Li, J.-B. Xia and S.-H. Wei, *Phys. Rev. Lett.*, 2009, **102**, 036402.
- 6 J. Kang, S. Tongay, J. Zhou, J. Li and J. Wu, Appl. Phys. Lett., 2013, 102, 012111.
- 7 V. Chakrapani, J. C. Angus, A. B. Anderson, S. D. Wolter, B. R. Stoner and G. U. Sumanasekera, *Science*, 2007, **318**, 1424.
- 8 A. K. Singh, K. Mathew, H. L. Zhuang and R. G. Hennig, *J. Phys. Chem. Lett.*, 2015, **6**, 1087.
- 9 S. Chen and L.-W. Wang, Chem. Mater., 2012, 24, 3659.
- 10 W. J. Ong, Front. Mater., 2017, 4, 11.
- D. Zeng, L. Xiao, W. J. Ong, P. Wu, H. Zheng, Y. Chen and
 D. L. Peng, J. Mater. Chem. A, 2017, 5, 16171–16178.
- 12 X. Lv, W. Wei, Q. Sun, F. Li, B. Huang and Y. Dai, *Appl. Catal.*, *B*, 2017, 217, 275.

- 13 H. L. Zhuang and R. G. Hennig, Chem. Mater., 2013, 25, 3232.
- 14 H. L. Zhuang and R. G. Hennig, *J. Phys. Chem. C*, 2013, **117**, 20440.
- 15 N. Singh, G. Jabbour and U. Schwingenschlögl, *Eur. Phys. J. B*, 2012, **85**, 392.
- 16 G. Guo and W. Liang, J. Phys. C: Solid State Phys., 1986, 19, 995.
- 17 P. D. Antunez, J. J. Buckley and R. L. Brutchey, *Nanoscale*, 2011, 3, 2399.
- 18 A. K. Singh and R. G. Hennig, *Appl. Phys. Lett.*, 2014, **105**, 042103.
- 19 Z. Ma, B. Wang, L. Ou, Y. Zhang, X. Zhang and Z. Zhou, *Nanotechnology*, 2016, 27, 415203.
- 20 L. Huang, F. Wu and J. Li, J. Chem. Phys., 2016, 144, 114708.
- 21 C. Xia, J. Du, W. Xiong, Y. Jia, Z. Wei and J. Li, *J. Mater. Chem. A*, 2017, 5, 13400.
- 22 A. Kudo and Y. Miseki, Chem. Soc. Rev., 2009, 38, 253.
- 23 Y. Hu, Y. Guo, Y. Wang, Z. Chen, X. Sun, J. Feng, T. M. Lu, E. Wertz and J. Shi, *J. Mater. Res.*, 2017, 325, 1.
- 24 J. Zhang, T. Song, Z. Zhang, K. Ding, F. Huang and B. Sun, J. Mater. Chem. C, 2015, 3, 4402.
- 25 R. Ai, X. Guan, J. Li, K. Yao, P. Chen, Z. Zhang, X. Duan and X. Duan, ACS Nano, 2017, 11, 3413.
- 26 I. Karbovnyk, I. Bolesta, I. Rovetskyi, V. Lesivtsiv, Y. Shmygelsky, S. Velgosh and A. Popov, *Low Temp. Phys.*, 2016, 42, 594.
- 27 J. Robertson, J. Phys. C: Solid State Phys., 1979, 12, 4753.
- 28 M. Zhong, S. Zhang, L. Huang, J. You, Z. Wei, X. Liu and J. Li, *Nanoscale*, 2017, **9**, 3736.

- 29 M. Zhong, L. Huang, H.-X. Deng, X. Wang, B. Li, Z. Wei and J. Li, J. Mater. Chem. C, 2016, 4, 6492.
- 30 P. Tyagi, A. G. Vedeshwar and N. C. Mehra, *Phys. B*, 2001, 304, 166.
- 31 I. A. Kariper, J. Mater. Res. Technol., 2016, 5, 77.
- 32 G. Kresse and J. Furthmüller, Phys. Rev. B, 1996, 54, 11169.
- 33 G. Kresse and J. Furthmüller, Comput. Mater. Sci., 1996, 6, 15.
- 34 P. E. Blöchl, Phys. Rev. B, 1994, 50, 17953.
- 35 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, 77, 3865.
- 36 J. Heyd, G. E. Scuseria and M. Ernzerhof, J. Chem. Phys., 2003, 118, 8207.
- 37 J. Heyd, J. E. Peralta, G. E. Scuseria and R. L. Martin, J. Chem. Phys., 2005, 123, 174101.
- 38 J. Heyd, G. E. Scuseria and M. Ernzerhof, *J. Chem. Phys.*, 2006, **124**, 9906.
- 39 S. Grimme, J. Comput. Chem., 2006, 27, 1787.
- 40 H. J. Monkhorst and J. D. Pack, Phys. Rev. B, 1976, 13, 5188.
- 41 İ. A. Kariper, J. Mater. Res. Technol., 2016, 5, 77.
- 42 Y. Wang, Y.-Y. Sun, S. Zhang, T.-M. Lu and J. Shi, *Appl. Phys. Lett.*, 2016, **108**, 013105.
- 43 J. Wang, Z. Guan, J. Huang, Q. Li and J. Yang, *J. Mater. Chem. A*, 2014, 2, 7960.
- 44 Q. Li, N. Zhang, Y. Yang, G. Wang and D. H. Ng, *Langmuir*, 2014, 30, 8965.
- 45 W. Jin, P.-C. Yeh, N. Zaki, D. Zhang, J. T. Sadowski, A. Al-Mahboob, A. M. van Der Zande, D. A. Chenet, J. I. Dadap, I. P. Herman, et al., Phys. Rev. Lett., 2013, 111, 106801.
- 46 R. Coehoorn, C. Haas, J. Dijkstra, C. Flipse, R. De Groot and A. Wold, *Phys. Rev. B*, 1987, 35, 6195.