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Introduction

Up-conversion is an anti-Stokes process in which a near infrared photon is converted into visible or even ultraviolet light owing to multi-photon action. Among various near infrared excitation lasers, the 980 nm one is the most commonly used excitation source generating up-conversion emission.¹ Actually there are many up-conversion processes with widely different conversion efficiencies. The processes require energy levels which are resonant with the incoming or outgoing radiation. Lanthani $de(m)$ ions (Ln³⁺) are very suitable as the emitting centers owing to their numerous electronic levels and narrow emission bands. Among the Ln(m) ions Ho³⁺ can be applied to obtain red and green up-conversion luminescence and $Tm³⁺$ acts as a blue

Up-conversion white emission and other luminescence properties of a YAG:Yb₂O₃ · Tm₂O₃ · Ho₂O₃@SiO₂ glassnanocomposite†

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We report on a glass-nanocomposite material consisting of yttrium aluminum garnet (Y₃Al₅O₁₂, YAG) nanocrystals co-doped with Yb³⁺, Tm³⁺ and Ho³⁺ ions as well as entrapped into a SiO₂ xerogel. This $94YAG·5Yb_2O_3·0.8Tm_2O_3·0.2Ho_2O_3@SiO_2$ (abbr. YAG:YbTmHo@SiO₂) nanocomposite material has been prepared by sol–gel procedure. Its structure and morphology has been characterized by means of X-ray diffraction (XRD) and scanning electron microscope (SEM) techniques as well as energy dispersive X-ray (EDX), X-ray photoelectron (XPS) and luminescence spectroscopies. The luminescent glassnanocomposite exhibited an up-conversion effect under $\lambda_{\text{exc}} = 980$ nm and emission when excited under 355 nm in steady-state conditions. Then time-resolved luminescence emission was observed, when the sample was excited at 290 and 355 nm by a pulse laser. Average decay times for the $SiO₂$ matrix and for some transitions of the Tm^{3+} and Ho³⁺ dopants present in the YAG:YbTmHo@SiO₂ material have been evaluated. The luminescent nanocomposite when excited under 290 or 355 nm wavelengths in both conditions emits blue light. However, the nanocomposite is promising as a singlesource white-light phosphor owing to its up-conversion luminescence under 980 nm excitation. Such optical features make the studied material an alternative phosphor. **PAPER**
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emitting activator under infrared (980 nm) excitation. While, Yb^{3+} can act as the sensitizer enhancing the emission efficiency due to appropriate structure of energy levels and longer decay lifetime of the excited states.²

The efficiency of the system depends strongly on the choice of the host matrix. It is known that oxides are less suitable than fluorides, since lifetimes in the former are shorter than in the latter due to stronger interaction between the luminescent ion and its coordination environment. Since in the up-conversion oxide material co-doped, e.g. with Yb^{3+} and Er^{3+} the lifetime of the intermediary ${}^{4}I_{11/2}$ $(Er³⁺)$ is shorter, in consequence the total efficiency of the up-conversion process will be lower, in contrast to fluoride material.³ Irrespective of this fact there are many papers devoted to up-conversion in oxide materials, among them also works considering yttrium aluminum garnet (e.g. ref. 2 and 4–11).

Yttrium aluminum garnet (YAG) is a prominent crystalline material of the garnet group. The choice of YAG relies on this material exhibiting exceptional optical transparency, having a cubic crystallographic structure that can well accommodate such dopants as lanthanide ions.^{12,13} YAG nanocrystals doped with Ln^{3+} ions were prepared using many different methods.^{7,14,15}

A nanocomposite is defined as a material consisting of at least two different solid phases. In our case glassy silica as the

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host phase of larger extent plays a matrix role. Then the nanosized YAG guest phase doped with $Ln(m)$ ions is dispersed into the matrix. Obviously, the constituent phases of the composite are separated by distinct interface. In the nanocomposites, the guest nanoparticles have dimension in the range 1–100 nm. In this type of nanocomposite materials glass is an excellent host owing to its several inherent superior properties compared to those of other encapsulating hosts. There are many types of glass-nanocomposites depending on the compositions, sizes and shapes of the nano-guests as well as processing parameters.¹⁶

One of the encapsulating methods used in nanocomposite preparation is sol–gel processing. This method involves the generation of colloidal suspensions (sols) which are subsequently converted to viscous gels and then to rigid xerogels.¹⁷ Thus, by sol–gel procedure crystalline particles can be immobilized in a solid xerogel. Wide range of inorganic and hybrid organic–inorganic composite materials share this common preparation procedure. In contrast to bulk materials prepared by the melting method, sol–gel materials can be obtained as thin films and coatings. If these materials are transparent, then they gain a competitive advantage with respect to bulk glasses and crystals. It is important especially in glass-nanocomposite luminescent materials doped with such luminescent species as Ln^{3+} ions, where the nanosize of the Ln^{3+} doped crystals is crucial to avoid light scattering. The nanocrystals play important role, when activated by lanthanide ions. The presence of the crystalline environment around the lanthanide ion allows high absorption and emission cross sections as well as tailoring of the ion–ion interaction by the control of the lanthanide ion partition. This last point is crucial and still an object of intense experimental and theoretical studies. Paper

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The integration of luminescent nanocrystals into vitreous matrix can lead to new optical nanocomposites. There is growing interest to integrate up-conversion nanocrystals with an optical matrix. Not only glasses but also rigid vitreous xerogels can serve as indispensable optical matrix materials owing to their high optical transparency, thermomechanical strength and the ability (especially in the case of xerogels) to be shaped into an almost unlimited range of geometrical structures.¹⁸ However, it remains a preparation challenge to integrate up-conversion nanoparticles with tailored nanophotonic properties into a vitreous matrix to create a new hybrid optical material. Wet chemistry is the satisfied synthetic technique in the case of up-conversion nanocrystals because it enables control over crystallite phase, size, shape, composition and nanostructure.¹⁹⁻²¹ This control gives a clear advantage over in situ glass ceramic procedure by synthesizing nanocrystals and glass separately, and then integrate them into a hybrid material. This preparation method can overcome the limitations of the glass-ceramic processing. Among the limitations is an extremely high risk of completely dissolving the nanosized crystals at higher temperatures.

Up-conversion nanocrystals containing thousands of $Ln³⁺$ ions emit high brightness up-conversion by utilizing highirradiance excitation to enrich the effective emitters $19,20,22$ or clustering Yb^{3+} sensitizers in arrays at sublattice level to promote localized excited states.²³

In our study glass-nanocomposites consisting of the $Ln³⁺$ codoped YAG nanocrystals incorporated into glassy silica matrix by sol–gel procedure were synthesized. We demonstrate that this approach allows incorporation of a controlled amount of YAG:Ln up-conversion nanocrystals with well-defined concentrations of Ln^{3+} ions into amorphous xerogel. The vitreous materials doped with up-conversion nanocrystals exhibit high transparency and identical glassy xerogel structure as a blank silica xerogel. Aim of the study is to compare luminescence of the glass nanocomposite excited in different conditions, viz. in steady-state ($\lambda_{\rm exc}$ = 355 and 980 nm) or by pulse laser ($\lambda_{\rm exc}$ = 290 and 355 nm) spectroscopy.

Experimental

Synthetic procedures

Pure $Y_3AI_5O_{12}$ (YAG) and co-doped $94Y_3AI_5O_{12}$ 6Ln₂O₃ $(in \, mol\%)$ materials (abbr. YAG:Ln, where $Ln = Yb$, Tm and Ho), were synthesized by glycol modification of the sol-gel procedure after Veith et $al.^{24}$ The pure YAG as well as doubly co-doped $94YAG \cdot 5Yb_2O_3 \cdot 1Ho_2O_3$ (abbr. YAG:YbHo) and triply co-doped $94YAG \cdot 5Yb_2O_3 \cdot 0.8Tm_2O_3 \cdot 0.2Ho_2O_3$ (abbr. YAG:YbTmHo) agglomerated crystals were powdered in an agate mortar and then pulverized in methanol slurry to nanosize particles by an ultrasonic technique for 3 h.²⁵ The chemicals for the doped YAG synthesis such as yttrium trioxide (Y_2O_3) , aluminum nitrate $\left[\text{Al}(\text{NO}_3)_3 \cdot 5\text{H}_2\text{O} \right]$ and appropriate lanthanide salts $\text{Ln}(\text{NO}_3)_3 \cdot 5\text{H}_2\text{O}$ (where $Ln = Yb$, Tm and Ho) of p.a. grade were purchased from Aldrich Co. While acetic acid (CH₃COOH), methanol (CH₃OH) and 1,2-ethanediol (glycol, $HOCH_2CH_2OH$) of p.a. grade were from Polish Chemical Reagents.

The co-doped YAG:YbHo and YAG:YbTmHo nanopowders were immobilized into silica xerogel by sol–gel processing in a one-pot method.¹⁷ 2.5 cm³ of tetramethoxysilane (TMOS was) dissolved in 5 cm^3 of methanol and 1 cm^3 of water was added. The mixture was stirred and 1.2 cm^3 of NH₃(aq) catalyst dropwise was added. After that the MeOH sol consisting of the co-doped YAG nanoparticles together with the sol–gel mixture was stirred for 15 min. The molar ratio of TMOS : H_2O : MeOH was established as $1:4:8$. After two days, the co-doped YAG nanocrystals YAG:YbHo@SiO₂ and YAG:YbTmHo@SiO₂ incorporated into rigid wet gels with the materials were obtained. The nanocomposites were dried at 200 $^{\circ}$ C, then crushed and sieved $(0.25-0.50 \text{ mm})$.²⁵ After that they were calcined at 200, 600, 800 and 1000 $^{\circ}$ C for 3 hours. The chemicals for the silica matrix preparation such as tetramethoxysilane (TMOS) was purchased from Aldrich Co. but methanol $(CH₃OH)$ and ammonia NH₃(aq) were from Polish Chemical Reagents.

Apparatus

The ultrasonic equipment for the crystals crushing to nanoparticles was Sonorex Super RK 103H. After straining and drying at 120 \degree C the nanocrystals were thermally treated at 200 \degree C in a programmable oven equipped with an SM-946 temperature controller.

Powder X-ray diffraction (XRD) analysis of the YAG, doped and co-doped YAG:Ln nanocrystals as well as YAG:Ln@SiO₂ nanocomposites was made using a Bruker diffractometer of the D2Phaser model employing CuKa radiation operated at 30 kV and 0.2 mA. The XRD patterns were collected using scanning steps of 0.02 deg. and counting time 0.4 s per step. The average crystalline size as well as percentage of crystallinity were calculated using Diff. EVA software provided by Bruker.

The topography of the samples was characterized using scanning electron microscope SEM S-3400N (Hitachi, Japan), equipped with a tungsten source. A secondary electron (SE) detector was used and analysis was performed under 20 kV accelerating voltage.

Chemical composition of the investigated materials was examined by means of energy dispersive X-ray spectroscopy (EDX) equipped with ThermoFisher Scientific detector attached to (SEM) and X-ray photoelectron spectroscopy (XPS) fitted with ThermoFisher Scientific: Escalab 250Xi utilizing monochromatic Al-Ka source with charge neutralization implemented by means of flood gun. High-resolution XPS spectra were recorded at energy step size of 0.1 eV at pass energy 10 eV. In order to normalize the spectroscopic measurements, the X axis (binding energy) from XPS spectrum was calibrated for the peak characteristics for neutral carbon C1s (284.6 eV) .²⁶ Data analysis was performed using Avantage software provided by the manufacturer. **Exchauses**
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Photoluminescence spectra were recorded using a monochromators SPM2 produced by Carl Zeiss Jena for excitation and emission. An excitation source in this system was Osram 250 W ozone-free Xenon lamp and laser produced by Optoelectronics Tech. Co., Ltd MDL-III-980L-1000 mW. Fluorescence intensity was measured using a Hamamatsu R928 Side-on photomultiplier. All spectra were recorded at 300 K.

To follow the luminescence kinetics, we used a system consisting of a YAG:Nd (PL 2143 A/SS) laser and system OPG (optical paramagnetic generators). This system can generate 30 ps laser pulses, with the frequency of 10 Hz with wavelengths ranging 290 and 355 nm. The emission signal was analyzed by a Bruker Optics 2501S spectrometer and the Hamamatsu Streak Camera model C4334-01 with a final spectral resolution 0.47 nm. Luminescence spectra were collected by integration of the streak camera pictures over time intervals, whereas luminescence decays were obtained by the integration of streak camera images over the wavelength intervals.²⁷

Results and discussion

XRD patterns

XRD patterns of the pure YAG (curve a), YAG doped with 6 mol% Yb^{3+} (YAG:Yb, curve b) and YAG triply co-doped with Yb^{3+} , Tm³⁺ and Ho $^{3+}$ (YAG:YbTmHo, curve c) dried at 200 $^{\circ}$ C are presented in Fig. S1.† All the patterns are consistent with the standard data corresponding to cubic $Y_3Al_5O_{12}$ nanocrystals (PDF 01-082-3500), i.e. no secondary crystalline phase is observed and presence of the dopants did not change the crystalline structure. The calculated by the Scherrer formula average size of the nanocrystals was equal to 27 ± 4 nm.

XRD patterns of the thermally treated YAG:YbTmHo@SiO₂ nanocomposite is shown in Fig. S2.† The sample was dried at 200 $^{\circ}$ C and then calcined at 600, 800 and 1000 $^{\circ}$ C. As previously, the XRD patterns are in accordance with the YAG standard data. The average crystallite sizes of the triply co-doped YAG nanoparticles in the nanocomposite collected in Table S1† are within experimental error close to the crystallite size of the un-doped, doubly and triply co-doped YAG nanoparticles. The crystallinity percentage of the samples were estimated. Its values change a little with temperature of the thermal treatment (see Table S1†).

SEM images

SEM micrographs of the pure YAG and doped YAG:Yb powder samples presented in Fig. 3S(a and b,† respectively) exhibit agglomerated crystals. Then in images of greater enlargement in insets are shown nanoparticles which average size could be estimated to ca. 30 nm. It is within experimental error close to the size calculated by the Scherrer formula (vide supra). On the other hand, SEM image of the YAG:YbTmHo@SiO₂ nanocomposite powder thermally treated at 600 $^{\circ}$ C shows agglomerates (Fig. S4†). Image of greater enlargement in inset shows the sphere-like nanoparticles consisted of the doped YAG nanocrystals capped with amorphous $SiO₂$. Size of the nanoparticles is near two-fold greater than size of the YAG and YAG:Yb nanocrystals.

EDS and XPS measurements

EDS measurement confirm the composition of the YAG and Ybdoped YAG nanocrystals (Fig. S5,† inset). Whereas highresolution XPS spectra of the YAG and YAG:Yb demonstrated in Fig. S6† allow to calculate the Y/Al/O ratio for both the samples. The results give a close match to $Y_3Al_5O_{12}$ (3.0/6.4/11.6) and 3.0/6.1/11.9 for YAG and YAG:Yb samples, respectively) proving their high purity (see Table S2†).

Luminescence spectroscopy

Steady-state emission spectra. The normalized upconversion emission spectra recorded at room temperature under near infrared (NIR) excitation at 980 nm for the co-doped YAG nanocrystals and the glass-nanocomposites with silica are compared in Fig. 1(A and B). In general, the emission spectra of the co-doped YAG:YbHo and YAG:YbTmHo nanocrystals (curves a) did not change after their incorporation into silica xerogel (compare with curves b, *i.e.* the YAG:YbHo@SiO₂ and YAG:YbTmHo@SiO₂ nanocomposite spectra, respectively).

The up-conversion emissions shown in Fig. 1A are generated by two- and three-photon processes such as energy transfer (ET) and excited state absorption (ESA) illustrated in the known schematic energy level diagrams presented in Fig. 2 for $Ho³⁺$ ions. In the first step Yb^{3+} sensitizer is excited under 980 nm from the ground state ${}^{2}F_{7/2}$ to the ${}^{2}F_{5/2}$ one. Then, occurs the Yb³⁺:²F_{5/2} \rightarrow $Ho^{3+}:^5I_6$ transfer. In this situation either the non-radiative transition Ho^{3+} ⁵ $\text{I}_6 \rightarrow {}^{5}\text{I}_7$ takes place and subsequently by step (2) the Ho³⁺ is excited to the ⁵F₅ state or directly occurs the Ho³⁺:⁵I₆ \rightarrow ⁵F₅ 5S₅ transition by two-photon process (step 2) from Vb³⁺₅ In F_4 , ⁵S₂ transition by two-photon process (step 2) from Yb³⁺. In

Fig. 1 Up-conversion emission spectra of the samples: (A) YAG:YbHo nanocrystals (curve a) and YAG:YbHo@SiO₂ glass nanocomposite (curve b) as well as (B) YAG:YbTmHo nanocrystals (curve a) and YAG:YbTmHo@SiO₂ glass nanocomposite (curve b). The samples dried at 200 °C were excited under $\lambda_{\text{exc}} = 980$ nm by the laser in steadystate condition.

Fig. 2 Scheme of energy levels of Yb^{3+} (sensitizer) as well as Tm^{3+} and $Ho³⁺$ (activators) in YAG:YbTmHo@SiO₂ nanocomposite. The scheme demonstrates up-conversion mechanisms of the emissions under 980 nm laser excitation. GSA means ground state absorption, ESA – excited state absorption, ET – energy transfer and NRT – non-radiative transition. After ref. 22 and 28.

former case the red emission about 670 nm $(14\,870\,$ $\mathrm{cm}^{-1})$ corresponds to the ${}^{5}F_{5} \rightarrow {}^{5}I_{8}$ transition. Whereas in the latter case can be present two transitions, $viz. \, {}^5F_4$, ${}^5S_2 \rightarrow {}^5I_8$ and 5F_1 ${}^5S_2 \rightarrow {}^5I_1$ Owing these transitions are observed such emis- F_4 , ${}^5S_2 \rightarrow {}^5I_7$. Owing these transitions are observed such emission bands as green one at about 550 nm $(18\ 180\ \text{cm}^{-1})$ and near-infra red (NIR) band centered at 764 nm (13 004 $\rm cm^{-1}$), respectively. However, after the ${}^{5}F_{5} \rightarrow {}^{5}I_{5}$ non-radiative decay three-photon process (step 3) takes place the ${}^{5}F_{5} \rightarrow {}^{5}G_{4}$, ${}^{3}K_{7}$ transition. The ⁵G₄, ³K₇ \rightarrow ⁵I₆ relaxation generates an emission band peaked at *ca.* 650 nm (15 380 cm⁻¹).^{4,22}

Emission spectra shown in Fig. 1B are superpositions of spectra consisting of lines assigned to transitions of Ho^{3+} and Tm^{3+} dopants in YAG:YbTmHo and its glass nanocomposite counterpart YAG:YbTmHo@SiO₂. The observed emission bands of Tm^{3+} are due to two-, three- and four-photon up-conversion processes. As previously, the up-conversion mechanism starts when the excited Yb^{3+} ion transfers its energy to the Tm³⁺ and the latter ion reaches the excited ${}^{3}\text{H}_5$ level. As in the case of Ho $^{3+}$ it is a phonon assisted energy transfer because there is energy difference between the excited state of Yb^{3+} and the excited ${}^{3}H_{5}$ level of Tm^{3+} (see left part of the energy diagram in Fig. 7). Then after the non-radiative relaxation Tm^{3+} ; $^{3}\text{H}_{5} \rightarrow {}^{3}\text{F}_{4}$ takes place ET from the Yb³⁺ (step 2) that causes the ³F₄ \rightarrow ³F_{2,3} transition of Tm^{3+} . From this energy level the Tm^{3+} ions can firstly nonradiatively relax to ${}^{3}H_4$ and then radiatively relax to the ground state $(^{3}H_{6})$. The latter transition generates the NIR emission band centered at 826 nm $(11 990 \text{ cm}^{-1})$ in the spectra shown in Fig. 1B. Step (3) allows to excite Tm^{3+} to the $^1\text{G}_4$ state. From this level start three emission transitions of Tm^{3+} , viz. ${}^{1}G_{4}$ \rightarrow ${}^{3}H_{5}$, ${}^{1}G_{4}$ \rightarrow ${}^{3}F_{4}$ and ${}^{1}G_{4}$ \rightarrow ${}^{3}H_{6}$ which correspond to bands at 810 nm (12 $\rm 250\,cm^{-1}$), 637 nm (15 $\rm 635\,cm^{-1}$) and 484 nm $(20 640 \text{ cm}^{-1})$, respectively. After the next ESA (step 4) the Tm³⁺ ions are transferred to the ${}^{1}D_{2}$ state from which they can relax by photon emissions of wavelengths of 792 nm (12540 cm^{-1}) and 470 nm $(21 \ 220 \text{ cm}^{-1})$. In this case the emission lines are related to ${}^{1}D_{2} \rightarrow {}^{3}F_{2,3}$ and ${}^{1}D_{2} \rightarrow {}^{3}F_{4}$ transitions, respectively.²² Paper

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The up-conversion mechanism in the triply co-doped nanocrystals and nanocomposite could be more complicated because such pairs of energy levels as $(Ho^{3+1.5}I_7$ and $Tm^{3+1.3}F_4$), $\rm Ho^{3+}:^5I_6$ and $\rm Tm^{3+}:^3H_5$, $\rm Ho^{3+}:^5F_5$ and $\rm Tm^{3+}:^3F_{2,3}$ as well as $\text{Ho}^{3+}: \mathbb{F}_1$ and $\text{Tom}^{3+}: \mathbb{G}_4$ possessing small energy differences between them (compare in Fig. 2). Thus, a quasi-resonant ET $\text{Ho}^{3+} \rightarrow \text{Tm}^{3+}, e.g. \text{Ho}^{3+}$:⁵F₅ $\rightarrow \text{Tm}^{3+}$:³F_{2,3} and next multiphonon non-radiative transitions (NRTs) to the Ho^{3+} ⁵I₇ and Tm³⁺:³H₄: levels could be taken into account (see right part of the diagram in Fig. 2). In this case population of the $\text{Tm}^{3+1}\text{G}_4$ state can be reached in the three-photon processes as follows: $\text{Yb}^{3+}:^2\text{F}_{7/2} \rightarrow \text{GSA Yb}^{3+}:^2\text{F}_{5/2} \rightarrow \text{ESA Ho}^{3+}:^5\text{I}_6 \rightarrow \text{NRT Ho}^{3+}:^5\text{I}_7 \rightarrow$ ESA $\text{Ho}^{3+}:^5\text{F}_5 \rightarrow \text{ET} \text{ Tm}^{3+}:^3\text{F}_{2,3} \rightarrow \text{NRT} \text{ Tm}^{3+}:^3\text{H}_4 \rightarrow \text{ESA}$ $\text{Tm}^{3\text{+}}:\text{G}_4$. The other observed bands are assigned to the same Ho^{3+} electronic transitions as in the doubly co-doped samples YAG:YbHo and YAG:YbHo@SiO₂ (Fig. 1A).²⁸

The material consisting of the triply co-doped YAG nanocrystals (YAG:YbTmHo) embedded into silica (YAG:YbTmHo@SiO₂) was excited under 355 nm, *i.e.* the excitation wavelength typical for Tm^{3+} , Ho³⁺ ions and additionally for silica matrix (see Fig. 3). The appropriate emission spectra of the YAG:YbTmHo@SiO $_2$ phosphor dried at 200 $^{\circ}$ C and then calcined at 600, 800 and 1000 $^{\circ}$ C are demonstrated in Fig. 3. Characteristic feature of the spectra is a band wing penetrated deeply from UV into the visible range. This wing is the long wavelength part of the spectrum (Fig. 3, inset) attributed to

Fig. 3 Steady-state emission spectra of the YAG:YbTmHo@SiO₂ nanocomposite excited under 355 nm after drying at (a) 200 $^{\circ} \mathsf{C}$ and calcined at: (b) 600, (c) 800 and (d) 1000 °C for 3 h. Inset presents emission spectrum of SiO $_2$ xerogel dried at 200 $^\circ$ C and excited under 355 nm.

electronic transitions between structural defect levels of amorphous SiO_2 ^{29,30} Emission lines of Tm³⁺ and Ho³⁺ are observed against the wing background. The distinct Ho^{3+} bands attributed to $^{5}S_{2}$, $^{5}F_{4} \rightarrow \ ^{5}I_{8}$ (peaked at 541 nm, 18 484 cm $^{-1}$) and $^5{\rm G}_4 \rightarrow {^5{\rm I}_6}$ (at 585 nm, 17 094 cm $^{-1})$ transitions are present only in the spectra of the material thermally treated in the 200–800 $^{\circ}\mathrm{C}$ range (curves a–c). While the Ho^{3+} bands centered at 466 nm $(21\ 459\ {\rm cm}^{-1})$ and 771 nm $(12\ 970\ {\rm cm}^{-1})$ ascribed to $^3{\rm K}_8\rightarrow\ ^5{\rm I}_8$ and ${}^5S_2, {}^5F_4 \rightarrow {}^5I_7$ transitions, respectively, are present also after calcination at 1000 °C (curve d). The Tm³⁺ bands corresponding to ${}^{1}D_{2} \rightarrow {}^{3}F_{4}$ (466 nm, 21 459 cm⁻¹), ${}^{1}G_{4} \rightarrow {}^{3}H_{6}$ (491, 20 367),
¹G₁ \rightarrow ³E₁ (608, 16 447) and ³H₁ \rightarrow ³H₁ (832, 12 019) transitions $G_4 \rightarrow {^{3}F_4}$ (608, 16 447) and ${^{3}H_4} \rightarrow {^{3}H_6}$ (832, 12 019) transitions are also present in all the spectra (a–d). One can assume that the faint red emission related to the ${}^{5}S_{2}$, ${}^{5}F_{4} \rightarrow {}^{5}I_{7}$ transition of Ho³⁺ is present by additional population for the $(^5S_2,~^5F_4)$ levels by energy transfer from Tm^{3+} to Ho^{3+} ions.³¹ However, the Ho^{3+} band quenching in range between 530 and 600 nm signifies that the $Tm^{3+} \rightarrow Ho^{3+}$ energy transfer in the material thermally treated at 1000 °C did not exist in this case.^{32,33} The set of the emission spectra in Fig. 3 suggests that 600 $^{\circ}\mathrm{C}$ as the lowest calcination temperature applied for 3 h allows to obtain optimal luminescence properties of the nanocomposite material. **PSC Advances**
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Thus, the luminescence excitation spectrum of the YAG:YbTmHo@SiO $_2$ nanocomposite calcined at 600 $^{\circ}$ C for 3 h is demonstrated in Fig. 4. The spectrum monitored at 466 nm (corresponding to Tm³⁺:¹D₂ \rightarrow ³F₄ and Ho³⁺:³K₈ \rightarrow ⁵I₈ transitions as well as simultaneously transition in the amorphous silica) consists of lines centered at: 284 nm $(35\ 211\ {\rm cm}^{-1},$ corresponding to transitions $\text{Ho}^{3+,5}\text{I}_8 \rightarrow {}^3\text{G}_3, {}^3\text{L}_8$ and $\text{Tom}^{3+,3}\text{H}_6 \rightarrow {}^3\text{P}_0$), 290 (34 483, $\text{Ho}^{3+}:^5\text{I}_8 \rightarrow {}^{3}\text{D}_3$), 319 (31 348, $\text{Ho}^{3+}:^5\text{I}_8 \rightarrow {}^{5}\text{G}_2$), 343 (29 155, $\rm{Ho}^{3+}{:}^{5}I_8 \rightarrow {^{5}G_3}$), 355 (28 169, $\rm{Ho}^{3+}{:}^{5}I_8 \rightarrow {^{3}H_6}$ and $\text{Trm}^{3^{\text{+}}}:H_6 \rightarrow {}^{1}\text{D}_2$), 384 (26 042, $\text{Ho}^{3^{\text{+}}}:^5\text{I}_8 \rightarrow {}^{5}\text{G}_4$), 399 (25 063, $\text{Ho}^{3+}:^5\text{I}_8 \rightarrow {^3\text{K}_7}$, 415 (24 096, $\text{Ho}^{3+}:^5\text{I}_8 \rightarrow {^5\text{G}_5}$), 425 (23 529, $\rm Ho^{3+}\!\!:\!\!^5I_8\rightarrow \!^5G_6)$ and 435 (22 989, $\rm Ho^{3+}\!\!:\!\!^5I_8\rightarrow \!^5F_1$). In the inset is presented excitation spectrum of glassy $SiO₂$. This spectrum is attributed to many transitions between defects states in the amorphous silica matrix.

Fig. 4 Excitation spectrum of the YAG:YbTmHo QSiO_2 nanocomposite monitored at 466 nm. The nanocomposite was calcined at 600 \degree C for 3 h. The inset demonstrates excitation spectrum of pure $SiO₂$ glass monitored at 440 nm. The arrows indicate the excitation wavelengths (290 and 355 nm) used in steady-state condition.

Time-resolved luminescence

To obtain time-resolved spectra, the YAG:YbTmHo@SiO₂ nanocomposite was excited using of a YAG:Nd laser and system OPG (see Apparatus). The apparatus allows to integrate luminescence in the 0-200 µs time interval. The spectra are collected in Fig. 5. Obviously, they are superpositions consisting of Tm^{3+} and Ho^{3+} bands. When the sample is excited under 290 nm (see curve a), among the lines there are four most intensive bands, viz. the resultant emission band peaked at 371 nm (26 963 cm⁻¹) attributed to Tm³⁺:¹D₂ \rightarrow ³H₆ and Ho³⁺:³H₆ \rightarrow ⁵I₈ transitions. The next band placed at 466 nm $(21\ 459\ cm^{-1})$ is a result of overlapping three bands related to both transitions of Tm³⁺:¹D₂ \rightarrow ³F₄ and ¹G₄ \rightarrow ³H₆ as well as the Ho³⁺:³K₈ \rightarrow ⁵I₈ transition. Whereas lines centered at 652 $(15\ 337\ cm^{-1})$ and 736 nm (13 592 cm⁻¹) are owing to the ¹G₄ \rightarrow ³F₄ transition of Tm³⁺ and the ⁵S₂, ⁵F₄ \rightarrow ⁵I₇ transitions of Ho³⁺, respectively.

Fig. 5 Time-resolved emission spectra of the YAG:YbTmHo@SiO₂ calcined at 600 \degree C for 3 h. The material was excited by the pulse laser under: (a) 290 and (b) 355 nm and the luminescence was integrated for the time interval $0-200 \text{ }\mu\text{s}$

Similar bands or their traces are present in the emission spectrum excited under 355 nm (see Fig. 5, curve b). However, mutual intensity ratios of the bands are quite other and two lines centered at 393 (25 434) and 407 (24 593) as well as a shoulder at *ca.* 615 nm (16 345 cm^{-1}) are absent in comparison to the spectrum excited under 290 nm (curve a). Moreover, the spectrum (curve b) differs from the emission spectrum excited under the same wavelength but in steady-state condition (see Fig. 3, curve b).

Photoluminescence lifetime measurements for YAG:YbTmHo@SiO₂ were performed by the pulsed laser under excitation at 290 nm (see decay curves in Fig. 6). The average lifetime values for electronic transitions related to the four most intense lines mentioned above (see Fig. 5) are presented in Table 1. The average decay times τ_{av} were evaluated using the equation:³⁴

$$
\tau_{\rm av} = \frac{\int tI(t)dt}{\int I(t)dt} \tag{1}
$$

where $I(t)$ is the emission intensity at time t.

Calculated average decay times τ_{av} of the above mentioned transitions are listed in Table 1 for $\lambda_{\text{exc}} = 290$ and 355 nm. One notes that independently on the excitation wavelength the decay times demonstrate approximate values for the observed transitions in both lanthanide dopants of YAG nanoparticles incorporated into silica matrix.

In much shorter time interval (0–50 ns) and under 290 nm excitation was observed for the YAG:Yb,TmHo@SiO₂ only a time-resolved spectrum of the vitreous silica matrix (Fig. 7). The spectrum consists of a strong and broad emission band centered at *ca.* 440 nm (22 738 cm^{-1}). Its origin is attributed to photon recombinations from numerous defects (as dangling bonds) in the amorphous silica.^{29,30}

The measured decay profile of the silica luminescence under excitation by the pulse laser at 290 nm is shown in Fig. 8. The curve allows to estimate average emission decay time value for silica equal to 5.6 ± 0.6 ns.

Fig. 6 Decay luminescence curves for electronic transitions related to four emission lines of Tm^{3+} and Ho^{3+} dopants in YAG nanocrystals incorporated into SiO $_2$ matrix (YAG:YbTmHo@SiO $_2$) calcined at 600 $^{\circ}$ C for 3 h and excited under 290 nm. The lines are presented in the spectra in Fig. 3 (curve b). For details vide supra.

Table 1 Selected electronic transitions of the co-doping Tm^{3+} and Ho^{3+} ions in YAG nanocrystals entrapped in silica matrix $(YAG:YbTmHo@SiO₂)$ and their average decay times

		Average decay time $\tau_{av}/\mu s$ Excitation wavelength $\lambda_{\rm exc}/\rm{nm}$	
Tm^{3+} Tm^{3+} Ho^{3+} Ho^{3+}	${}^{1}D_{2} \rightarrow {}^{3}H_{6}$ ${}^{1}G_{4} \rightarrow {}^{3}H_{6}$ ${}^5G_4 \rightarrow {}^5I_5$ ${}^5S_2, {}^5F_4 \rightarrow {}^5I_7$	41 ± 3 36 ± 3 33 ± 2 30 ± 2	30 ± 2 28 ± 2 36 ± 3 41 ± 3

Fig. 7 Time-resolved emission spectrum of the YAG:YbTmHo@SiO₂ nano-composite calcined at 600 $^{\circ}$ C for 3 h and excited at 290 nm by the pulse laser. The luminescence was integrated in the time interval 0–50 ns.

Chromaticity

The emission spectra of the YAG:YbTmHo@SiO₂ glassnanocomposite under different excitation wavelengths and luminescence techniques (such as the steady state and pulse laser measurements) are demonstrated in Fig. 9. Under the 355 nm excitation (curve a) in the steady-state condition, the nanocomposite material emits simultaneously reddish, green

Fig. 8 Decay of the luminescence excited under 290 nm and integrated in the band 450-470 nm for SiO₂ matrix in YAG:YbTmHo@SiO₂ material calcined at 600 °C for 3 h. Average lifetime $\tau_{\text{av}} = 5.6 \pm 0.6$ ns.

Fig. 9 Comparison of the steady-state emission spectra excited under: (a) 355 nm and (b) 980 nm (up-conversion) as well as excited by the pulse laser under: 355 (c) and 290 nm (d) for the <code>YAG:YbTmHo@SiO</code>2 nanocomposite calcined at 600 $^{\circ}$ C for 3 h. The luminescence in the cases (c) and (d) was integrated in the time interval $0 - 200$ us.

and blue (RGB) luminescent light (see Fig. 3 curve b).³⁵ The emissions are ascribed to the transitions of Tm^{3+} and Ho^{3+} ions, while blue emission is especially intense due to electronic transition in silica (see Fig. 3, inset). The up-conversion spectrum (curve b) obtained under 980 nm excitation by the laser in steady-state condition supply with red, green and bluish green light owing to the transitions of Tm^{3+} and Ho^{3+} ions (compare Fig. 1B). On the other hand, under excitations 355 and 290 nm by the pulse laser (curves c and d, respectively) also the both lanthanide ions are responsible for the emission in visible range. However, the 355 nm excitations delivers only red and blue emission (Fig. 5, curve b), while under 290 nm excitation is observed light consisting of emissions of primary colours (RGB) (Fig. 5, curve a).

The CIE chromaticity coordinates and diagram of YAG:YbTmHo@SiO₂ under the different excitation wavelengths and luminescent techniques are presented in Table 2 and Fig. 10, respectively. The luminescent nanocomposite emits cold white light with bluish tone (point A), when excited by 355 nm in steady state-condition. In this case the correlated colour temperature (T_{cc}) is rather high (ca. 8100 K).³⁶ Whereas the up-conversion emission (point B) excited at 980 nm in the same condition is white and its $T_{cc} \approx 4800$ K is in between temperatures of mean noon sunlight and household tungsten-

Table 2 The CIE coordinates of the YAG:YbTmHo α SiO₂ glassnanocomposite calcined at 600 $^{\circ}$ C for 3 h

Sample	Х		T_{cc} K	λ_{exc} nm
A	0.300	0.307	8100	355
B	0.359	0.351	4800	980
C	0.194	0.191	$-$ ^a	355^b
D	0.260	0.282	14 000	290^b

^{*a*} Impossible to evaluate due to high T_{cc} . b Excited by the pulse laser.

Fig. 10 The CIE chromaticity diagram of the YAG:YbTmHo@SiO₂ nanocomposite calcined at 600 °C for 3 h. For details see Table 2.

filament light bulb (100 W). $35,36$ Both the emissions excited under 290 and 355 nm by pulse laser emitted blue light which $T_{\rm cc}$ values in the former case (point D) is equal to 14 000 K but in the latter case (point C) T_{cc} is too high to evaluate them.

Conclusions

The glass nanocomposite, i.e. nanocrystalline YAG co-doped with Yb^{3+} , Tm^{3+} and Ho^{3+} entrapped in glassy silica $(YAG:YbTmHo@SiO₂)$, was successively prepared by sol-gel method. Presence of the specific lanthanide ions such as Yb^{3+} , Tm^{3+} and Ho³⁺ in the nanocomposite allows to observe upconversion effect, when the composite sample was excited under 980 nm in steady-state condition. Other steady-state emission was also recorded, when $\lambda_{\text{exc}} = 355$ nm. The nanocomposite material excited by pulse laser shows dissimilar time-resolved luminescence spectra for different excitation wavelengths (λ_{exc} = 290 and 355 nm). Luminescence property based on up-conversion effect of the glass-nanocomposite excited under 980 nm in steady-state condition allows to recognize the material as a white phosphor candidate.

Conflicts of interest

There are no conflicts to declare.

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References

- 1 J. Xu, M. Sun, Y. Kuang, H. Bi, B. Liu, D. Yang, R. Lv, S. Gai, F. He and P. Yang, Dalton Trans., 2017, 46, 1495–1501.
- 2 J. Jin, K. Yang, J. Su and Z. Si, J. Lumin., 2015, 159, 178–182.
- 3 G. Blasse and B. C. Grabmaier, Luminescent Materials, Springer-Verlag, Berlin, 1994, ch. 10.
- 4 C. S. Lim, V. V. Atuchin, A. S. Aleksandrovsky, M. S. Molokeev and A. S. Oreshonkov, J. Alloys Compd., 2017, 695, 737–746.
- 5 H. W. Zheng, S. J. Liu, G. S. Yin, W. C. Wang, C. L. Diao, Y. Z. Gu and W. F. Zhang, J. Sol-Gel Sci. Technol., 2011, 59, 290–296.
- 6 X. Mateos, M. C. Pujol, F. Güell, R. Sole, J. Gavalda, J. Massons, M. Aguilo and F. Diaz, Opt. Mater., 2004, 27, 475–479.
- 7 M. Nakielska, J. Sarnecki, M. Malinowski and R. Piramidowicz, J. Alloys Compd., 2008, 451, 190–193.
- 8 J. Zhang and J. Jia, J. Lumin., 2016, 174, 1–5.
- 9 F. G. Yang, F. P. Yan, Z. Y. You, C. Y. Tu, C. L. Sun, Y. Wang, Z. J. Zhu and J. F. Li, Laser Phys. Lett., 2010, 7, 867–869.
- 10 L. Marciniak, A. Bednarkiewicz and W. Stręk, J. Lumin., 2017, 184, 179–184.
- 11 S. Hu, J. Yang, C. Li and J. Lin, Mater. Chem. Phys., 2012, 133, 751–756.
- 12 A. Ikesue, T. Kinoshita, K. Kamata and K. Yoshida, J. Am. Chem. Soc., 1995, 78, 1033–1040.
- 13 S. Pfeifer, M. Bischoff, R. Niewa, B. Clauss and M. R. Buchmeister, J. Eur. Ceram. Soc., 2014, 34, 1321–1328.
- 14 X.-X. Ge, Y.-H. Sun, C. Liu and W.-K. Qi, J. Sol-Gel Sci. Technol., 2009, 52, 179–187.
- 15 R. Muenchausen, L. G. Jacobsohn, B. L. Bennett, E. A. McKigney, J. F. Smith, J. A. Valdez and D. W. Cooke, J. Lumin., 2007, 126, 838–842.
- 16 B. Karmakar, Fundamentals of Glass and Glass Nanocomposites in Glass Nanocomposites, Synthesis, Properties and Applications, ed. B. Karmakar, K. Rademann, A. L. Stepanov, Elsevier, Amsterdam, 2016.
- 17 C. J. Brinker and G. W. Scherer, Sol-Gel Science: The Physics and Chemistry of Sol-Gel Processing, Academic Press, Boston, 1990.
- 18 J. Zhao, X. Zheng, E. P. Schartner, P. Ionescu, R. Zhang, T.-L. Nguyen, D. Jin and H. Ebendorff-Heidepriem, Adv. Opt. Mater., 2016, 4, 1507–1517.
- 19 J. Zhao, D. Jin, E. P. Schartner, Y. Lu, Y. Liu, A. V. Zvyagin, L. Zhang, J. M. Dawes, P. Xi, J. A. Piper, E. M. Goldys and T. M. Monro, Nat. Nanotechnol., 2013, 8, 729–734.
- 20 D. J. Gargas, E. M. Chan, A. D. Ostrowski, S. Aloni, M. V. P. Altoe, E. S. Barnard, B. Sanii, J. J. Urban, D. J. Milliron, B. E. Cohen and P. J. Schuck, Nat. Nanotechnol., 2014, 9, 300–305. Paper Vescott, Computer 1991. Zhang, D. J. Samer, J. G. Samer, J. C. Samer, J. C. H. Samer, J. Samer, J. D. H. Samer, J. D. H. Samer, J. D. H. Samer, J. D. H.
	- 21 D. Liu, X. Xu, Y. Du, X. Qin, Y. Zhang, C. Ma, S. Wen, W. Ren, E. M. Goldys, J. Pieper, S. Dou, X. Liu and D. Jin, Nat. Commun., 2016, 7, 10432–10438.
	- 22 J. Zhou, G. Chen, Y. Zhu, L. Huo, W. Mao, D. Zou, X. Sun, E. Wu, H. Zeng and J. Zhang, J. Mater. Chem. C, 2015, 3, 364–369.
	- 23 J. Wang, R. Deng, M. A. MacDonald, B. Chen, J. Yuan, F. Wang, D. Chi, T. S. Andy Hor, P. Zhang, G. Liu, Y. Han and X. Liu, Nat. Mater., 2014, 13, 157–162.
	- 24 M. Veith, S. Mathur, A. Kareiva, M. Jilavi, M. Zimmer and V. Huch, J. Mater. Chem., 1999, 9, 3069–3079.
	- 25 M. Kubus, H.-J. Meyer, L. Kienle and A. M. Kłonkowski, J. Non-Cryst. Solids, 2009, 355, 1333–1337.
	- 26 D. Pawlak, K. Wozniak, Z. Frukacz, T. L. Barr, D. Fiorentino and S. Seal, J. Phys. Chem. B, 1999, 103, 1454–1461.
	- 27 A. Kubicki, P. Bojarski, M. Grinberg, M. Sadownik and B. Kukliński, Opt. Commun., 2006, 263, 275-280.
	- 28 J. Zmojda, D. Dorosz, M. Kochanowicz, P. Miluski and ˙ J. Dorosz, Acta Phys. Pol., A, 2013, 124, 598–601.
	- 29 L. N. Skuja, J. Non-Cryst. Solids, 1992, 149, 77–95.
	- 30 A. M. Kłonkowski, W. Wiczk, J. Ryl, K. Szczodrowski and D. Wilenska, J. Alloys Compd., 2017, 724, 649-658.
	- 31 A. C. Yanes, J. J. Velazquez, J. del-Castillo, J. Mendez-Ramos and V. D. Rodriguez, J. Sol-Gel Sci. Technol., 2009, 51, 4–9.
	- 32 D. Wang, Y. Guo, G. Sun, J. Li, L. Zhao and G. Xu, J. Alloys Compd., 2008, 451, 122–124.
	- 33 D. Chen, Y. Wang, Y. Yu, P. Huang and F. Weng, J. Solid State Chem., 2008, 181, 2763–2769.
	- 34 S. Shinoya and W. M. Yen, Phosphor Handbook, CRC Press, Boca Raton, 1999.
	- 35 R. J. D. Tilley, Colour and the optical properties of materials, J. Wiley, Chichester, 2 edn, 2011, ch 3.
	- 36 http://www.vendian.org/mncharity/dir3/blackbody/Unstable URLs/bbr_color.html.