

Cite this: *RSC Adv.*, 2018, 8, 10848Received 28th January 2018
Accepted 13th March 2018

DOI: 10.1039/c8ra00861b

rsc.li/rsc-advances

Metallic VO₂ monolayer as an anode material for Li, Na, K, Mg or Ca ion storage: a first-principle study

Yusheng Wang,^{id}*^{ac} Nahong Song,^{bc} Xiaoyan Song,^a Tianjie Zhang,^a Qiaoli Zhang^a and Meng Li^d

Using density functional theory (DFT), we assess the suitability of monolayer VO₂ as promising electrode materials for Li, Na, K, Mg and Ca ion batteries. The metallic VO₂ monolayer can offer an intrinsic advantage for the transportation of electrons in materials. The results suggest that VO₂ can provide excellent mobility with lower diffusion barriers of 0.043 eV for K, 0.119 eV for Li, 0.098 eV for Na, 0.517 eV for Mg, and 0.306 eV for Ca. The specific capacities of Li, Na and Mg can reach up to 968, 613 and 815 mA h g⁻¹ respectively, which are significantly larger than the corresponding value of graphite. Herein, with high open-circuit voltage the VO₂ sheet could be a promising candidate for the anode material in battery applications.

1. Introduction

The increase of the energy demands and the concentration of carbon dioxide in the atmosphere signify the search for sustainable and renewable energy sources. Li-ion battery (LIB) is an attractive power source and has been widely applied in portable electronic devices and electric vehicles.^{1–7} However, several issues including cost, safety, and capacity restrict the further development in large scale energy storage applications of LIB.^{8,9} Therefore, the development of new types of batteries is urgently needed. Batteries based on alternative carrier ions, such as Na,^{10–13} K,^{14,15} Mg^{16,17} or Ca,¹⁷ may be more suitable for large-scale energy storage systems. It is well known that the performance of ion batteries relies greatly on the properties of electrode materials. Therefore, great effort is being expended on finding novel and efficient electrode materials for next generation of renewable energy technologies. So far, two-dimensional (2D) materials, such as graphene,¹⁸ porous graphene,¹⁹ silicene,²⁰ transition-metal dichalcogenides (TMDS),^{21–24} MXenes,²⁵ black phosphorus,^{26,27} and borophene,^{6,28–30} have been widely studied because of their particular structural, physical, and chemical properties, which may enable the high capacity and the fast metal ion diffusivity.

Recently, a new 2D material, namely vanadium dioxide (VO₂) has attracted great attention owing to its semiconductor to

metal transition around 340 K.^{31–34} The existence of monolayer VO₂ and its metallic properties were already theoretically predicted.³⁵ Beyond these novel properties, how VO₂ behave as battery electrodes is still unknown. More importantly, the monolayer VO₂ shows metallic characteristic with light weight and large surface area, suggesting its potential applications for metal ion storage. It is well known that the high capacity, electric conduction and excellent ion mobility are all requirements for an anode material. This work therefore aims to explore the applicability of monolayer VO₂ for Li, Na, K, Mg and Ca ions storage. From our calculation results, we show that the monolayer VO₂ can be taken as a promising anode material for storing Li, Na and Mg.

2. Computational method

First-principles density functional theory (DFT) calculations in the present study were carried out using Vienna *ab initio* simulation package (VASP)^{36–38} with a generalized gradient approximation (GGA) in the Perdew–Burke–Ernzerhof (PBE).³⁹ Since the van der Waals (vdW) interaction plays a crucial role in the adsorption of adatoms on the substrate, a non-local vdW density functional in the form of Becke88 optimization (optB88)⁴⁰ is adopted, which has been evidenced to give a much improved description for various systems including heterostructures,^{41–43} lithium carbides⁴⁴ and the molecules/2D materials systems.^{45,46} A kinetic energy cutoff of the plane-wave basis set was used to be 500 eV. A *k*-point sampling of 5 × 5 × 1 Monkhorst–Pack grids⁴⁷ in the first Brillouin zone of the 4 × 4 supercell is used in the calculation. An inter-layer vacuum space larger than 20 Å was kept to avoid mirror interactions. The convergence criteria of total energy and Hellmann–Feynman force were set as 10^{−4} eV and 0.01 eV Å^{−1}, respectively. The

^aSchool of Mathematics and Statistics, North China University of Water Resources and Electric Power, Zhengzhou, Henan, 450046, China. E-mail: wangyusheng@ncwu.edu.cn

^bCollege of Computer and Information Engineering, Henan University of Economics and Law, Zhengzhou, Henan, 450000, China

^cInternational Joint Research Laboratory for Quantum Functional Materials of Henan, School of Physics and Engineering, Zhengzhou University, Zhengzhou 450001, China

^dCollege of Science, Zhongyuan University of Technology, Zhengzhou 450007, China



optimized lattice constant of unit cell is 2.75 Å, the bond length of V–O is 1.95 Å, which are consistent with previous results.³⁵ For the evaluation of charge transfer between the ions and the VO₂, we performed Bader charge analysis.⁴⁸ To obtain diffusion pathways and the corresponding energy barriers, we performed standard nudged elastic band (NEB) calculations.⁴⁹

3. Results and discussion

3.1 Structures and electronic properties of metal adatoms adsorption on VO₂ monolayer

The geometric structure of monolayer VO₂ sheet is depicted in Fig. 1. Clearly, from the top view, the VO₂ sheet exhibits typical graphene-like hexagonal structure, and from the side view, the VO₂ sheet is a triple-layer structure, which can be viewed as the V-atomic layer being sandwiched by two O-atomic layers through strong V–O bonds. The unit cell of monolayer VO₂ sheet consists of two O atoms and one V atom. The optimized lattice constant and the bond length of V–O are 2.753 Å and 2.050 Å respectively, which agree with the previous results.⁵⁰ As depicted in Fig. 1(b), VO₂ exhibits metallic behavior, which can ensure a satisfied electrical conductivity during battery cycling.

We first examine the adsorption behavior of a single metal atom (Li, Na, K, Mg and Ca) on the monolayer VO₂. The adsorption energy is defined as the difference between the sum of the energies of the pristine VO₂ sheet and the metal atom and the total energy of the metal-adsorbed system:

$$E_a = E_M + E_{VO_2} - E_{M/VO_2} \quad (1)$$

where E_M , E_{VO_2} , and E_{M/VO_2} are the energies of the metal adatoms, VO₂ sheet and the metal atom adsorption on VO₂ sheet. By this definition, the larger E_a , the stronger the interaction between the metal adatom and the VO₂ sheet. It is well known that for a material to be suitable as an anode for a particular ion storage, a relatively large adsorption energy plays a fundamental role. The VO₂ with a 4 × 4 supercell structure is employed to avoid the interaction between the adjacent metal adatoms. We first determined the most favorable adsorption site for a single metal atom on VO₂ sheet. Herein, four different initial adsorption sites are considered: above the hollow center of hexagonal V–O ring (H), above the top of O atom (T1), above the

top of V atom (T2) and above the V–O bond (B) as shown in Fig. 1(a). The results suggest that all the above metal atoms prefer to be adsorbed above the H site. The corresponding adsorption energies are 3.85, 3.89, 4.42, 4.31, and 6.21 eV for Li, Na, K, Mg and Ca respectively. These large adsorption energies ensure that the metal atoms can bind strongly to the VO₂ sheet.

To obtain insight into the adsorption process, we investigated the charge transfers by performing Bader charge analysis.⁴⁸ The three-dimensional charge density difference of metal/VO₂ are illustrated in Fig. 2, which were obtained from the formula

$$\Delta\rho = \rho(\text{metal}/\text{VO}_2) - [\rho(\text{VO}_2) + \rho(\text{metal})]$$

where $\rho(\text{metal}/\text{VO}_2)$ and $\rho(\text{VO}_2)$ denote the total electron densities of the relaxed VO₂ sheet with and without metal atoms, respectively, and $\rho(\text{metal})$ is the total electron density of the metal atoms. The red and blue colors represent charge accumulation and depletion. It is found that all the metal atoms behave as charge donors. They donate 0.91e, 0.90e, 0.94e, 1.70e, and 1.64e to the VO₂ sheet for Li, Na, K, Mg and Ca respectively. To further understand the interaction between the metal atom and the VO₂ sheet, we have calculated the partial density of states (PDOS) of different systems and illustrated them in Fig. 3. Unlike many other two-dimensional materials that exhibit semi-conducting or semi-metallic characteristics, the VO₂ sheet as well as all the metal/VO₂ systems are all metallic states. Therefore, the metallic nature of the VO₂ sheet remains unchanged during the metal adsorbed on it, and an excellent electronic conductivity is maintained, which will significantly improve the performance of metal ion batteries. The metal atoms donate their electrons to the O atoms of the VO₂ so that the s and p orbitals of the Li, Na, K and Mg (s and d orbitals for Ca) partially appear above the Fermi level. Meanwhile, the orbital hybridization between the metal atoms and the O atoms are observed at the lower energy region, which contributes to the strong interaction between the metal atoms and the VO₂.

3.2 Metal atom diffusion on VO₂ monolayer

Another crucial characteristic to evaluate the suitability of an anode material for rechargeable batteries is mobility of metal

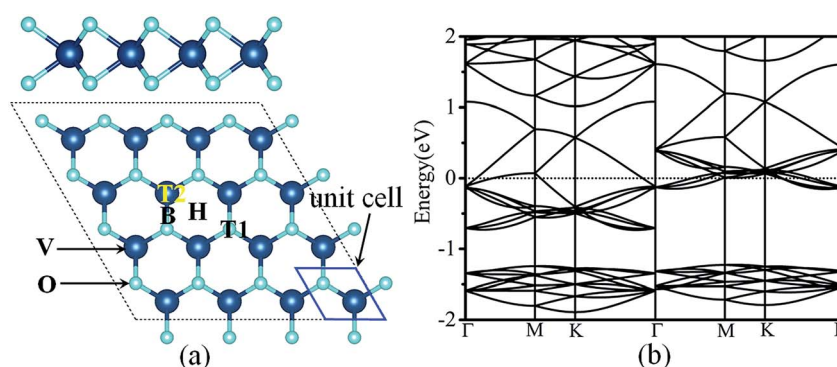


Fig. 1 (a) The optimized atomic structures and (b) the band structure of monolayer VO₂.



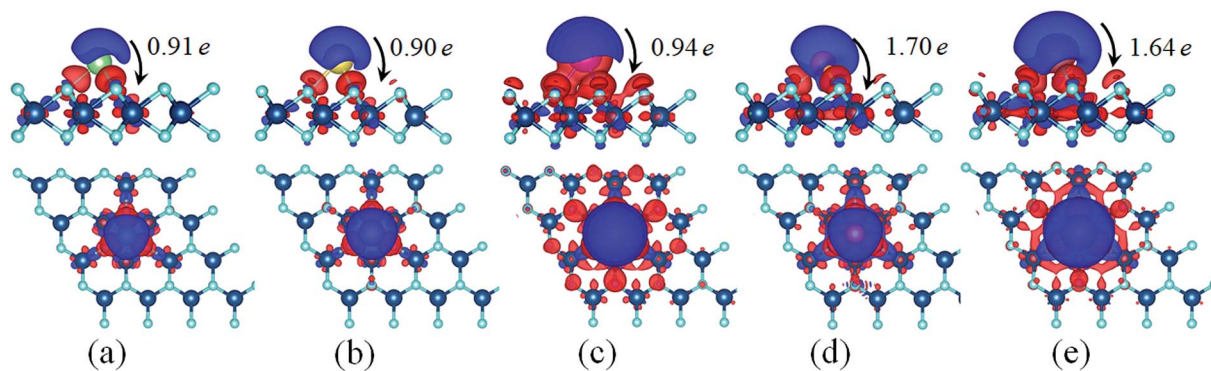


Fig. 2 Charge density difference plots for (a) Li/VO₂, (b) Na/VO₂, (c) K/VO₂, (d) Mg/VO₂, (e) Ca/VO₂. The red and blue iso-surface with an iso-value of 0.0015 e Å⁻³ corresponds to charge accumulation and depletion respectively.

atoms on the electrode material. From this point of view, the diffusion of metal atoms on the surface of VO₂ must be quantified. In this regard we investigate the diffusion barriers of metal atoms on the VO₂ sheet. As illustrated in Fig. 4(a), H-T-H pathway is the most favorable diffusion pathway for all metal atoms studied in this work. The corresponding diffusion energy profiles along the optimized pathway for different metal atoms are shown in Fig. 4(b). It is found that the diffusion barriers of the alkali atoms are lower than those of alkaline earth atoms. Among the three alkali atoms, K possesses the lowest diffusion barrier, with a value of 0.043 eV. The diffusion barriers of the Li and Na are 0.119 eV and 0.098 eV, respectively. The obtained low barrier for alkaline atoms on VO₂ is far superior to typical electrode materials, such as VS₂ (0.22 eV),²⁴ borophene (0.60 eV),⁵¹

porous graphene (0.37 eV),¹⁹ MoN₂ (0.49 eV),¹⁴ and so on.^{52–54} The low diffusion barriers of alkaline atoms on VO₂ implies the fast diffusion processes, especially for K ion. On the other hand, we find that the obtained diffusion barriers of Mg and Ca are 0.517 eV and 0.306 eV respectively, which are comparable to the energy barriers for Li diffusion on other anode materials.^{51,53,55} Therefore, we can conclude that the VO₂ monolayer will be easy for the above metal atoms to diffuse. In addition, the ZPE contributions are also considered. As a prototype model system, Li diffusion on VO₂ is tested. The diffusion barrier of Li on VO₂ with ZPE is 0.149 eV, which is slightly larger than that of the value 0.119 eV calculated without ZPE. Although the NEB method without ZPE underestimates the diffusion barrier, it is still comparable to that calculated with ZPE.

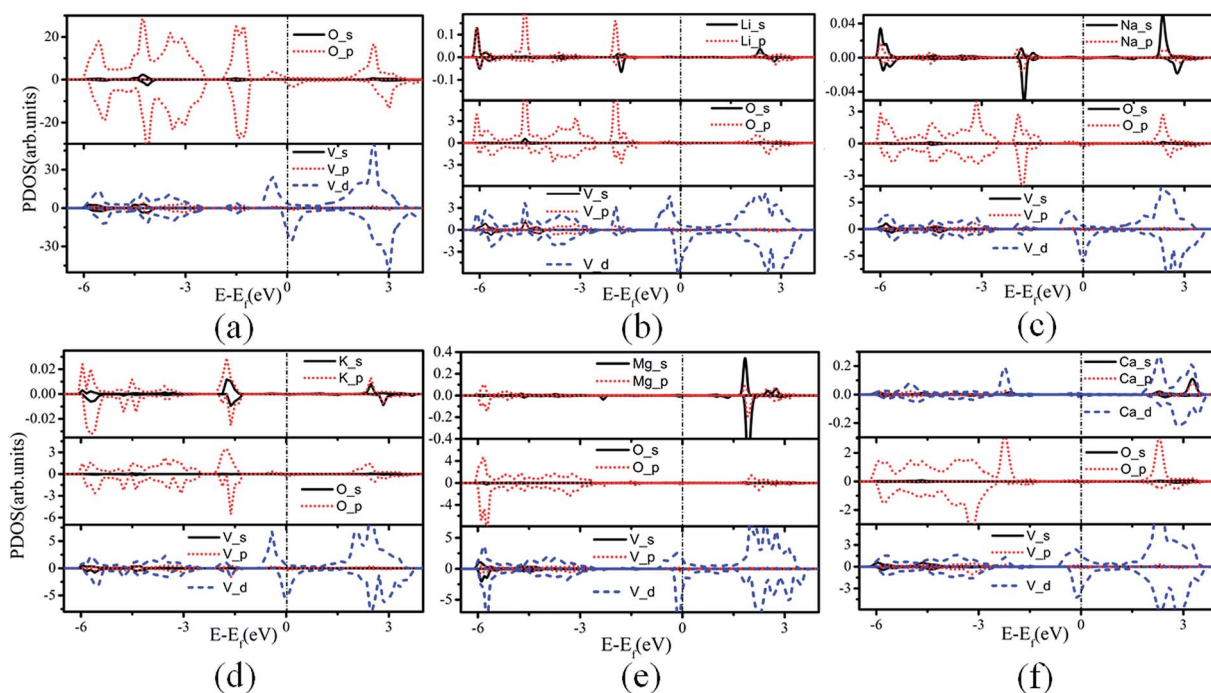


Fig. 3 PDOS for (a) VO₂, (b) Li/VO₂, (c) Na/VO₂, (d) K/VO₂, (e) Mg/VO₂, (f) Ca/VO₂. The Fermi level is set as zero.



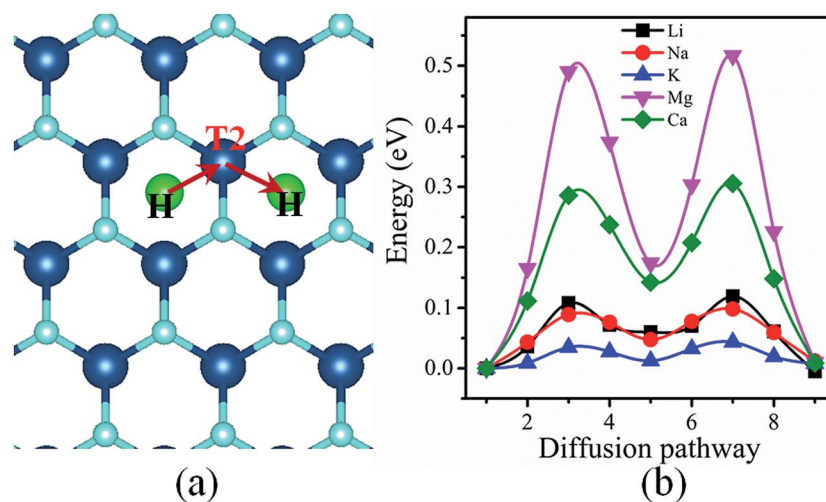


Fig. 4 (a) Schematic of the diffusion pathways of a metal atom along H–T2–H. (b) Energy profile corresponding to the minimum-energy pathway.

It should be noted that the charge states on metal atoms are positive 0.91, 0.90, 0.94, 1.70 and 1.64 for Li, Na, K, Mg and Ca respectively, which suggests that all the metal atoms exist in the cationic state. In the light of this idea, the diffusion of charged metal atoms on VO₂ sheet could be much easier with the external electric field during the processes. Thus, the ultrafast charge/discharge rate is expected for the VO₂ sheet as the metal-ion battery anode material.

3.3 Theoretical storage capacity of Li, Na, K, Mg and Ca on the VO₂ monolayer and open-circuit voltage

Next, we increased the coverage of adatoms over the VO₂ sheet. To this aim, we constructed a series of configurations with the chemical stoichiometry of M_xVO₂ ($x = 0.0625, 0.25, 0.5, 0.75, 1, 2, 4$) by placing different number of metal atoms on the single and both sides of the VO₂ sheet until the complete coverage was achieved. As a prototype model system, Fig. 5 displays the optimized structures for Li_xVO₂ with x varying from 0.0625 to 4. In the case of M₂VO₂, the VO₂ fully covered by metal atoms leads to x being 2 as illustrated in Fig. 5(f). In the case of K_{0.75}VO₂ and Ca_{0.75}VO₂, the average adsorption energies were estimated to be negative, this indicates that when the number of adsorbed of K and Ca atoms is larger than 0.75 in a unit cell, the systems are unstable. To test the feasibility of multilayer adsorption, we also studied the VO₂ with bilayer metal atoms adsorbing on its both sides. For adsorption of the second layer, it should be noted that once the adsorption sites for the first layer are fully occupied, subsequent metal atoms of the next layer will occupy T2 sites. In the case of Mg₄VO₂, the average adsorption energy is calculated to be negative. This indicates that the second layer of Mg atoms can not be adsorbed on the VO₂ with one layer of Mg atoms fully covering on both sides, because the extra atoms occupy a second layer and it is difficult for these atoms to be adsorbed on VO₂ sheet. Interestingly, bilayer Li and Na adsorptions on both sides of VO₂ sheet are feasible, corresponding to formulae of M₄VO₂, as shown in Fig. 5(g).

Open-circuit voltage data is another crucial parameter to characterize the performance of metal ion batteries. In theory, the open-circuit voltage for the metal atom coverages between $x_1 \leq x \leq x_2$ can be obtained by using the following equation:^{24,56–58}

$$V \approx -(E_{Mx_2VO_2} - E_{Mx_1VO_2} - (x_1 - x_2)E_M) / ((x_1 - x_2))$$

in which, $E_{Mx_2VO_2}$, $E_{Mx_1VO_2}$, and E_M are the energies of M_{x₂}VO₂, M_{x₁}VO₂ and metal atom, respectively. Fig. 6 shows the voltage profiles of the VO₂ sheet as a function of the metal atom coverage ratio. It is seen that the voltage profiles of Li_xVO₂ and Na_xVO₂ are all positive which additionally confirms the suitability of VO₂ as anode material. In the case of Li_xVO₂, there is a dramatic decrease from approximately 3.85 V when $x < 0.25$ to 2.08 V ($x < 0.75$), and then decreases again to 1.16 V. Next, the voltage values are fluctuating around 1.16 V ($0.5 < x < 2$). At last, the voltage value reduces to 0.31 V for the bilayer Li atoms adsorbed on both sides of the VO₂. The voltage profile of Na_xVO₂ is similar to that of Li_xVO₂. In the case of K_xVO₂ (Ca_xVO₂), when the K(Ca) atom coverage ratio is larger than 0.75 V, the open-circuit voltage becomes to be negative indicating that the VO₂ as K(Ca) ion battery node material is unstable for high coverage ratio of K(Ca). In the case of Mg_xVO₂, when the Mg atom coverage ratio is lower than 2, the Mg_xVO₂ is stable. Meanwhile, we calculated the average open-circuit voltage by numerically averaging the voltage profile, which are 1.56 V for Li ($0 < x < 4$), 0.98 V for Na ($0 < x < 4$), 2.36 V for K ($0 < x < 0.75$), 1.21 V for Mg ($0 < x < 1$), and 2.82 V for Ca ($0 < x < 0.75$). Thus, the VO₂ based metal ion battery can provide high charging voltage, which are proper for the practical anodic applications.^{57–59}

To designate VO₂ monolayer as anode material for metal ion battery, we calculated the specific capacities using the following equation^{54,58,60}

$$C = \frac{nN_A e}{3.6m_{M_xVO_2}}$$



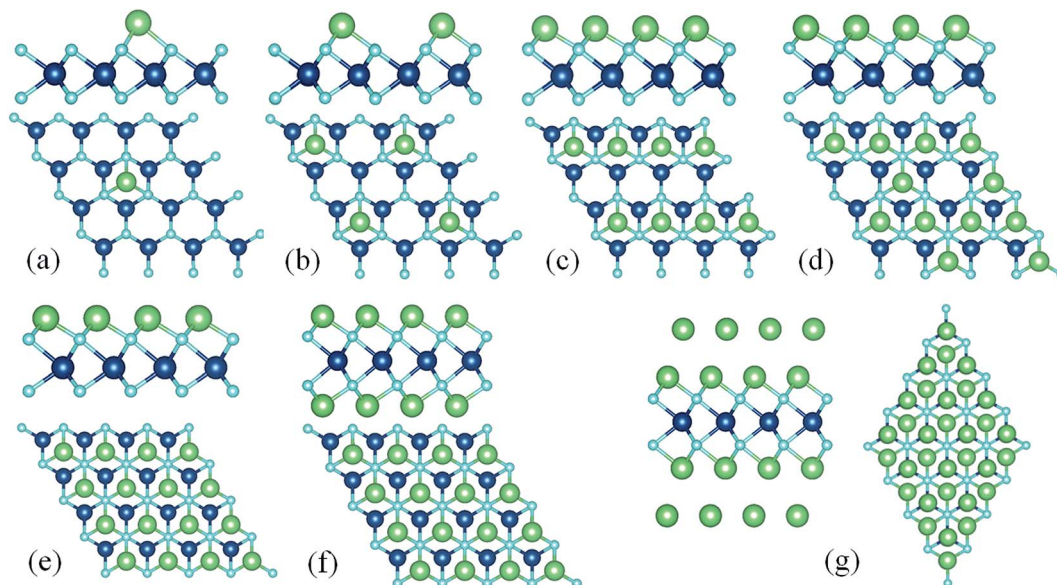


Fig. 5 The optimized atomic structures for (a) $\text{Li}_{0.625}\text{VO}_2$, (b) $\text{Li}_{0.25}\text{VO}_2$, (c) $\text{Li}_{0.5}\text{VO}_2$, (d) $\text{Li}_{0.75}\text{VO}_2$, (e) Li_1VO_2 , (f) Li_2VO_2 , (g) Li_4VO_2 .

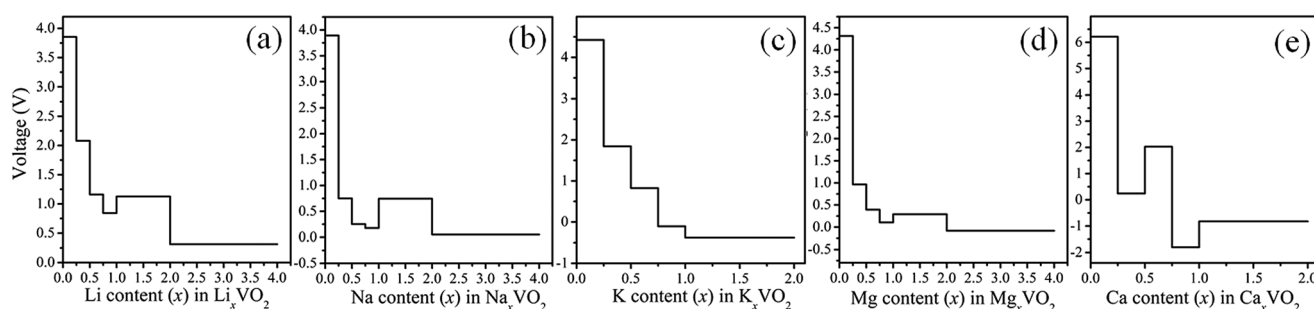


Fig. 6 Predicted open-circuit voltage profiles as a function of metal atoms content x . (a) Li_xVO_2 , (b) Na_xVO_2 , (c) K_xVO_2 , (d) Mg_xVO_2 , (e) Ca_xVO_2 .

in which n is the adsorbed electric charge per mole VO_2 , N_A is the Avogadro constant, e is the elementary charge, and $m_{\text{M}_x\text{VO}_2}$ represents the molar mass of M_xVO_2 . For bilayer Li or Na adsorption on both sides of the VO_2 sheet, each formula unit of VO_2 sheet can accommodate up to 4Li or 4Na atoms. Corresponding theoretical capacities are found to be about 968 and 613 mA h g^{-1} , respectively, which are much larger than that for most of 2D materials. For example, the theoretical capacity of Li on graphene is 372 mA h g^{-1} , Li and Na on Sc_2C are 462 and 362 mA h g^{-1} ,⁵⁸ Li and Na on Mo_2C are 526 and 132 mA h g^{-1} ,⁶⁰ and Li and Na on Mo_2N are 432 and 864 mA h g^{-1} .¹⁴ The maximum Mg content is 1 resulting the theoretical capacity of 815 mA h g^{-1} , which is also higher than conventional electrode capacity. Regarding to K and Ca ions, the predicted capacities are 131 and 260 mA h g^{-1} respectively, which are distinctly lower than that for Li, Na and Mg. On the basis of the above analyses, we can conclude that the VO_2 monolayer should be a good candidate for the application of an electrode material in metal ion batteries, especially for Li, Na and Mg ion batteries. Also, it is equally important to explore the properties of the VO_2 with the substrate it was formed on in the future.

4. Conclusion

In summary, we have demonstrated the application of the monolayer VO_2 as anode material for metal (Li, Na, K, Mg and Ca) ion batteries by DFT-based first-principles calculations. The results shown that the metal adatoms prefer to occupy the center of hexagonal V-O ring and play the role of charge donors. The atoms adsorption on VO_2 exhibit the diffusion barriers as low as 0.043 eV for K, 0.119 eV for Li, 0.098 eV for Na, 0.517 eV for Mg, and 0.306 eV for Ca, suggesting fast charging/discharging processes especially for K, Li and Na. The metallic nature of the VO_2 sheet remains unchanged during the metal adsorbed on it, which is essential for the use as an electrode. In addition, the theoretical specific capacity of VO_2 can be as much as 968, 613 and 815 mA h g^{-1} for Li, Na and Mg respectively, which are significantly larger than that of graphene. The results also shown that the VO_2 based metal ion battery can provide high charging voltage. Given these advantages, we can conclude that the monolayer VO_2 has fascinating potential as the high capacity and fast charging-discharging anode material for metal ion batteries.



Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was supported by the NSF of China (Grant No. 11404112), the Key Scientific Research Project of Colleges and Universities of Henan Province (Grant No. 16A140024) and Research in Cutting-edge Technologies of Zhengzhou (Grant No. 141PRKXF622).

References

- H. Wang, H. Feng and J. Li, *Small*, 2014, **10**, 2165–2181.
- J. M. Tarascon and M. Armand, *Nature*, 2001, **414**, 359–367.
- H. Li, Z. Wang, L. Chen and X. Huang, *Adv. Mater.*, 2009, **21**, 4593–4607.
- M. Liang and L. Zhi, *J. Mater. Chem.*, 2009, **19**, 5871–5878.
- J. Dahn, T. Zheng, Y. Liu and J. Xue, *Science*, 1995, **270**, 590–593.
- H. R. Jiang, Z. Lu, M. C. Wu, F. Ciucci and T. S. Zhao, *Nano Energy*, 2016, **23**, 97–104.
- Q. Zhang, J. Wang, J. Dong, F. Ding, X. Li, B. Zhang, S. Yang and K. Zhang, *Nano Energy*, 2015, **13**, 77–91.
- J.-M. Tarascon, *Nat. Chem.*, 2010, **2**, 510.
- J. B. Goodenough and K.-S. Park, *J. Am. Chem. Soc.*, 2013, **135**, 1167–1176.
- B. Xiao, Y.-c. Li, X.-f. Yu and J.-b. Cheng, *ACS Appl. Mater. Interfaces*, 2016, **8**, 35342–35352.
- L. Shi and T. Zhao, *J. Mater. Chem. A*, 2017, **5**, 3735–3758.
- N. Yabuuchi, K. Kubota, M. Dahbi and S. Komaba, *Chem. Rev.*, 2014, **114**, 11636–11682.
- F. Legrain and S. Manzhos, *J. Power Sources*, 2015, **274**, 65–70.
- X. Zhang, Z. Yu, S.-S. Wang, S. Guan, H. Y. Yang, Y. Yao and S. A. Yang, *J. Mater. Chem. A*, 2016, **4**, 15224–15231.
- A. Eftekhari, *J. Power Sources*, 2004, **126**, 221–228.
- D. Aurbach, Z. Lu, A. Schechter, Y. Gofer, H. Gizbar, R. Turgeman, Y. Cohen, M. Moshkovich and E. Levi, *Nature*, 2000, **407**, 724–727.
- R. Y. Wang, C. D. Wessells, R. A. Huggins and Y. Cui, *Nano Lett.*, 2013, **13**, 5748–5752.
- K. Persson, Y. Hinuma, Y. S. Meng, A. Van der Ven and G. Ceder, *Phys. Rev. B*, 2010, **82**, 125416.
- Y. Wang, Q. Zhang, M. Jia, D. Yang, J. Wang, M. Li, J. Zhang, Q. Sun and Y. Jia, *Appl. Surf. Sci.*, 2016, **363**, 318–322.
- G. A. Tritsarlis, E. Kaxiras, S. Meng and E. Wang, *Nano Lett.*, 2013, **13**, 2258–2263.
- D. B. Putungan, S.-H. Lin and J.-L. Kuo, *ACS Appl. Mater. Interfaces*, 2016, **8**, 18754–18762.
- C. N. R. Rao, U. Maitra and U. V. Waghmare, *Chem. Phys. Lett.*, 2014, **609**, 172–183.
- Y. Jing, E. O. Ortiz-Quiles, C. R. Cabrera, Z. Chen and Z. Zhou, *Electrochim. Acta*, 2014, **147**, 392–400.
- Y. Jing, Z. Zhou, C. R. Cabrera and Z. Chen, *J. Phys. Chem. C*, 2013, **117**, 25409–25413.
- D. Sun, Q. Hu, J. Chen, X. Zhang, L. Wang, Q. Wu and A. Zhou, *ACS Appl. Mater. Interfaces*, 2016, **8**, 74–81.
- Q.-F. Li, C.-G. Duan, X. G. Wan and J.-L. Kuo, *J. Phys. Chem. C*, 2015, **119**, 8662–8670.
- L.-Q. Sun, M.-J. Li, K. Sun, S.-H. Yu, R.-S. Wang and H.-M. Xie, *J. Phys. Chem. C*, 2012, **116**, 14772–14779.
- H. R. Jiang, W. Shyy, M. Liu, L. Wei, M. C. Wu and T. S. Zhao, *J. Mater. Chem. A*, 2017, **5**, 672–679.
- B. Mortazavi, A. Dianat, O. Rahaman, G. Cuniberti and T. Rabczuk, *J. Power Sources*, 2016, **329**, 456–461.
- S. Banerjee, G. Periyasamy and S. K. Pati, *J. Mater. Chem. A*, 2014, **2**, 3856–3864.
- B.-J. Kim, Y. W. Lee, B.-G. Chae, S. J. Yun, S.-Y. Oh, H.-T. Kim and Y.-S. Lim, *Appl. Phys. Lett.*, 2007, **90**, 023515.
- J. Liang, W. Li, J. Liu and M. Hu, *Mater. Lett.*, 2016, **184**, 92–95.
- L. Liu, T. Yao, X. G. Tan, Q. H. Liu, Z. Q. Wang, D. C. Shen, Z. H. Sun, S. Q. Wei and Y. Xie, *Small*, 2012, **8**, 3752–3756.
- S. Long, H. Zhou, S. Bao, Y. Xin, X. Cao and P. Jin, *RSC Adv.*, 2016, **6**, 106435–106442.
- C. Ataca, H. Şahin and S. Ciraci, *J. Phys. Chem. C*, 2012, **116**, 8983–8999.
- G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1993, **47**, 558–561.
- G. Kresse and J. Furthmüller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169.
- G. Kresse and J. Furthmüller, *Comput. Mater. Sci.*, 1996, **6**, 15–50.
- J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865–3868.
- A. D. Becke, *Phys. Rev. A: At., Mol., Opt. Phys.*, 1988, **38**, 3098–3100.
- W. Yu, Z. Zhu, S. Zhang, X. Cai, X. Wang, C.-Y. Niu and W.-B. Zhang, *Appl. Phys. Lett.*, 2016, **109**, 103104.
- L. Zhang, P. Bampoulis, A. N. Rudenko, Q. Yao, A. van Houselt, B. Poelsema, M. I. Katsnelson and H. J. W. Zandvliet, *Phys. Rev. Lett.*, 2016, **116**, 256804.
- J. E. Padilha, A. Fazzio and A. J. R. da Silva, *Phys. Rev. Lett.*, 2015, **114**, 066803.
- Y. Lin, T. A. Strobel and R. E. Cohen, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2015, **92**, 214106.
- Y. Cai, Q. Ke, G. Zhang and Y.-W. Zhang, *J. Phys. Chem. C*, 2015, **119**, 3102–3110.
- T. P. Kaloni, G. Schreckenbach and M. S. Freund, *J. Phys. Chem. C*, 2014, **118**, 23361–23367.
- H. J. Monkhorst and J. D. Pack, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1976, **13**, 5188–5192.
- G. Henkelman, A. Arnaldsson and H. Jonsson, *Comput. Mater. Sci.*, 2006, **36**, 354–360.
- G. Henkelman, B. P. Uberuaga and H. Jónsson, *J. Chem. Phys.*, 2000, **113**, 9901.
- Z.-K. Tang, X.-B. Li, D.-Y. Zhang, Y.-N. Zhang and L.-M. Liu, *J. Mater. Chem. C*, 2015, **3**, 3189–3197.
- X. Zhang, J. Hu, Y. Cheng, H. Y. Yang, Y. Yao and S. A. Yang, *Nanoscale*, 2016, **8**, 15340–15347.
- T. H. Osborn and A. A. Farajian, *J. Phys. Chem. C*, 2012, **116**, 22916–22920.



- 53 C. Sun and D. J. Searles, *J. Phys. Chem. C*, 2012, **116**, 26222–26226.
- 54 C. C. Leong, H. Pan and S. K. Ho, *Phys. Chem. Chem. Phys.*, 2016, **18**, 7527–7534.
- 55 C. Uthaisar and V. Barone, *Nano Lett.*, 2010, **10**, 2838–2842.
- 56 M. K. Aydinol, A. F. Kohan, G. Ceder, K. Cho and J. Joannopoulos, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1997, **56**, 1354–1365.
- 57 B. Mortazavi, M. Shahrokhi, M. Makaremi and T. Rabczuk, *Applied Materials Today*, 2017, **9**, 292–299.
- 58 X. Lv, W. Wei, Q. Sun, L. Yu, B. Huang and Y. Dai, *ChemPhysChem*, 2017, **18**, 1627–1634.
- 59 H. Zheng, K. Jiang, T. Abe and Z. Ogumi, *Carbon*, 2006, **44**, 203–210.
- 60 Q. Sun, Y. Dai, Y. Ma, T. Jing, W. Wei and B. Huang, *J. Phys. Chem. Lett.*, 2016, **7**, 937–943.

