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Enhanced thermoelectric performance of CNT/P3HT composites with low CNT content†

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Carbon nanotubes (CNTs) have emerged as one of the leading additives for improving the thermoelectric properties of organic materials due to their unique structure and excellent electronic transport properties. However, since as-grown CNTs generally possess different diameters, it is of high interest to determine the influence of the diameter of carbon nanotubes on the thermoelectric properties of CNT/poly(3-hexylthiophene) (P3HT) composite films. Herein, we prepared CNT/P3HT composite films with diameters of <8 nm, 8–15 nm, 20–30 nm, 30–50 nm and >50 nm and studied their thermoelectric properties. It was found that the diameter of CNTs had an important influence on the TE performance of the composite films. The P3HT- d_{CNT} (<8 nm) and P3HT- d_{CNT} (8–15 nm) composite films exhibited almost the same thermoelectric performance and almost more than double that of the other three composite films with increased CNT diameter. The different mass fractions of CNT/P3HT composite films have also been investigated. The maximum TE power factor of CNT ($d < 8$ nm)/P3HT composite films reached $49.0 \mu\text{W mK}^{-2}$ at the mass fraction of 95 wt% P3HT, that is, 5 wt% CNTs. This superior TE power factor of CNT ($d < 8$ nm)/P3HT composite films can be ascribed to the fully connected interlayer of the P3HT polymer and also the heterogeneous dispersion of short-length CNTs.

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1. Introduction

Thermoelectric (TE) materials have great potential in applications such as power generation and solid-state heating/cooling, realizing direct conversion between thermal energy and electrical energy without involving moving mechanical components or hazardous working fluids. The performance of TE materials is characterized by a dimensionless figure of merit $ZT = S^2\sigma T/\kappa$, where S , σ , κ and T are the Seebeck coefficient, electrical conductivity, thermal conductivity and absolute temperature, respectively, and $S^2\sigma$ is the power factor. Conventional inorganic semiconductors such as Bi_2Te_3 , PbTe and SiGe have received the most attention as TE materials because of their high power factors.^{1–5} However, the relatively high cost of raw materials and the lack of scalability of production for these conventional TE materials have created a great need for inexpensive and easily processed materials for expanding TE applications.

In recent years, polymer-based composite TE materials have attracted increasing attention in thermoelectric research

because of their low cost, mechanical flexibility and generally low thermal conductivity.^{6–18} Carbon nanotubes (CNTs) are the most commonly used and effective material among numerous fillers. CNTs can provide conductive paths when embedded in polymer matrix because CNTs possessed excellent electrical conductivity and high mobility. For example, Bounioux *et al.* used single-walled carbon nanotubes (SWNTs) as fillers to prepare SWNTs/P3HT composites and obtained a high electrical conductivity of $\sim 103 \text{ S cm}^{-1}$ as the SWNTs formed a three-dimensional conductive network, resulting in a large TE power factor in the in-plane direction of the composites, $\sim 95 \mu\text{W mK}^{-2}$ at room temperature.¹² Meanwhile, quasi one-dimensional structure of CNTs could be used as templates to induce highly ordered polymer molecular chain structures through the strong π - π conjugated interactions between CNTs and conjugated conducting polymers.^{19–22} Du *et al.*¹⁹ fabricated P3HT/CNT films by oxidative polymerization of 3-hexylthiophene in chloroform solution containing dispersed CNT. The P3HT chains were orderly attached onto the surface of the CNTs by the π - π conjugated interactions between CNTs and P3HT, which resulted in a reasonable thermoelectric performance with CNT content of 5 wt%. Yao *et al.*^{23,24} fabricated SWNTs/PANI composite TE materials and found that PANI molecular chains were orderly arranged on the surface of SWNTs by the π - π conjugated interactions between SWNTs and PANI, which facilitated the increase of carrier mobility and thereby improved the electronic transport properties of SWNTs/PANI composites.

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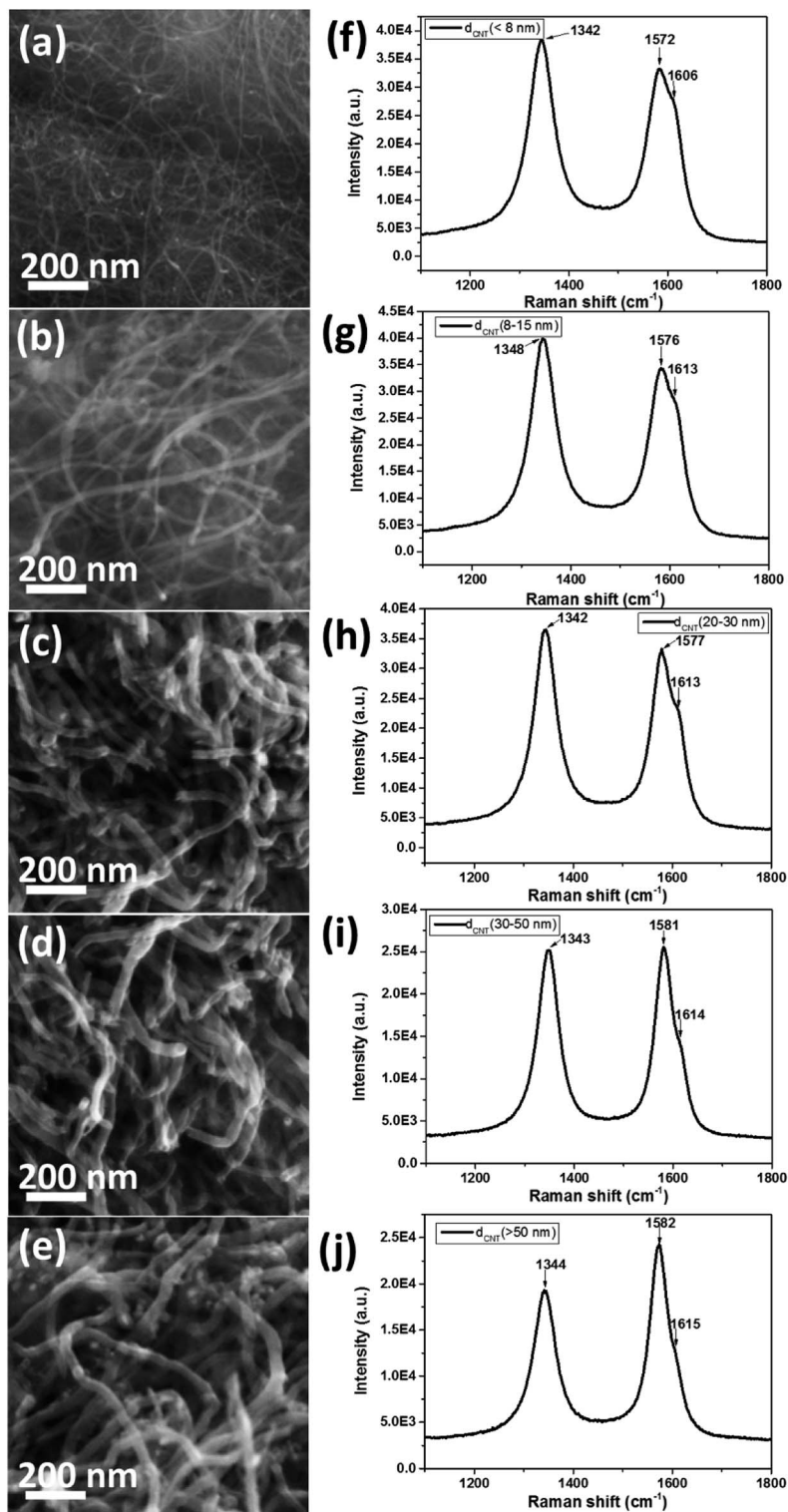


Fig. 1 SEM images of CNTs with diameter of (a) < 8 nm; (b) 8–15 nm; (c) 20–30 nm; (d) 30–50 nm; (e) > 50 nm and Raman shift of CNTs with diameter of (f) < 8 nm; (g) 8–15 nm; (h) 20–30 nm; (i) 30–50 nm; (j) > 50 nm.

Although the TE properties of organic materials can be largely enhanced by compounding SWNTs with polymer matrix and numerous CNTs/polymer composite TE materials have been prepared by all kinds of methods, the effect of physical

properties of carbon nanotubes on the thermoelectric properties of CNTs/polymer composites have been rarely studied. Previously, Wang²⁵ compared the TE properties of CNTs/polymer composites using metallic CNT and semi-conducting



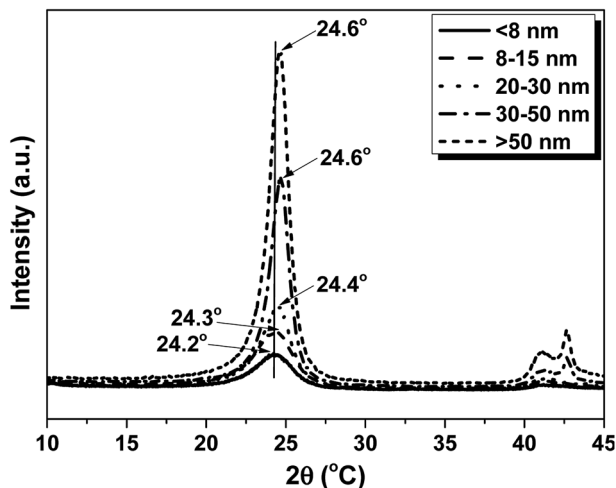


Fig. 2 X-ray diffraction of CNTs with diameter of <8 nm; 8–15 nm; 20–30 nm; 30–50 nm; >50 nm.

CNT, confirming that the semi-conducting CNT/polymer composites exhibited a much higher Seebeck coefficient than that of metallic CNT/polymer composites, attributed to the effective energy filtering effect at the interfaces between semi-conducting CNT and polymer. For carbon nanotubes, the diameter is another important physical properties. Blackburn and Ferguson studied the properties of semiconducting CNTs with different diameter by conjugated polymer wrapping process. The wrapping conjugated polymers are fluorine-based polymers. However, the CNTs extracted by polymers still contain a broad diameter range, while the diameter of the

separated CNT was determined by average energy of the extracted CNTs calculated from absorbance spectra.^{26,27} Therefore, it is still not so clear about how the effect of CNT diameter on the thermoelectric properties of CNTs/polymer composites.

Generally, the chirality and diameter of SWNTs are uniquely specified by the vector, $C_h = ma_1 + na_2$, where m , n are integers and a_1 , a_2 are the unit vectors of graphite, as shown in Fig. S1.^{†28–31} CNTs can be deemed to form by crimping a graphene sheet along the direction of the vector C_h . The diameter (d) and chiral angle (θ) can be calculated from the indices of n and m using the following equation:

$$d = \frac{a}{\pi} \sqrt{n^2 + m^2 + nm}$$

$$\cos \theta = \frac{2n + m}{2\sqrt{n^2 + m^2 + nm}}$$

Therefore, the different diameter of CNTs would lead to diverse electronic transport properties for CNTs/polymer composites. Thus, it is vital to clarify the effect of diameter of CNTs on the TE properties of CNTs/polymer composites.

Herein, we fabricated CNT/P3HT composite films with five different diameter of CNTs ($d < 8$ nm, 8–15 nm, 20–30 nm, 30–50 nm and >50 nm) and their TE properties were investigated. It was found that the diameter of CNTs had an important influence on the TE performance of the composite films. The P3HT- $d_{\text{CNT}} (<8$ nm) and P3HT- $d_{\text{CNT}} (8-15$ nm) composite films exhibited almost the same thermoelectric performance and almost more than twice times of the other three composite films with increased CNT

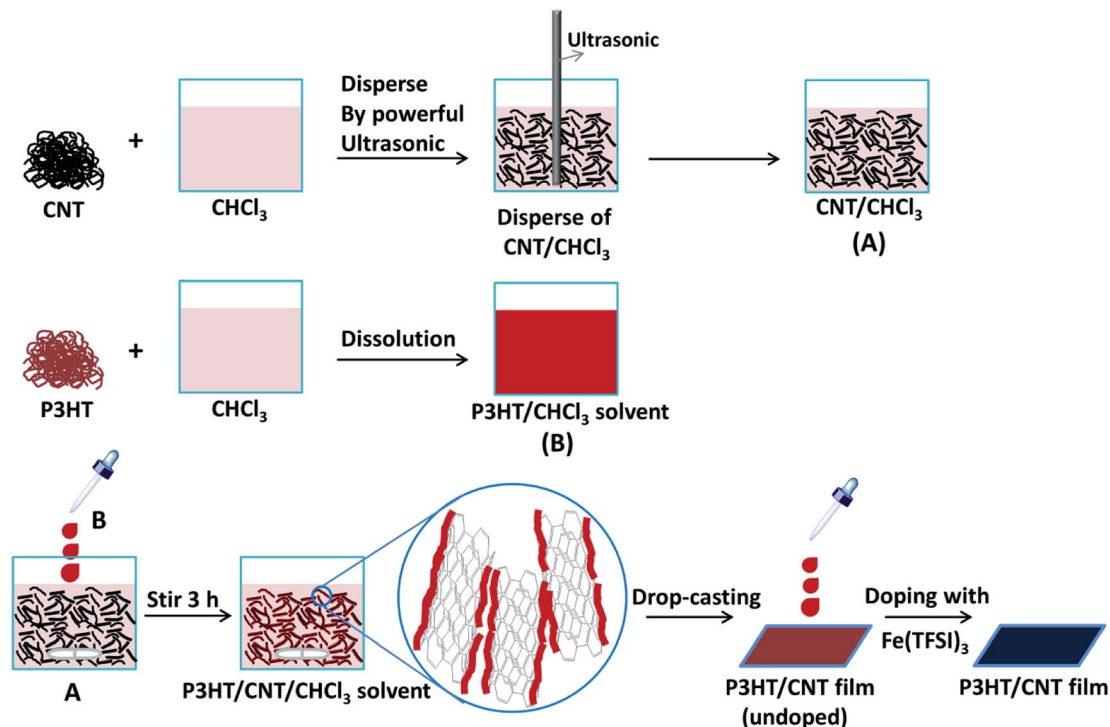


Fig. 3 Schematic illustration of the synthesis procedure of P3HT/CNT composite films with different diameter or CNT content.



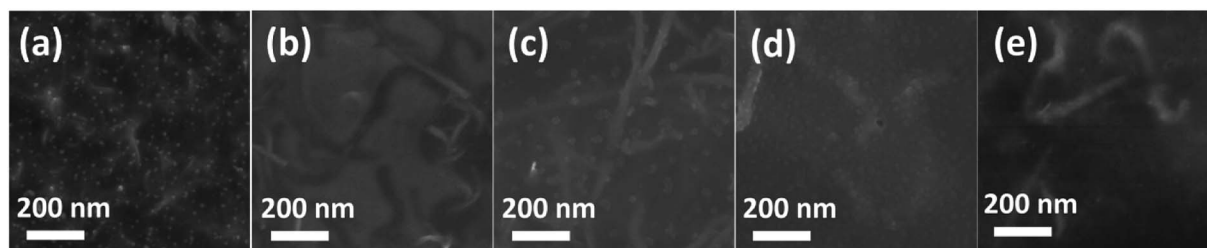


Fig. 4 SEM images of P3HT/CNT composite films with diameter of (a) $< 8\text{ nm}$; (b) 8–15 nm; (c) 20–30 nm; (d) 30–50 nm; (e) $> 50\text{ nm}$.

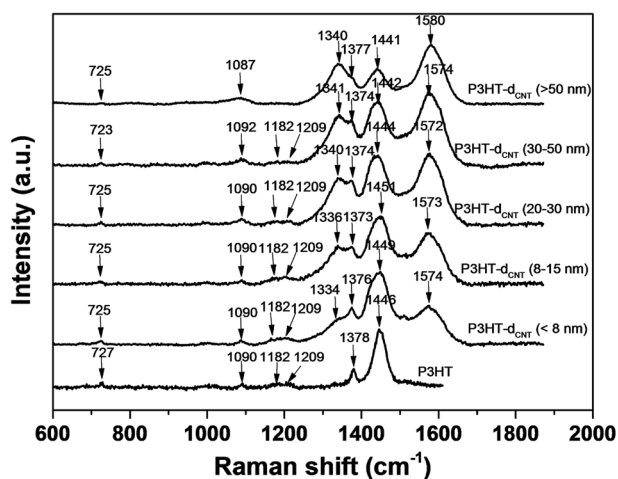


Fig. 5 Raman shift of P3HT/CNT composite films with diameter of $< 8\text{ nm}$; 8–15 nm; 20–30 nm; 30–50 nm; $> 50\text{ nm}$.

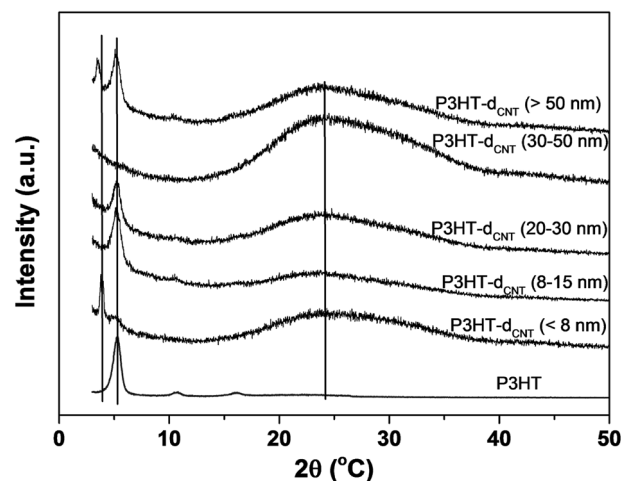


Fig. 6 X-ray diffraction of P3HT/CNT composite films with diameter of $< 8\text{ nm}$; 8–15 nm; 20–30 nm; 30–50 nm; $> 50\text{ nm}$.

Table 1 The assignment of the typical Raman spectra peaks for P3HT

| Raman bands (cm ⁻¹) | Assignment |
|---------------------------------|---|
| ~1445 | Symmetric C=C stretch mode |
| ~1381 | C-C intra-ring stretch mode |
| ~1208 | Inter-ring C-C stretch mode |
| ~1180 | C-H bending mode with C-C inter-ring stretch mode |
| ~1090 | C-H bending |
| ~728 | C-S-C deformation mode |

diameter. The different mass fraction of CNT/P3HT composite films have also been investigated. The maximum TE power factor of CNT ($d < 8\text{ nm}$)/P3HT composite films reached $49.0\ \mu\text{W mK}^{-2}$ at the mass fraction of 95 wt% P3HT, that is, 5 wt% CNT. This superior TE power factor of CNT ($d < 8\text{ nm}$)/P3HT composite films can be ascribed to the fully connected interlayer of P3HT polymer and also the heterogeneous dispersion of short-length CNTs.

2. Experiments

2.1. Materials

The CNTs (MWNT) with length of $\sim 1\ \mu\text{m}$ and diameters of (1) $< 8\text{ nm}$, (2) 8–15 nm, (3) 20–30 nm, (4) 30–50 nm, (5) $> 50\text{ nm}$, were purchased from DK nano (Beijing, China). Regioregular

P3HT (HT-P3HT, the percentage of molecules with the head-to-tail configuration is up to 98%) ($M_w = 87\ \text{kg mol}^{-1}$) was purchased from Sigma-Aldrich (Shanghai, China) and used as received. Trifluoromethanesulfonimide ($\text{CF}_3\text{SO}_2)_2\text{NH}$ (95%) was purchased from Aladdin (Shanghai, China). Analytical grade nitromethane (CH_3NO_2) was dried over CaCl_2 before use. The dopant ferric bis(trifluoromethanesulfonyl)imide ($\text{Fe}(\text{TFSI})_3$) was synthesized by firstly treating $\text{Fe}_2(\text{SO}_4)_3$ aqueous solution with NaHCO_3 aqueous solution to get the precipitates, and then the precipitates was added into acid $(\text{CF}_3\text{SO}_2)_2\text{NH}$ in anhydrous CH_3NO_2 to obtain $\text{Fe}(\text{TFSI})_3/\text{CH}_3\text{NO}_2$ solvent.

2.2. Synthesis of CNTs/P3HT composite films

First, carbon nanotubes (CNTs) were crudely dispersed in chloroform (CHCl_3) by sonication (650 w) for 30 min. Then, regioregular P3HT was dissolved in CHCl_3 ($10\ \text{mg ml}^{-1}$) and different volume of P3HT solution was gradually added into the CNTs/ CHCl_3 ($0.05\ \text{mg ml}^{-1}$) dispersion. After addition step, the mixed CNTs/polymer/ CHCl_3 dispersion was stirred for another 3 h at room temperature, resulting in homogeneous dispersion of CNTs. Subsequently, the well dispersed suspensions were dropped on glass substrates and dried at room temperature. Finally, the CNT/P3HT composite films were immersed into $\text{Fe}(\text{TFSI})_3/\text{CH}_3\text{NO}_2$ for 1 h. The doped CNT/P3HT composite films with different diameters and weight content of CNTs were obtained.



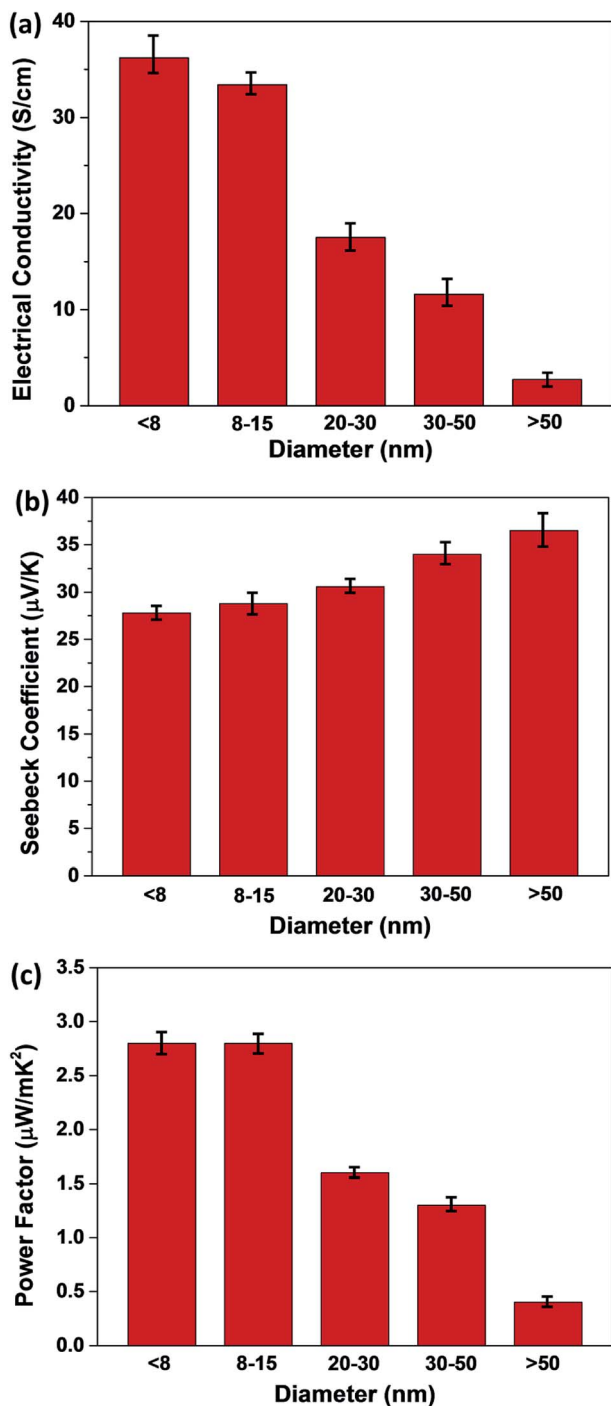


Fig. 7 (a) Electrical conductivity, (b) Seebeck coefficient and (c) power factor of P3HT/CNT (40 wt% CNT) composite films with diameter of <8 nm, 8–15 nm, 20–30 nm, 30–50 nm and >50 nm.

2.3. Characterizations

The morphologies of all the CNTs, CNT/polymer composite films were characterized by scanning electron microscopy (SEM) (FEI Magellan 400). Raman spectra were collected to determine the ordering of CNTs from a Raman spectrometer (Thermo Scientific DXR) equipped with two different laser wavelengths (532 nm and 780 nm). The Grazing incidence X-ray

diffraction (GIXRD) patterns of the films were examined by an X-ray diffractometer (Bruker D8 Advance) with an incidence angle of 1° . The in plane electrical conductivity and Seebeck coefficient of all films were measured at room temperature using a ZEM-5 instrument system (Advance Riko). The thickness of the films were detected on a Dektak profilometer. Typical thickness of the samples was 0.5–1 μm . The power factor was calculated from the corresponding Seebeck coefficient and electrical conductivity at room temperature.

3. Results and discussion

3.1. Characterization of pure CNTs

Fig. 1a–e shows the SEM images of pure CNTs with different diameter. The pure CNTs were uniformly separated with nearly few aggregation. The diameter of the CNTs was gradually increased from Fig. 1a–e. The CNT bundles showed partially disordered structures with diameter of less than 15 nm (Fig. 1a and b), while for the CNTs with diameter higher than 20 nm, the CNT bundles were displaying ordered and clear structure as shown in Fig. 1c–e. Fig. 1f–j presents the Raman spectra of the pure CNTs with different diameter. The spectra of pure CNTs show two strong peaks at $\sim 1580\text{ cm}^{-1}$ and $\sim 1340\text{ cm}^{-1}$, assigned to the G-band (in-plane stretching E_{2g} mode) and D-band (defect or disorder induction mode), respectively.³²

The G-band spectra of all the CNTs consisted of one main peak and one shoulder peak. The main peak with lower frequency attributed to vibrations along the circumferential direction and the acromion peak with upper frequency associated with vibrations along the direction of the nanotube axis. The D-band attributed to defect or disorder induction mode was also observed in the Raman spectra. However, the G : D ratio was gradually increased from Fig. 1f–j, indicating increased ordering and less defects with increased diameter of the CNT.

X-ray diffraction (XRD) of CNTs were characterized to give further evidence of the change in crystallinity for the CNTs with different diameter, as shown in Fig. 2. The samples show typical peak at $2\theta = 24.2^\circ, 24.3^\circ, 24.4^\circ, 24.6^\circ$ and 24.6° for the $d_{\text{CNT}} (<8\text{ nm})$, $d_{\text{CNT}} (8\text{--}15\text{ nm})$, $d_{\text{CNT}} (20\text{--}30\text{ nm})$, $d_{\text{CNT}} (30\text{--}50\text{ nm})$ and $d_{\text{CNT}} (>50\text{ nm})$, respectively. This peak was assigned to lattice plane (002), similar with the peak of graphite.³³ The gradually upper shift of the peaks indicated that the CNTs have decreased interlayer spacing with increasing CNT diameter. Besides, the full width at half-maximum (FWHM) gradually decrease with increasing CNT diameter. This results also display increased ordering of the CNTs with increasing CNT diameter, as the FWHM always relates with the material ordering. Another typical peak at $2\theta = 42^\circ$ can also be observed for all the samples, which was assigned to lattice plane (100).³³

3.2. Preparation and characterization of CNT/P3HT composite films with different CNT diameter

In order to improve the connectivity between the CNTs, we induced poly(3-hexylthiophene) (P3HT) to bridge the uniformly dispersed CNT bundles and CNT/P3HT composite films with different CNT diameter ($d_{\text{CNT}} (<8\text{ nm})$, $d_{\text{CNT}} (8\text{--}15\text{ nm})$, $d_{\text{CNT}} (20\text{--}30\text{ nm})$, $d_{\text{CNT}} (30\text{--}50\text{ nm})$ and $d_{\text{CNT}} (>50\text{ nm})$).



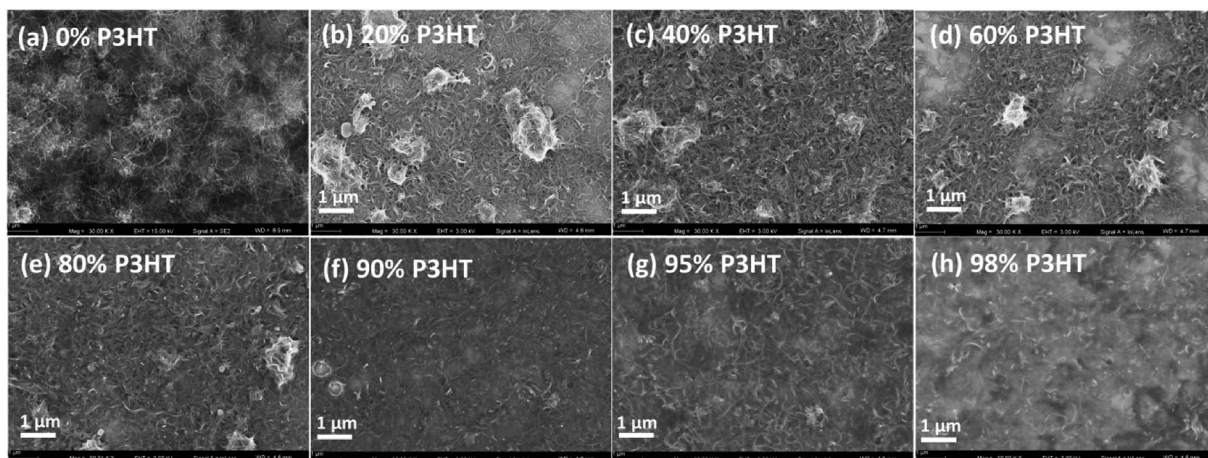


Fig. 8 SEM images of P3HT/CNT ($d < 8$ nm) composite films with P3HT content (wt%) of 0 wt%, 20 wt%, 40 wt%, 60 wt%, 80 wt%, 90 wt%, 95 wt% and 98 wt%.

30 nm), d_{CNT} (30–50 nm) and d_{CNT} (>50 nm)) were prepared. The CNT/P3HT composite films were synthesized based on the procedures described in Fig. 3. During this process, first the CNT was homogeneously dispersed in chloroform (CHCl_3) through powerful sonication (650 w) for 0.5 h, and then the P3HT dissolved in CHCl_3 was slowly added into the CNT/ CHCl_3 dispersion, keeping stirring for another 3 h. The color of the mixture solvents gradually change from dark to deep red when P3HT was added into the CNT dispersions. The P3HT chains were preliminary oriented on the surface of CNT by the π - π conjugation interactions between P3HT and CNT, providing less contact area between CNTs, therefore the CNT/P3HT can steadily dispersed in CHCl_3 solvent. Subsequently, the well dispersed P3HT/CNT/ CHCl_3 suspensions were drop-casted onto the glass substrate and dried at room temperature. Finally, the CNT/P3HT composite films were immersed into solution of the dopant $\text{Fe}(\text{TFSI})_3/\text{CH}_3\text{NO}_2$ solvent for 1 h. The doped CNT/P3HT composite films with different diameters and weight content of CNTs were obtained. The SEM images of CNT/P3HT composite films with different CNT diameter were displayed in Fig. 4 (40 wt% CNT was selected). All the CNTs were well dispersed in the composite films. Meanwhile, the CNTs can be clearly observed in the SEM images and show gradually increased diameters.

The Raman spectra of the CNT/P3HT composite films with different CNT diameter were measured, and the spectra of pure P3HT was also measured as comparison, as shown in Fig. 5. The assignments of all typical Raman peaks are listed in Table 1.³⁴ For the pure P3HT film, all the typical peaks of the pure P3HT were observed. The spectra of the CNT/P3HT films show almost all the peaks of the pure P3HT film and the additional peak at ~ 1340 cm^{-1} and ~ 1570 cm^{-1} of the CNT can also be observed. The peak intensities at ~ 1340 cm^{-1} and ~ 1570 cm^{-1} of the CNT gradually increase with the increase of the CNT diameter, while the intensities at ~ 1182 cm^{-1} , ~ 1209 cm^{-1} and ~ 1445 cm^{-1} decrease on the other hand, suggesting an increase in the in-plane stretching of CNT and a decrease of C–C inter-ring stretch and symmetric C=C stretch of P3HT.

XRD spectra of P3HT/CNT composite films with different CNT diameter has also been measured and displayed in Fig. 6. The pure P3HT show three sharp peaks at $2\theta = 5.3^\circ$, 10.7° and 16.1° , which are assigned to the repeat unit of hexyl chains ((100), (200) and (300) lattice planes). The broad peak at $2\theta = 24.2^\circ$ assigned to the repeat unit of π - π stacking of pure P3HT was not obvious in the figure. After composite of P3HT with different CNTs, there appears an obvious peak at $2\theta = \sim 24^\circ$, ascribed to the CNT peaks. However, the peak at $2\theta = 42^\circ$ of CNTs was not observed in all the P3HT/CNT composites. The peak intensities of P3HT- $d_{\text{CNT}} (< 8$ nm) and P3HT- $d_{\text{CNT}} (30\text{--}50$ nm) composite films were low compared to other samples, ascribed to the high intensity of the (002) lattice planes of CNT. New peaks at lower 2θ degree were observed in the P3HT- $d_{\text{CNT}} (< 8$ nm) and P3HT- $d_{\text{CNT}} (> 50$ nm) composite films, which is due to two possible factors. One factor is the strong π - π interaction between the thiophene rings and CNT surface of the composite films. The π - π interaction between the thiophene rings and $d_{\text{CNT}} (< 8$ nm) is because of the small diameter increased the van der Waals force between the two subject. While the strong interaction between the thiophene rings and $d_{\text{CNT}} (> 50$ nm) is possibly because of the increased ordering of CNT surface, therefore improved the π - π conjugation length. The other possible factor is the effect of the side chains of the polymer. The interaction between P3HT and CNTs is strongly depends on CNT diameter during the polymer wrapping process.

The electrical conductivity (σ) and Seebeck coefficient (S) of P3HT/CNT composite films with different CNT diameters at room temperature were measured, and the results are shown in Fig. 7, with error bars. The average σ of the composite films are showing remarkable difference with different CNT diameters. The σ of P3HT- $d_{\text{CNT}} (< 8$ nm), P3HT- $d_{\text{CNT}} (8\text{--}15$ nm), P3HT- $d_{\text{CNT}} (20\text{--}30$ nm), P3HT- $d_{\text{CNT}} (30\text{--}50$ nm) and P3HT- $d_{\text{CNT}} (> 50$ nm) composite films are 36.2 S cm^{-1} , 33.4 S cm^{-1} , 17.5 S cm^{-1} , 11.6 S cm^{-1} and 2.7 S cm^{-1} , respectively. The σ of the composite films are decreased with increasing CNT diameter. The S of the composite films are 27.8 $\mu\text{V K}^{-1}$, 28.8 $\mu\text{V K}^{-1}$, 30.6 $\mu\text{V K}^{-1}$, 34.0 $\mu\text{V K}^{-1}$ and 36.5 $\mu\text{V K}^{-1}$, respectively, showing slight



increasement with increasing CNT diameter. As a result, the power factor (PF) of the composite films are 2.8, 2.8, 1.6, 1.3 and $0.4 \mu\text{W mK}^{-2}$, respectively. The P3HT- $d_{\text{CNT}} (<8 \text{ nm})$ and P3HT- $d_{\text{CNT}} (8\text{--}15 \text{ nm})$ composite films exhibited almost the same thermoelectric performance, while the other three composite films show decreased thermoelectric performance with increased CNT diameter, owing to the deterioration of the electrical conductivity.

3.3. Preparation and characterization of CNT/P3HT composite films with different mass fraction of P3HT content

The CNT/P3HT composite films with 40 wt% CNT showed relatively low thermoelectric performance possibly because of insufficient electrical connection due to less composition of P3HT. Therefore, the different mass fraction of CNT/P3HT composite films have been prepared. In order to get improved properties of composites, we chose the CNT ($d < 8 \text{ nm}$) and prepared the P3HT/CNT composite films with different P3HT content (0 wt%, 20 wt%, 40 wt%, 60 wt%, 80 wt%, 90 wt%, 95 wt% and 98 wt%). The SEM images of the composite films with different P3HT content (wt%) are shown in Fig. 8. The pure CNT film with 0 wt% P3HT was effectively dispersed with little connection. After composite with 20 wt% P3HT, some part showed aggregation of CNT bundles with P3HT, while the other part showed dispersed CNTs without enough connection. When the introduced polymer P3HT was increased to 40 wt%, 60 wt%, 80 wt% and 90 wt%, the aggregation bundles started to decrease gradually and showed better connection and more compact structure of the film. When further increased P3HT content to 95 wt% and 98 wt%, the CNTs were observed to be homogeneously dispersed into the composite films and with compact structure. The increasing content of P3HT improved the connectivity and accessory of the CNTs with short length.

Fig. 9 presents the thermoelectric properties of P3HT/CNT composite films as a function of the mass fraction of P3HT at room temperature. The electrical conductivity increased with the increase of P3HT content, reaching the peak value of 579 S cm^{-1} for the sample with 95 wt% P3HT, then slightly decreased to 555 S cm^{-1} for the sample with 98 wt% P3HT. The electrical conductivity of pure P3HT was decreased to 84.1 S cm^{-1} . While Seebeck coefficient of composites slightly increased from $21.2 \mu\text{V K}^{-1}$ to $29.1 \mu\text{V K}^{-1}$ and suddenly dropped to $21.5 \mu\text{V K}^{-1}$ at 98 wt% P3HT. The Seebeck coefficient of pure P3HT was increased to $52.9 \mu\text{V K}^{-1}$. The suddenly drops of Seebeck coefficient and electrical conductivity is possibly because of the less connection of the CNTs decrease the carrier mobility of the composites. Consequently, the highest power factor of the P3HT/CNT composite film was $49.0 \mu\text{W mK}^{-2}$ at the mass fraction of 95 wt% P3HT as shown in Fig. 9.

The theoretical electrical conductivity and Seebeck coefficient value of composite system are generally estimated based on either the series or parallel connected mixture models.^{35,36} For the P3HT/CNT system, the volume fraction of CNT and P3HT can be substituted by the weight fraction due to the similar density of CNT and P3HT. Thus, the electrical

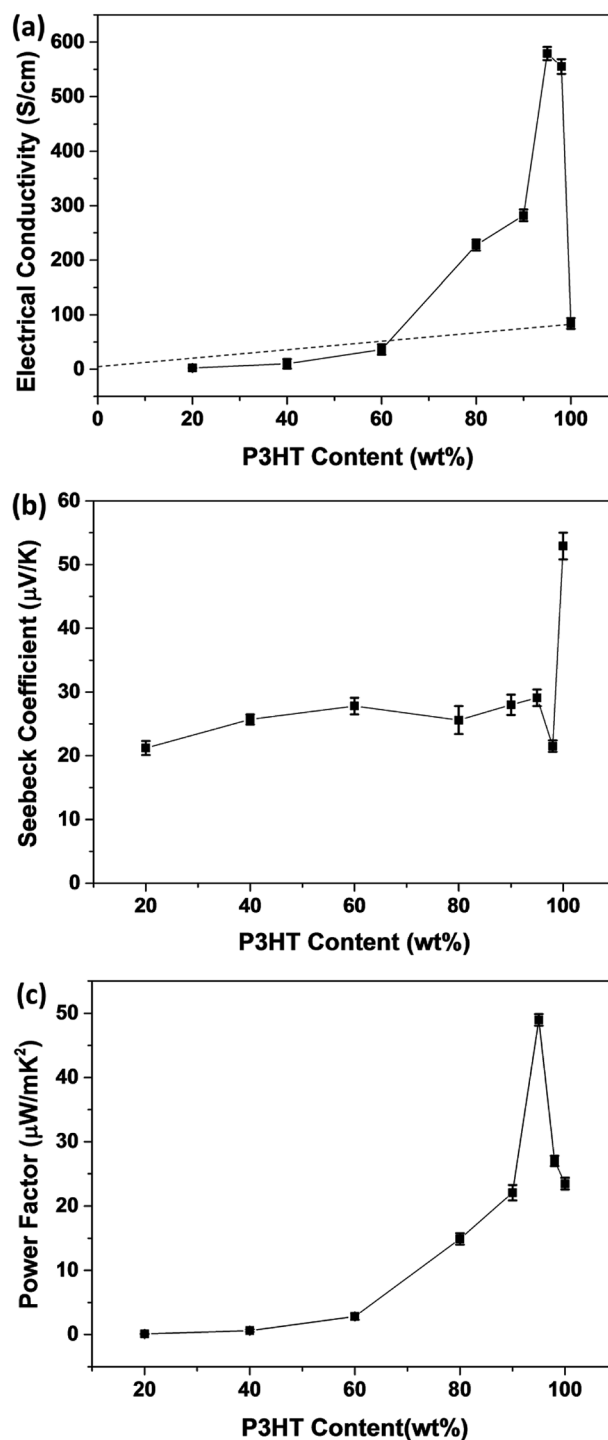


Fig. 9 (a) Electrical conductivity, (b) Seebeck coefficient and (c) power factor of P3HT/CNT ($d < 8 \text{ nm}$) composite films with P3HT content (wt%) of 20 wt%, 40 wt%, 60 wt%, 80 wt%, 90 wt%, 95 wt%, 98 wt% and 100 wt%. The dashed line is the estimated electrical conductivity based on the parallel connected mixture model.

conductivity and Seebeck coefficient of the P3HT/CNT composite films can be calculated using the following expressions:³⁷

$$\sigma (\text{parallel}) = \sigma_{\text{P3HT}} V_{\text{P3HT}} + \sigma_{\text{CNT}} (1 - V_{\text{P3HT}})$$



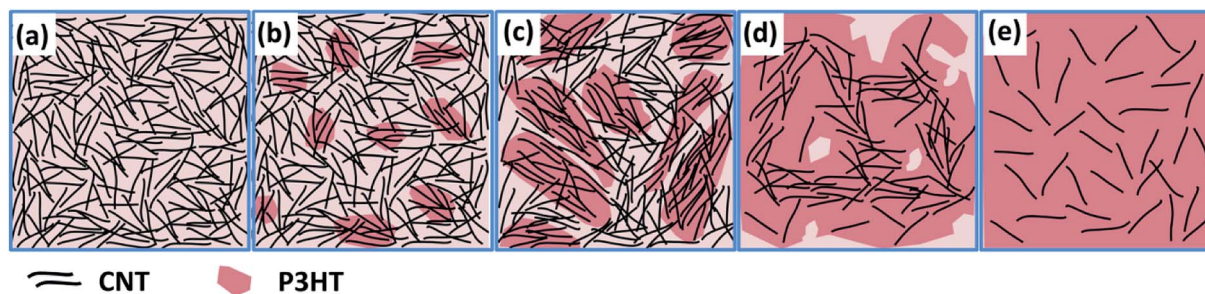


Fig. 10 Illustration scheme of electron transport in composite film: (a) pure CNT; (b) low P3HT content (<20 wt%); (c) medium P3HT content (20–60 wt%); (d) high P3HT content (60–95 wt%); (e) ultra high P3HT content (98 wt%).

$$\sigma (\text{series}) = \frac{\sigma_{\text{P3HT}} \sigma_{\text{CNT}}}{\sigma_{\text{P3HT}} (1 - V_{\text{P3HT}}) + \sigma_{\text{CNT}} V_{\text{P3HT}}}$$

$$S (\text{parallel}) = \frac{S_{\text{P3HT}} \sigma_{\text{P3HT}} V_{\text{P3HT}} + S_{\text{CNT}} (1 - V_{\text{P3HT}})}{\sigma_{\text{P3HT}} V_{\text{P3HT}} + \sigma_{\text{CNT}} (1 - V_{\text{P3HT}})}$$

$$S (\text{series}) = \frac{S_{\text{P3HT}} k_{\text{CNT}} V_{\text{P3HT}} + S_{\text{CNT}} k_{\text{P3HT}} (1 - V_{\text{P3HT}})}{k_{\text{CNT}} V_{\text{P3HT}} + k_{\text{P3HT}} (1 - V_{\text{P3HT}})}$$

where σ (parallel) and S (parallel) are the calculated electrical conductivity and Seebeck coefficient of the P3HT/CNT composite system based on the parallel connected mixture model. σ_{P3HT} and S_{P3HT} are the electrical conductivity and Seebeck coefficient of the P3HT content. σ_{CNT} and S_{CNT} are the electrical conductivity and Seebeck coefficient of the CNT content. V_{P3HT} is the volume fraction of the P3HT.

The percolation theory is generally a reasonable tool to describe the insulator-to-conductor transitions in composites. The estimated electrical conductivity based on the parallel connected mixture model has been shown in the dashed line in Fig. 9a. Interestingly, it is shown that the experimental conductivity values of the P3HT/CNT composite film are firstly lower than those of calculated values when the P3HT content is smaller than 60 wt%, and then dramatically higher than those of calculated values with further increasing P3HT content.

In general, the percolation threshold is no more than 10 vol% in the CNT-based system due to its high aspect ratio for the long-length CNTs. However, in this CNT/P3HT composite films, the CNT is with short length of about only 1 μm . The CNT film without adding P3HT is of individual CNT or nanotube clusters isolated like the island, which prevents contact between the CNTs, showing insulator character of the pure CNT film, as seen in Fig. 10a. After introducing 20 wt% polymer P3HT into the CNTs, the electrical conductivity of the composite film can be observed but with low value. This is possibly because the aggregation of CNT bundles with P3HT only connect some part of the CNTs, while most of the other CNTs are still remained insulated clusters. When increasing the P3HT content to 40 wt% and 60 wt%, the electrical conductivity of composite films increased slowly, which is possibly because of the

aggregation of CNT bundles with P3HT, letting the interlayer of P3HT polymer and CNT cannot be fully connected, as shown in Fig. 10c. When further increasing the P3HT content to 80 wt%, 90 wt% and 95 wt%, there is enough polymer chains to wrap around the CNTs, therefore the CNTs can be sufficiently debundle and disperse with P3HT wrapping around its tubes. The prepared CNT/P3HT composite films are therefore composed of fully dispersed CNTs with P3HT wrapping around its tubes, without aggregated CNT bundles, which can be proved by the SEM images shown in Fig. 8. The interlayer of P3HT polymer and CNT become fully connected into a network, which make the delocalization carriers spread over the entire network, enabling the rapid propagation of electrons along the π - π conjugated region, therefore the electrical conductivity was dramatically increased to the maximum value. And then there is a little bit decrease of the electrical conductivity probably because of the less connection again due to the decreased amount of CNTs.

4. Conclusions

In summary, we have demonstrated influence of diameter of CNTs on the TE performance of the composite films. The SEM, Raman and XRD analysis results showed that CNTs display increased ordering with increasing CNT diameter. Consequently, the P3HT- d_{CNT} (<8 nm) and P3HT- d_{CNT} (8–15 nm) composite films exhibited almost the same thermoelectric performance and almost more than twice times of the other three composite films with increased CNT diameter. The CNT/P3HT composite films with different mass fraction of CNT have been prepared. With the increasing of bridging material P3HT, the electrical conductivity increased gradually for the P3HT content from 40 wt% to 60 wt%, and then dramatically increased from 60 wt% to 95 wt%, reaching the maximum value of 579 S cm^{-1} for the sample with 95 wt% P3HT, that is, 5 wt% CNT. The interlayer of P3HT polymer and CNT become fully connected into a network, which make the delocalization carriers spread over the entire network, enabling the rapid propagation of electrons along the π - π conjugated region, therefore the power factor was dramatically increased to the maximum value of 49.0 $\mu\text{W mK}^{-2}$ at the mass fraction of 5 wt% CNT.



Conflicts of interest

There are no conflicts to declare.

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References

- 1 L. Bell, *Science*, 2008, **321**, 1457–1461.
- 2 A. Ioffe, *Semiconductor Thermoelements and Thermo-electric Cooling*, Infosearch Limited, London, UK, 1957.
- 3 R. Mehta, Y. Zhang, C. Karthik, B. Singh, R. Siegel, T. Tasciuc and G. Ramanath, *Nat. Mater.*, 2012, **11**, 233–240.
- 4 X. Tang, W. Xie, H. Li, W. Zhao, Q. Zhang and M. Niino, *Appl. Phys. Lett.*, 2007, **90**, 012102.
- 5 H. Liu, X. Shi, F. Xu, L. Zhang, W. Zhang, L. Chen, Q. Li, C. Uher, T. Day and G. J. Snyder, *Nat. Mater.*, 2012, **11**, 422–425.
- 6 Q. Zhang, Y. Sun, W. Xu and D. Zhu, *Adv. Mater.*, 2014, **26**, 6829–6851.
- 7 Y. Chen, Y. Zhao and Z. Liang, *Energy Environ. Sci.*, 2015, **8**, 401–422.
- 8 K. Shi, F. Zhang, C. Di, T. Yan, Y. Zou, X. Zhou, D. Zhu, J. Wang and J. Pei, *J. Am. Chem. Soc.*, 2015, **137**, 6979–6982.
- 9 H. Ju and J. Kim, *ACS Nano*, 2016, **10**, 5730–5739.
- 10 N. Toshima, K. Oshima, H. Anno, T. Nishinaka, S. Ichikawa, A. Iwata and Y. Shiraishi, *Adv. Mater.*, 2015, **27**, 2246–2251.
- 11 R. Tian, C. Wan, Y. Wang, Q. Wei, T. Ishida and A. Yamamoto, *J. Mater. Chem. A*, 2016, **5**, 564–570.
- 12 C. Bounioux, P. Diaz-Chao, M. Campoy-Quiles, M. S. Martin-Gonzalez, A. R. Goni, R. Yerushalmi-Rozen and C. Muller, *Energy Environ. Sci.*, 2013, **6**, 918–925.
- 13 Y. Wang, S. Zhang and Y. Deng, *J. Mater. Chem. A*, 2016, **4**, 3554–3559.
- 14 R. Wu, H. Yuan, C. Liu, J. Lan, X. Yang and Y. Lin, *RSC Adv.*, 2018, **8**, 26011.
- 15 W. Shi, S. Qu, H. Chen, Y. Chen, Q. Yao and L. Chen, *Angew. Chem., Int. Ed.*, 2018, **57**, 8037–8042.
- 16 D. Huang, C. Wang, Y. Zou, X. Shen, Y. Zang, H. Shen, X. Gao, Y. Yi, W. Xu, C. Di and D. Zhu, *Angew. Chem., Int. Ed.*, 2016, **55**, 10672–10675.
- 17 C. Wan, R. Tian, M. Kondou, R. Yang, P. Zong and K. Koumoto, *Nat. Commun.*, 2017, **8**, 1024–1032.
- 18 S. Qu, Q. Yao, L. Wang, Z. Chen, K. Xu, H. Zeng, W. Shi, T. Zhang, C. Uher and L. Chen, *NPG Asia Mater.*, 2016, **8**, e292.
- 19 Y. Du, S. Shen, W. Yang, S. Chen, Z. Qin, K. Cai and P. Casey, *J. Electron. Mater.*, 2012, **41**, 1436–1441.
- 20 Y. Du, S. Shen, W. Yang, K. Cai and P. Casey, *Synth. Met.*, 2012, **162**, 375–380.
- 21 C. T. Hong, W. Lee, Y. H. Kang, Y. Yoo, J. Ryu, S. Y. Cho and K.-S. Jang, *J. Mater. Chem. A*, 2015, **3**, 12314.
- 22 W. Lee, C. T. Hong, O. H. Kwon, Y. Yoo, Y. H. Kang, J. Y. Lee, S. Y. Cho and K.-S. Jang, *ACS Appl. Mater. Interfaces*, 2015, **7**, 6550–6556.
- 23 Q. Yao, L. Chen, W. Zhang, S. Liufu and X. Chen, *ACS Nano*, 2010, **4**, 2445–2451.
- 24 Q. Yao, Q. Wang, L. Wang and L. Chen, *Energy Environ. Sci.*, 2014, **7**, 3801–3807.
- 25 L. Wang, Q. Yao, S. Qu, W. Shi and L. Chen, *Org. Electron.*, 2016, **9**, 146–152.
- 26 A. D. Avery, B. H. Zhou, J. Lee, E. Lee, E. M. Miller, R. Ihly, D. Wesenberg, K. S. Mistry, S. L. Guillot, B. L. Zink, Y.-H. Kim, J. L. Blackburn and A. J. Ferguson, *Nat. Energy*, 2016, **1**, 1–9.
- 27 B. A. MacLeod, N. J. Stanton, I. E. Gould, D. Wesenberg, R. Ihly, Z. R. Owczarczyk, K. E. Hurst, C. S. Fewox, C. N. Folmar, K. H. Hughes, B. L. Zink, J. L. Blackburn and A. J. Ferguson, *Energy Environ. Sci.*, 2017, **10**, 2168–2179.
- 28 R. Saito, M. Fujita, G. Dresselhaus and M. S. Dresselhaus, *Appl. Phys. Lett.*, 1992, **60**, 2204–2206.
- 29 S. Reich and C. Thomsen, *Phys. Rev. B*, 2000, **62**, 4273–4276.
- 30 M. S. Strano, C. A. Dyke, M. L. Usrey, P. W. Barone, M. J. Allen, H. W. Shan, C. Kittrell, R. H. Hauge, J. M. Tour and R. E. Smalley, *Science*, 2003, **301**, 1519–1522.
- 31 H. Wang and Z. Bao, *Nano Today*, 2015, **10**, 737–758.
- 32 M. S. Dresselhaus, G. Dresselhaus, A. Jorio, A. G. Souza and R. Saito, *Carbon*, 2002, **40**, 2043–2061.
- 33 S. Basu-Dutt, M. L. Minus, R. Jain, D. Nepal and S. Kumar, *J. Chem. Educ.*, 2012, **89**, 221–229.
- 34 W. Tsoi, D. James, J. Kim, P. Nicholson, C. Murphy, D. Bradley, J. Nelson and J. Kim, *J. Am. Chem. Soc.*, 2011, **133**, 9834–9843.
- 35 Y. Gelbstein, *J. Appl. Phys.*, 2009, **105**, 023713.
- 36 N. E. Coates, S. K. Yee, B. McCulloch, *et al.*, *Adv. Mater.*, 2013, **25**, 1629–1633.
- 37 L. Wang, Q. Yao, W. Shi, S. Qu and L. Chen, *Mater. Chem. Front.*, 2017, **1**, 741–748.

