



Cite this: *RSC Adv.*, 2019, 9, 7432

# Structural exploration of $Au_xM^-$ ( $M = Si, Ge, Sn; x = 9-12$ ) clusters with a revised genetic algorithm†

Ping Huang,<sup>a</sup> Yan Jiang,<sup>a</sup> \*<sup>a</sup> Tianquan Liang,<sup>b</sup> Enhui Wu,<sup>a</sup> Jun Li<sup>a</sup> and Jing Hou<sup>c</sup>

We used a revised genetic algorithm (GA) to explore the potential energy surface (PES) of  $Au_xM^-$  ( $x = 9-12$ ;  $M = Si, Ge, Sn$ ) clusters. The most interesting finding in the structural study of  $Au_xSi^-$  ( $x = 9-12$ ) is the 3D ( $Au_9Si^-$  and  $Au_{10}Si^-$ )  $\rightarrow$  quasi-planar 2D ( $Au_{11}Si^-$  and  $Au_{12}Si^-$ ) structural evolution of the Si-doped clusters, which reflects the competition of Au–Au interactions (forming a 2D structure) and Au–Si interactions (forming a 3D structure). The  $Au_xM^-$  ( $x = 9-12$ ;  $M = Ge, Sn$ ) clusters have quasi-planar structures, which suggests a lower tendency of  $sp^3$  hybridization and a similarity of electronic structure for the Ge or Sn atom.  $Au_9Si^-$  and  $Au_{10}Si^-$  have a 3D structure, which can be viewed as being built from  $Au_8Si^-$  and  $Au_9Si^-$  with an extra Au atom bonded to a terminal gold atom, respectively. In contrast, the quasi-planar structures of  $Au_xM^-$  ( $x = 9-12$ ;  $M = Ge, Sn$ ) reflect the domination of the Au–Au interactions. Including the spin–orbit (SO) effects is very important to calculate the simulated spectrum (structural fingerprint information) in order to obtain quantitative agreement between theoretical and future experimental PES spectra.

Received 8th February 2019  
 Accepted 27th February 2019

DOI: 10.1039/c9ra01019j

[rsc.li/rsc-advances](http://rsc.li/rsc-advances)

## 1. Introduction

The properties of nanoclusters with several hundreds of atoms and molecules strongly depend on both their size and shape. The discovery of catalytic activities of supported gold nanoparticles has attracted intense research interest in the structures and properties of gaseous gold clusters,<sup>1</sup> which provides an atomic-level understanding of catalytic mechanisms of gold nanoparticles. Various techniques have been used to investigate the structures of gold nanoclusters, such as photoelectron spectroscopy (PES),<sup>2</sup> ion mobility,<sup>3,4</sup> infrared multiphoton dissociation spectroscopy,<sup>5,6</sup> and trapped ion electron diffraction.<sup>7</sup> All of these techniques are quite powerful, especially when combined with *ab initio* calculations in obtaining structural information. Many of the unusual properties of the gold clusters derive from the strong relativistic effects that reduce the Au 5d–6s energy gap and enhance s–d hybridization.<sup>8–10</sup> Small-sized anion gold clusters ( $Au_n^-$ ) up to 12 atoms have been found to have 2D planar structures and are understood in terms

of the s–d hybridization.<sup>3,10,11</sup> In the medium-sized gold clusters,  $Au_{16}^-$  to  $Au_{18}^-$  clusters have been found to possess highly stable hollow-cage structures,<sup>7,12</sup> and both  $Au_{20}^-$  and  $Au_{20}$  clusters have tetrahedral structures.<sup>6,13</sup>  $Au_{24}^-$  possibly possesses tubular structure,<sup>14,15</sup> and other large gold clusters, such as  $Au_{25}^-$ ,  $Au_{32}^-$ ,  $Au_{34}^-$ , and  $Au_{55}^-$  to  $Au_{64}^-$ ,<sup>15–23</sup> show core shell structures.

In comparison to pure gold clusters, doped gold clusters have received increasing attention because that the properties of gold clusters can be greatly influenced by the presence of impurities.<sup>24–32</sup> For example, both W and V atoms can change the structure of the pure  $Au_{12}^-$  cluster to form an endohedral structure.<sup>33–35</sup> The pure gold clusters doped with transition-metal atom also have been studied through experiment and theory.<sup>36–38</sup> However, doping an isoelectronic substitution atom, such as Ag and Cu, in a pure gold cluster has a little effect on the electronic and geometrical structure of gold cluster anions.<sup>39–42</sup>

Few studies have been undertaken on the gold clusters doped with group-14 atoms.<sup>26,43–62</sup> Small Si-doped gold clusters, such as  $Au_4Si^{-1/0}$ ,  $Au_4Si_2^{-1/0}$ ,  $Au_2Si_2^{-1/0}$ , and  $Au_3Si_3^{-1/0/+1}$ , have been investigated by Wang *et al.*<sup>26,43,44</sup> The Au/H analogy has been found in the tetrahedral  $Au_4Si^{-1/0}$  cluster. For Ge or Sn dopant, they also affect the structure of gold clusters by different ways relative to the Si dopant. For example, the theoretical calculation shows that tetrahedral structure does not exist in  $Au_4Ge^-$  and  $Au_4Sn^-$ , and they have square-pyramidal structures.<sup>46</sup> The tetrahedral structure derives from the  $sp^3$  hybridization in the  $Au_4Si^-$ . The results suggest that Si atom has stronger tendency to form directional bonding by  $sp^3$  hybridization than Ge and Sn. In order to obtain different isomers of

<sup>a</sup>Panzhuhua International Research Institute of Vanadium and Titanium, Panzhuhua University, Panzhuhua 617000, People's Republic of China. E-mail: [jiangyanzky@163.com](mailto:jiangyanzky@163.com)

<sup>b</sup>College of Environment and Planning, Liaocheng University, Liaocheng 252059, People's Republic of China

<sup>c</sup>College of Vanadium and Titanium, Panzhuhua University, Panzhuhua 617000, People's Republic of China

† Electronic supplementary information (ESI) available: Predicted low-energy structures of the  $Au_xM^-$  ( $x = 9-12$ ;  $M = Si, Ge, Sn$ ) clusters using the revised genetic algorithm combined with density functional theory. See DOI: 10.1039/c9ra01019j



atom cluster system, the genetic algorithm (GA) combined with density functional theory has been used to search the potential energy surface (PES). This method has been proven as an effective stochastic global search algorithm.<sup>63,64</sup> To compare with the future experimental data, we computed the density of state (DOS) spectra by including the relativistic and spin-orbit effects for several lowest-lying isomers. The simulated DOS spectra for the isomers of different systems will be compared with future experimental data, providing considerable credence for the identified isomers for these clusters.

In this paper, the revised genetic algorithm was used to explore the possible low-lying cluster structures of  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ;  $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) systems. We want to explore the structure evolution of different Si-, Ge-, and Sn-doped gold anion clusters and find out the different effect of Si, Ge, and Sn atoms on pure gold clusters ( $\text{Au}_x$ ;  $x = 9-12$ ).

## II. Theoretical methods

A revised genetic algorithm (GA), which is a global searching method based on biological evolution theory, was used to search global minimum of  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ;  $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) systems.<sup>64</sup> The GA method contains several similar procedures as Deaven used: initial population, fitness function evaluation, selection, crossover, mutation and selecting new population.<sup>63</sup> A new subprocess was added after crossover by our GA. The new process was called “exchange operation”, switching Cartesian coordinates of “M” and “Au” element in cluster ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ). Fig. 1 shows the basic operator of mating by cut and swap. Due to the complex PES of multicomponent cluster, the swap operator is very important for multicomponent searching to

effectively explore the local minimum around PES basin of a fixed cluster architecture. The initial population containing 20 individuals were generated by random process, and their coordinates were optimized by quasi-Newton L-BFGS routine.<sup>65</sup> The fitness was calculated by individual energy in population, and the fitness functions can be represented as below:

$$P = (E_i - E_{\min}) / (E_{\max} - E_{\min})$$

where  $E_{\max}$  and  $E_{\min}$  are the highest and lowest energy clusters in the current population, respectively. The selection of two parent individuals was accomplished by roulette wheel, and the low energy individual will have high probability to be selected for crossover. The crossover was achieved by a random plane passing through the center of mass of parent clusters, and the child was generated by collecting two slices atoms which lay above and below the plane of parent 1 and parent 2. The swap operation, which has been executed by exchanging coordinates between different elements based on local optimization structure, was performed for each new produced child. This operation was carried out by a fixed number of times set up by user (10 times for current research), and the lowest energy structure was selected as a new child to replace the child generated by crossover. The mutation is very important for keeping population diversity and also prevents searching trapped in local basin on PES. Ten percent of the offspring were selected for mutation, and the mutation was accomplished by the perturbation of each atomic coordinates with randomly values ranging from  $-1$  to  $1$ . Relaxation to local minimum was performed using L-BFGS method after each operation.<sup>65</sup> All processes of crossover, mutation and selection were then

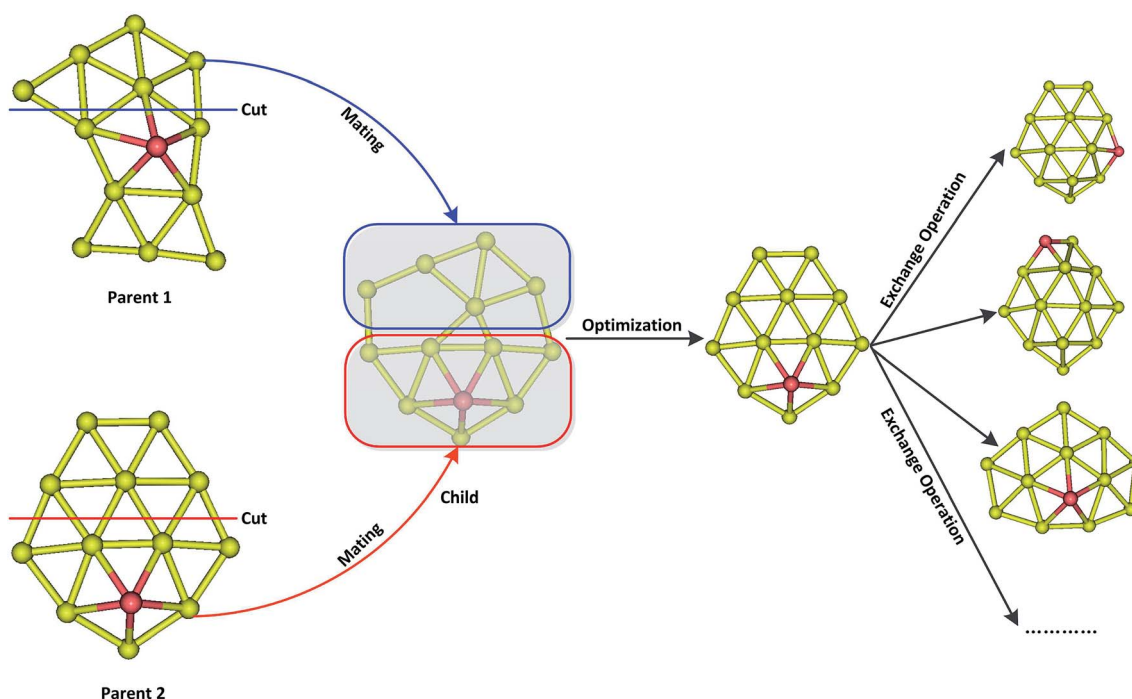


Fig. 1 The schematic diagram of mating operator by cut and swap operator via exchanging two different elements after mating.



repeated for a specified number of generations. Fig. 2 shows the total energy of  $\text{Au}_{12}\text{Ge}^-$  and  $\text{Au}_{12}\text{Si}^-$  clusters against genetic generation by our revised GA.

To search for low-lying structures of  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ;  $M = \text{Si, Ge, Sn}$ ) systems, we used the revised GA global optimization technique combined with DFT for geometry optimization. Generalized gradient approximation in the Perdew–Burke–Ernzerhof (PBE)<sup>66</sup> functional and the double-numerical polarized (DNP) basis set with effective core potentials (ECPs), implemented in the DMol<sup>3</sup> code,<sup>67</sup> were chosen for structure optimization of  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ;  $M = \text{Si, Ge, Sn}$ ) systems. To generate enough isomers in GA searches, a medium level convergence criterion was chosen such that the optimization gradient convergence was less than  $4 \times 10^{-3}$  hartree per Å and the optimization energy convergence was  $2 \times 10^{-5}$  hartree in the DFT calculations. When more than 200 generations were collected, the top-5 lowest-energy structures for last generation were reoptimized using the hybrid functional PBE0 and CRBNBL basis set for all of elements (PBE0/CRBNBL) in the NWChem software package.<sup>68</sup> The isomers were ranked according to their relative energies. The relative energies of the top-5 isomers of the  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ;  $M = \text{Si, Ge, Sn}$ ) systems were further calculated using the PBEPBE/Def2-TZVPPD functional/basis set, PBE1PBE/Def2-TZVPPD functional/basis set, B3LYP/Def2-TZVPPD functional/basis set, and MP2/Def2-TZVPPD second-order Møller–Plesset perturbation theory/basis set, implemented in Gaussian 09 software package.<sup>69</sup> In order to further confirm cluster structure information and compare with future experiment data, we simulated the photoelectron spectroscopy of  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ;  $M = \text{Si, Ge, Sn}$ ) systems. It is very important to include the spin-orbit (SO) effects in the calculation of the density of states in order to obtain quantitative agreement between theoretical and experimental PES spectra.<sup>58</sup> We calculated simulated spectra of all candidate isomers at the PBE0/CRBNBL(SO) level using NWChem software package.<sup>68</sup> Photoelectron spectra were calculated using PBE0 functional and CRBNBL basis set for Si, Ge, and Sn, CRBNBL basis set for Au with spin-orbit effects

included. The first vertical detachment energies (VDEs) of the anion clusters were calculated at PBE0/CRENBL(SO) level as the energy difference between the optimization anion isomer and the neutral at the corresponding anion geometry. The binding energies of deeper orbitals were then added to the first VDE to approximate the higher binding energy detachment features. Each peak was then fitted with a 35 meV-wide Gaussian curve to simulate the DOS spectra.

### III. Theoretical results and discussions

The top-5 lowest-energy structures of the  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ;  $M = \text{Si, Ge, Sn}$ ) systems are depicted in Fig. S1–S4 in the ESI† respectively, together with their simulated PES spectra and coordinates. Tables 1–4 list the relative energies of the corresponding isomers at several levels of theory (see table titles). The simulated spectra of the primary structures of the  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ;  $M = \text{Si, Ge, Sn}$ ) systems are depicted in Fig. 3–6.

#### $\text{Au}_9\text{M}^-$

For  $\text{Au}_9\text{Si}^-$ , PBE0 (NWChem), PBEPBE, and PBE1PBE computations all predict isomer 1 as the global minimum structure (Fig. S11a in the ESI†), which can be viewed as the lowest-energy structure of  $\text{Au}_8\text{Si}^-$  cluster bonded by another Au atom through a terminal gold atom.<sup>58</sup> The B3LYP and MP2 calculations seem to prefer isomer 5 (Fig. S11e in the ESI†) as the lowest energy structure. In order to further identify the primary structure information for future PES experiment spectrum of  $\text{Au}_9\text{Si}^-$  cluster, the simulated spectra of the top-5 lowest-energy structures of  $\text{Au}_9\text{Si}^-$  cluster (Fig. S11a–e in the ESI†) have been calculated at the PBE0/CRENBL(SO) level. It is very important for theoretical spectrum simulation of Au atom to consider spin-orbital interaction by previous study.<sup>58</sup> For  $\text{Au}_9\text{Ge}^-$ , PBE0 (NWChem) computation predicts isomer 1 (Fig. S11a in the ESI†) as the global minimum structure, while the PBEPBE, PBE1PBE and B3LYP computations suggest isomer 2 (Fig. S11b in the ESI†) as the global minimum structure. The MP2 calculations seem to prefer isomer 5 (Fig. S11e in the ESI†) as the

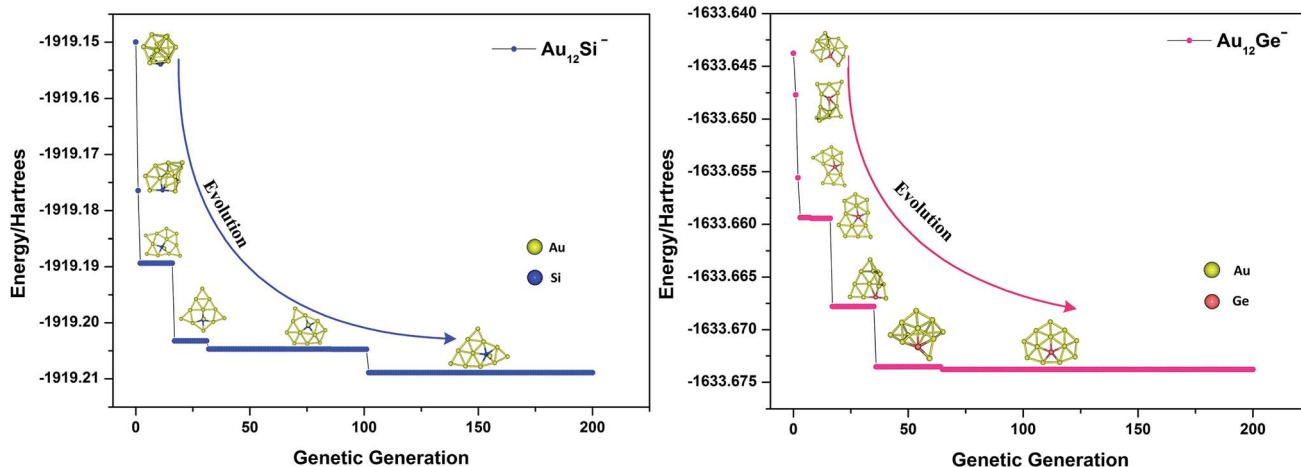


Fig. 2 The total energy against genetic generation of  $\text{Au}_{12}\text{Si}^-$  and  $\text{Au}_{12}\text{Ge}^-$  clusters using revised genetic algorithm.



**Table 1** Relative energies of five low-lying isomers of  $Au_9M^-$  ( $M = Si, Ge, Sn$ ) at PBE0/CRENBL (NWChem) and PBEPBE/Def2-TZVPPD (G09), PBE1PBE/Def2-TZVPPD (G09), B3LYP/Def2-TZVPPD (G09), and MP2/Def2-TZVPPD (G09) levels of theory and basis set<sup>a</sup>

Relative energies (eV)						
Isomer		PBE0	PBEPBE	PBE1PBE	B3LYP	MP2
$Au_9Si^-$	1	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	0.040	0.448
	2	0.182	0.174	0.265	0.315	0.972
	3	0.186	0.141	0.252	0.195	0.908
	4	0.186	0.072	0.171	0.160	0.802
	5	0.379	0.170	0.127	<b>0.000</b>	<b>0.000</b>
$Au_9Ge^-$	1	<b>0.000</b>	0.034	0.122	0.150	0.422
	2	0.005	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	0.019
	3	0.047	0.001	0.104	0.097	0.409
	4	0.069	0.123	0.205	0.336	0.828
	5	0.104	0.069	0.074	0.069	<b>0.000</b>
$Au_9Sn^-$	1	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	0.034
	2	0.090	0.140	0.015	0.147	<b>0.000</b>
	3	0.270	0.295	0.206	0.330	0.322
	4	0.271	0.232	0.129	0.229	0.112
	5	0.273	0.324	0.233	0.379	0.372

<sup>a</sup> Isomers are ranked according to their relative energies at the PBE0/CRENBL level of theory.

lowest energy structure. For  $Au_9Sn^-$ , PBE0 (NWChem), PBEPBE, PBE1PBE, and B3LYP computations all predict isomer 1 as the global minimum structure (Fig. S1ra in the ESI†), while the MP2 computation suggests isomer 2 (Fig. S1rb in the ESI†) as the global minimum structure. Due to the similar chemical properties of Ge and Sn atom, they should have similar lowest-energy structure as previous study. But the  $Au_9Ge^-$  and  $Au_9Sn^-$  have different lowest energy structures by our calculation. Fig. 3 shows the simulated spectra of the primary structures of  $Au_9M^-$

**Table 2** Relative energies of five low-lying isomers of  $Au_{10}M^-$  ( $M = Si, Ge, Sn$ ) at PBE0/CRENBL (NWChem) and PBEPBE/Def2-TZVPPD (G09), PBE1PBE/Def2-TZVPPD (G09), B3LYP/Def2-TZVPPD (G09), and MP2/Def2-TZVPPD (G09) levels of theory and basis set<sup>a</sup>

Relative energies (eV)						
Isomer		PBE0	PBEPBE	PBE1PBE	B3LYP	MP2
$Au_{10}Si^-$	1	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	0.017
	2	0.010	0.113	0.249	0.281	0.789
	3	0.032	0.011	0.048	0.044	0.182
	4	0.069	0.075	0.084	0.152	0.185
	5	0.072	0.098	0.144	0.045	<b>0.000</b>
$Au_{10}Ge^-$	1	<b>0.000</b>	0.133	0.082	0.321	0.647
	2	0.015	<b>0.000</b>	0.019	0.144	0.526
	3	0.066	0.099	0.069	0.196	0.388
	4	0.101	0.060	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>
	5	0.140	0.176	0.147	0.291	0.483
$Au_{10}Sn^-$	1	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	0.098
	2	0.078	0.103	0.060	0.063	<b>0.000</b>
	3	0.112	0.161	0.146	0.187	0.291
	4	0.149	0.151	0.185	0.190	0.353
	5	0.175	0.213	0.186	0.205	0.192

<sup>a</sup> Isomers are ranked according to their relative energies at the PBE0/CRENBL level of theory.

**Table 3** Relative energies of five low-lying isomers of  $Au_{11}M^-$  ( $M = Si, Ge, Sn$ ) at PBE0/CRENBL (NWChem) and PBEPBE/Def2-TZVPPD (G09), PBE1PBE/Def2-TZVPPD (G09), B3LYP/Def2-TZVPPD (G09), and MP2/Def2-TZVPPD (G09) levels of theory and basis set<sup>a</sup>

Relative energies (eV)						
Isomer		PBE0	PBEPBE	PBE1PBE	B3LYP	MP2
$Au_{11}Si^-$	1	<b>0.000</b>	0.001	0.031	0.022	0.243
	2	0.082	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>
	3	0.106	0.074	0.089	0.259	0.506
	4	0.114	0.136	0.159	0.188	0.394
	5	0.130	0.415	0.459	0.951	1.658
$Au_{11}Ge^-$	1	<b>0.000</b>	0.009	0.065	0.186	0.470
	2	0.035	0.238	0.161	0.828	1.264
	3	0.051	0.043	0.112	0.247	0.643
	4	0.125	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>
	5	0.134	0.047	0.116	0.122	0.210
$Au_{11}Sn^-$	1	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	0.032	0.226
	2	0.009	0.037	0.050	0.106	0.410
	3	0.100	0.022	0.143	0.071	0.474
	4	0.144	0.075	0.088	<b>0.000</b>	<b>0.000</b>
	5	0.196	0.183	0.257	0.200	0.452

<sup>a</sup> Isomers are ranked according to their relative energies at the PBE0/CRENBL level of theory.

( $M = Si, Ge, Sn$ ) clusters. The primary structure of  $Au_9Si^-$  cluster was dominated by  $Au_8Si^-$  cluster, reflecting that the electronic structure of  $Au_9Si^-$  cluster should be similar with  $Au_8Si^-$  cluster. The result shows that the lowest-energy structure of  $Au_9Si^-$  cluster has a strong tendency of forming tetrahedral bonding structure by  $sp^3$  hybridization for Si atom. However, Ge- and Sn-doped gold anion cluster have a quasi-planar structure differing from  $Au_9Si^-$  cluster. It suggests the lower tendency of  $sp^3$  hybridization for Ge and Sn atoms. The primary

**Table 4** Relative energies of five low-lying isomers of  $Au_{12}M^-$  ( $M = Si, Ge, Sn$ ) at PBE0/CRENBL (NWChem) and PBEPBE/Def2-TZVPPD (G09), PBE1PBE/Def2-TZVPPD (G09), B3LYP/Def2-TZVPPD (G09), and MP2/Def2-TZVPPD (G09) levels of theory and basis set<sup>a</sup>

Relative energies (eV)						
Isomer		PBE0	PBEPBE	PBE1PBE	B3LYP	MP2
$Au_{12}Si^-$	1	<b>0.000</b>	<b>0.000</b>	0.021	0.150	0.640
	2	0.084	0.069	<b>0.000</b>	0.192	0.343
	3	0.117	0.075	0.003	<b>0.000</b>	<b>0.000</b>
	4	0.131	0.268	0.249	0.942	2.037
	5	0.142	0.099	0.137	<b>0.000</b>	0.741
$Au_{12}Ge^-$	1	<b>0.000</b>	0.016	0.031	0.147	0.511
	2	0.005	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	0.102
	3	0.135	0.175	0.127	0.124	<b>0.000</b>
	4	0.157	0.151	0.152	0.142	0.217
	5	0.182	0.357	0.241	0.680	0.764
$Au_{12}Sn^-$	1	<b>0.000</b>	<b>0.000</b>	<b>0.000</b>	0.092	0.332
	2	0.089	0.111	0.062	0.138	0.167
	3	0.090	0.038	0.033	<b>0.000</b>	<b>0.000</b>
	4	0.163	0.324	0.204	0.821	1.140
	5	0.216	0.323	0.185	0.735	0.856

<sup>a</sup> Isomers are ranked according to their relative energies at the PBE0/CRENBL level of theory.



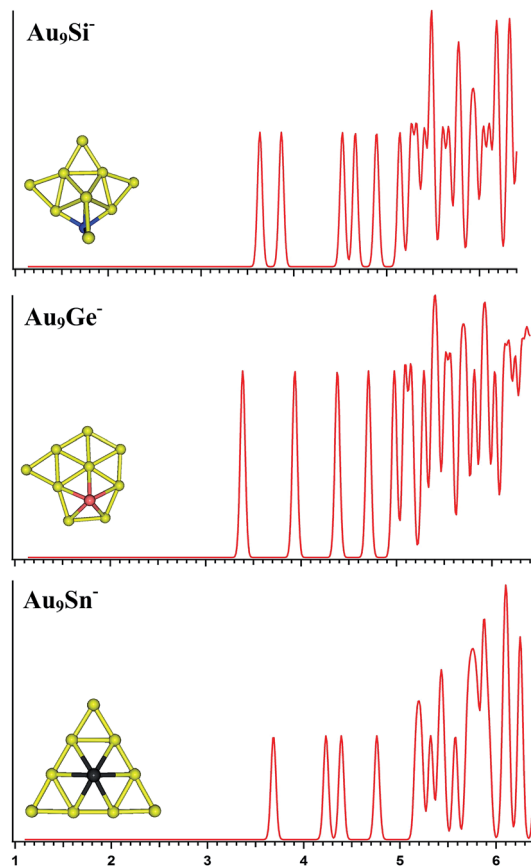


Fig. 3 The simulated photoelectron spectra of  $\text{Au}_9\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) using PBE0/CRENBL (NWChem) level of theory and basis set including the relativistic and spin-orbit effects. The insets show the corresponding structures. The dopant atoms are shown in color (Si in blue, Ge in red, and Sn in black).

structures of  $\text{Au}_9\text{Ge}^-$  and  $\text{Au}_9\text{Sn}^-$  are mainly dominated by Au-Au interactions.

### $\text{Au}_{10}\text{M}^-$

The PBE0 (NWChem), PBEPBE, PBE1PBE, and B3LYP calculations predict isomers 1 (Fig. S2la in the ESI<sup>†</sup>) as the global minimum structure for  $\text{Au}_{10}\text{Si}^-$ . Isomer 1 can be viewed based on the isomer 1 (Fig. S1la in the ESI<sup>†</sup>) of  $\text{Au}_9\text{Si}^-$  cluster by adding an Au atom. However, the MP2 calculation reveals isomer 5 as the lowest energy structure, which has a quasi-planar structure differing from isomer 1 (Fig. S2la in the ESI<sup>†</sup>). For  $\text{Au}_{10}\text{Ge}^-$ , The PBE0 (NWChem) computation (Table 2) indicates that isomer 1 (Fig. S2ma in the ESI<sup>†</sup>) has the global minimum, while PBEPBE calculation seems to prefer isomer 2 (Fig. S2mb in the ESI<sup>†</sup>) as the lowest energy structure, and PBE1PBE, B3LYP, and MP2 calculations reveal isomer 4 as the lowest energy structure. The isomer 1 of  $\text{Au}_{10}\text{Ge}^-$  is also a quasi-planar structure and can be viewed as being derived from isomer 1 of  $\text{Au}_9\text{Ge}^-$  by adding one Au atom. For  $\text{Au}_{10}\text{Sn}^-$ , PBE0 (NWChem), PBE1PBE, PBEPBE and B3LYP calculations reveal that isomer 1 is the lowest-energy structure, while MP2 calculation prefers isomer 2 (Fig. S2rb in the ESI<sup>†</sup>) as the lowest

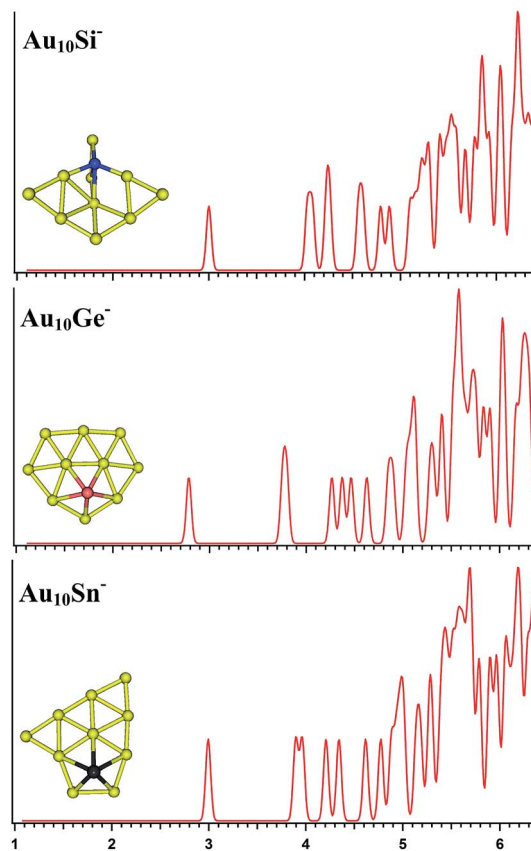


Fig. 4 The simulated photoelectron spectra of  $\text{Au}_{10}\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) using PBE0/CRENBL (NWChem) level of theory and basis set including the relativistic and spin-orbit effects. The insets show the corresponding structures. The dopant atoms are shown in color (Si in blue, Ge in red, and Sn in black).

energy structure. Different theory level computations produce different energy order (Table 2). The isomer 1 and isomer 2 are the primary structures of  $\text{Au}_{10}\text{Sn}^-$  and can be viewed as being derived from isomer 2 of  $\text{Au}_9\text{Sn}^-$ . Fig. 4 shows the simulated spectra of the primary structures of  $\text{Au}_{10}\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) clusters.  $\text{Au}_{10}\text{Si}^-$  has a 3D structure, and shows a strong tendency of forming tetrahedral bonding structure by  $\text{sp}^3$  hybridization for Si atom.  $\text{Au}_{10}\text{M}^-$  ( $\text{M} = \text{Ge}, \text{Sn}$ ) systems show the quasi-planar structures, reflecting that the structure is mainly dominated by Au-Au interactions.

### $\text{Au}_{11}\text{M}^-$

For  $\text{Au}_{11}\text{Si}^-$ , the PBE0 (NWChem) calculation (Table 3) reveals isomer 1 (Fig. S3la in the ESI<sup>†</sup>) as the ground-state structure. Isomer 1 can be viewed as being derived from isomer 5 (Fig. S2le in the ESI<sup>†</sup>) of  $\text{Au}_{10}\text{Si}^-$  by attaching one Au atom to form triangle structure, and has a quasi-planar structure differing from the 3D structures of  $\text{Au}_9\text{Si}^-$  and  $\text{Au}_{10}\text{Si}^-$  because of  $\text{sp}^3$  hybridization for Si atom. Other computations suggest that isomer 2 (Fig. S3lb in the ESI<sup>†</sup>) is the lowest energy structure. The simulated spectrums of 5-top isomers of  $\text{Au}_{11}\text{Si}^-$  cluster can be found in ESI<sup>†</sup> (Fig. S3la-e in the ESI<sup>†</sup>). For  $\text{Au}_{11}\text{Ge}^-$ , PBE0 (NWChem) calculation (Table 3) indicates that isomer 1



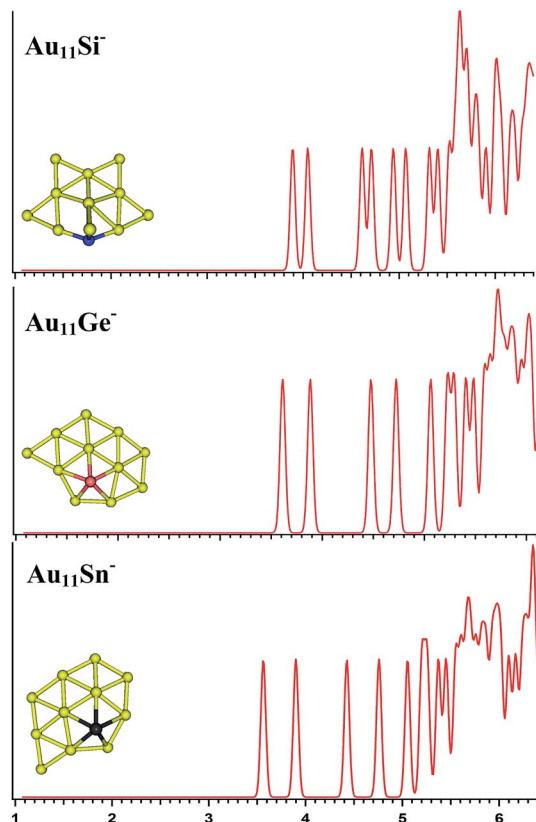


Fig. 5 The simulated photoelectron spectra of  $\text{Au}_{11}\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) using PBE0/CRENBL (NWChem) level of theory and basis set including the relativistic and spin-orbit effects. The insets show the corresponding structures. The dopant atoms are shown in color (Si in blue, Ge in red, and Sn in black).

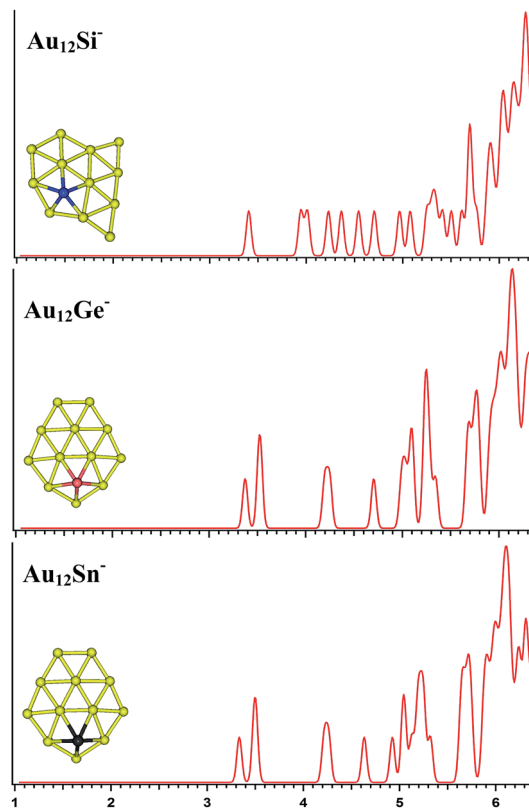


Fig. 6 The simulated photoelectron spectra of  $\text{Au}_{12}\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) using PBE0/CRENBL (NWChem) level of theory and basis set including the relativistic and spin-orbit effects. The insets show the corresponding structures. The dopant atoms are shown in color (Si in blue, Ge in red, and Sn in black).

(Fig. S3ma in the ESI<sup>†</sup>) is the global minimum structure and can be viewed as being derived from isomer 3 (Fig. S2mc in the ESI<sup>†</sup>) of  $\text{Au}_{10}\text{Ge}^-$  by adding one Au atom to form triangle structure. Other computations suggest isomer 2 (Fig. S3mb in the ESI<sup>†</sup>) as the lowest energy structure. For  $\text{Au}_{11}\text{Sn}^-$ , the lowest-energy structure is found to be isomer 1 (Fig. S3ra in the ESI<sup>†</sup>), which has the same structure as  $\text{Au}_{11}\text{Ge}^-$  (Fig. S3ma in the ESI<sup>†</sup>) for PBE0 (NWChem) calculation. B3LYP and MP2 calculations reveal that isomer 4 is the lowest-energy structure, and can be viewed as being derived from isomer 2 (Fig. S2rb in the ESI<sup>†</sup>) of  $\text{Au}_{10}\text{Sn}^-$  by attaching one Au atom to form triangle structure. Fig. 5 shows the simulated spectra of the primary structures of  $\text{Au}_{11}\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) clusters. All of the  $\text{Au}_{11}\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) systems show the same quasi-planar structure, reflecting that the structure is mainly dominated by Au–Au interactions.

### $\text{Au}_{12}\text{M}^-$

The PBE0 (NWChem) and PBEPBE calculations indicate that isomer 1 (Fig. S4la in the ESI<sup>†</sup>, a quasi-planar structure) is the lowest-energy structure for  $\text{Au}_{12}\text{Si}^-$  cluster, and other calculations reveal isomer 2 and isomer 3 as the ground state structure. For  $\text{Au}_{12}\text{Ge}^-$ , PBE0 (NWChem) calculation (Table 4) reveals isomer 1 (Fig. S4ma in the ESI<sup>†</sup>) as the lowest energy structure,

while PBE1PBE, PBEPBE, and B3LYP computations predict isomer 2 (Fig. S4mb in the ESI<sup>†</sup>) as the global minimum structure. Isomer 2 has very similar quasi-planar structure with isomer 1 of  $\text{Au}_{12}\text{Si}^-$  cluster. The MP2 calculations reveal isomer 3 has the global minimum, which can be viewed evolved from the isomer 1 of  $\text{Au}_{11}\text{Ge}^-$  by attaching one Au atom and subsequent structural relaxation. For  $\text{Au}_{12}\text{Sn}^-$ , PBE0 (NWChem), PBEPBE, and PBE1PBE calculations predict isomer 1 (Fig. S4ra in the ESI<sup>†</sup>) as the global minimum, which has a similar structure with isomer 1 of  $\text{Au}_{12}\text{Ge}^-$  cluster. However, the B3LYP and MP2 calculation reveals isomer 3 (Fig. S4rc in the ESI<sup>†</sup>) as the lowest energy structure. Fig. 6 shows the simulated spectra of the primary structures of the  $\text{Au}_{12}\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) clusters. The  $\text{Au}_{12}\text{M}^-$  ( $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) clusters have quasi-planar structures, reflecting that their structures are dominated by Au–Au interactions.

To summarize, the structure evolutions of  $\text{Au}_x\text{M}^-$  ( $x = 9-12$ ,  $\text{M} = \text{Si}, \text{Ge}, \text{Sn}$ ) systems have been explored by revised GA method. The most intriguing finding is the structure change of 3D to quasi-planar for  $\text{Au}_x\text{Si}^-$  ( $x = 9-12$ ). The result suggests the competition between the tendency of forming tetrahedral bonding structure by  $\text{sp}^3$  hybridization for Si atom (Au–M interaction) and the tendency of forming planar structures by Au–Au interactions for small anion gold clusters. The Si-doped



gold clusters easily form local tetrahedral bonding structures because of the strong  $sp^3$  hybridization of the Si atom. But the structure changing from 3D to quasi-planar is found at  $x = 11$  for  $Au_xSi^-$  ( $x = 9-12$ ) systems, which reflects that their structures are dominated by Au–Au interactions when the cluster size increases for  $Au_xSi^-$  ( $x = 9-12$ ) systems. However, the Ge- or Sn-doped gold clusters exhibit different structural images from pure gold anion clusters, which reflects its lower tendency of  $sp^3$  hybridization in the doped clusters. The structural evolutions of  $Au_xM^-$  ( $x = 9-12$ ;  $M = Si, Ge, Sn$ ) systems reflect that the Au–Au interactions become competitive in deciding the cluster structures when the cluster size increases for dopant clusters of gold. Including the spin-orbit (SO) effects are very important to calculate the density of states in order to obtain quantitative agreement between theoretical and future experimental PES spectra.

## IV. Conclusion

In conclusion, a detailed investigation of the structures and isomers of  $Au_xM^-$  ( $x = 9-12$ ;  $M = Si, Ge, Sn$ ) clusters using revised genetic algorithm combined with density functional theory were reported. The subprocess of exchange operation was added to explore local basin on PES for multicomponent systems in GA method. The  $Au_9Si^-$  and  $Au_{10}Si^-$  clusters have a tetrahedral-based 3D structure, and  $Au_{11}Si^-$  and  $Au_{12}Si^-$  clusters have a quasi-planar structure. The results show that the Si-doped gold clusters, which often form 3D direction bonding structure owing to the strong  $sp^3$  hybridization of Si atom, will be replaced by Au–Au interactions to form quasi-planar structure when the cluster size increases for dopant clusters of gold. However, the quasi-planar structures of Ge- or Sn-doped gold clusters reflect that the Ge and Sn atoms have the lower tendency of s–p hybridization in the doped clusters. For  $Au_xM^-$  ( $x = 9-12$ ;  $M = Ge, Sn$ ) systems, quasi-planar structures tend to be formed due to the domination of Au–Au interactions, suggesting that the Ge and Sn atoms have the lower tendency of s–p hybridization in the doped clusters.

## Conflicts of interest

There are no conflicts of interest to declare.

## Acknowledgements

The theoretical work was supported by the National Natural Science Foundation of China (Grant No. 51601189), and the foundation from Sichuan science and technology program (Grant No. 2017JY0117). Some of the computations were performed at the Supercomputing Center of the Chinese Academy of Sciences.

## References

- 1 M. Haruta, *Catal. Today*, 1997, **36**, 153–166.
- 2 K. J. Taylor, C. L. Pettiettehall, O. Cheshnovsky and R. E. Smalley, *J. Chem. Phys.*, 1992, **96**, 3319–3329.

- 3 F. Furche, R. Ahlrichs, P. Weis, C. Jacob, S. Gilb, T. Bierweiler and M. M. Kappes, *J. Chem. Phys.*, 2002, **117**, 6982–6990.
- 4 P. Weis, *Int. J. Mass Spectrom.*, 2005, **245**, 1–13.
- 5 A. Fielicke, A. Kirilyuk, C. Ratsch, J. Behler, M. Scheffler, G. von Helden and G. Meijer, *Phys. Rev. Lett.*, 2004, **93**, 023401–023405.
- 6 P. Gruene, D. M. Rayner, B. Redlich, A. F. G. van der Meer, J. T. Lyon, G. Meijer and A. Fielicke, *Science*, 2008, **321**, 674–676.
- 7 X. Xing, B. Yoon, U. Landman and J. H. Parks, *Phys. Rev. B*, 2006, **74**, 165423.
- 8 P. Pyykko, *Chem. Rev.*, 1988, **88**, 563–594.
- 9 P. Pyykko, *Inorg. Chim. Acta*, 2005, **358**, 4113–4130.
- 10 H. Hakkinen, M. Moseler and U. Landman, *Phys. Rev. Lett.*, 2002, **89**, 033401–033405.
- 11 H. Hakkinen, B. Yoon, U. Landman, X. Li, H. J. Zhai and L. S. Wang, *J. Phys. Chem. A*, 2003, **107**, 6168–6175.
- 12 S. Bulusu, X. Li, L.-S. Wang and X. C. Zeng, *Proc. Natl. Acad. Sci. U. S. A.*, 2006, **103**, 8326–8330.
- 13 J. Li, X. Li, H. J. Zhai and L. S. Wang, *Science*, 2003, **299**, 864–867.
- 14 B. Yoon, P. Koskinen, B. Huber, O. Kostko, B. von Issendorff, H. Hakkinen, M. Moseler and U. Landman, *ChemPhysChem*, 2007, **8**, 157–161.
- 15 S. Bulusu, X. Li, L.-S. Wang and X. C. Zeng, *J. Phys. Chem. C*, 2007, **111**, 4190–4198.
- 16 M. Ji, X. Gu, X. Li, X. G. Gong, J. Li and L. S. Wang, *Angew. Chem., Int. Ed.*, 2005, **44**, 7119–7123.
- 17 A. F. Jalbout, F. F. Contreras-Torres, L. A. Perez and I. L. Garzon, *J. Phys. Chem. A*, 2008, **112**, 353–357.
- 18 A. Lechtken, D. Schooss, J. R. Stairs, M. N. Blom, F. Furche, N. Morgner, O. Kostko, B. von Issendorff and M. M. Kappes, *Angew. Chem., Int. Ed.*, 2007, **46**, 2944–2948.
- 19 X. Gu, S. Bulusu, X. Li, X. C. Zeng, J. Li, X. G. Gong and L.-S. Wang, *J. Phys. Chem. C*, 2007, **111**, 8228–8232.
- 20 I. E. Santizo, F. Hidalgo, L. A. Perez, C. Noguez and I. L. Garzon, *J. Phys. Chem. C*, 2008, **112**, 17533–17539.
- 21 I. L. Garzon, K. Michaelian, M. R. Beltran, A. Posada-Amarillas, P. Ordejon, E. Artacho, D. Sanchez-Portal and J. M. Soler, *Phys. Rev. Lett.*, 1998, **81**, 1600–1603.
- 22 H. Hakkinen, M. Moseler, O. Kostko, N. Morgner, M. A. Hoffmann and B. von Issendorff, *Phys. Rev. Lett.*, 2004, **93**, 093401–093405.
- 23 W. Huang, M. Ji, C.-D. Dong, X. Gu, L.-M. Wang, X. G. Gong and L.-S. Wang, *ACS Nano*, 2008, **2**, 897–904.
- 24 T. K. Ghanty, A. Banerjee and A. Chakrabarti, *J. Phys. Chem. C*, 2010, **114**, 20–27.
- 25 D. Yuan, X. Gong and R. Wu, *Phys. Rev. B*, 2008, **78**, 035441.
- 26 X. Li, B. Kiran, L. F. Cui and L. S. Wang, *Phys. Rev. Lett.*, 2005, **95**, 253401.
- 27 K. Koszinowski, D. Schroder and H. Schwarz, *ChemPhysChem*, 2003, **4**, 1233–1237.
- 28 A. M. Joshi, W. N. Delgass and K. T. Thomson, *J. Phys. Chem. B*, 2006, **110**, 23373–23387.
- 29 S. Pande, T. Jian, N. S. Khetrpal, L. S. Wang and X. C. Zeng, *J. Phys. Chem. C*, 2018, **122**, 6947–6954.



- 30 N. S. Khetrpal, T. Jian, G. V. Lopez, S. Pande, L. S. Wang and X. C. Zeng, *J. Phys. Chem. C*, 2017, **121**, 18234–18243.
- 31 N. S. Khetrpal, T. Jian, R. Pal, G. V. Lopez, S. Pande, L. S. Wang and X. C. Zeng, *Nanoscale*, 2016, **8**, 9805–9814.
- 32 Y. R. Zhao, Y. Qian, M. G. Zhang and Y. F. Hu, *Mol. Phys.*, 2015, **113**, 3598–3605.
- 33 P. Pyykko and N. Runeberg, *Angew. Chem., Int. Ed.*, 2002, **41**, 2174–2176.
- 34 X. Li, B. Kiran, J. Li, H. J. Zhai and L. S. Wang, *Angew. Chem., Int. Ed.*, 2002, **41**, 4786–4789.
- 35 H. J. Zhai, J. Li and L. S. Wang, *J. Chem. Phys.*, 2004, **121**, 8369–8374.
- 36 Y. Gao, S. Bulusu and X. C. Zeng, *J. Am. Chem. Soc.*, 2005, **127**, 15680–15681.
- 37 Y. Gao, S. Bulusu and X. C. Zeng, *ChemPhysChem*, 2006, **7**, 2275–2278.
- 38 S. Neukermans, E. Janssens, H. Tanaka, R. E. Silverans and P. Lievens, *Phys. Rev. Lett.*, 2003, **90**, 033401–033405.
- 39 L.-M. Wang, S. Bulusu, H.-J. Zhai, X.-C. Zeng and L.-S. Wang, *Angew. Chem., Int. Ed.*, 2007, **46**, 2915–2918.
- 40 L.-M. Wang, R. Pal, W. Huang, X. C. Zeng and L.-S. Wang, *J. Chem. Phys.*, 2009, **130**, 051101–051105.
- 41 L.-M. Wang, R. Pal, W. Huang, X. C. Zeng and L.-S. Wang, *J. Chem. Phys.*, 2010, **132**, 114306–114314.
- 42 R. Pal, L.-M. Wang, W. Huang, L.-S. Wang and X. C. Zeng, *J. Chem. Phys.*, 2011, **134**, 054306–054307.
- 43 B. Kiran, X. Li, H. J. Zhai, L. F. Cui and L. S. Wang, *Angew. Chem., Int. Ed.*, 2004, **43**, 2125–2129.
- 44 B. Kiran, X. Li, H.-J. Zhai and L.-S. Wang, *J. Chem. Phys.*, 2006, **125**, 133204.
- 45 P. Pyykko and Y. F. Zhao, *Chem. Phys. Lett.*, 1991, **177**, 103–106.
- 46 R. Pal, S. Bulusu and X. C. Zeng, *J. Comput. Methods Sci. Eng.*, 2007, **7**, 185–193.
- 47 C. Majumder, A. K. Kandalam and P. Jena, *Phys. Rev. B*, 2006, **74**, 205437.
- 48 C. Majumder, *Phys. Rev. B*, 2007, **75**, 235409.
- 49 M. Walter and H. Hakkinen, *Phys. Chem. Chem. Phys.*, 2006, **8**, 5407–5411.
- 50 Q. Sun, Q. Wang, G. Chen and P. Jena, *J. Chem. Phys.*, 2007, **127**, 214706.
- 51 L.-M. Wang, S. Bulusu, W. Huang, R. Pal, L.-S. Wang and X. C. Zeng, *J. Am. Chem. Soc.*, 2007, **129**, 15136.
- 52 M. Abe, T. Nakajima and K. Hirao, *J. Chem. Phys.*, 2002, **117**, 7960–7967.
- 53 R. Pal, L.-M. Wang, W. Huang, L.-S. Wang and X. C. Zeng, *J. Am. Chem. Soc.*, 2009, **131**, 3396–3404.
- 54 S. J. Lu, X. L. Xu, H. G. Xu and W. J. Zheng, *J. Chem. Phys.*, 2018, **148**, 244306.
- 55 Y. Zhao, F. Zhao and L. H. Kang, *J. Mol. Model.*, 2018, **24**, 61–68.
- 56 K. Joshi and S. Krishnamurthy, *J. Phys. Chem. C*, 2017, **121**, 17514–17522.
- 57 E. M. Dore and J. T. Lyon, *J. Cluster Sci.*, 2016, **27**, 1365–1381.
- 58 Y. R. Liu, T. Huang, Y. B. Gai, Y. Zhang, Y. J. Feng and W. Huang, *Sci. Rep.*, 2015, **5**, 17738–17745.
- 59 S. Gautam, N. Goel and K. Dharamvir, *RSC Adv.*, 2014, **4**, 13927–13933.
- 60 T. Quoc Tri, S.-J. Lu, L.-J. Zhao, X.-L. Xu, H.-G. Xu, T. Van Tan, J. Li and W. J. Zheng, *J. Phys. Chem. A*, 2018, **122**, 3374–3382.
- 61 S.-J. Lu, X.-L. Xu, G. Feng, H.-G. Xu and W.-J. Zheng, *J. Phys. Chem. C*, 2016, **120**, 25628–25637.
- 62 S. J. Lu, L. R. Hu, X. L. Xu, H. G. Xu, H. Chen and W. J. Zheng, *Phys. Chem. Chem. Phys.*, 2016, **18**, 20321–20329.
- 63 D. M. Deaven and K. M. Ho, *Phys. Rev. Lett.*, 1995, **75**, 288–291.
- 64 R. L. Johnston, *Dalton Trans.*, 2003, 4193–4207.
- 65 R. H. Byrd, P. H. Lu, J. Nocedal and C. Y. Zhu, *SIAM J. Sci. Comput.*, 1995, **16**, 1190–1208.
- 66 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865–3868.
- 67 B. Delley, *J. Chem. Phys.*, 1990, **92**, 508–517.
- 68 R. A. Kendall, E. Apra, D. E. Bernholdt, E. J. Bylaska, M. Dupuis, G. I. Fann, R. J. Harrison, J. L. Ju, J. A. Nichols, J. Nieplocha, T. P. Straatsma, T. L. Windus and A. T. Wong, *Comput. Phys. Commun.*, 2000, **128**, 260–283.
- 69 R. Ferrando, J. Jellinek and R. L. Johnston, *Chem. Rev.*, 2008, **108**, 845–910.

