



Cite this: *RSC Adv.*, 2019, 9, 35039

# p-MoS<sub>2</sub>/n-InSe van der Waals heterojunctions and their applications in all-2D optoelectronic devices†

Pan Li,<sup>ab</sup> Kai Yuan,<sup>a</sup> Der-Yuh Lin,<sup>c</sup> Tingting Wang,<sup>ab</sup> Wanying Du,<sup>a</sup> Zhongming Wei,<sup>id d</sup> Kenji Watanabe,<sup>id e</sup> Takashi Taniguchi,<sup>e</sup> Yu Ye<sup>id \*abf</sup> and Lun Dai<sup>id \*ab</sup>

A library of two-dimensional (2D) semiconductors with different band gaps offers the construction of van der Waals (vdWs) heterostructures with different band alignments, providing a new platform for developing high-performance optoelectronic devices. Here, we demonstrate all-2D optoelectronic devices based on type-II p-MoS<sub>2</sub>/n-InSe vdWs heterojunctions operating at the near infrared (NIR) wavelength range. The p–n heterojunction diode exhibits a rectification ratio of  $\sim 10^2$  at  $V_{ds} = \pm 2$  V and a turn-on voltage of  $\sim 0.8$  V. Under a forward bias exceeding the turn-on voltage and a proper positive back-gate voltage, the all-2D vdWs heterojunction diode exhibits an electroluminescence with an emission peak centered at  $\sim 1020$  nm. Besides, this p-MoS<sub>2</sub>/n-InSe heterojunction shows a photoresponse at zero external bias, indicating that it can serve as a photodiode working without an external power supply. The as-demonstrated all-2D vdWs heterojunction which can work as both a light-emitting diode and a self-powered photodetector may find applications in flexible wear, display, and optical communication fields, etc.

Received 24th August 2019  
 Accepted 20th October 2019

DOI: 10.1039/c9ra06667e

[rsc.li/rsc-advances](http://rsc.li/rsc-advances)

Optoelectronic devices, including light-emitting diodes (LEDs), lasers, solar cells, and photodetectors, have been extensively studied, because of their potential applications in illumination, communication, energy, and healthcare. With the growing demand for advanced functional optoelectronic devices, such as flexible wear and display, faster communications between devices on microchips, etc., next-generation optoelectronic devices require new optoelectronic materials with characteristics superior to those currently in use. Candidate materials must be flexible for wearable devices, transparent for interactive displays, and efficient for light-emitting and/or solar cells. The emergence of two-dimensional (2D) materials, with atomic layers combined by van der Waals (vdWs) force, has brought a new era to the development of optoelectronic devices, because of their unique structural and physical properties, such as strong light–matter interactions, flexibility, vdWs assembly,

etc.<sup>1–6</sup> In addition, the extensive library of 2D materials offers the formation of vdWs heterostructures with a broad range of band alignments, and resultant interfacial physical properties, providing a new platform for developing high-performance optoelectronic devices.<sup>7–12</sup> Specifically, heterojunctions with a type-II band alignment, where the valence band maximum and the conduction band minimum reside in two separate materials, have enabled efficient electron–hole separation at the interface. Therefore type-II 2D heterojunctions have been used in high-performance 2D optoelectronic devices such as photodetectors, photovoltaic devices, and LEDs.<sup>6,13,14</sup> The near-infrared (NIR) optoelectronic devices can be used in a wide range of applications, including night vision, optical communication, and computing. However, due to the limit of 2D materials with a band gap in the NIR range, NIR optoelectronic devices based on type-II vdWs heterojunction are rarely studied.

InSe, a group III–VI layered semiconductor with a direct band-gap of 1.25 eV for few-layer one ( $>10$  L),<sup>15</sup> has gained various interests due to its unique optical, electronic, and mechanical properties, targeting the applications in memory devices, optical sensors, and thermoelectric implements.<sup>16–19</sup> As a direct band-gap semiconductor, which has an efficient electron–hole radiative recombination rate, the few-layer InSe is also suitable for making NIR optoelectronic devices. Recently, high-performance broadband few-layer InSe photodetectors from the visible to NIR region have been realized.<sup>17–19</sup> Few-layer MoS<sub>2</sub> is with an indirect band-gap of  $\sim 1.28$  eV, which can form a type-II band alignment with few-layer InSe. Under a forward bias, the injected carriers recombine to emit photons in the InSe region that has a narrower direct band gap.

<sup>a</sup>State Key Laboratory for Artificial Microstructure & Mesoscopic Physics, School of Physics, Peking University, Beijing 100871, China. E-mail: [ye\\_yu@pku.edu.cn](mailto:ye_yu@pku.edu.cn); [lundai@pku.edu.cn](mailto:lundai@pku.edu.cn)

<sup>b</sup>Collaborative Innovation Center of Quantum Matter, Beijing 100871, China

<sup>c</sup>Department of Electronics Engineering, National Changhua University of Education, Changhua 50007, Taiwan

<sup>d</sup>State Key Laboratory of Superlattices and Microstructures, Institute of Semiconductors, Chinese Academy of Sciences, Beijing 100083, China

<sup>e</sup>National Institute for Materials Science, 1-1 Namiki, Tsukuba 305-0044, Japan

<sup>f</sup>Frontiers Science Center for Nano-optoelectronics, Peking University, Beijing 100871, China

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c9ra06667e



Here, we present type-II vdWs heterojunctions based on the n-type few-layer InSe and p-type MoS<sub>2</sub> nanoflakes (Fig. 1). The heterostructure was stacked and encapsulated by top and bottom hexagonal boron nitride (hBN) nanoflakes using a spatially controlled dry transfer method,<sup>20</sup> which guarantees the clean heterojunction interfaces and prevents contaminations from any polymers or solutions. The p-MoS<sub>2</sub>/n-InSe heterojunctions exhibit typical diode characteristic with a rectification ratio of  $\sim 10^2$  at  $V_{ds} = \pm 2$  V, and a low turn-on voltage of  $\sim 0.8$  V. Under a forward bias exceeding the turn-on voltage, the vdWs heterojunction shows strong NIR electroluminescence (EL) centered at  $\sim 1020$  nm, corresponding to the band-edge emission of the InSe. Under the simulated solar illumination, the heterojunction also exhibits typical photovoltaic behavior. Photocurrent mapping at zero external bias indicates that the photoresponse of the diode is from the heterojunction region. Our work provides a new feasible structure of all-2D NIR optoelectronic devices for flexible and microchip communication applications.

## Results and discussion

In this work, the p-type MoS<sub>2</sub> and n-type InSe bulk crystals were grown *via* a chemical vapor transport (CVT) method. To p-type doping the MoS<sub>2</sub>, we used Nb as the dopant during the CVT process. To make a high-performance p-MoS<sub>2</sub>/n-InSe vdWs heterojunction, we first exfoliated the n-type InSe nanoflakes and p-type MoS<sub>2</sub> nanoflakes on Si/SiO<sub>2</sub> substrates in a glove box. Then, the heterostructure was stacked and encapsulated by top and bottom hBN nanoflakes using a spatially controlled dry transfer method,<sup>20</sup> which guarantees the clean heterojunction interfaces and prevents contaminations from any polymers or solutions. After that, the top layer hBN was selectively etched and the Ti/Au (10/70 nm) contact electrodes were fabricated on both the InSe and MoS<sub>2</sub> (Fig. 1) (see Method for details). hBN encapsulation avoids the degradation of InSe. To characterize the electrical properties of both p-MoS<sub>2</sub> and n-InSe, as well as the contact behaviors of the Ti/Au electrodes, we fabricated field-effect transistors (FETs) of both p-MoS<sub>2</sub> and n-InSe nanoflakes, with thickness of  $\sim 15$  nm and  $\sim 15$  nm, respectively (insets of Fig. 2b and d). The  $I$ - $V$  curves of the few-layer InSe FET above turn-on back gate voltage ( $V_{gs}$ ) show linear behavior at low source-drain bias ( $V_{ds}$ ) (Fig. 2a), confirming the ohmic contact

characteristic between the low-work function Ti/Au ( $\Phi_{Ti} = -4.33$  eV) electrode and the n-type InSe. The source-drain current ( $I_{ds}$ ) saturates at a larger source-drain voltage (0.45 V). The channel current increases with increasing  $V_{gs}$ , indicating the n-type nature of the InSe. From the transfer curve (Fig. 2b), we can see that the threshold voltage and on/off ratio are about 22 V and  $>10^7$  (see ESI S1†), respectively. The electron mobility ( $\mu_e$ ) and concentration ( $n_{2D}$ ) are estimated to be about  $285 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $\sim 10^{12} \text{ cm}^{-2}$ , respectively, when the channel was turned on (see ESI S1†). Besides, the on-current of few-layer InSe increase by about two to three orders of magnitude compared with non-encapsulation ones (see ESI S3†). The  $I$ - $V$  curves of the few-layer MoS<sub>2</sub> FET show near-linear behaviors at different  $V_{gs}$  (Fig. 2c), because there exists kind of Schottky contact barrier at the interface of Ti/Au and the p-type MoS<sub>2</sub>. The channel current decreases with increasing  $V_{gs}$ , indicating the p-type nature of the MoS<sub>2</sub>. The channel cannot be fully depleted even when  $V_{gs}$  is as high as +80 V (Fig. 2d), indicating the heavily p-type doping nature of the MoS<sub>2</sub>. The hole mobility ( $\mu_h$ ) and concentration ( $p_{2D}$ ) is estimated to be about  $1.7 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $\sim 3.03 \times 10^{13} \text{ cm}^{-2}$ , respectively (see ESI S2†).

LED is a semiconductor device that converts electrical energy into light energy. The light emission mechanism from the p-MoS<sub>2</sub>/n-InSe LED can be qualitatively understood by studying the energy band diagram of the vdWs heterojunction, in which the Fermi levels of n-InSe and p-MoS<sub>2</sub> coincide with each other in equilibrium state (Fig. 3a). Based on the depletion approximation, depletion layers with abrupt boundaries cause a built-in potential, while the semiconductors away from the depletion region boundaries are assumed to be electrically neutral. The determination of MoS<sub>2</sub>/InSe band alignment is based on the reported band gaps, affinities, and calculated work functions of MoS<sub>2</sub> and InSe nanoflakes (see ESI S4†). Due to the difference in their band gaps ( $E_g$ ) and electron affinities ( $\chi$ ), large conduction-band barrier ( $\Delta E_c \sim 1.16$  eV) and valence-band barrier ( $\Delta E_v \sim 1.10$  eV) are formed between MoS<sub>2</sub> and InSe. Because the conduction band barrier is larger, with forward bias increasing, the valence band barrier will disappear first (Fig. 3b). At this stage, large amount holes will flow from the p-MoS<sub>2</sub> to the n-InSe and recombine with the electrons at the InSe region. Therefore, the EL spectrum is determined by the band-edge emission of the InSe. Due to the degenerate p-doping nature of the MoS<sub>2</sub>, the EL intensity is limited by the electron concentration in InSe. The  $I$ - $V$  curves of the vdWs heterojunction diode (Fig. 3d) show a clear rectifying behavior when the voltage changes from  $-2$  V to  $2$  V. The turn-on voltage is  $\sim 0.8$  V. Due to the modest electron concentration of the InSe at zero gate voltage, the turn-on current of the diode is limited. When the back gate voltage increases from  $0$  V to  $20$  V, the on-current of the InSe/MoS<sub>2</sub> heterojunction at  $V_{ds} = 2$  V increases from  $2.78 \mu\text{A}$  to  $19.56 \mu\text{A}$ , while the reverse current at  $V_{ds} = -2$  V increases from  $4.71 \text{ nA}$  to  $78.23 \text{ nA}$ , resulting the rectification ratio at  $\pm 2$  V from  $5.9 \times 10^2$  to  $2.5 \times 10^2$ . The  $I$ - $V$  relationship of a p-n junction can be expressed as  $I = I_0 \left[ \exp\left(\frac{qV}{nkT}\right) - 1 \right]$ , where  $I_0$  is the reverse saturation current,  $k$  is the Boltzmann constant,  $T$  is the temperature,  $q$  is the

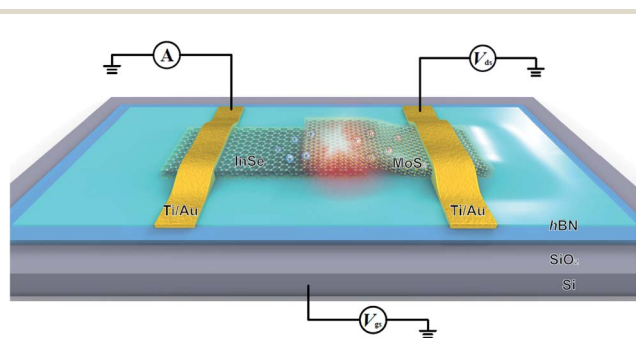
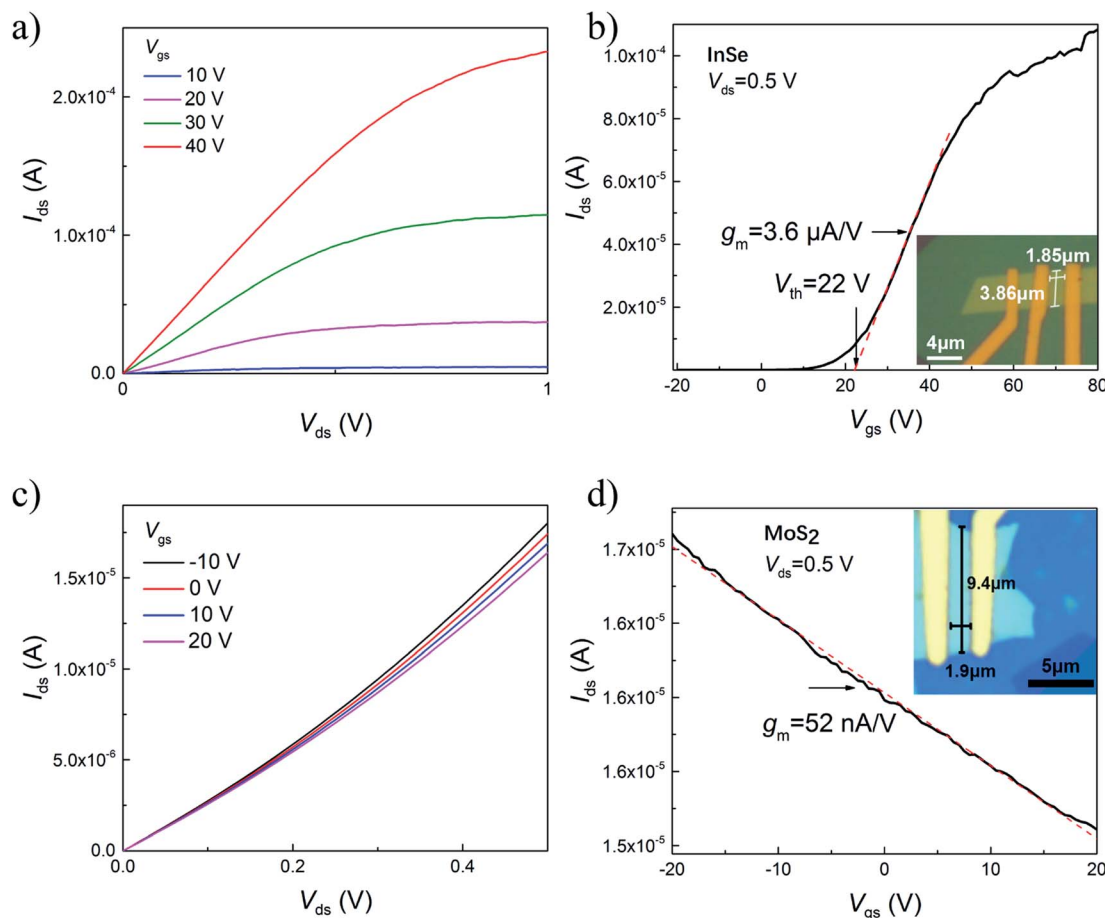


Fig. 1 Schematic illustration of the p-MoS<sub>2</sub>/n-InSe vdWs heterostructure.





**Fig. 2** Electrical characterizations of the InSe and MoS<sub>2</sub> nanoflakes. (a) The  $I_{ds}$  versus  $V_{ds}$  curves for the InSe FET under different  $V_{gs}$ . (b) The  $I_{ds}$  versus  $V_{gs}$  curve of the InSe FET at  $V_{ds} = 0.5$  V. The red straight dash line shows the fitting result with channel transconductance ( $g_m$ ) of about  $3.6 \mu\text{A V}^{-1}$  and turn-on threshold voltage ( $V_{th}$ ) of 22 V. Inset: the optical image of the measured InSe FET with a channel length and width of  $1.85 \mu\text{m}$  and  $3.86 \mu\text{m}$ , respectively. (c) The  $I_{ds}$  versus  $V_{ds}$  curves for the MoS<sub>2</sub> FET under different  $V_{gs}$ . (d) The  $I_{ds}$  versus  $V_{gs}$  curve of the MoS<sub>2</sub> FET at  $V_{ds} = 0.5$  V. The red straight dash line shows the fitting result with  $g_m$  of about  $52 \text{ nA V}^{-1}$ . Inset is the optical image of the measured MoS<sub>2</sub> FET with a channel length and width of  $1.9 \mu\text{m}$  and  $9.4 \mu\text{m}$ , respectively.

electron charge, and  $n$  is the ideality factor.<sup>21</sup> For a p–n junction, the ideality factor  $n$  equals 2 when the recombination current dominates, and  $n$  equals 1 when the diffusion current dominates.<sup>21</sup> By fitting the measured  $I$ – $V$  curve at  $V_{gs} = 20$  V with the above equation, we obtain  $n = 2.58$ . This ideality factor larger than 2 indicates that there exists recombination from other transport channels, which may be caused by the defects and impurities arising from the heavy doping.<sup>22</sup> When the back-gate voltage is higher than 20 V and the source–drain bias is above the turn-on voltage, we can observe strong EL from the vdWs diode. Fig. 3e shows the room-temperature EL spectra under a forward bias of 9 V at various back-gate voltages. Under a positive back gate voltage, more electrons are injected from InSe and accumulated at the heterojunction interface to recombine with the injected holes, resulting in an enhanced EL emission. The EL spectra are dominated by an emission peaked around  $\sim 1020$  nm, coincident with the band-edge emission of the InSe nanoflake (see ESI S4†).

Photodetectors, which convert optical signals to electrical signals, are basic components in imaging techniques, optical

communications, and future intra-chip optical interconnects.<sup>23</sup> Herein, the p-MoS<sub>2</sub>/n-InSe heterojunction can serve as a photodiode, which can work without an external power supply. Upon light illumination, the photo-generated holes and electrons are driven towards the p-MoS<sub>2</sub> and n-InSe, respectively, by the built-in field, resulting in a short-circuit current or a voltage drop on the external resistor. Under a standard solar illumination with a power density of  $100 \text{ mW cm}^{-2}$ , the MoS<sub>2</sub>/InSe heterojunction exhibits good photovoltaic behavior (Fig. 4a), with a typical open-circuit voltage ( $V_{oc}$ ), short-circuit current density ( $J_{sc}$ ), fill factor (FF), and energy conversion efficiency ( $\eta = FF J_{sc} V_{oc} / P_{in}$ , where  $P_{in}$  is the illumination light power density) to be about 0.2 V,  $1.17 \text{ mA cm}^{-2}$ , 50.4%, and 0.59%, respectively. The photosensitivity, defined as  $(I_{photo} - I_{dark}) / I_{dark}$ , where  $I_{photo}$  is the photocurrent and  $I_{dark}$  is the dark current, is measured to be about  $4.42 \times 10^3$  at zero external bias. The photocurrent mapping at zero external bias under 633 nm laser excitation shows a uniform photoresponse at the heterostructure region (Fig. 4c), indicating the high-quality vdWs interface at the vertical stacked heterojunction.



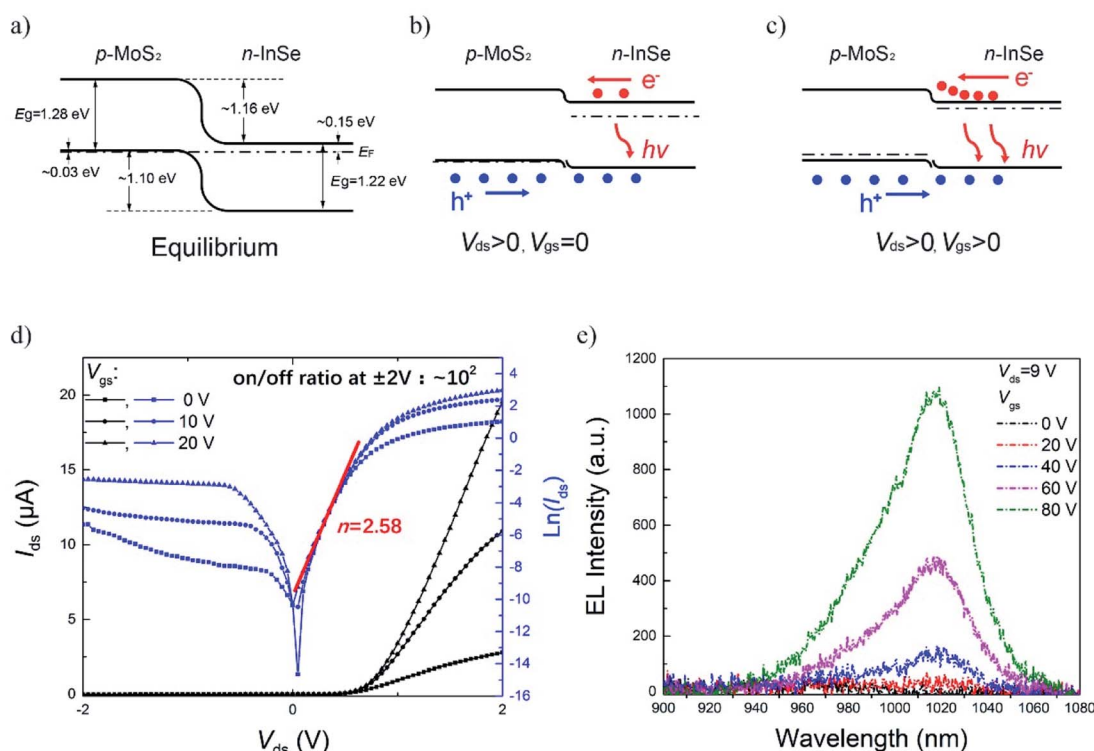


Fig. 3 (a, b and c) The energy band diagrams of the p-MoS<sub>2</sub>/n-InSe heterojunction diode in equilibrium, under a forward bias and zero back gate voltage, and under a forward bias and positive back gate voltage, respectively. (d) Room-temperature  $I$ - $V$  characteristics of the heterojunction diode depicted in the linear scale (black) and semi-log scale (blue), respectively. The red straight line shows the fitting result with an ideality factor  $n$  of about 2.58. (e) Room-temperature EL spectra under a forward source-drain bias of 9 V at various back-gate voltages. The EL intensity increases with the positive back-gate voltage, due to the increase of injected electrons from InSe.

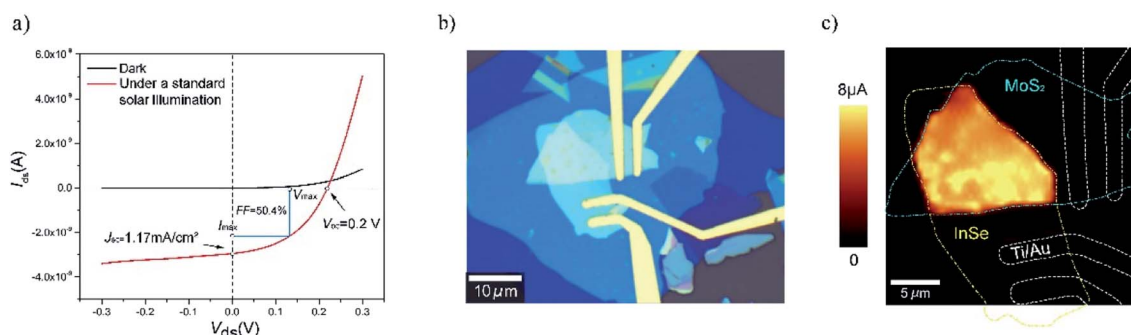


Fig. 4 (a) The  $I$ - $V$  curves of this p-n vdWs heterojunction diode in dark and under a standard solar illumination with a power density of 100 mW cm<sup>-2</sup>. (b) Optical image of the fabricated p-n vdWs heterojunction diode. (c) Photocurrent mapping of the p-n vdWs heterojunction under the excitation of a 633 nm laser with a power of 0.5 mW.

## Conclusion

In conclusion, we have fabricated all-2D vdWs hBN/p-MoS<sub>2</sub>/n-InSe/hBN vertical stacking heterostructure by the dry transfer method in a glove box. The type-II p-MoS<sub>2</sub>/n-InSe heterojunction diode exhibits a good rectification behavior and shows strong light emission at wavelength of  $\sim 1020$  nm under a forward bias exceeding the turn-on voltage and a positive back-gate voltage higher than 20 V. The EL intensity can be tuned by the back-gate voltage. Under simulated solar

illumination, the vdWs p-n heterojunction exhibits a good photovoltaic behavior, with a typical  $V_{oc}$ ,  $J_{sc}$ , FF, and energy conversion of about 0.2 V, 1.17 mA cm<sup>-2</sup>, 50.4%, and 0.59%, respectively. The photosensitivity at zero external bias is about  $4.42 \times 10^3$ . The photocurrent mapping shows a uniform photoresponse at the heterostructure region. The as-demonstrated all-2D vdWs heterojunction, which can work as both a light-emitting diode and a self-powered photodetector, may find applications in flexible wear, display, and optical communication fields, *etc.*





## Methods

### Synthesis of p-type MoS<sub>2</sub> bulk

To p-type dope the MoS<sub>2</sub>, we used Nb as the dopant during the CVT process. A horizontal three-zone furnace was utilized where the high (low) temperature zone was used as the reaction (growth) zone. Prior to the crystal growth, a quartz tube containing the iodine transport agent and the elements required for the doped crystals was evacuated to below 10<sup>−6</sup> Torr and securely sealed. The purity of the source materials was Mo 99.99%, S 99.99%, and Nb 99.99%, with a molar ratio of Mo : S : Nb equal to 1 : 2 : 0.005 for a nominal 0.5% Nb doping. The sealed quartz tube was then inserted into the three-zone furnace. In the beginning, the temperature of the growth zone was set to 1000 °C for 12 h to remove any possible residues. The process of the CVT was initiated with a temperature gradient varying from 1050 °C in the reaction zone to 935 °C in the growth zone over a distance of ~30 cm, and the system was held in this condition over 500 h in order to provide sufficient diffusion and equilibration for the single-crystal synthesis.<sup>24</sup>

### Dry transfer method for assembly the hBN encapsulated p-MoS<sub>2</sub>/n-InSe heterostructures

The hBN/MoS<sub>2</sub>/InSe/hBN vertical stacked heterostructure was fabricated as follows: first, the hBN, MoS<sub>2</sub>, and InSe nanoflakes with appropriate thicknesses were exfoliated on Si/SiO<sub>2</sub> substrates (the SiO<sub>2</sub> layer was about 285 nm) (Fig. S3a–c in ESI S5†). Then the exfoliated hBN and MoS<sub>2</sub> nanoflakes were picked up in sequence by a polydimethylsiloxane (PDMS) stamp covered by a layer of polypropylene carbonate (PPC) (Fig. S4a, b in ESI S6†). To avoid the degradation of InSe, the InSe related exfoliation and transfer processes were done in a glove box (O<sub>2</sub>, H<sub>2</sub>O < 0.1 ppm). The selected InSe nanoflake was picked up by the previously picked hBN/MoS<sub>2</sub> heterostructure on PPC (Fig. S4c in ESI S6†) and transferred onto another hBN nanoflake on Si/SiO<sub>2</sub> substrate (Fig. S4d in ESI S6†). The hBN encapsulated heterostructure was detached from PDMS/PPC at 60 °C and left on the device substrate (Fig. S4e in ESI S6†).

### Fabrication of the electrodes

The contact areas were patterned by electron-beam lithography, and the exposed top hBN layer was subsequently etched by reactive-ion etching using trifluoromethane and oxygen with flow rates of 20 sccm and 2 sccm for 2 min, respectively. Electron beam evaporation, and lift-off processes were used to fabricate the Ti/Au (10/70 nm) electrodes on both the InSe and MoS<sub>2</sub> nanoflakes.

### Electrical measurements

All the electrical measurements were conducted with a semiconductor characterization system (Keithley 4200) under a probe station.

### PL and EL spectra characterizations

PL and EL spectra were recorded by a micro-zone confocal Raman spectroscopy (HORIBA Jobin Yvon, LabRam HR 800) equipped with a color charge-coupled device (CCD) at room temperature. In the PL measurement, a 633 nm laser was used as the excitation and focused on the sample by a 50× objective (Olympus LMPlanFI IR). The EL and PL from devices were collected by the same objective.

### Photoresponse measurements

The photovoltaic measurements were conducted with a semiconductor characterization system (Keithley 4200) under a standard simulated solar illumination with a power density of 100 mW cm<sup>−2</sup>. The photocurrent mapping was recorded under the 633 nm laser illumination with a power of 0.5 mW using a micro-zone confocal Raman spectroscopy (WITec alpha 300R). The photocurrent signal was collected and converted to a voltage signal by an amplifier with a magnification of 10<sup>6</sup>, and recorded simultaneously with the scanning position information during the mapping process.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

This work was supported by the National Natural Science Foundation of China (No. 61521004, 61874003, and 61875001), and the National Key R&D Program of China (Grant No. 2018YFA0306900). D.-Y. L. acknowledges the financial support from the Ministry of Science and Technology of Taiwan, Republic of China under contract No. MOST 108-2221-E-018-010. K. W. and T. T. acknowledge support from the Elemental Strategy Initiative conducted by the MEXT, Japan and the CREST (JPMJCR15F3), JST.

## References

- 1 K. S. Novoselov, V. Fal, L. Colombo, P. Gellert, M. Schwab and K. Kim, *Nature*, 2012, **490**, 192.
- 2 N. O. Weiss, H. Zhou, L. Liao, Y. Liu, S. Jiang, Y. Huang and X. Duan, *Adv. Mater.*, 2012, **24**, 5782–5825.
- 3 F. Bonaccorso, Z. Sun, T. Hasan and A. Ferrari, *Nat. Photonics*, 2010, **4**, 611.
- 4 F. Schwierz, *Nat. Nanotechnol.*, 2010, **5**, 487.
- 5 P. Avouris, *Nano Lett.*, 2010, **10**, 4285–4294.
- 6 L. Britnell, R. Ribeiro, A. Eckmann, R. Jalil, B. Belle, A. Mishchenko, Y.-J. Kim, R. Gorbachev, T. Georgiou and S. Morozov, *Science*, 2013, **340**, 1311–1314.
- 7 H. Fang, S. Chuang, T. C. Chang, K. Takei, T. Takahashi and A. Javey, *Nano Lett.*, 2012, **12**, 3788–3792.
- 8 D. J. Groenendijk, M. Buscema, G. A. Steele, S. Michaelis de Vasconcellos, R. Bratschitsch, H. S. van der Zant and A. Castellanos-Gomez, *Nano Lett.*, 2014, **14**, 5846–5852.



- 9 J. S. Ross, P. Klement, A. M. Jones, N. J. Ghimire, J. Yan, D. Mandrus, T. Taniguchi, K. Watanabe, K. Kitamura and W. Yao, *Nat. Nanotechnol.*, 2014, **9**, 268.
- 10 Y. Zhang, T. Oka, R. Suzuki, J. Ye and Y. Iwasa, *Science*, 2014, **344**, 725–728.
- 11 Y. Zhou, Y. Nie, Y. Liu, K. Yan, J. Hong, C. Jin, Y. Zhou, J. Yin, Z. Liu and H. Peng, *ACS Nano*, 2014, **8**, 1485–1490.
- 12 J. Pu, Y. Yomogida, K.-K. Liu, L.-J. Li, Y. Iwasa and T. Takenobu, *Nano Lett.*, 2012, **12**, 4013–4017.
- 13 X. Hong, J. Kim, S.-F. Shi, Y. Zhang, C. Jin, Y. Sun, S. Tongay, J. Wu, Y. Zhang and F. Wang, *Nat. Nanotechnol.*, 2014, **9**, 682.
- 14 Y. Gong, J. Lin, X. Wang, G. Shi, S. Lei, Z. Lin, X. Zou, G. Ye, R. Vajtai and B. I. Yakobson, *Nat. Mater.*, 2014, **13**, 1135.
- 15 G. W. Mudd, S. A. Svatek, T. Ren, A. Patanè, O. Makarovskiy, L. Eaves, P. H. Beton, Z. D. Kovalyuk, G. V. Lashkarev and Z. R. Kudrynskiy, *Adv. Mater.*, 2013, **25**, 5714–5718.
- 16 D. A. Bandurin, A. V. Tyurnina, L. Y. Gelang, A. Mishchenko, V. Zólyomi, S. V. Morozov, R. K. Kumar, R. V. Gorbachev, Z. R. Kudrynskiy and S. Pezzini, *Nat. Nanotechnol.*, 2017, **12**, 223.
- 17 Z. Chen, J. Biscaras and A. Shukla, *Nanoscale*, 2015, **7**, 5981–5986.
- 18 G. W. Mudd, S. A. Svatek, L. Hague, O. Makarovskiy, Z. R. Kudrynskiy, C. J. Mellor, P. H. Beton, L. Eaves, K. S. Novoselov and Z. D. Kovalyuk, *Adv. Mater.*, 2015, **27**, 3760–3766.
- 19 S. R. Tamalampudi, Y.-Y. Lu, R. U. Kumar, R. Sankar, C.-D. Liao, K. B. Moorthy, C.-H. Cheng, F. C. Chou and Y.-T. Chen, *Nano Lett.*, 2014, **14**, 2800–2806.
- 20 L. Wang, I. Meric, P. Huang, Q. Gao, Y. Gao, H. Tran, T. Taniguchi, K. Watanabe, L. Campos and D. Muller, *Science*, 2013, **342**, 614–617.
- 21 S. Sze and C. A. Mead, *Analog VLSI and Neural Systems*, Addison-Wesley, Reading, MA, 1989, p. 36.
- 22 P. Gehring, R. Urcuyo, D. L. Duong, M. Burghard and K. Kern, *Appl. Phys. Lett.*, 2015, **106**, 233110.
- 23 Z. Gao, W. Jin, Y. Zhou, Y. Dai, B. Yu, C. Liu, W. Xu, Y. Li, H. Peng and Z. Liu, *Nanoscale*, 2013, **5**, 5576–5581.
- 24 J. Suh, T.-E. Park, D.-Y. Lin, D. Fu, J. Park, H. J. Jung, Y. Chen, C. Ko, C. Jang and Y. Sun, *Nano Lett.*, 2014, **14**, 6976–6982.

