# **Nanoscale** Advances

# PAPER

Cite this: Nanoscale Adv., 2020, 2, 2352

Received 21st February 2020 Accepted 11th April 2020

DOI: 10.1039/d0na00147c

rsc.li/nanoscale-advances

## Introduction

Two-dimensional transition metal dichalcogenides (TMDs) have attracted great interest due to their excellent electrical and optical properties.<sup>1</sup>–<sup>5</sup> The semiconducting nature of TMDs allows overcoming the shortcomings of zero-bandgap graphene. In fact, semiconducting TMDs are also attractive as a potential complement to graphene for constructing digital circuits.<sup>6-8</sup>

One of the main advantages of two-dimensional (2D) semiconducting materials comes from the indirect-to-direct bandgap transition at the single layer limit.<sup>9,10</sup> This particular effect leads to a strong interaction with light, which can pave the way for envisioning the next generation of optoelectronic devices<sup>11</sup>–<sup>14</sup> and integration with existing silicon photonic technology.15,16 Furthermore, due to the breaking of inversion symmetry and strong spin–orbital coupling in monolayer TMDs, the electronic states of the two valleys have different chiralities.<sup>17</sup> This motivates significant interest in manipulating the spin and valley degrees of freedom for the development of novel spintronic and valleytronic devices.<sup>18-21</sup>

The most versatile and efficient technique for the synthesis of TMDs is chemical vapor deposition  $(CVD)$ .<sup>2,22</sup> Several studies have demonstrated the large area CVD synthesis of different

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# Influence of organic promoter gradient on the MoS<sub>2</sub> growth dynamics†

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Chemical vapor deposition has been demonstrated to be the most efficient, versatile and reliable technique for the synthesis of monolayers of transition metal dichalcogenides. The use of organic promoters during the growth process was a turning point in order to increase the monolayer lateral size or to obtain complete coverage of the growth substrate. In this work we clarify the influence of the promoter gradient on the growth dynamics of MoS<sub>2</sub>. In particular, we place a sacrificial substrate covered with a promoter (a low sublimation-temperature perylene-based compound) downstream with respect to the growth substrate in order to maximize its gradient on the growth substrate through upstream diffusion. We demonstrate that the morphology and the number of layers of  $MOS<sub>2</sub>$  are drastically affected by the distance of the growth substrate from the promoter sacrificial substrate. The farthermost area from the promoter substrate presents micrometric MoS<sub>2</sub> triangular monolayers and large low hierarchy dendritic multi-layer structures. On the contrary the closest area reveals an almost continuous polycrystalline  $MoS<sub>2</sub>$ monolayer, with bilayer terraces, with a lateral dimension up to hundreds of micrometers. **PAPER**<br> **PAPER**<br> **IMPLINENCE Of Organic promoter gradient on the state of the** 

TMDs on different insulating substrates, such as silicon dioxide or sapphire.<sup>23-27</sup> Recently, a CVD-based approach compatible with "back-end-of-line" for integration with complementary metal-oxide-semiconductor (CMOS) technology has been reported in the case of two-dimensional tungsten based TMDs following the strict restrictions for technological compatibility.<sup>28,29</sup>

In the last few years, the CVD approach with gaseous metalorganic precursors, such as molybdenum or tungsten hexacarbonyl (MoCO<sub>6</sub> or WCO<sub>6</sub>) for the metal and  $H_2S$  for the sulfur, is gaining increasing interest for the possible wafer-scale synthesis of different TMDs.<sup>30,31</sup>

An efficient approach to increase the grain size and crystalline quality with the MOCVD technique is the pre-exposure of the growth substrate to alkali metal halides (KI or NaCl) before the growth stage, in order to suppress nucleation and to increase the lateral size of single-crystal domains from a few hundred nanometers to tens of microns.<sup>32</sup>

An additional novel approach for the large-area growth of TMDs is to utilize water-soluble metal precursors.<sup>33</sup> In this approach a solution with a metal precursor and a seeding salt, such as NaOH or NaCl, is spin-coated directly on the growth substrate, allowing a wafer scale synthesis.<sup>34</sup>

The main drawback of the use of salts in the growth process is the formation of a sodium/potassium metal oxide layer below the TMDC monolayer, as a byproduct of the growth reaction.<sup>33</sup> The combination of organic promoters, namely metalloporphyrins, and water-soluble metal precursors has been employed in order to achieve large-area synthesis with concurrent doping of  $MoS<sub>2</sub>$  monolayers.<sup>35</sup>



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<sup>†</sup> Electronic supplementary information (ESI) available. See DOI: 10.1039/d0na00147c

In the case of CVD with solid precursors, such as  $MoO<sub>3</sub>$  and  $WO<sub>3</sub>$  powders, an efficient approach to increase the grain size and substrate coverage is the employment of organic promoters.<sup>36</sup>

In this perspective a turning point was the work of Ling et al.<sup>37</sup> in which the authors were able to obtain a continuous film of monolayer  $MoS<sub>2</sub>$ , by introducing an organic promoter during the CVD process. Another advantage of seeding promoters is the possible decrease of the growth temperature without affecting the quality of two-dimensional materials. Among them, the most widely employed promoters are perylene based compounds, namely perylene-3,4,9,10-tetracarboxylic dianhydride (PTCDA) and perylene-3,4,9,10-tertracarboxylic acid tetrapotassium salt (PTAS).<sup>25,37,38</sup> PTCDA has the advantage of a lower sublimation temperature (450  $^{\circ}$ C) compared to PTAS  $(600 °C)$ ,<sup>37</sup> leading to less amorphous carbon or promoter residuals on growth substrates.

Organic promoters have also been used in the synthesis of large-area van der Waals heterostructures, in particular  $MoS<sub>2</sub>/$  $WS_2$  and  $MoSe_2/WSe_2$  lateral structures<sup>39</sup> or the synthesis of  $MoS<sub>2</sub>$  on top of other 2-dimensional materials,<sup>37</sup> like graphene or h-BN.

In addition, it has been demonstrated that organic promoters can affect the optical and electrical properties of monolayer  $MoS<sub>2</sub>$ . For instance, the PTAS promoter allows the growth of a large-area monolayer with excellent optical quality and high electron mobility. Meanwhile the employment of crystal violet as a growth promoter results in the synthesis of highly defective MoS<sub>2</sub> monolayers.<sup>40</sup>

In this work, we studied the growth dynamics of  $MoS<sub>2</sub>$  twodimensional crystals in relation with the gradient of the organic promoter (PTCDA), that sublimates during the heating ramp and diffuses upstream in the gas phase. The promoter sacrificial substrate is placed downstream with respect to the growth substrate: during the heating ramp the PTCDA backdiffuses in the gas phase towards the growth substrate, resulting in an inhomogeneous distribution. This PTCDA inhomogeneity, in conjunction with flake growth-promoting properties, causes two distinct  $MoS<sub>2</sub>$  growth dynamics in relation with the distance from the promoter sacrificial substrate. In fact, the farthermost area from the promoter substrate presents micrometric  $MoS<sub>2</sub>$  triangular monolayers and large dendritic multilayer structures and meanwhile the closest area reveals an almost continuous polycrystalline monolayer  $MoS<sub>2</sub>$ , affected by bilayer terraces.

## Experimental section

Scanning electron microscopy (SEM) was performed to study the crystal morphology as a function of the distance from the PTCDA substrate. SEM analysis is carried out using a backscattered electron detector or a secondary electron detector in a Zeiss Merlin FEG-SEM microscope, with an accelerating voltage of 5 kV and an electron current of 12 pA, in order to minimize the damage due to electron beam irradiation.

Scanning Raman and photoluminescence (PL) spectroscopies are carried out with a Renishaw InVia system, equipped with a confocal microscope, a 532 nm excitation laser and a 1800 line per mm grating (spectral resolution 2  $\rm cm^{-1}$ ). All the analyses are performed with the following parameters: excitation laser power 500  $\mu$ W, acquisition time for each spectrum 4s, pixel size of 1  $\mu$ m  $\times$  1  $\mu$ m, and an 100 $\times$  objective (NA = 0.85).

 $MoS<sub>2</sub>$  flakes have been transferred from the growth samples onto standard copper grids for TEM investigations: the surface of the  $MoS<sub>2</sub>$ -covered  $SiO<sub>2</sub>$  substrate is coated with poly-methyl methacrylate (PMMA) and then the  $SiO<sub>2</sub>$  is dissolved using a high concentration KOH solution (30% at 80 °C).<sup>41,42</sup> After the etching, the floating PMMA membrane is scooped out, rinsed and positioned on top of the TEM grid. The PMMA is dissolved in acetone to free the  $MoS<sub>2</sub>$  flakes.

HAADF-STEM structural analysis was performed using a JEM-ARM200F electron microscope corrected for spherical aberration, operated at 200 kV.

## Results and discussion

 $MoS<sub>2</sub>$  flakes were synthesized at atmospheric pressure in an open tube using S and  $MoO<sub>3</sub>$  powders with nitrogen as the carrier gas, on 290 nm thick  $SiO<sub>2</sub>/silicon$  substrates. We use 50 mg of  $MoO<sub>3</sub>$  powder and 175 mg of Sulfur powder. Fig. 1(a) shows a schematic illustration of the CVD reactor, *i.e.* an open tube with a diameter of about 1 inch heated in a two zone furnace. A sulfur boat is positioned in a low temperature zone, while the  $MoO<sub>3</sub>$  and the growth substrate are placed in a high temperature zone. In order to favor the nucleation of the flakes, perylene-3,4,9,10-tetracarboxylic dianhydride (PTCDA) was used as the organic promoter.<sup>37</sup> PTCDA was dissolved in water and drop-casted with a concentration of 12.75 M on a separate Si substrate. The PTCDA solution was obtained by dissolving 100 mg of PTCDA in 20 cubic centimeter of deionized water. The SiO2/Si employed as the growth substrate and the PTCDA/Si are positioned on a graphite susceptor, with the PTCDA/Si adjacent and immediately downstream to the growth  $SiO<sub>2</sub>/Si$  substrate. Paper<br>
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> Fig. 1(b) summarizes the growth process, showing the precursor temperature and the  $N_2$  flow profile. In order to obtain a more controllable S flow during the  $MoS<sub>2</sub>$  synthesis, we initially heated the substrate, the MoO<sub>3</sub> powders at 750  $\degree$ C under a flow of 300 sccm of  $N_2$ , while keeping S at a low temperature (<70  $\degree$ C). The downstream position of the PTCDA/ Si substrate is chosen to maximize the gradient of the organic promoter: PTCDA decomposes during the heating ramp and diffuses upstream in the gas phase, depositing on  $SiO<sub>2</sub>$ . Due to the diffusion process, we expect a maximum PTCDA transfer to the growth substrate in the zone Z3, gradually decreasing to a minimum concentration in the zone  $Z1<sup>37</sup>$  After the growth, only carbon residues remain on the sacrificial substrate with PTCDA, as demonstrated by Raman analysis (Fig. S1 in the ESI†).

> Once the set point temperature has been reached, the  $N_2$ flow was stopped and the S powder was separately heated to 170  $\degree$ C, while maintaining a stagnant gas flow. This allowed limiting the S transport to the Mo powders and to the substrate during the temperature ramp; it was reported that this procedure prevents the uncontrollable reaction between the



Fig. 1 (a) Schematic illustration of the CVD reactor, with the configuration of the sulfur and Mo precursors and the disposition of the different substrates (b) Temperature (red lines) of the sulfur precursor (dashed line) and of the growth substrate (full line) and N<sub>2</sub> carrier flow (blue line) during the growth process (c) optical image of the growth substrate, a different contrast induced by the presence of MoS<sub>2</sub> is reported. The marker is 200  $\mu$ m. (d) Optical images of the three areas named Z1 to Z3 highlighted by the colored squares in the panel (c). All the markers are 20  $\mu$ m. The area Z1, the farthest from the PTCDA sacrificial substrate, shows the presence of MoS<sub>2</sub> multi-layered dendritic structures and micrometric triangular monolayers, and the optical image of Z2 shows the presence of triangular crystals affected by overgrowth. Z3, the closest to the PTCDA substrate shows the presence of the MoS<sub>2</sub> polycrystalline monolayer with high coverage, affected by a micrometric few-layer island, probably due to the seeding point of the crystal. (e) Large-area Raman spectroscopy spectra of the three areas in analysis. (f) Large area photoluminescence spectroscopy spectra of the same areas. The PL spectrum of Z3 is divided by 10 for the sake of clarity. (g) Normalized PL spectra.

precursors during the pre-growth stage.<sup>43</sup> The synthesis process lasts 10 minutes, and the system then cools down naturally to room temperature. In the optical image of the substrate after the growth process, a clear change in the optical contrast is evident from the closest to the farthermost areas with respect to the PTCDA substrate, red and blue squares respectively. The marker of the optical image is  $200 \mu m$ . A false color image in Fig. 1(c), is reported in Fig. S2.† Optical images at higher magnification are reported in Fig.  $1(d)$ , where the different areas are labelled Z1 to Z3, with a decreasing distance from the PTCDA substrate. The high magnification optical images clarify the different optical contrasts of the whole sample image; Z1 presents large dendritic  $MoS<sub>2</sub>$  multi-layer structures and

micrometric triangular  $MoS<sub>2</sub>$  flakes. In the Z2 area the lateral size of the MoS<sub>2</sub> flakes increases up to 20  $\mu$ m, although they are affected by dendritic overgrowth. In the Z3 area, the lateral size of the  $MoS<sub>2</sub>$  flakes increases up to hundreds of micrometers, with few-layer terraces, most likely in the seeding points of the flakes. The monolayer coverage, observed by optical microscopy, varies between 77% and 93% in the Z3 area. The optical microscope observations are confirmed by large-area Raman and PL spectroscopies. Fig. 1(e) shows the comparison of the Raman spectra obtained in the different areas, as expected the separation of the  $E_{2g}$  and  $A_{1g}$  modes of MoS<sub>2</sub>, defined as  $\Delta$ , changes in the different areas. In the case of Z1 and Z2,  $\Delta$  is equal to 25  $\mathrm{cm}^{-1}$ , the standard for bulk MoS<sub>2</sub>, and meanwhile  $\varDelta$ 

of Z3 is 21  $\text{cm}^{-1}$ , representative of the MoS<sub>2</sub> bilayer. Raman analysis reveals no detectable amorphous carbon after the  $MoS<sub>2</sub>$ growth process (Fig. S3†). Large-area spectroscopy spectra (Fig. 1(f)) show a peculiar behavior; Z3 shows the typical PL spectrum of the  $MoS<sub>2</sub>$  monolayer, with an intense peak at 672 nm (1.85 eV) related to the A exciton, $9$  and meanwhile the other areas present weak PL emissions. The Z3 spectrum is divided by 10 for a direct comparison with the PL spectra of the other areas. In order to have a clear comparison of the PL peak position, the normalized PL spectra are reported in Fig. 1(g). Z2 shows an important red-shift of the A exciton emission with the PL peak at 680 nm (1.82 eV). Z2 also presents a tail toward the low wavelength side due to the B exciton of  $MoS<sub>2</sub>$ , peaked at 632 nm (1.96 eV). The Z3 PL spectrum presents a faint broad peak centered at 665 nm (1.86 eV), probably related to the presence of micrometric monolayer flakes. In Fig. S4<sup>+</sup> the result of a growth process without the PTCDA sacrificial substrate is reported in order to highlight the effect of the Mo/S ratio on these growth parameters. Paper<br>
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Fig. 2 presents an accurate characterization of the morphological, optical and structural properties of the  $MoS<sub>2</sub>$  dendritic multilayer structures, representative of Z1. Fig. 1(a) and (b) report the secondary electron and back-scattered electron images, respectively. The SE image shows a three branched architecture, with oblong structures originating from a central backbone. It is worth noting that the dendritic structure originates from a central triangular pyramid. Aberration corrected TEM analysis of different pyramids is reported in Fig. S5 of the ESI† file, and it shows that the 3D-pyramids are deeply affected by extended defects. The defects in these pyramids can lead to the development of dendritic multi-layer structures. The BSE image shows a peculiar contrast behavior; some areas have a dark contrast in comparison with the  $SiO<sub>2</sub>$  substrate, while other areas present a bright contrast. This contrast inversion is probably due to the increase in the  $MoS<sub>2</sub>$  work function with the increasing number of layers.<sup>44</sup>

The Raman maps of  $\Delta$  and of the  $A_{1g}/E_{2g}$  intensity ratio are reported in Fig. 2(c) and (d), respectively. The  $\Delta$  Raman map reveals that the dendritic structure presents an average  $\Delta$  of



Fig. 2 (a) SE micrograph of the dendritic structures representative of Z1, (b) BSE micrograph, (c) Raman map of the separation between the  $A_{10}$ and  $E_{2g}$  modes, defined as  $\Delta$ , (d) Raman map of the A<sub>1g</sub> and  $E_{2g}$  mode intensity ratio, (e) PL intensity map of the MoS<sub>2</sub> A exciton, and (f) PL peak position map of the MoS<sub>2</sub> A exciton.

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27  $\rm cm^{-1}$ , indicating bulk MoS<sub>2</sub>, and meanwhile the edges have a  $\varDelta$  of 21  $\mathrm{cm}^{-1},$  indicative of MoS $_2$  bilayers. The intensity ratio of the  $A_{1g}$  and  $E_{2g}$  is employed as a benchmark for the presence of defects, in particular sulfur vacancies, as previously reported.45,46 In this regard, the intensity ratio map shows that the center and the edges of the dendritic structure have a higher intensity ratio, an indication that these areas are more defective than the body of each branch. However, the intensity ratio map is also sensitive to the number of layers, $47,48$  and it can give similar information to the  $\Delta$  Raman map. Indeed, the intensity ratio map, reported in Fig. 2(d) confirms the analysis of the  $\Delta$ Raman map (Fig. 2(c)). PL mapping has been carried out to clarify the optical properties of the dendritic structures. The intensity and peak position of the  $MoS<sub>2</sub>$  A exciton are presented in Fig. 2(e) and (f), respectively. As expected the central part of the dendritic structure presents a faint PL intensity of the  $MoS<sub>2</sub>$ A exciton, due to the indirect band-gap of bulk  $MoS<sub>2</sub>$ . However, the PL intensity increases at the edge confirming the ML nature of the edges of the dendritic structure. It is worth noting that the PL peak position drastically red shifts from 668 nm, at the edges Nanoscale Advances<br>
27 cm<sup>-1</sup>, indicating bulk MoS<sub>2</sub> and meanwhile the redposition of the survivace *Ng*, 50 presents are option and the survey of the

of the structure, up to 685 nm at the center of the dendritic structure. Fig. S6† presents the optical image of the structure reported in Fig. 2 and the representative PL spectra, extracted from the PL map.

The PL intensity and peak position maps perfectly overlap, demonstrating that the red-shift of the A exciton is related to the indirect-to-direct band gap transition of  $MoS<sub>2</sub>$  in the monolayer form. The AFM analysis of the dendritic structure is reported in Fig. S7.†

A similar analysis has been carried out for the Z2 area. The  $SE$  image (Fig. 3(a)) reveals flakes with a lateral size varying from 30 to 50 µm. The dendritic overgrowth occurs close to the center of the flakes, probably starting at the seeding point of the flake itself. The BSE image (Fig. 3(b)) shows the presence of few-layer terraces (area with a brighter contrast). It is possible to convey that these terraces are few-layers by comparing the SE and the BSE images; indeed the SE image does not show any morphological contrast and meanwhile the BSE image presents a clear contrast. Raman mapping permits the clarification of the number of layers of these particular few-layer terraces. Fig. 3(c)



Fig. 3 (a) SE micrograph of the MoS<sub>2</sub> flakes affected by dendritic overgrowth, representative of Z2, (b) BSE micrograph, (c) Raman map of the separation between the A<sub>1g</sub> and E<sub>2g</sub> modes, defined as  $\Delta$ , (d) Raman map of the A<sub>1g</sub> and E<sub>2g</sub> mode intensity ratio, (e) PL intensity map of the MoS<sub>2</sub> A exciton, and (f) PL peak position map of the  $MoS<sub>2</sub>$  A exciton.

reports the  $\Delta$  map, showing that the flakes are mainly monolayer ( $\Delta = 19 \text{ cm}^{-1}$ ) and the dendritic overgrowth is bulk with a value ranging from 23  $\rm cm^{-1}$  to 26  $\rm cm^{-1}.$  However, the  $\varDelta$  map reveals the presence of bilayer terraces, similar to that reported by the BSE imaging with  $\varDelta$  equal to 21 cm<sup>-1</sup>, representative of bilayer  $MoS<sub>2</sub>$ . Furthermore, all the few-layer areas, both the bilayer terraces and the dendritic overgrowth areas, present a lower intensity ratio of the Raman modes (Fig. 3(d)): this is a clear indication that the few-layer area is less defective than the MLs. Even in this case the intensity ratio map confirms the analysis of the  $\Delta$  Raman maps. In addition, the bilayer nature of such terraces is confirmed by PL intensity mapping (Fig.  $3(e)$ ): as a matter of fact, the terraces present a slightly local decrease in the PL intensity and the decrease is more enhanced in the dendritic overgrowth areas. In the case of the A exciton peak position map, the peak varies from 683 nm (1.815 eV) to 685 nm (1.81 eV) in the ML area. In the bilayer and dendritic overgrowth area the PL peak is at around 682 nm (1.82 eV). This is in agreement with the large area PL analysis presented in Fig. 1(g). The optical image of the structure reported in Fig. 3 and the Paper<br>
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representative PL spectra are shown in Fig. S8.† The AFM analysis of the overgrowth affected flake is reported in Fig.  $S9.†$ 

In Fig. 4 we present the morphological, structural and optical properties of large-area polycrystalline monolayer  $MoS<sub>2</sub>$ , representative of Z3. The SE image reveals that the lateral size of the  $MoS<sub>2</sub>$  is difficult to evaluate due to the coalescence to a polycrystalline quasi-continuous film. It is worth noting that the SE analysis reveals the presence of bright and dark round features probably related to residuals of organic promoters<sup>33,38</sup> and craters.<sup>33</sup> The comparison with the BSE images (Fig.  $4(b)$ ) shows the presence of few-layer terraces, similarl to the overgrowth affected flakes reported in Fig. 3, and of grain-boundaries, demonstrating the polycrystalline nature of the  $MoS<sub>2</sub>$  film. The  $\Delta$  Raman map confirms the monolayer nature of the edges of the flakes with an average  $\Delta$  equal to 18.5 cm<sup>-1</sup>; meanwhile the central part of the flake is composed of bilayer terraces, with an average  $\Delta$  equal to 20 cm<sup>-1</sup>. The intensity ratio map (Fig. 4(d)) reveals a faint contrast demonstrating a comparable concentration of defects of the monolayer and bilayer  $MoS<sub>2</sub>$ . The comparison between the  $\Delta$  Raman map and the intensity



Fig. 4 (a) SE micrograph of large polycrystalline MoS<sub>2</sub>, representative of Z3, (b) BSE micrograph, (c) Raman map of the separation between the  $A_{1g}$  and  $E_{2g}$  modes, defined as  $\Delta$ , (d) Raman map of the  $A_{1g}$  and  $E_{2g}$  mode intensity ratio, (e) PL intensity map of the MoS<sub>2</sub> A exciton, and (f) PL peak position map of the MoS<sub>2</sub> A exciton.



Fig. 5 Aberration corrected TEM analysis results of the flake transferred from Z3, (a) interface area between a MoS<sub>2</sub> monolayer and bilayer, (b) and (c) bilayer and monolayer atomically resolved TEM image, respectively.

ratio shows a slightly different spatial localization, revealing that probably, in this case, the different concentrations of point defects is more effective than the number of layers. The PL intensity map of the A exciton of  $MoS<sub>2</sub>$  (Fig. 4(e)) shows peculiar homogeneity; however in the case of the polycrystalline  $MoS<sub>2</sub>$ the PL peak position of the  $MoS<sub>2</sub>$  A exciton shows that the exciton in the bilayer terraces (685 nm) is red-shifted by 5 nm in comparison with the ML areas (680 nm). It is worth noting that the edges and the grain boundaries present a blue shift, revealing the PL peak position set at 678 nm. Some PL spectra, extracted from the map, are shown in Fig. S10.† The AFM analysis of the overgrowth affected flake is reported in Fig.  $S11.†$ In order to evaluate the crystalline quality of the  $MoS<sub>2</sub>$  film aberration corrected Scanning Transmission Electron Microscopy (STEM) has been carried out (Fig. 5). A similar Raman and PL analyses are reported in Fig. S12† for another area of the polycrystalline MoS<sub>2</sub>. Nanoscale Advances<br>  $\sum_{k=1}^{\infty} \frac{1}{k} \sum_{k=1}^{\infty} \frac{1}{k} \sum_{k=1}^{\in$ 

The  $MoS<sub>2</sub>$  flake has been transferred (see Methods) from the Z3 area of the sample where mainly mono and bilayer  $MoS<sub>2</sub>$  can be found, and in particular we focused our attention on the interface between the monolayer and bilayer highlighted in Fig. 4.

Two very high magnification atomic resolution images have been taken (reported in Fig. 5(b) and (c)) in the blue and green squared region located in the monolayer and bilayer respectively in order to confirm the crystalline quality of the sample and, at the same time, to count the number of layers. This last task can be easily achieved by analysing the atomic patterns. As a matter of fact, in the single layer  $MoS<sub>2</sub>$  the Mo and S atoms are alternately located at the six corners of an hexagon<sup>49</sup> and they show a different contrast in the image in Fig. 5(b), see ESI Fig. S13.†

On the other hand, in the bilayer  $MoS<sub>2</sub>$ , the second layer is rotated by  $60^\circ$  with respect to the first one, resulting in the superposition of the two atomic species. Consequently, all the atomic columns will contain the same number of atoms showing the same contrast in the HAADF image reported in image 5c.



Fig. 6 Schematic illustration of the growth dynamics of MoS<sub>2</sub> considering the gradient of the organic promoter (PTCDA) with representative SEM images of the different areas.

As a further confirmation, we performed multislice STEM simulations,<sup>50</sup> finding a perfect match between the experimental and the simulated images, which are reported as insets along with the crystal models.

Beside these extra layers, the  $MOS<sub>2</sub>$  flakes are defect free and they have very high crystallinity as apparent from the images of the atomic lattice. Nevertheless, some vacancies can be observed in Fig. 5(a) highlighted by white circles; however the formation of electron beam induced vacancies is a well reported phenomenon and<sup>51</sup> therefore these defects cannot be univocally ascribed to growth.<sup>52</sup>

Fig. 6 illustrates the  $MoS<sub>2</sub>$  growth dynamics in the presence of the two opposite flows, the flow of the sulfur and molybdenum precursors  $(Moo<sub>3</sub>)$  and the backflow of PTCDA. The latter results in a gradient of the growth-promoter on the  $SiO<sub>2</sub>$ substrate, causing a change in the  $MoS<sub>2</sub>$  structure morphology and lateral size along the length of the growth substrate. We expect a maximum promoter concentration in the zone Z3 and a minimum concentration in Z1. First, the size of the monolayer triangular flake increases laterally from few micrometers to several tens of micrometers, getting close to the sacrificial PTCDA substrate. The increase of the size of triangular monolayer crystals is the direct effect of the PTCDA, which acts as a growth promoter favoring the lateral growth of the flakes. In fact, as previously reported, by employing PTCDA it is possible to achieve the complete coverage of the growth substrate.<sup>37</sup> In addition, the presence of dendritic structures decreases. Albeit the growth of dendritic structures of transition metal dichalcogenides previously reported, the origin of this particular effect is not completely clarified. The Kinetic Monte Carlo (KMC) method was used to simulate the growth behavior during van der Waals epitaxy growth of WSe2 and other atomically thin TMD films.<sup>53,54</sup> The shape transition from triangular crystals to fractal and dendritic structures corresponds to the change in the growth parameters.55,56 In particular, the simulation shows that one important parameter in order to achieve dendritic and fractal structures is the ratio of the chalcogen and metal flow<sup>53</sup> as experimentally verified in recent studies.<sup>55,57,58</sup> Paper<br>
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Another possible origin of the dendritic and fractal structures is their formation at the initial stage of the growth of a high-quality TMD triangular monolayer. In fact, the high growth rate is predicted to occur through a two-step process: after the formation of fractal/dendritic structures, they can eventually evolve into highly crystalline triangular structures. This possible origin was experimentally addressed in the case of growth of  $WS_2$  triangular flakes on SiO<sub>2</sub>/Si substrates.<sup>59</sup> One important aspect of the dendritic TMD structure is the hierarchy level. A high level of hierarchy is normally reported in the case of six branched structures due to the presence of twin boundaries.<sup>57,58</sup>

Most of the studies presenting dendritic structures do not employ a growth promoter except for the work of J. Wang et al.<sup>58</sup> in which the authors employed tape glue residuals to promote CVD growth. It is noticeable that the dendritic structure morphologies depend on which tape was used in the growth process.

It is worth noting that in this work we never obtain dendritic structures with a high hierarchy level. We report three branched structures, however, in our case the central part of the dendritic structure is multi-layer, as demonstrated by Raman and PL mappings. We assume that the dendritic structures probably originate from three-dimensional defective pyramids (Fig. S5†) found at the center of the dendritic structures, as suggested by the multi-layer central part of the structure (Fig. 2 and S7†). The crystalline disorder due to the extended defects in the central pyramid can cause the growth of the dendritic structures. The synthesis of the three dimensional pyramidal seeds can be due to inhomogeneity of the PTCDA growth promoter on  $SiO<sub>2</sub>$ boosting three dimensional growth. These effects are supported by a recent study,<sup>60</sup> where by utilizing a porphyrin based promoter (5,10,15,20-tetrakis(4-hydroxyphenyl)-21H, 23Hporphyrin) large-scale vertically oriented  $MoS<sub>2</sub>$  nanoflakes were obtained. In particular, the thickness of the promoter layer was important to achieve large scale homogeneity.

## **Conclusions**

In conclusion we clarify the effect of gradient of an organic promoter on the growth dynamics of molybdenum disulphide. We employ a perylene-based growth promoter (PTCDA) that sublimates during the heating ramp and diffuses upstream in the gas phase. A substrate covered with the promoter is placed downstream with respect to the growth substrate; the promoter is transferred to the growth substrate by upstream flow diffusion, thus creating a gradient of its distribution. The farthest zone from the sacrificial promoter substrate presents micrometric triangular monolayers and dendritic multi-layer  $MoS<sub>2</sub>$ branched structures. This growth dynamics demonstrates that a low amount of organic promoter leads to reduced lateral growth and boosting of the vertical growth. We suppose that the dendritic structures originate from three-dimensional defective pyramids found at the center of these structures. In the zone closest to the promoter substrate, a quasi-continuous  $MoS<sub>2</sub>$ polycrystalline film has been found, demonstrating that a high quantity of growth promoter enhances the lateral growth, suppressing the vertical one. However, the presence of bilayer terraces is reported. The transition zone presents monolayer crystals with an average lateral size of 50 micrometers affected by multilayer dendritic overgrowth at the flake center, supporting the direct connection between the flow gradient and the  $MoS<sub>2</sub>$  growth dynamics.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

F. F. wants to thank Dr Pasqualantonio Pingue for the help in the analysis of the AFM images. Dr Corneliu Ghica (National Institute of Materials Physics, Magurele, Ro) is acknowledged for the help in the TEM analysis.

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