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Absorption, fluorescence, Raman spectroscopic and density functional theoretical studies on the singlet and triplet excited state decay of 3-amino-5-mercapto-1,2,4-triazole†

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The excited state decay process of N-heterocyclic compounds is attracting increasing attention due to their fundamental applications in pharmaceutical and biological sciences. In this study, 3-amino-5-mercapto-1,2,4-triazole (AMT) was investigated in solid, protic, and aprotic solvents using vibrational and electronic spectroscopies combined with density functional theory (DFT) and time-dependent density functional theory (TDDFT) calculations. The steady absorption and resonance Raman spectra indicated that the AMT structure was quite sensitive to the polarity and proton of the solvent, and the pH environments. The intermolecular hydrogen bonding may contribute significantly to the decay channels of the singlet excited $S_2(\pi\pi^*)$ state process. Moreover, ns-transient absorption spectroscopy detected the short-time triplet species with ~ 200 ns lifetime in solvents. The DFT and TDDFT calculations interpreted the photophysical and photochemical process from the excited $S_2(\pi\pi^*)$ state, including the singlet and triplet decay mechanisms.

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Introduction

1,2,4-Triazole and its derivatives are widely used most notably as antifungals, pharmaceuticals and organic electroluminescent devices.^{1–3} Furthermore, 1,2,4-triazoles containing nitrogen atoms are easily coordinated with metals. Thus, apart from the biological significance,⁴ 1,2,4-triazoles are also applied to the protection of metals and alloys, such as corrosion inhibition.^{5–7} Important structured features, such as the planarity (π) and lone pair of electrons on the hetero-atom, are attributed to their strong coordination on the metal surface.^{8–13} In particular, 3-amino-5-mercapto-1,2,4-triazole (AMT) can be considered as one pull–push electronic structure due to both electron donor ($-\text{NH}_2$) and acceptor ($-\text{S}=\text{C}$ or $-\text{SH}$) groups bridged by one aromatic 1,2,4-triazole ring. Due to the donor– π -acceptor (D– π -A) structure, AMT demonstrates distinguished physical and chemical behaviour besides the above-mentioned characteristics. Moreover, AMT contains both donor and acceptor of hydrogen, which provides the prerequisite to the proton transfer reaction. Proton transfer has long been an issue of fundamental importance. In particular, the excited state proton transfer induces a unique large fluorescence Stokes shift

emission due to the dramatic alteration of the structure. The shifts usually can reach as large shift as $6000\text{--}1200\text{ cm}^{-1}$.¹⁴ Albinsson *et al.* reported that the adenosine derivative N^6,N^6 -dimethyladenosine (DMA) had a double fluorescence band located at 330 and 500 nm corresponding to the normal and isomer structures, respectively, in a series of different polar solvents.¹⁵ Similarly, 7-azaindole was reported to have double fluorescent bands in water. Shen and Wu¹⁶ elaborated that the fluorescence band of the isomer was red-shifted as the pH increased.¹⁷ Generally, popular hydrogen bonds (H-bonds) involve a combination of an O–H or an N–H proton donor and an N or an O proton acceptor, which demonstrates the double fluorescent phenomenon and large Stokes shift. However, there are only a few reports on H-bonds that involve a sulfur atom as either the proton donor ($-\text{SH}$) or acceptor (S). The concept of “sulfur participating in an H-bond” appeared in early reports despite the relatively low electronegativity (2.58) for sulfur on the Pauling scale.^{18–21} Recent quantum chemical calculations in conjunction with high-resolution spectroscopic evidence have revealed that a sulfur-containing H-bond can be as strong as conventional H-bonds.^{22–24} In biological systems, methionine-containing dipeptides have been reported to form amide–N–H \cdots S intermolecular H-bonds that are stronger than amide–N–H \cdots O–C H-bonds.²⁵ Additionally, the strengths of the intermolecular H-bonds in thiobase pairs are reported to be of the same magnitude with the canonical nucleobase pairs.²⁶ As a result, the intermolecular sulfur H-bonds play an important role in artificial DNA or RNA research, and have been applied to

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RNAi resistant genes²⁷ and proteins. Recently, Chou *et al.* reported an important 7-hydroxy-2,2-dimethyl-2,3-dihydro-1*H*-indene-1-thione (DM-7HIT) molecule with O–H···S H-bond formation and its excited-state intramolecular H-bond on/off reaction by room-temperature phosphorescence (RTP). The results concluded that DM-7HIT possessed a strong polar (hydroxy)-dispersive (thione)-type H-bond, which opened a new chapter in the chemistry of hydrogen bonds, especially the generality of the sulfur H-bond on/off switching mechanism. Moreover, the tautomerism between 2-thiopyridone (2TP) and 2-mercaptopyridine (2MP) concluded that the excited state proton transfer (ESPT) reaction of 2TP in solvents led to the formation of 2MP transients by nanosecond time-resolved resonance Raman spectroscopy.²⁸ In early 2006, Fu *et al.* reported the tautomerization of 2-pyridinethione in water, and found that water promoted the proton transfer in the ground state *via* a bridged hydrogen bonding.²⁹ Recently, we also reported several solvent-dependent hydrogen bonding molecules, *i.e.*, 5-(methylthio)-1,3,4-thiadiazole-2(3*H*)-thione (MTTN), 3-amino-1,2,4-triazole (3AT), 4,6-dimethyl-2-mercaptopyrimidine (DMMP), and 2-mercaptobenzothiazole by Raman spectroscopy and density functional theory calculations.^{30–34} These results revealed that solvent hydrogen bonding can stabilize the thione isomers and the thione ↔ thiol tautomerization may occur in a higher excited state, which was greatly controlled by the solvent character and UV light wavelength.

In this paper, we reported the ground structure and excited state decay dynamics of AMT in different environments using resonance Raman spectra, steady and nanosecond transient absorption spectroscopies combined with theoretical calculation. The decay channels of the singlet excited $S_2(\pi\pi^*)$ state were investigated by the DFT and TD-DFT calculations.

Experimental and computational methods

3-Amino-5-mercapto-1,2,4-triazole (AMT, solid, purity, 99%) was purchased from J&K Scientific Ltd., China and employed without further purification. UV absorption spectra in water, methanol, and acetonitrile were measured using an ultraviolet/visible spectrometer at a concentration of $\sim 10^{-5}$ mol L⁻¹. The Fourier transform (FT) Raman and FT-IR spectra were obtained using a Thermo Nicolet FT-Raman 960 spectrometer and a PerkinElmer 1 FT-IR spectrometer with 2 cm⁻¹ resolution. For the resonance Raman experiments, the harmonics of a nanosecond Nd:YAG laser and their hydrogen Raman-shifted laser lines were utilized to generate the 239, 245, 252, 266, and 273 nm excitation wavelengths, which have been described previously.^{28,35,36} The excitation laser beam used a ~ 100 μ J pulse energy loosely focused to a 0.5–1.0 mm diameter spot size onto a flowing and circulating liquid stream of sample. Solution-phase samples were used at a concentration of $\sim 10^{-3}$ mol L⁻¹ AMT (99% purity) in water and methanol ($\geq 99.5\%$ purity). The Raman signal was recorded for about 100 s before being read out to an interfaced PC computer; 80–100 of these readouts were combined to obtain a resonance Raman spectrum with a 2 cm⁻¹

resolution. The Raman shifts were calibrated using the known vibrational frequencies of the solvents (*i.e.*, methanol and water).

The complete geometry optimization and vibration analysis were carried out at the B3LYP-D3(BJ)/6-311++G(d,p) level, and the normal-mode analysis of AMT was carried out using the VEDA program.³⁷ The vertical electronic transition energies were performed at the B3LYP(TD)-D3(BJ)/6-311++G(d,p) level. The vertical $S_0 \rightarrow S_n$ transition energy and the S_n minimum were estimated at the B3LYP-TD-D3(BJ)/6-31G* level by employing a self-consistent reaction field (SCRF) and the polarized continuum (overlapping spheres) model (PCM). The ground state and vertical transition computations were obtained using the Gaussian 09 program.³⁸ The minimum energies and optimized geometric structures for S_0 , S_1 , S_2 , and the conical intersection points were calculated using the single-component artificial force-induced reaction (SC-AFIR) method as implemented in the global reaction route mapping (GRRM) program,³⁹ which runs in combination with Gaussian 09. Moreover, to avoid the erroneous description of weak interactions, DFT including long-range dispersion effects (DFT-D3) by Grimme^{40–42} are involved in all optimized geometries.

Results and discussion

Vibrational spectra

The crystalline AMT structure was initially confirmed by FT-IR and FT-Raman spectra, as shown in Fig. 1(a) and (c). To further assign the experimental spectra, the possible tautomers were calculated at the B3LYP-D3(BJ)/6-311++G(d, P) level. As shown in Fig. S1,† eight monomeric isomers were found and thiol structure (a) is the lowest one in energy. Five possible thiol structures were also obtained, which were lower than thiones (b) and (c) in energy except for structure (a). The vibrational

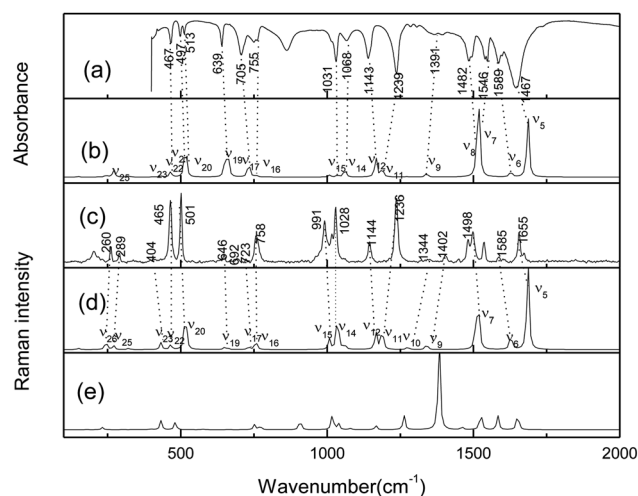


Fig. 1 (a) FTIR spectrum of the molecule of neat powder in KBr pellet. (b) Calculated IR spectrum of gas thione-AMT monomer (1,4-*H*-AMT in Fig. S1†). (c) FT-Raman spectrum. (d) Calculated Raman spectrum of gas thione-AMT monomer (1,4-*H*-AMT in Fig. S1†) and (e) thiol (structure A in Fig. S2†) at the B3LYP-D3(BJ)/6-311++G(d, p) level.



frequencies were further analysed for the optimized structures at the same level. However, only the calculated Raman activity spectra of the most stable monomer are compared in Fig. 1.

As shown in Fig. 1(c), two strong bands were observed at 465 and 501 cm^{-1} in the 100–920 cm^{-1} spectral region, which correlate reasonably well to the calculated (scaled) modes at 467 cm^{-1} (ν_{22}) and 518/513 cm^{-1} (ν_{20}/ν_{21}). Based on the calculated results, both bands corresponded to the C=S stretching vibration, which indicates that AMT existed as a thione structure in the crystal state. The bands with medium intensities observed at 260, 289 and 758 cm^{-1} correlate reasonably well to 244 (ν_{26}), 271 (ν_{25}), and 757 cm^{-1} (ν_{16}) calculated at the B3LYP level. Three weak bands observed at 404, 646 and 723 cm^{-1} correlate reasonably well to the calculated modes at 432 (ν_{23}), 649 (ν_{19}), and 732 cm^{-1} (ν_{17}). In the 920–1700 cm^{-1} region, the FT-Raman spectra display three strong bands at 991, 1028, and 1236 cm^{-1} , in good agreement with the calculated modes ν_{15} at 1008 cm^{-1} , ν_{14} at 1035 cm^{-1} and ν_{11} at 1188 cm^{-1} . Other observed bands at 1144, 1344, 1402, 1498, 1585 and 1655 cm^{-1} should correspond to the calculated 1166 (ν_{12}), 1275 (ν_{10}), 1340 (ν_9), 1517 (ν_7), 1628 (ν_6), and 1686 (ν_5) cm^{-1} , corresponding to the thione monomer in Fig. 1(d). Thus, solid-state AMT can be attributed to exist as a thione monomer in the solid form. Complete assignments are given for the computed and experimental bands in Table S1.†

Resonance Raman experiments were also carried out in different solvents with selected excitation wavelengths. The excitation wavelength was determined according to the ultraviolet absorption band, and the low solubility of AMT in acetonitrile did not meet the requirements of resonance Raman detection, so we did not obtain any spectra in acetonitrile. Fig. 2 displays the 266 nm resonance Raman spectra in water and methanol. Compared to the FT-Raman spectra, the intensities of some Raman bands were greatly altered. In addition, some bands became very broad and could not be resolved due to the effects of the solvents. Fig. 2 compares the 266 nm resonance Raman spectra of AMT in water and methanol (figure shows

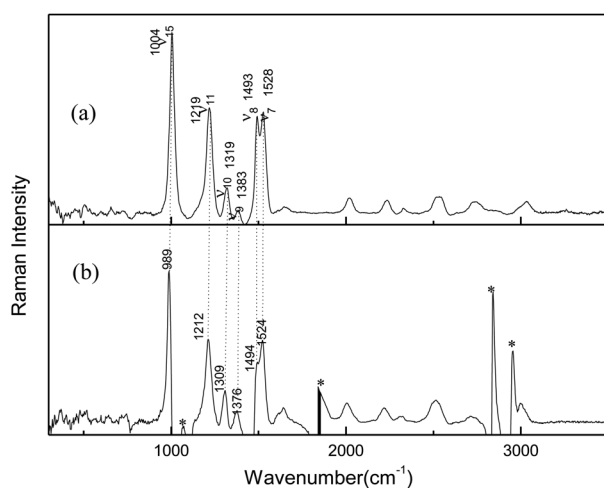


Fig. 2 266 nm resonance Raman spectra in (a) H_2O , (b) CH_3OH (asterisks mark regions where solvent subtraction artifacts are present).

only the main vibration modes). Calculated vibrational frequencies at the B3LYP-D3(BJ)/6-311++G(d, p) level using the H_2O PCM model and assignments of thione AMT in solvents are shown in Table S2.† In methanol and water, the obvious Raman shifts were observed between corresponding peaks. In particular, strong bands at 1004 (ν_{15}), 1219 (ν_{11}), 1493 (ν_8), and 1528 (ν_7) cm^{-1} dominated the Raman spectra in water. Conversely, in methanol, strong bands were observed at 989 (ν_{15}), 1212 (ν_{11}), 1494 (ν_8), and 1524 (ν_7) cm^{-1} . The intense bands were observed at 1319 (ν_{10}) and 1383 (ν_9) cm^{-1} in water, while 1309 (ν_{10}) and 1376 (ν_9) cm^{-1} were the weak bands in methanol. The Raman shifts of these modes were caused by the hydrogen bond between the solvent and AMT molecules. Furthermore, we observed a big difference in water and methanol for ν_8 , which was assigned to the $\text{N}_2\text{C}_3/\text{N}_4\text{C}_3$ stretching combined with $\text{H}_{11}\text{N}_4\text{C}_3/\text{H}_8\text{N}_7\text{C}_3$ in-plane bending. Therefore, we confirmed that the hydrogen bond between the solvent and the AMT molecule should be mainly located on both $-\text{N}_4\text{H}_{11}\text{O}$ and $\text{H}_9\text{N}_7\text{H}_8\text{O}$ groups. In our recent reports, methanol was shown to form stronger hydrogen bonds with solute molecules than water.^{31,32} Thus, the weakness of ν_8 in methanol may be the result of stronger solute–solvent hydrogen bonding. This suggests that structural dynamics are solvent and hydrogen-bonding dependent along the reaction coordinates in the Franck–Condon region.

Fig. 3 presents an expanded view of the 252 nm resonance Raman spectra with the tentative vibrational assignments indicated above the spectra. Most of the 252 nm resonance Raman spectra can be assigned to the fundamentals, overtones, and combination bands of about nine Franck–Condon active vibrational modes: the $\text{H}_6\text{N}_1\text{N}_2\text{C}_3$ tors/ $\text{C}_3\text{N}_4\text{C}_5\text{N}_1$ tors/ $\text{S}_{10}\text{N}_1\text{N}_4\text{C}_5$ out ν_{18} , the $\text{N}_2\text{C}_3\text{N}_4\text{C}_5$ tors/ $\text{N}_7\text{N}_4\text{N}_2\text{C}_3$ out ν_{17} , N_4C_3 stretch/ $\text{N}_2\text{C}_3\text{N}_4$ bend ν_{15} , the N_1C_5 stretch/ S_{10}C_5 stretch/ $\text{H}_{11}\text{N}_4\text{C}_3$ bend/ $\text{N}_4\text{C}_5\text{N}_1$ bend ν_{11} , the N_1C_5 stretch/ $\text{H}_6\text{N}_1\text{N}_2$ bend ν_{10} , the N_4C_5 stretch/ $\text{H}_6\text{N}_1\text{N}_2$ bend/ $\text{H}_{11}\text{N}_4\text{C}_3$ bend ν_9 , the N_2C_3 stretch/ N_4C_3 stretch/ $\text{H}_{11}\text{N}_4\text{C}_3$ bend/ $\text{H}_8\text{N}_7\text{C}_3$ bend ν_8 , the N_1C_5 stretch/ $\text{H}_6\text{N}_1\text{N}_2$

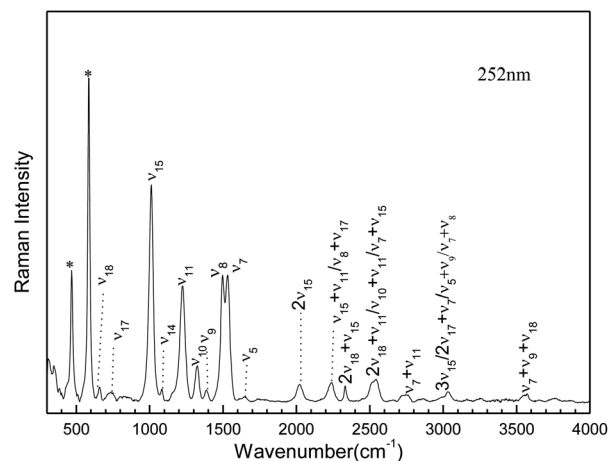


Fig. 3 Expanded view of 252 nm resonance Raman spectra of AMT in water with the tentative vibrational assignments indicated above the spectrum.



bend ν_7 , and the N_2C_3 stretch/ N_7C_3 stretch/ $\text{H}_9\text{N}_7\text{H}_8$ bend ν_5 , as shown in Table S2.†

Steady absorption spectra

Fig. 4 shows the UV absorption spectra of AMT in acetonitrile, water and methanol solvents. In acetonitrile, the experimental UV spectrum for AMT exhibits one strong absorption band with the maxima at approximately 266 nm ($f = 0.1626$, $\epsilon = 0.985 \times 10^4 \text{ L mol}^{-1} \text{ cm}^{-1}$). In methanol, the band absorption appears at 260 nm ($f = 0.2953$, $\epsilon = 1.4472 \times 10^4 \text{ L mol}^{-1} \text{ cm}^{-1}$), and it is blue-shifted 6 nm relative to that observed in acetonitrile. In water, the maximum band absorption was blue-shifted to 254 nm ($f = 0.3312$, $\epsilon = 1.1876 \times 10^4 \text{ L mol}^{-1} \text{ cm}^{-1}$) relative to that in acetonitrile. The blue-shifts of this band indicated that there was an intermolecular hydrogen bond between the protic solvent molecules and AMT, in addition to the solvent polarity. To simulate the solvent effects on the UV spectra of AMT in solvents, the B3LYP(TD)-D3(BJ)/6-311++G(d, p) calculations employed explicit and implicit solvent polarizable continuum models (PCM).

Table 1 lists the computed electronic transition energies, oscillator strengths and their corresponding orbitals of AMT with the PCM model. The calculated S_2 electronic transition bands at 252.91 nm ($f = 0.3534$) in acetonitrile, 252.77 nm ($f = 0.3353$) in methanol and 252.68 nm ($f = 0.3650$) in water, respectively, all of which were assigned to $\pi \rightarrow \pi^*$ electronic transitions, were in good agreement with the bands observed experimentally. Next, we discussed the change of the absorption spectra at different pH environments in methanol and water. Fig. 5 and S3† show the absorption spectra of AMT in H_2O and CH_3OH with different volumes of $[(\text{Bu})_4\text{N}]^+\text{OH}^-$. At pH = 6.25 in H_2O without $[(\text{Bu})_4\text{N}]^+\text{OH}^-$, AMT exhibited a strong band at around 254 nm, which represents the normal species. As the pH increased from 6.25 to 7.44, the band at 254 nm was gradually weakened and even disappeared, while the band at 242 nm gradually increased, and no change occurred when $\text{pH} > 8.59$.

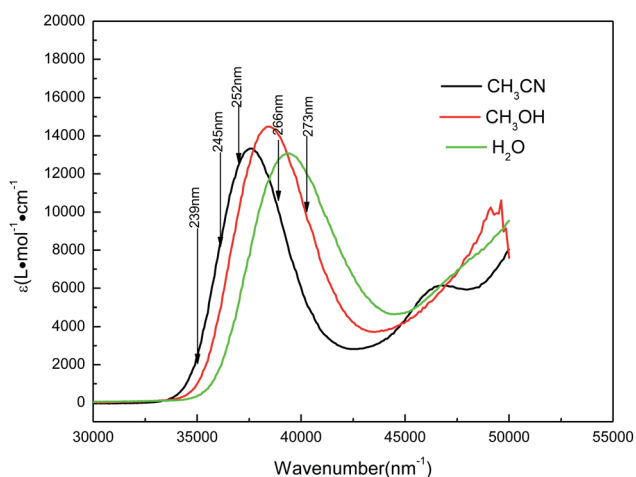


Fig. 4 UV absorption spectra of AMT in water (green line), methanol (red line), and acetonitrile (black line) with the excitation wavelengths labeled for the resonance Raman spectra.

Wrzosek and Bukowska proved that the single nitrogen atom in the triazole ring of AMT remained protonated even in $\text{pH} = 12$ solution.^{43–45} Therefore, we speculated that the band at 242 nm may be attributed to the single detached proton, and only the other kept on the triazole ring, which led to the red-shifts of the absorption band. So, we concluded that the AMT is deprotonated in basic solution to form anion monomers.

Transient absorption spectra

Fig. 6 shows the transient absorption spectra obtained at 266 nm excitation wavelength in CH_3CN , CH_3OH and H_2O . Two absorption bands at about 330 and 600 nm appeared in all three solvents. As shown in Fig. 6, the absorption band at ~ 600 nm became broad as the polarity of the solvent increased. According to the TD-DFT calculation, two absorption bands corresponded to the lowest triplet $^3\pi\pi^*$ transition, as shown in Table S3.† Fig. S4† shows the kinetics of AMT measured at 550 and 600 nm in CH_3CN , CH_3OH and H_2O . The decay time for AMT in acetonitrile was 274 ns, 140 ns in methanol and 153 ns in water. AMT has a ~ 120 ns longer lifetime in acetonitrile than in water and methanol. Therefore, in protic solvents, besides the intersystem jumping back to the ground state, there should be other main decay pathways.













Table 2 lists the computed electronic transition energies, oscillator strengths and their corresponding orbitals of AMT with the PCM model. The calculated electronic transition-allowed band occurred at 565.44 nm ($f = 0.1494$) in acetonitrile, and the calculated electronic transition-allowed band occurred at 564.82 nm ($f = 0.1489$) in water, which corresponds to the triplet state transition. This is basically consistent with the experimentally observed absorption of ~ 600 nm in solvents.

Excited $S_2(\pi\pi^*)$ state decay mechanism

To elaborate the experimental results, ground state structures were initially calculated at the B3LYP-D3(BJ)/6-31G* level. At this level, the ground-state AMT–thiol tautomer is 13.8 kcal mol^{-1} higher in energy than the AMT–thione tautomer, which is connected *via* transition state with a barrier of 66.3 kcal mol^{-1} . Therefore, the ground-state proton transfer reaction between AMT–thione and AMT–thiol was inhibited as a result of the high barrier, which is similar to those previously reported for DMMP.³² According to the UV absorption spectrum, the ground AMT molecule was excited to the excited $S_2(\pi\pi^*)$ state in the Franck–Condon region. Subsequently, it evolved to the lower excited state *via* internal conversion (IC) or illumination. To further compare the experimental results, excited state decay potential surface profiles on the excited state were calculated, including the excited singlet triplet states as shown in Fig. 7. AMT in the $S_2(\pi\pi^*)$ state of the Franck–Condon region with an energy of 120.2 kcal mol^{-1} vibrationally reached $S_{2 \text{ min}}(\text{n}\pi^*/\pi\pi^*)$ with 96.4 kcal mol^{-1} or $S'_{1 \text{ min}}(\pi\pi^*)$ *via* CI $S_1/S_2(^1\pi\pi^*/^1\pi\text{Ryd})$ with 119.2 kcal mol^{-1} . This is only 1.0 kcal mol^{-1} energy difference between them, as shown in Fig. 7. Fig. S5† showed the calculated energies of S_0 and $^1(\pi\pi^*)$ excited states on the selected ground state IRC points at the B3LYP(TD)-D3(BJ)/6-31G* level. The results show the higher



Table 1 Experimental and calculated singlet electronic transition energies, corresponding orbitals and oscillator strengths with the electronic transition character for the optimized ground AMT at B3LYP(TD)-D3(BJ)/6-311++G(d, p) level

States	Orbitals (coefficient)	Electronic transition	Transition energy (nm)		Oscillator strength (f)	
			Calc.	Expt	Calc.	Expt
AMT in CH₃CN						
T ₁	30 → 31 (0.68938)	π → π*	337.62			
T ₂	29 → 31 (0.70333)	n → π*	280.57			
S ₁	29 → 31 (0.70488)	n → π*	269.69		0.0002	
T ₃	30 → 34 (0.56829)	π → π*	269.59			
T ₄	30 → 32 (0.67340)	π → Ryd ₁	254.49			
S ₂	30 → 31 (0.66706)	π → π*	252.91	266.0	0.3534	0.1626
	30 → 32 (-0.15659)	π → Ryd ₁				
	30 → 34 (0.10341)	π → π*				
S ₃	30 → 31 (0.14705)	π → π*	251.90		0.0294	
	30 → 32 (0.66903)	π → Ryd ₁				
	30 → 33 (0.15603)	π → Ryd ₂				
S ₄	29 → 32 (0.11813)	n → Ryd ₁	231.36		0.0250	
	30 → 32 (-0.14267)	π → Ryd ₁				
	30 → 33 (0.66957)	π → Ryd ₂				
						
AMT in H₂O						
T ₁	30 → 31 (0.68932)	π → π*	337.41			
T ₂	29 → 31 (0.70340)	n → π*	279.92			
T ₃	30 → 34 (0.56828)	n → π*	269.16			
S ₁	29 → 31 (0.70496)	n → π*	269.15		0.0002	
T ₄	30 → 32 (0.67497)	π → Ryd ₁	253.94			
S ₂	30 → 31 (0.67880)	π → π*	252.68	254.0	0.3650	0.3312
	30 → 34 (0.10503)	π → π*				
S ₃	30 → 32 (0.68195)	π → Ryd ₁	251.28		0.0172	
	30 → 33 (0.15238)	π → Ryd ₂				
S ₄	29 → 32 (0.11076)	n → Ryd ₁	230.80		0.0249	
	30 → 32 (-0.14095)	π → Ryd ₁				
	30 → 33 (0.67140)	π → Ryd ₂				
						

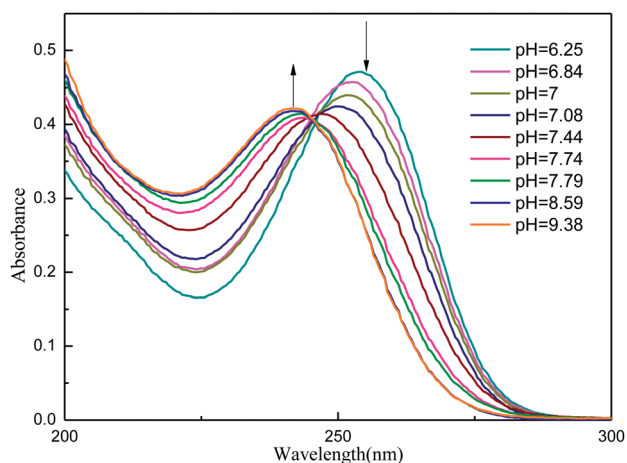


Fig. 5 UV for AMT ($3.95 \times 10^{-5} \text{ mol L}^{-1}$) in H₂O with different volumes of [(Bu)₄N]⁺OH⁻ ($1.53 \times 10^{-2} \text{ mol L}^{-1}$).

barrier on the ¹(ππ*) potential energy surface profile from thione to thiol, which indicated that the thione → thiol isomerization occurred on the CI (S₁/S₂) point. Table S4† lists the orbitals and transition nature of the 11 structures along the IRC path in detail.

Moreover, the spin-orbit coupling (SOC) constants were calculated between the higher singlet and lower triplet states using the single-particle BP operator with an effective charge approximation in PySOC.⁴⁶ The calculated SOC value results are shown in Table 3. For ground S₀, there is an obvious intersystem crossing (ISC) character due to large SOC constants for S₀/T₂, S₁/T₁ and S₂/T₂. Even for S_{1 min} and S_{2 min}, the ISC process became much easier than ground state. Based on the above-mentioned results, the decay path on the triplet state was calculated to elaborate the decay process on the decay mechanism of the higher excited state S₂(ππ*), according to the electronic spectroscopic results. As shown in Fig. 7, AMT was initially excited to S₂(ππ*) and then ISC to T₂(nπ*) or vibrational relaxation to the



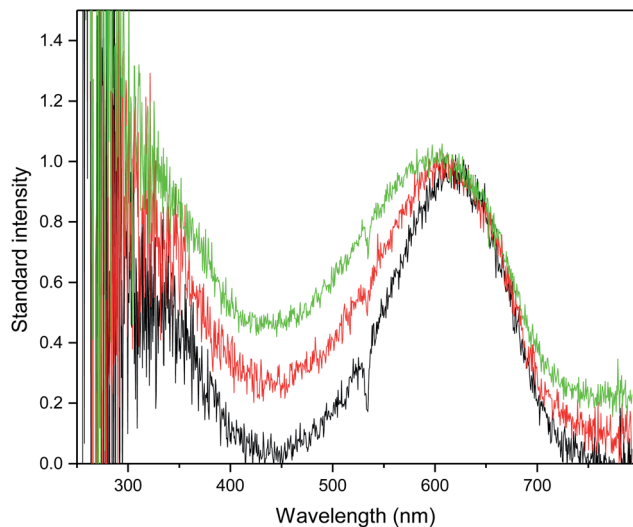


Fig. 6 Nanosecond transient absorption at 266 nm excitation of AMT in CH_3CN (in green), CH_3OH (in red) and H_2O (in black).

dark $S_1(n\pi^*)$ via a fast internal conversion in Franck-Condon region. Then, both $T_2(n\pi^*)$ and $S_1(n\pi^*)$ decayed to the $T_1(\pi\pi^*)$ state. The calculated spin-orbit coupling constant SOC of $S_1/T_1(1n\pi^*/^3n\pi^*)$ was 118 cm^{-1} , which also confirmed that intersystem crossings are quite favorable on the ground state. Another possibility was that the decay occurred from $S_1(n\pi^*)$ to $T_2(n\pi^*)$ states passing through the $S_1/T_2(1n\pi^*/^3n\pi^*)$ intersection. Then, the $T_2(n\pi^*)$ vibrationally relaxed to $T_1(\pi\pi^*)$. However, on the T_1 potential energy surface profile shown in Fig. 7, the energy of the thiol $T_1'_{\text{min}}(\pi\pi^*)$ was 38 kcal mol^{-1} higher than the energy of the thione $T_1_{\text{min}}(\pi\pi^*)$, so we believe the photophysical process occurred easily from T_1 to the ground state. The second path was that after AMT reaches the $S_2(\pi\pi^*)$ state, an intersystem crossing to the $T_2(n\pi^*)$ state occurred immediately due to the 133 cm^{-1} SOC constants for $S_2/T_2(1\pi\pi^*/^3n\pi^*)$ calculated at S_0 . Then, $T_2(n\pi^*)$ vibrationally

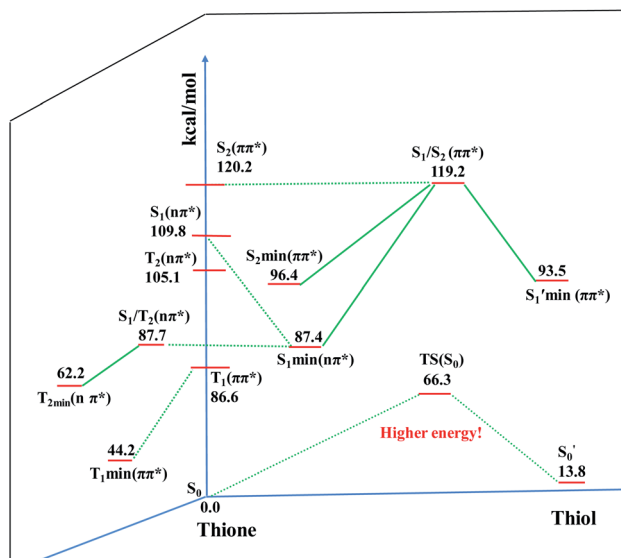


Fig. 7 Possible decay pathways of excited state AMT calculated at TD-B3LYP-D3(BJ)/6-31G* level using the H_2O PCM model.

Table 3 Calculated spin-orbit constants (SOC, cm^{-1}) of optimized S_0 , S_1_{min} and S_2_{min} structures

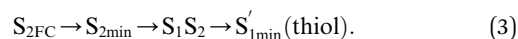
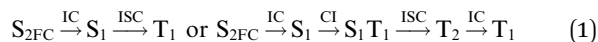
	S_0	S_1_{min}	S_2_{min}
S_0/T_1	1	119	117
S_0/T_2	93	148	143
S_1/T_2	4	134	110
S_1/T_1	118	78	72
S_2/T_2	133	82	75
S_2/T_1	4	115	121

relaxed to $T_1(\pi\pi^*)$. The third path was for the excited state proton transfer on the S_2 state, where a photochemical reaction might have occurred. In summary, three possible decay channels on S_2 are proposed as:

Table 2 Experimental and calculated triplet electronic transition energies, corresponding orbitals and oscillator strengths with the electronic transition character for the optimized T_1_{min} at the B3LYP(TD)-D3(BJ)/6-31G* level

States	Orbitals (coefficient)	Electronic transition	Transition energy (nm)		Oscillator strength (<i>f</i>)
			Cal.	Exp.	Cal.
AMT in CH_3CN					
T_1	29 → 30 (0.99635)	$\pi \rightarrow \pi^*$	1125.57		0.0000
T_2	28 → 30 (0.97857)	$\pi \rightarrow \pi^*$	565.44	613	0.1494
AMT in H_2O					
T_1	29 → 30 (0.99632)	$\pi \rightarrow \pi^*$	1112.52		0.0000
T_2	28 → 30 (0.97863)	$\pi \rightarrow \pi^*$	564.82	600	0.1489

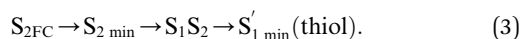
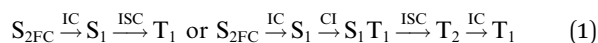




However, these possible pathways require further experimental support, such as transient fluorescence.

Conclusion

In conclusion, resonance Raman spectroscopy was employed to assign the vibrational spectra of 3-amino-5-mercapto-1,2,4-triazole (AMT) in methanol and water by a combination of FT-IR, FT-Raman, UV-vis spectroscopy, and density functional theoretical (DFT) calculations. The FT-Raman spectra confirmed that AMT formed a monomer in its neat solid form. In methanol and water, we speculated that the hydrogen bond between the solvent and the AMT molecule might be at $-N_4H_{11}O$ and $H_9N_7H_8O$, and the Franck-Condon region structural dynamics along various reaction coordinates are solvent dependent. After AMT is excited, we propose three decay channels on the S_2 state.



The triplet reaction channel may be applied to photodynamic therapy in disease diagnosis and treatment.

Conflicts of interest

There are no conflicts to declare.

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