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Electronic spectra of ytterbium fluoride from relativistic electronic structure calculations†

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We report an investigation of the low-lying excited states of the YbF molecule—a candidate molecule for experimental measurements of the electron electric dipole moment—with 2-component based multi-reference configuration interaction (MRCI), equation of motion coupled cluster (EOM-CCSD) and the extrapolated intermediate Hamiltonian Fock-space coupled cluster (XIHFS-CCSD). Specifically, we address the question of the nature of these low-lying states in terms of configurations containing filled or partially-filled Yb 4f shells. We show that while it does not appear possible to carry out calculations with both kinds of configurations contained in the same active space, reliable information can be extracted from different sectors of Fock space—that is, by performing electron attachment and detachment IHFS-CCSD and EOM-CCSD calculation on the closed-shell YbF⁺ and YbF[−] species, respectively. From these calculations we predict $\Omega = 1/2, 3/2$ states, arising from the $4f^{13}5d^1\sigma_{6s}^2$, $4f^{14}5d^1/6p^1$, and $4f^{13}5d^1\sigma_{6s}^1$ configurations to be able to interact as they appear in the same energy range around the ground-state equilibrium geometry. As these states are generated from different sectors of Fock space, they are almost orthogonal and provide complementary descriptions of parts of the excited state manifold. To obtain a comprehensive picture, we introduce a simple adiabaticization model to extract energies of interacting $\Omega = 1/2, 3/2$ states that can be compared to experimental observations.

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1 Introduction

In a previous paper,¹ we introduced all-electron relativistic basis sets for the lanthanides (La–Lu) and discussed their performance for the determination of spectroscopic constants for the ground state of ytterbium fluoride (YbF), an open-shell molecule with a $2\Sigma^+$ ground state. This molecule has received a fair amount of experimental and theoretical attention because of its potential application in the observation of parity-violating interactions^{2–4} via determination of the electric dipole moment of the electron (eEDM)—see, for instance ref. 4–11 and references therein). There is also some interest in the Yb atom, cation and dimer in connection to ultracold

physics.^{12–14} An example¹⁵ is the parity violation observed in the Yb atom.

A noteworthy finding in our previous work was the sensitivity of coupled cluster calculations to the basis set in use and, indirectly, to the amount of electron correlation recovered in the calculations. We observed a spike in the values of the T_1 diagnostic around the ground state equilibrium geometry, so it appears that the perturbative treatment of triple excitations in the CCSD(T) calculations breaks down in this region of the potential energy curves unless there is enough flexibility in the correlation treatment. The same was recently observed by Pasteka *et al.*¹⁶ for the nuclear quadruple coupling constant. This suggests the existence of a low lying perturbing state, which we want to investigate further in the current work.

Experimental^{17,18} and previous theoretical^{19–21} investigations suggest that in the ground state the unpaired electron is located in a σ_s orbital with dominant contributions from the 6s orbital of Yb, corresponding to a Yb($4f^{14}5d^1\sigma_{6s}^1$)F configuration. This $2\Sigma_{1/2}$ state ground state was studied in greater detail by combining microwave and optical spectroscopy for the odd ^{171}Yb isotope.²²

Experimentally,¹⁸ the lowest excited state observed is assigned as $2\Pi_{1/2}$, with an energy of $18\,106\text{ cm}^{-1}$, while the $2\Pi_{3/2}$ component is found at $19\,471\text{ cm}^{-1}$, yielding a spin-orbit splitting of 1365 cm^{-1} of this spin-orbit split A 2Π state. The lower component will be denoted $3_{1/2}$ in the current work.

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Experiments indicate a perturbation of its vibrational levels,^{18,23–25} which was attributed to the presence of a perturbing state (denoted by $4_{1/2}$ here) found at $18\,705\text{ cm}^{-1}$.¹⁸ This perturbing $4_{1/2}$ state is sometimes referred to as $[18.6]_{1/2}$ by experimentalists^{18,23} (energy in cm^{-1} divided by 10^3 in the square brackets, and Ω -value as subscript). The mixing of these two $\Omega = 1/2$ states gives rise to states designated as $[557]$ and $[561]$ (the values in square brackets referring to transition energies in Thz from the vibronic ground state) with transition energies of $18\,574$ and $18\,699\text{ cm}^{-1}$, respectively.^{24,25} These two states are of importance for laser cooling schemes that have been investigated²⁶ and tested²⁵ with the purpose of realizing high-accuracy measurements of YbF at very low temperatures. Besides these first excited states, Smallman²⁶ investigated also two not yet fully characterized mixed states, $[574]$ ($\approx 19\,150\text{ cm}^{-1}$) and $[578]$ ($\approx 19\,280\text{ cm}^{-1}$) at higher energies. These can be compared with the $\Omega = 3/2$ state at $19\,471\text{ cm}^{-1}$ found earlier by Dunfield,¹⁸ which will be denoted $2_{3/2}$ in the current work. Uttam *et al.*^{27,28} furthermore measured additional unidentified higher bands at about $23\,035$, $23\,256$ and $26\,015\text{ cm}^{-1}$, which they denoted as C_1 , C_2 and D , respectively.

Theoretically, excited states arising from the $\text{Yb}(4f^{14}6p^1)\text{F}$ and $\text{Yb}(4f^{14}5d^1)\text{F}$ configurations were considered by Nayak and Chaudhuri⁶ with RAS-CI based on 4-component spinors, yielding the $A^2\Pi_{1/2}$ ($3_{1/2}$), $A^2\Pi_{3/2}$ ($2_{3/2}$), and a $2^2\Sigma_{1/2}$ state. Earlier multi-reference CI calculations by Dolg *et al.*¹⁹ furthermore indicate the possibility of low-lying $\Omega = 1/2$ states arising from the $\text{Yb}(4f^{13}[F_{7/2}^\circ]\sigma_{6s}^2)\text{F}$ or $\text{Yb}(4f^{13}[F_{7/2}^\circ]5d^1\sigma_{6s}^1)\text{F}$ configurations, to be lying below or close to the $\text{Yb}(4f^{14}6p^1)\text{F}$ states. This was also found in the DFT calculations of Liu *et al.*²⁰ who place excited states arising from the $\text{Yb}(4f^{13}[F_{7/2}^\circ]\sigma_{6s}^2)\text{F}$ configuration in the range from 9000 to 15000 cm^{-1} relative to the $\text{Yb}(4f^{14}\sigma_{6s}^1)\text{F}$ ground state. These findings make it of interest to explicitly consider the configuration interaction between the f^{13} and f^{14} configurations in the Yb atom.²⁹

The vibronic states are additionally split due to hyperfine interactions. In atomic experiments they were measured for the ground and excited states^{9,22,30–33} also using Zeeman spectroscopy.³⁴ The hyperfine interaction of the atom^{35–37} and molecule³⁸ were studied theoretically, and should have similar uncertainties to the contribution of the eEDM to the spectrum due to the similarity of the matrix elements. Recently, uncertainties of the hyperfine constants arising in relativistic coupled cluster computations have been studied.³⁹

It is clear from the above that a proper description of the Yb atom and the YbF molecule requires an accurate treatment of both spin-orbit coupling and electron correlation, for ground as well as excited states. A popular approach is the so-called two-step approach to spin-orbit coupling (SOC), in which electron correlation methods based on non-relativistic or scalar relativistic Hamiltonians are used to obtain excited state energies, that are in turn used to dress a spin-orbit configuration interaction (SOC-ICI) matrix. This approach can yield quite accurate spin-orbit coupled states, but results are particularly sensitive to the number of spin-free states serving as a basis for the SOC-ICI

step.^{40–46} An alternative is to include SOC already at the mean-field level, and use fully SO-coupled molecular spinors to construct the correlated wave functions.⁴⁷ This can be done with four-component Hamiltonians, as done for the ground¹ and excited states^{6,37,48,49} of YbF, or with more computationally efficient two-component Hamiltonians based on the eXact 2-Component (X2C) approach,^{50–57} in which a transformation to decouple the positive and negative energy states of the Dirac Hamiltonian can be carried out in matrix form, yielding the same positive energy spectrum as the original 4-component Hamiltonian. More details can be found in the recent review by Liu.⁴⁷ Among the different X2C flavors, we can distinguish two main strategies for the decoupling, which is performed based on: (i) the one-electron Dirac Hamiltonian prior to the mean-field step,^{54,57} and for which two-electron spin-orbit contributions due to the untransformed two-electron potential are included *via* atomic mean-field contributions calculated with the AMFI code^{41,58,59} (X2C-AMFI); (ii) after a converged 4-component mean-field calculation on atoms^{52,53,56} or molecules⁵⁵ ($^2\text{DC}^M$). Recent benchmarks show that $^2\text{DC}^M$ calculations closely reproduce equivalent 4-component ones for valence⁶⁰ or core⁶¹ states.

Moreover, the aforementioned calculations for the excited states of YbF have mostly employed multireference CI (MRCI) approaches. While these can provide great flexibility in capturing static correlation, it remains the case that dynamical correlation is better accounted for with coupled cluster approaches. Among the coupled cluster singles and doubles (CCSD) approaches for excited states, we have the equation of motion (EOM-CCSD) method as well as Fock-space (IHFS-CCSD) methods,⁶² of which the single electron attachment, detachment, and singly excited states variants are the most commonly used. The two approaches have been found to yield very accurate results in general and in particular for calculations with relativistic Hamiltonians as discussed elsewhere (see ref. 60 and references therein).

The first goal of this work is therefore to go beyond the investigations performed to date in the literature, and apply the relativistic EOM-CCSD and IHFS-CCSD approaches to describe the low-lying excited states of YbF. For such states, where the most important excited state configurations appear have a single open-shell character ($4f^{14}\sigma_{6s}^1$, $4f^{13}[F_{7/2}^\circ]\sigma_{6s}^2$, $4f^{14}5d^1$, $4f^{13}[F_{7/2}^\circ]5d^1\sigma_{6s}^1$, $4f^{14}6p^1$), these coupled cluster approaches are in principle applicable, provided one starts from closed shell configurations such as $\text{Yb}(4f^{14}\sigma_{6s}^2)\text{F}^-$ or $\text{Yb}(4f^{14})\text{F}^+$. Additionally, we assess the performance of relativistic MRCI with respect to the coupled cluster methods. Our second goal is to confirm whether any low-lying state is close enough to the ground state to perturb the latter, and explain the anomalous behavior observed in the open-shell ground-state calculations in the literature.

2 Computational details

All relativistic electronic structure calculations were performed with a development version of the Dirac program suite⁶³



(revision 6e10c5d3), employing for Yb the valence double-zeta (24s19p13d9f2g), triple-zeta (30s24p18d14f4g2h) and quadruple-zeta (35s30p19d16f6g4h2i) basis sets from the previous work,¹ along with the matching augmented correlation-consistent (aug-cc-pVnZ, $n = 2, 3, 4$) basis sets of Dunning⁶⁴ for F. All basis sets were kept uncontracted, with the small component basis generated by restricted kinetic balance. In addition to these individual basis sets, we have used the calculations with triple- and quadruple-zeta sets to construct extrapolations to the complete basis set limit (E_∞) for the underlying potential energy curves, using the relation⁶⁵

$$E_\infty(\mathbf{R}) = \frac{4^3 E_4(\mathbf{R}) - 3^3 E_3(\mathbf{R})}{4^3 - 3^3} \quad (1)$$

where the subscripts denote the cardinal numbers for the basis sets and $E_n(\mathbf{R})$ the energy for a given geometry and electronic structure method for a basis of cardinal number $n (= 2, 3, 4)$.

In the coupled cluster computations the ²DC^M Hamiltonian^{54,55} was applied, all two-electron integrals over small component (S) basis sets (*i.e.* the so-called (SS|SS)-type integrals) appearing in the SCF step have been replaced by a simple correction.⁶⁶ In order to account for spin-orbit coupling and other relativistic effects the X2C-AMFI Hamiltonian was employed for the Kramers-restricted configuration interaction (KRCI) method.

Spectroscopic constants (r_e , D_e , ω_e and $\omega_e \chi_e$) were determined from a Morse potential fit in the vicinity of the potential energy minima. The potential energy curves were determined for bond lengths between 1.6 Å and 2.3 Å spaced by 0.02 Å and additional points with larger spacing up to 3.5 Å. In the calculation of D_e the asymptotic dissociation limit is calculated from the energies of the isolated neutral atoms, F in the ²P_{3/2} state and Yb in the ¹S₀ state.

The dataset associated with this manuscript (outputs from calculations, codes to extract and process information from these, and code to obtain the spectroscopic constants) is provided in ref. 67.

2.1 Kramers-restricted configuration interaction

For YbF we first consider Kramers-restricted configuration interaction (KRCI) based on an average-of-configuration Hartree-Fock approach (AOC-SCF).⁶⁸ This method was employed in order to treat the open shells, where one or two valence electrons were distributed over the s- and d-orbitals and the f-shell was either completely filled or contained one hole, depending on the states of interest. The AOC-SCF reference wave function in the KRCI computation is occupied according to a definition given by a generalized active space (GAS).⁶⁹ In this approach the Hamiltonian is computed for all allowed configurations and then diagonalized. The GAS space was defined by a f-shell which was completely filled or contained one hole and one or two electrons distributed over 29 orbitals.

2.2 Equation-of-motion coupled cluster

The first approach we use to describe the dynamical correlation that is largely missing in KRCI is EOM-CCSD, which can give

access to electronic states of different kinds, depending on the single determinant wave function that is chosen as the starting point. In it, the CCSD amplitudes are determined for the chosen ground state in the first step, and subsequently the similarity transformed Hamiltonian is constructed using these amplitudes and the desired states are generated by an operator that either removes or adds an electron.

The first set of states was obtained by electron attachment on Yb(4f¹⁴F⁺ ion, where the HOMO ($\sigma_{6s,1/2}$) of YbF was initially empty. This computation on the (0h,1p) sector of Fock space yielded states with 4f¹⁴ and a valence electron in the σ_{6s} , d or p orbital. This means that, in the process of obtaining the potentials for the ground and excited states of YbF, we immediately obtain energies of CCSD quality for YbF⁺, and therefore vertical ionization potentials (IP) at each geometry.

Another set of states was obtained by ionizing the Yb(4f¹⁴ σ_{6s}^2)F⁻ anion, where the HOMO ($\sigma_{6s,1/2}$) of YbF was initially doubly occupied. States arising from the Yb(4f¹⁴ σ_{6s}^1)F, Yb(4f¹³[F_{7/2}] σ_{6s}^2)F, Yb(4f¹³[F_{5/2}] σ_{6s}^2)F and Yb(4f¹⁴ σ_{6s}^2)F(2p⁵) configurations were obtained by considering the (1h,0p) sector of Fock space. This means that, in the process of obtaining the potentials for the ground and excited states of YbF, we immediately obtain energies of CCSD quality for YbF⁻, and therefore vertical electron affinities (EA) at each geometry. We note that states arising from the (2h,1p) and (1h,2p) manifolds are also accessible from EOM-IP and EOM-EA calculations, though the energy of electronic states determined by such configurations will be less accurate than states dominated by single detachment or attachment configurations.

The EOM-CCSD electronic states are obtained by an iterative diagonalization (Davidson) procedure in which only the energies of a certain number of the lowest states are determined. For the IP-EOM-CCSD we obtained 16 $\Omega = 1/2$, 8 $\Omega = 3/2$, 6 $\Omega = 5/2$ and 2 $\Omega = 7/2$ states, whereas for EA-EOM-CCSD we obtained 8 $\Omega = 1/2$, 6 $\Omega = 3/2$, 4 $\Omega = 5/2$ and 2 $\Omega = 7/2$ states.

As transition moments are not yet available for the EOM-CCSD implementation in Dirac, we have only obtained the potential energy curves. These are nevertheless useful since, by not requiring the definition of model spaces or the use of an extrapolation procedure, they serve as a cross validation of the IHFS-CCSD calculations below.

2.3 Fock-space coupled cluster

Fock-space coupled cluster⁷⁰ (FS-CCSD) is our second approach to include dynamical correlation in the electronically excited states. Here it was employed in a similar fashion to EOM-CCSD, starting from YbF⁺ or YbF⁻ and proceeding to the (0h,1p) and (1h,0p) sectors of Fock space, respectively. For FS-CCSD a model space is defined by selecting a number of occupied and virtual orbitals and how many electrons are added and removed. The matrix for this subspace is constructed and subsequently diagonalized, thus yielding all states within the chosen model space, in this case states arising from single electron attachment (EA) or single electron detachment (IP). This method requires solving first for the underlying sectors,



starting with (0h,0p), which corresponds to CCSD. Due to computational constraints, we have truncated the virtual space so that 117, 230 and 296 orbitals were used in the double-, triple- and quadruple-zeta CCSD calculations, respectively.

The separation into a model and external space leads to the appearance of the so-called intruder states, a well-known difficulty with Fock-space coupled cluster and other effective Hamiltonian approaches, that can be dealt with in many cases by the intermediate Hamiltonian (IH) Fock-space coupled cluster (IHFS-CCSD) method.^{71,72}

The IH approach was employed to compute $\text{Yb}(4f^{14}\{\sigma_{6s}^1 6p, 5d, \dots\}^1)\text{F}$ states starting from YbF^+ . The active P space in such calculations contained about 50 spinors varying slightly with bond distance and basis set. Of these 26 spinors are always present in the model (P_m) space, whereas the remaining active spinors are placed in the intermediate (P_I) space. Due to using the (0h,1p) sector for the cation, states arising from configurations where the Yb 4f shell is partially filled (such as $\text{Yb}(4f^{13}[\text{F}_{7/2}^{\circ}]\sigma_{6s}^2)\text{F}$, $\text{Yb}(4f^{13}[\text{F}_{7/2}^{\circ}]5d^1\sigma_{6s}^1)\text{F}$ or $\text{Yb}(4f^{13}[\text{F}_{7/2}^{\circ}]\sigma_{6s}^1 6p^1)\text{F}$) are not accessible in this calculation.

The approach outlined above was, however, not enough to avoid divergence in the computation of the (1h,0p) sector using the anion as a reference. Therefore, the extrapolated intermediate Hamiltonian (XIH) Fock-space coupled-cluster approach⁷³ (XIHFS-CCSD) was applied using the same shifts as in ref. 73. Values of 0.1 and 0.2 Hartree were selected if one of the holes is not in the model space. These shifts were doubled for two holes outside the model space. Using the determined energies an extrapolation to the system without shifts was performed. The model (P_m) space in these computations contained 22 spinors, the intermediate (P_I) space about 24 spinors depending on the bond distance and basis set. Since we start out from the anion and only allow holes, only $\text{Yb}(4f^{14}\sigma_{6s}^1)\text{F}$, $\text{Yb}(4f^{13}[\text{F}_{7/2}^{\circ}]\sigma_{6s}^2)\text{F}$, $\text{Yb}(4f^{13}[\text{F}_{5/2}^{\circ}]\sigma_{6s}^2)\text{F}$, and $\text{Yb}(4f^{14}\sigma_{6s}^2)\text{F}(2p^5)$ configurations are accessible in this computation.

Combining the two sectors allows us to get different excited states of YbF, although there are limitations. Firstly, the interactions between configurations with open f-shell and the ones with an electron in the p- or d-shell are not included, since they will be obtained for different sectors of Fock space. This interaction will nevertheless be treated with a simple adiabaticization approach, described in Sections 2.4 and 3.6. Secondly, configurations such as $\text{Yb}(4f^{13}[\text{F}_{7/2}^{\circ}]\sigma_{6s}^1 6p^1)\text{F}$ or $\text{Yb}(4f^{13}[\text{F}_{7/2}^{\circ}]5d^1\sigma_{6s}^1)\text{F}$ are not included in the current treatment. This limitation is not as significant because these states have higher energies than the ones we are interested in. Both of these problems could be dealt with by using the (1h,1p) sector, but this goes beyond the current work as convergence is very unstable for this sector and it requires the use of an open-shell reference.

2.4 Adiabaticization of electronic states

As we separated the computations of states with $4f^{13}$ and $4f^{14}$ character, these states cannot interact with each other, and states with the same Ω value cross although they should have

an avoided crossing. In order to correct this deficiency we considered a simple adiabaticization model, in which we set up and diagonalize the following matrix for each Ω value:

$$M = \begin{pmatrix} E(f^{13}) & C \\ C & E(f^{14}) \end{pmatrix} \quad (2)$$

where C is a matrix where every entry is a coupling constant (whose value is kept constant for all states and geometries considered; we have investigated values of 0.01, 0.001 and 0.0001 a.u.), E are matrices with the eigenvalues of the different electronic states on the diagonal. The potential curves were computed for different coupling constants and the results are shown in Section 3.6. We note that since the ground state energy, associated with the $\text{Yb}(4f^{14}\sigma_{6s}^1)\text{F}$ configuration, appears in both coupled cluster approaches, we have only considered one such energy. As we shall see in the discussion, this is valid in the region between 1.8 Å and 2.5 Å, since for these distances the ground-state energies from IP and EA calculations are nearly identical.

3 Results and discussion

We start our discussion with the electronic transitions of the atomic Yb^+ cation, before moving on to the YbF molecule. This is because the cation's electronic structure is similar to the Yb in YbF since, due to the large electronegativity of fluorine, one electron is almost completely removed from the Yb atom.

Subsequently, the potential energy curves for Kramers-restricted configuration interaction are presented. A discussion of the coupled cluster approaches follows, with a focus on the comparison of the coupled cluster results for the Fock space and equation-of-motion approach. This section is followed by a presentation of the spectroscopic data for the ground and excited states. In the last part we take a closer look at the mixing of states at around 18 000 cm^{-1} and apply the adiabaticization procedure.

3.1 Ytterbium cation

As discussed in the introduction, states from both the $4f^{13}$ and the $4f^{14}$ configurations are of importance. This is difficult to realize in a balanced manner when using one set of orbitals to describe all states. Any change in the occupation of the 4f-shell will alter the screening of the 5s and 5p orbitals of Yb, resulting in differences between orbital sets optimized for a $4f^{13}$ or a $4f^{14}$ configuration. Additionally, the 4f orbitals are very compact and since they are the first f-shell there are no orthogonality conditions limiting the radial expansion or contraction of the orbital. Depictions of the orbitals for both configurations can be found in Table S1 in the ESI.[†]

These observations help to understand why it turned out to be very difficult to treat both sets of states in the same calculation, which we attempted to do from AOC-SCF on the Yb^+ . We started out by performing AOC-SCF computations on the atom, based on the 4f closed shell configuration. While we obtained the correct ground state configuration, the $^2\text{F}_{7/2}^{\circ}$ has an energy of about 46 000 cm^{-1} (over two times higher than the



Table 1 Kramers-restricted configuration interaction transition energies (in cm^{-1}), squared transition dipole moments (TDM^2), and line strength (S) for the Yb^+ cation, the latter two are in atomic units ($e^2 a_0^2$). Reference values and notation have been taken from the NIST database.⁷⁴ For the $4f^{13}$ configurations, energies relative to the $^2F_{7/2}^{\circ}$ state are also given. 2z, 3z, 4z, and extr. indicate double, triple, quadruple zeta and extrapolated results, respectively

State	Conf.	NIST ⁷⁴		2z		3z		4z		extr.
		E	S	E	TDM ²	E	TDM ²	E	TDM ²	E
$^2S_{1/2}$	$4f^{14}6s^1$	0		0		0		0		0
$^2D_{3/2}$	$4f^{14}5d^1$	22 961		23 322	0.0	22 802	0.0	23 606	0.0	24 192
$^2D_{5/2}$	$4f^{14}5d^1$	24333		23 882	0.0	23 321	0.0	24 117	0.0	24 698
$^2P_{1/2}^{\circ}$	$4f^{14}6p^1$	27062	6.1	25 210	3.5	24 533	3.8	25 331	3.6	25 914
$^2P_{3/2}^{\circ}$	$4f^{14}6p^1$	30392	11.4	28 104	16.9	27 385	18.9	28 153	17.4	28 712

State	Conf.	E	ΔE				
$^2F_{7/2}^{\circ}$	$4f^{13}6s^2$	21 419	0	0	0	0	0
$^3[3/2]_{5/2}^{\circ}$	$4f^{13}5d^16s^1$	26 759	5340	4260	5538	4618	3946
$^3[3/2]_{3/2}^{\circ}$	$4f^{13}5d^16s^1$	28 758	7339	6387	7822	7123	6613
$^3[11/2]_{9/2}^{\circ}$	$4f^{13}5d^16s^1$	30 224	8806	8214	9325	8314	7576
$^3[11/2]_{11/2}^{\circ}$	$4f^{13}5d^16s^1$	30 563	9144	8320	9431	8447	7729

experimental value), and the wrong order for the hole states is observed. If the wave function is optimized for a $4f^{13}$ configuration, one obtains the $^2F_{7/2}^{\circ}$ as the lowest state and finds the true ground state more than $20\,000\text{ cm}^{-1}$ higher. Because of these difficulties, the KRCI calculations discussed below all follow the strategy of different orbital sets that is also employed in the subsequent coupled cluster calculations.

Table 1 contains KRCI values of electronic transitions for the cation. The transition energies show deviations of about 10% and the spin-orbit splitting is underestimated for states with a $4f^{14}$ configuration. The squared transition dipole moment (TDM^2) of the $^2P_{1/2}^{\circ}$ state is underestimated by about 13%, the one for the $^2P_{3/2}^{\circ}$ state overestimated by about 47%. The second set of states with a hole in the f-shell and different distributions of the 2 valence electrons are given in the lower part of Table 1, the energies are relative to the $^2F_{7/2}^{\circ}$ state. In this case the two valence electrons are distributed over the s- and d-shell. The lowest state with a $4f^{13}[F_{7/2}^{\circ}]6s^16p^1$ configuration has a transition energy of 47912.31 cm^{-1} and was not included in the current treatment.

The excited states of the Yb atom for the $4f^{14}$ configuration have already been investigated by relativistic Fock-space coupled cluster^{36,37,75,76} as well as for the cation,⁷⁵ including the transition moment of magnetic transitions.⁷⁶ With our current calculations we can go beyond these studies and investigate the f^{13} configurations as well. Before discussing our IHFS-CCSD calculations for Yb^+ , we focus on the EOM-CCSD excitation energies, shown in Table 2. The EOM-IP-CCSD energies of $4f^{13}$ states obtained from the extrapolation to the complete basis set limit underestimate the experimental transition energies by around 3000 cm^{-1} , whereas the values for $4f^{14}$ states, obtained with EOM-EA-CCSD are within 1000 cm^{-1} of the experimental values, which yields a quantitative improvement over the KRCI ones for both configurations, even though qualitatively the two methods provide a similar picture. From that and the preceding discussion, we attribute the relatively lower accuracy for the $4f^{13}$ to arise from the incomplete account of the relaxation

Table 2 Transition energies (in cm^{-1}) for the Yb^+ cation, obtained for different basis set with EOM-IP-CCSD ($4f^{13}$) and EOM-EA-CCSD ($4f^{14}$), except for the ground state, for which both methods yield the same configuration and total energy. 2z, 3z, 4z, and extr. indicate double, triple, quadruple zeta and extrapolated results, respectively. Reference values were obtained from the NIST database⁷⁴

State	Conf.	NIST ⁷⁴	2z	3z	4z	Extr.
$^2S_{1/2}$	$4f^{14}6s^1$	0	0	0	0	0
$^2F_{7/2}^{\circ}$	$4f^{13}6s^2$	21 419	12 054	13 524	16 092	17 966
$^2F_{5/2}^{\circ}$	$4f^{13}6s^2$	31 568	22 629	24 139	26 655	28 491
$^2D_{3/2}$	$4f^{14}5d^1$	22 961	24 073	24 209	24 060	23 951
$^2D_{5/2}$	$4f^{14}5d^1$	24 333	25 351	25 457	25 341	25 257
$^2P_{1/2}^{\circ}$	$4f^{14}6p^1$	27 062	27 539	27 780	27 857	27 913
$^2P_{3/2}^{\circ}$	$4f^{14}6p^1$	30 392	30 954	31 246	31 323	31 380

of the wave function upon the creation of the hole in the f shell. Beyond the states presented in Table 2, which are dominated by single electron attachment and detachment, we are able to access states with significant (1h,2p) and (2h,1p) character with EOM-CCSD. These states, available in Table S14 in the ESI,[†] are about $10\,000\text{ cm}^{-1}$ higher in energy than the experimental ones.

Finally, our IHFS-CCSD results are presented in Table 3. The transition energies for $4f^{14}$ configuration reproduce well the experimental ones, with errors below 6%, and only show a small dependence on the basis set. The states arising from the $4f^{13}$ configuration ($4f^{13}[F_{7/2}^{\circ}]s^2$ etc.), in contrast, show a significant dependence on the basis and a rather slow convergence and underestimate the value by about 30%, which makes them less accurate than the EOM-CCSD ones. This lower accuracy is a consequence of the reduced flexibility in the model spaces, due to the need of adding the 5p-shell just below the 4f-shell to the intermediate space, in order to achieve convergence. These results are in line with the observations of Shee *et al.*,⁶⁰ in that the formal equivalence between EOM-CCSD and IHFS-CCSD for the sectors of Fock space considered depends, in fact, on the flexibility of the main model space.



Table 3 Transition energies for the Yb⁺ cation. Reference values have been obtained from the NIST database,⁷⁴ the computed values were obtained for different basis set sizes with Fock-space coupled cluster

State	Conf.	NIST ⁷⁴	2z	3z	4z	Extr.	DCB ⁷⁵
² S _{1/2}	4f ¹⁴ 6s ¹	0	0	0	0	0	0
² F _{7/2}	4f ¹³ 6s ²	21 419	11 087	12 390	13 618	14 514	
² F _{5/2}	4f ¹³ 6s ²	31 568	21 631	22 976	24 170	25 042	
² D _{3/2}	4f ¹⁴ 5d ¹	22 961	24 058	24 223	24 059	23 938	23 720
² D _{5/2}	4f ¹⁴ 5d ¹	24 333	25 336	25 469	25 340	25 246	24 998
² P _{1/2}	4f ¹⁴ 6p ¹	27 062	27 518	27 774	27 851	27 907	27 870
² P _{3/2}	4f ¹⁴ 6p ¹	30 392	30 934	31 241	31 316	31 371	31 312

Furthermore, the removal of the 5p spinors from the main model space underscores the importance of the 5p for the energetics of the states with a hole in the 4f shell, since by doing so, we undress the contributions from the 5p configurations, and thus prevent them from interacting effectively with 4f¹³ determinants.

3.2 Kramers-restricted configuration interaction potential energy curves

In Fig. 1 the potential energy curves obtained by an approach corresponding to the one used for Yb⁺ are shown for f¹⁴ and f¹³ configurations. For separate potential energy curves and transition dipole moments we refer the interested reader to the ESI[†] (Fig. S2–S7).

The potential energy curves have been determined up to 15 Å for the closed and open f-shells. The energy difference in the atom between these states is 21418.75 cm⁻¹. Accordingly, the PECs for the hole states were shifted to obtain this separation at this distance. There is still some interaction between ytterbium and fluorine at 15 Å, but the long range behaviour can be expected to be similar for the two configurations (this assumption was checked, see Fig. S8 in the ESI[†] for further details). Taking into account the position of the minima, the curvatures, spin-orbit splitting, the avoided crossings and asymptotes the states can be assigned to a dominant configuration, shown in Fig. 1.

Regarding the 4f¹⁴ manifold, the lowest two excited states in the figure belong to the Yb(4f¹⁴6p¹)F configuration, but approach asymptotically the ²D_{3/2} state. The asymptote of the next three states is ²D_{5/2} corresponding to the Yb(4f¹⁴5d¹)F configuration for smaller internuclear separations. For Ω = 5/2 the transition dipole moment with the ground state is zero, for the other four the values are shown in Fig. S3 in the ESI[†]. The first Ω = 3/2 and the third Ω = 1/2 state have a larger transition dipole moment close to equilibrium, but get close to each other at the largest internuclear separations.

Regarding states of the 4f¹³ manifold, the lowest four states belonging to the Yb(4f¹³[F_{7/2}^o]σ_{6s}²)F configuration are well separated from a dense region with a lot of states about 12 000 cm⁻¹ higher. Most of these states are of the Yb(4f¹³[F_{7/2}^o]5d¹σ_{6s}¹)F configuration, with the Yb(4f¹³[F_{5/2}^o]σ_{6s}²)F state slightly higher

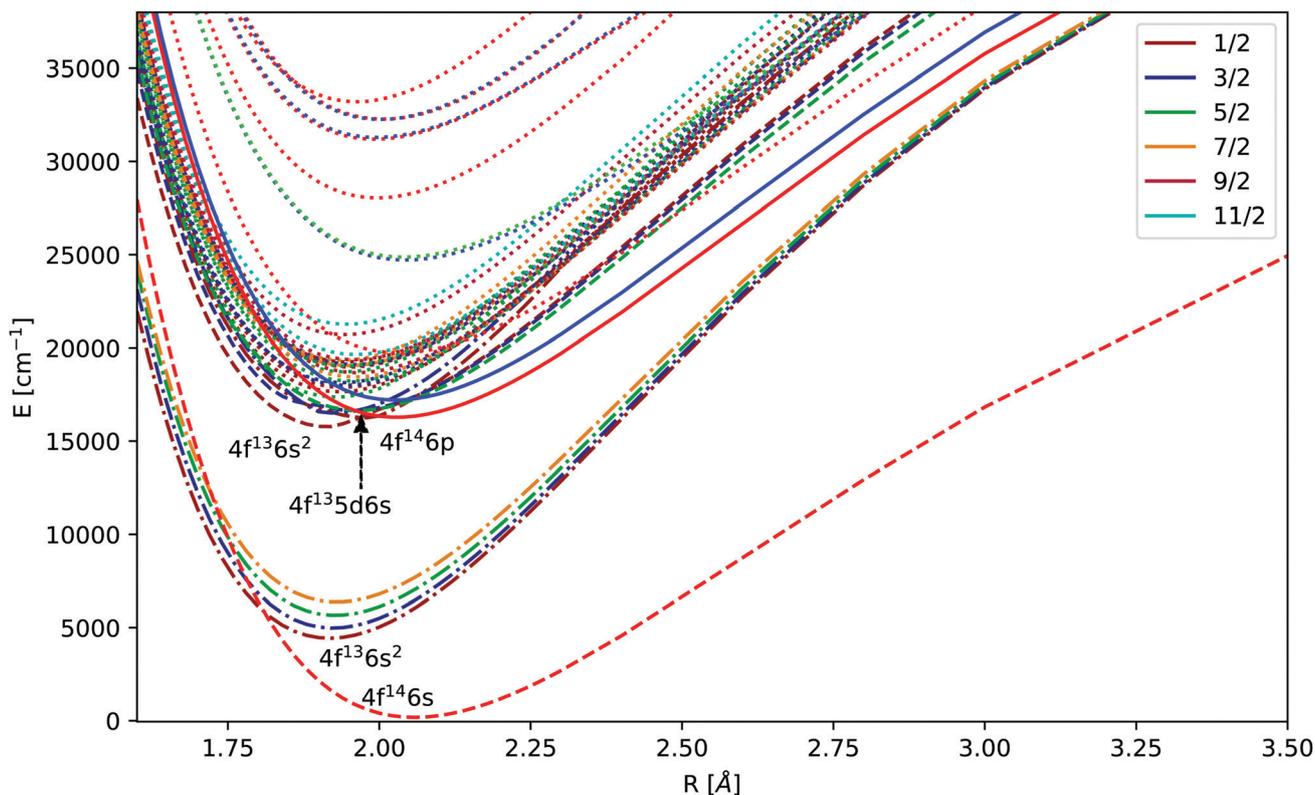


Fig. 1 Combination of the sets of KRCI potentials obtained by extrapolating triple and quadruple zeta basis sets. The lowest Ω = 1/2 states are denoted by their dominant configuration.



in energy asymptotically and more strongly bound, resulting in several avoided crossings. For each of the four $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]\sigma_{6s}^2)F$ states the transition dipole moments with higher excited states of varying Ω are plotted in Fig. S6 in the ESI.† The transition dipole moments are substantially smaller than the ones for the closed 4f-shells but some of them are non-zero.

An alternative to AOC-SCF for obtaining orbitals for several configurations is multiconfigurational SCF, but similar difficulties as for AOC-SCF in obtaining a balanced description of the $4f^{14}$ and $4f^{13}$ states are observed: either the wrong ground state is obtained (if only the hole states are optimized in MCSCF), or the hole states are too high in energy by about $20\,000\text{ cm}^{-1}$ (if the ground state is optimized). We also made attempts using state-averaged MCSCF in a non-relativistic quantum chemistry code and observed the same difficulties (see dataset⁶⁷). If the $4f^{13}$ configurations are excluded one obtains meaningful results, but at the expense of obtaining a $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]\sigma_{6s}^2)F$ states too high in energy. If all the states are included, the wrong ground state is obtained.

3.3 Coupled cluster potential energy curves

The potential energy curves of excited states obtained by the equation-of-motion and Fock space methods are displayed in Fig. 2, the values for the complete basis set limit are shown. The basis set dependence in the molecule is similar to the one observed for Yb^+ : energies for $4f^{14}$ states depend only weakly on the basis set, while the gap between the ground state and the excited states corresponding to the $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]\sigma_{6s}^2)F$ configuration increases upon improving the basis sets.

While the EOM-CCSD excitations energies of Yb^+ are closer to the experimental ones, the $\text{Yb}(4f^{13}[F_{5/2}^{\circ}]\sigma_{6s}^2)F$ states are too high to perturb the $\text{Yb}(4f^{14}6p^1/5d^1)F$ PECs. From the extended potential energy curves provided in the ESI† (Fig. S15 and S16), we can observe that the ground state of the non-interacting system ($\text{Yb}(4f^{14}\sigma_{6s}^2)F(2p^5)$) is repulsive and has a high energy at the equilibrium distance. This results in several avoided crossings being observed at 3, 3.5, and 5 Å.

Since the $\text{Yb}(4f^{14}\sigma_{6s}^2)F$ ground state is accessible for both sectors employed in the coupled cluster calculations (Fock space as well as EOM), we can assess the compatibility of the two separate sets of calculations (in the sense of having comparable accuracies) by looking more closely at the differences between the ground states in Fig. 2. From that, we can see that the EOM-CCSD and IHFS-CCSD approaches the curves are on top of each other from the smallest considered internuclear separation up to about 2.8 Å. This assures us that there should not be artifacts in putting together and comparing the calculations on the two sectors.

3.4 Dissociation and ionization energies

Since the (1h,0p) and (0h,1p) sectors have been considered in our EOM-CCSD and IHFS-CCSD calculations, we have as a by-product of our calculations the ionization potentials (IP) and electron affinities (EA) for YbF for all computed distances. Therefore these quantities are presented first in Table 4, before proceeding to the spectroscopic constants.

Unlike coupled cluster calculations, for KRCI a consistent definition of active spaces is difficult, and its lack of size-consistency results in large deviations from experiment and from the coupled cluster values. For adiabatic electron affinities, for which to the best of our knowledge there are no experimental values, the extrapolated values are 8393 and 8197 cm^{-1} for EOM-CCSD and IHFS-CCSD, respectively. For a distance of 6.5 Å a value of $28\,651\text{ cm}^{-1}$ was obtained, which is reasonably close to the electron affinity of fluorine ($27\,432\text{ cm}^{-1}$).⁷⁷ Corresponding results for the atoms are listed in the table, which allow to calculate the dissociation energies (D_e). They deviate from the experimental values of $43\,600 \pm 800\text{ cm}^{-1}$ by Kaledin *et al.*⁷⁸ and $43\,260 \pm 800\text{ cm}^{-1}$ by Yokozeki and Menzinger.⁷⁹ The ionization potentials in Table 4 show acceptable agreement with experimental values.

3.5 Spectroscopic constants

The spectroscopic constants for the ground state are now considered. In Table 5 our results are summarized, along those from the literature.

We observe that the extrapolated KRCI bond distances, at about 2.058 Å , are significantly longer (by around 0.04 Å) than experiment,³⁰ whereas the coupled cluster calculations show differences from experiment smaller than 0.01 Å , with EOM-CCSD showing slightly larger discrepancies than IHFS-CCSD. Between the extrapolated EOM-CCSD and IHFS-CCSD, we also see small differences between the $4f^{14}$ and $4f^{13}$ for EOM-CCSD these differ by around 0.001 Å whereas for IHFS-CCSD the difference is slightly under 0.002 Å , with the $4f^{13}$ configuration yielding a slightly underestimated value, compared to experiment, something that can be traced back to the differences in model spaces for this configuration.

Our results for harmonic frequencies further indicate that KRCI seems to underestimate the bonding strength in YbF , as the harmonic vibrational frequency is smaller (491 cm^{-1}) than experiment (between 505.5 and 506.7 cm^{-1} depending on the experiment). The coupled cluster results, on the other hand, show the typical $5\text{--}6\text{ cm}^{-1}$ overestimation of the harmonic frequencies with respect to experiment (something also seen for the anharmonic constants), which can be attributed to lack of triples in the EOM or FS treatment, that would introduce further orbital relaxation. This can be seen in comparison to the unrestricted coupled cluster calculations of Gomes *et al.*,¹ which in spite of the large value of the T_1 diagnostic, reproduce well the experimental bond lengths, harmonic frequencies and anharmonic constants.

Taken together, our 2-component CCSD-based calculations and the 4-component ones of Gomes *et al.*¹ compare consistently better to experiment than the other theoretical works for bond lengths, vibrational frequencies and anharmonic constants. For the dissociation energies, on the other hand, the extrapolated calculations presented here do not provide a significant improvement over the results of prior theoretical investigations (quadruple zeta values are closer to the experimental ones for this quantity, see Table 4). Especially, electron attachment values are off, which might be related to



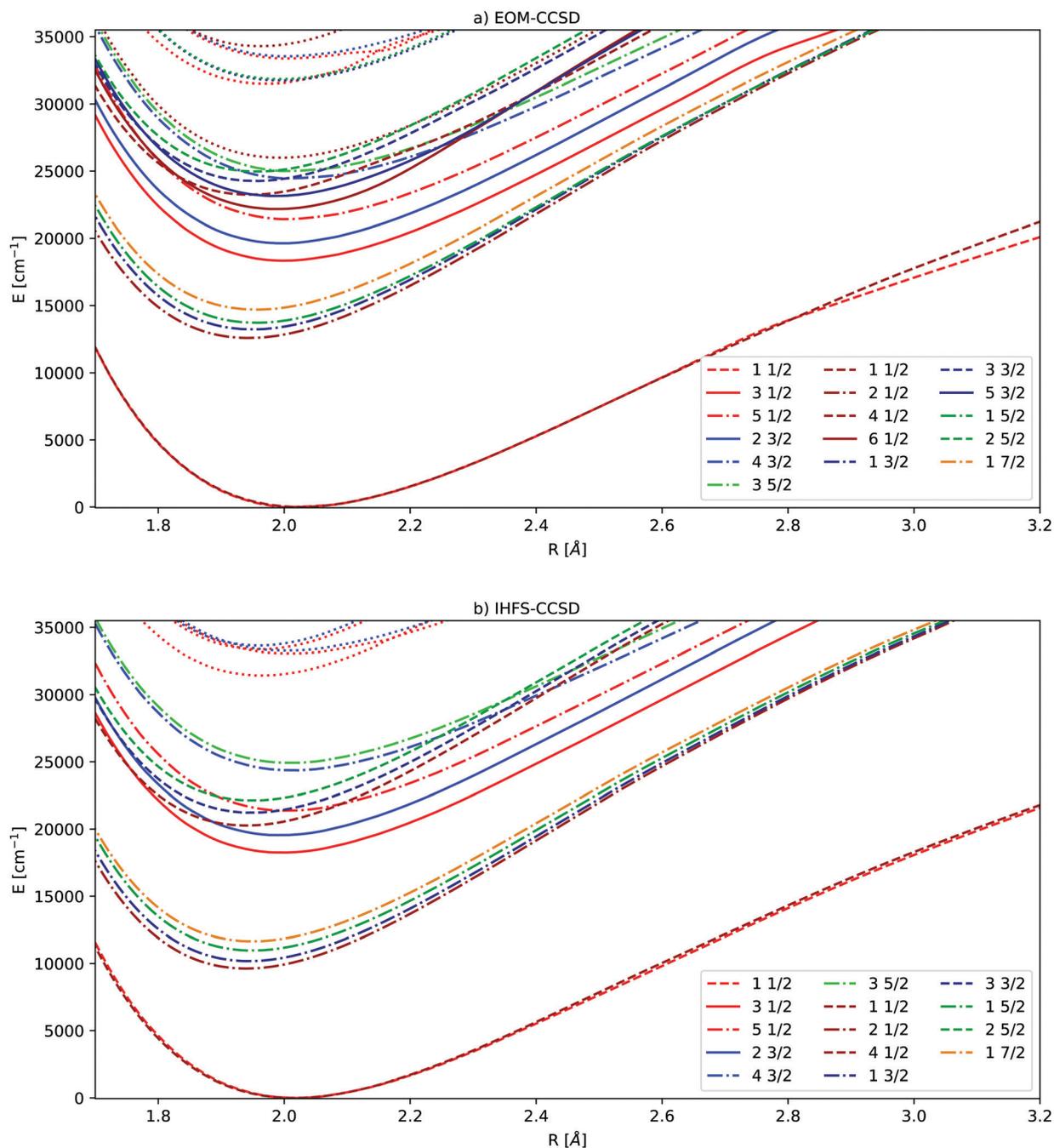


Fig. 2 Potential energy curves obtained by extrapolating triple and quadruple zeta basis sets. EOM-CCSD results in the upper part, IHFS-CCSD in the lower part. The Ω values of 1/2, 3/2, 5/2 and 7/2 are indicated by the colors red, blue, green and orange. The light colors are used for the (0h,1p) sector, dark ones for (1h,0p). For the states below 30 000 cm⁻¹ we employ the same color coding and state notation as in Fig. 1.

the absence of the configuration with a hole in the p orbitals of fluorine.

Moving now to excited states, we start by considering the four lowest excited states, which belong to the $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]\sigma_{6s}^2)F$ configuration. These states are well separated from the ground state (the lowest excited state is about 10 000 cm⁻¹ above the ground state) and higher excited states. That such states are quite well separated from the ground state would, in our view, tend to exclude the interaction with a low-lying

excited state as an explanation for the appearance of the large T_1 diagnostic values observed by Gomes *et al.*¹ From their spectroscopic constants, presented in Table 6, we see that with the exception of DFT all methods yield similar level splittings of about 500, 1200, and 2000 cm⁻¹. To the best of the authors knowledge there is no experimental data available for these states, due to their negligible transition dipole moments for dipole excitations (see for instance Fig. S6 in the ESI†) and small Franck-Condon factors due to the difference in bond



Table 4 Ionization potential (IP), electron affinity (EA) and dissociation energy (D_e) of Yb, F, and YbF. All values in cm^{-1} . They are listed for a quadrupole zeta basis set and a basis set extrapolation. The minimum of the potentials were determined using a Morse fit and used to compute the adiabatic values listed here

Quant.	System	KRCI		EOM-CCSD		IHFS-CCSD		Experiment
		4z	Extr.	4z	Extr.	4z	Extr.	
IP	Yb	38 406	39 128	50 735	50 822	50 740	50 837	50 443 ⁷⁴
IP2	Yb	90 581	90 926	97 919	98 035	97 918	98 040	98 232 ⁷⁴
IP	F	127 131	126 617	144 153	143 321	144 076	144 703	140 525 ⁷⁴
EA	F	13 326	11 978	27 279	27 740	27 246	27 759	27 432 ⁷⁷
IP	YbF	58 478	56 884	48 471	48 578	48 426	49 901	47 700 ⁷⁸
EA	YbF	7423	7326	9713	9876	9579	8197	
$D_e(\text{IP})$	YbF	26 059	25 887	43 824	47 782	40 591	40 931	43 260 ⁷⁹
$D_e(\text{EA})$	YbF	40 394	39 660	45 534	49 629	40 430	49 053	43 260 ⁷⁹

Table 5 Spectroscopic constants for ground state parameters for different approaches. Dissociation energies (D_e), harmonic frequencies (ω_e) and anharmonicity constants ($\omega_e\chi_e$) are given in cm^{-1} , the equilibrium bond distances (r_e) in Å. For the theoretical results we listed the values obtained by extrapolating triple and quadruple zeta basis sets (CBS)

Method	Ref.	r_e	ω_e	$\omega_e\chi_e$	D_e
KRCI	YbF	2.0829	465	2.40	39 660
EOM-CCSD	YbF ⁺	2.0230	511	2.80	49 629
	YbF ⁻	2.0250	508	2.53	47 782
IHFS-CCSD	YbF ⁺	2.0176	515	2.82	49 053
	YbF ⁻	2.0159	513	2.42	40 931
CCSD ¹	YbF	2.0174	507.6	2.357	40 904
CCSD(T) ¹	YbF	2.0289	528.2	1.939	41 156
CCSD ²¹		2.0127	566.8	3.7885	55 650
RASCI ⁶		2.051	529		
CCSD(T) ⁸⁰		2.03			38 900
CISD ⁸¹		2.034	502		42 100
DFT ²⁰		1.987	532		45 000
Exp. ⁷⁹					43 260
Exp. ¹⁸		2.0158	506.6674	2.2452	
Exp. ²⁷			505.5	1.9	
Exp. ⁷⁸					43 600
Exp. ³⁰		2.016514			
Exp. ²⁴		2.0195	506.616	2.235	

distances between these states and the ground state, see Section 3.6.

The smallest equilibrium distance was obtained for the $2_{1/2}$ state with 1.94 Å for the coupled cluster methods and 0.02 Å less for KRCI. The vibrational frequencies are between 570 and 600 cm^{-1} for the coupled cluster methods and about 30 cm^{-1} higher for KRCI.

For higher excited states, as apparent from the figures in the previous section, the identification and assignment of states gets more difficult and there are differences between the methods. We have nevertheless provided in Table 7 the spectroscopic constants for excited with Ω values of 1/2, 3/2, and 5/2, respectively.

The comparison with experimental results allows assignment of the lowest excited state reported in experiments and give some indications for higher states. The lowest $\Omega = 1/2$ state observed in experiment can be identified as the $3_{1/2}$ state ($5_{1/2}$ for KRCI). Spectroscopic parameters agree well with the ones obtained by fitting to the $A^2\Pi_{1/2}$ in experiments. A bond distance of 1.9935 Å

Table 6 Spectroscopic constants for the lowest excited states Yb($4f^{13}[F_{7/2}^2\sigma_{6s}^2]F$) for different wave function methods using the potential energy curves extrapolated to the basis set limit. In the case of KRCI and MRCI¹⁹ the ground state is not included in the computation and absolute transition energies are not available. The transition energy (T_e), level splitting (T_{rel} , energy relative to $2_{1/2}$), harmonic frequencies (ω_e) and anharmonicity constants ($\omega_e\chi_e$) are given in cm^{-1} , the equilibrium bond distances (r_e) in Å

State	Method	T_e	r_e	ω_e	$\omega_e\chi_e$	T_{rel}
$2_{1/2}$	KRCI		1.9200	631	2.51	
	EOM-CCSD	12 568	1.9432	591	2.59	
	IHFS-CCSD	9627	1.9396	599	2.79	
	DFT ²⁰	3790	1.9570	561		
$1_{3/2}$	MRCI ¹⁹		1.9480	600		
	KRCI		1.9253	628	2.50	540
	EOM-CCSD	13 211	1.9494	588	2.61	643
	IHFS-CCSD	10 180	1.9438	595	2.79	553
$1_{5/2}$	DFT ²⁰	9520	1.9440	597		5730
	MRCI ¹⁹		1.951 0	598		428
	KRCI		1.9296	622	2.45	1223
	EOM-CCSD	13 703	1.9553	582	2.61	1135
$1_{7/2}$	IHFS-CCSD	10 968	1.9493	589	2.78	1341
	DFT ²⁰	10 970	1.9360	598		7180
	MRCI ¹⁹		1.9540	594		1021
	KRCI		1.9315	616	2.43	1933
	EOM-CCSD	14 685	1.9556	577	2.62	2117
	IHFS-CCSD	11 645	1.9496	583	2.77	2018
	DFT ²⁰	16 530	1.936	592		12 740
	MRCI ¹⁹		1.954	589		1709

obtained by fitting to the same states in ref. 34 agrees well with the coupled cluster values for the $3_{1/2}$ state, the vibrational constant of about 540 cm^{-1} is close to the experimental value of ref. 18. Similarly, the lowest $\Omega = 3/2$ state reported by Dunfield *et al.*¹⁸ can be identified as the $2_{3/2}$ state ($4_{3/2}$ for KRCI), see Table 7.

The lowest states with $\Omega = 1/2$ and $\Omega = 3/2$ for this energy range approach asymptotically a state with a Yb($4f^{14}5d^1$)F configuration, but if one analyses the EOM-CCSD and IHFS-CCSD orbital composition, significant contributions of the atomic 6p are identified. The $\Omega = 3/2$ state is dominated (97%) by a single configuration, corresponding to a HOMO($\sigma_{6s,1/2}$) \rightarrow LUMO+1 where the latter is made up of a mixture of $6p_\pi$ and $5d_\pi$ orbitals (the $6p_\pi$ contributions being the dominant – $\approx 80\%$ – in the reference YbF⁺ orbitals). The few other significant configurations arise from excitations to higher-lying orbitals with increasingly



Table 7 Spectroscopic constants for selected excited states (complete list presented in Table S17 in the ESI) with $\Omega = 1/2, 3/2, 5/2$, starting from 18 000 cm^{-1} for different methods using the values after extrapolation to the basis set limit. Transition energy (T_e), vibrational constant (ω_e), and anharmonicity constant ($\omega_e x_e$) are given in cm^{-1} , the equilibrium bond distance (r_e) in Å. Experimental transitions that were not assigned (n.a.) are also listed. Labels of experimental results are defined in the introduction

Ω	Method	State	Configuration	T_e	r_e	ω_e	$\omega_e x_e$
1/2	KRCI ^a	3	$4f^{13}\sigma_{6s}^2$	15 572	1.9038	655	13.57
		5	$4f^{14}6p^1$	16 189	2.0504	496	2.38
		10	$4f^{14}5d^1$	19 631	2.0552	490	2.49
	EOM-CCSD	3	$4f^{14}6p^1$	18 373	2.0004	536	2.72
		4	$4f^{14}5d^1$	21 448	2.0079	532	2.78
		6	$4f^{13}\sigma_{6s}^2$	23 241	1.9432	582	4.06
		IHFS-CCSD	3	$4f^{13}\sigma_{6s}^2$	18 249	1.9953	539
	4		$4f^{13}\sigma_{6s}^2$	20 267	1.9397	597	2.78
	5		$4f^{14}5d^1$	21 375	2.0032	533	2.73
	MRCI ¹⁹		$4f^{13}\sigma_{6s}^2$		1.948	600	
		Exp. ^{18 a}	3		18106.20		537
	Exp. ¹⁸	4	[18.6] _{1/2}		18705.06		
[557]				18 574	1.9656	502.15	
[561]				18 699	1.9571		
3/2	KRCI ^a	2	$4f^{13}\sigma_{6s}^2$	16 206	1.9331	711	8.45
		4	$4f^{14}6p^1$	17 123	2.0473	499	2.37
		> 7	$4f^{14}5d^1$	24 583	2.0669	470	2.50
	EOM-CCSD	2	$4f^{14}6p^1$	19 672	1.9971	540	2.72
		4	$4f^{13}\sigma_{6s}^2$	24 251	1.9537	584	2.64
		5	$4f^{14}5d^1$	24 468	2.0177	509	2.78
	IHFS-CCSD	2	$4f^{14}6p^1$	19 543	1.9920	542	2.63
		3	$4f^{13}\sigma_{6s}^2$	21 222	1.9480	591	2.80
		4	$4f^{14}5d^1$	24 363	2.0120	512	2.73
	MRCI ¹⁹		$4f^{13}\sigma_{6s}^2$		1.953	596	
		Exp. ¹⁸	2		19471.49		
	5/2	KRCI ^a	3	$4f^{13}\sigma_{6s}^2$	17 063	1.9302	635
> 6			$4f^{14}5d^1$	24 744	2.0639	474	2.48
EOM-CCSD		2	$4f^{13}\sigma_{6s}^2$	24 957	1.9536	577	2.62
		3	$4f^{14}5d^1$	25 023	2.0146	513	2.80
IHFS-CCSD		2	$4f^{13}\sigma_{6s}^2$	22 127	1.9499	584	2.77
		3	$4f^{14}5d^1$	24 919	2.0089	515	2.77
MRCI ¹⁹			$4f^{13}\sigma_{6s}^2$		1.954	590	
		n.a.	Exp. ²⁶	[574]	19 150		
Exp. ²⁶			[578]	19 280			
Exp. ²⁷	C ₁		23035.3		523	2	
Exp. ²⁷	C ₂		23256.0		507	2	
Exp. ²⁷	D		26014.8		574.6	2.8	

^a KRCI transition energies for the $4f^{13}$ sector were obtained by adding 4144 cm^{-1} , an estimate for the energy of the lowest state in this manifold.

large ($\geq 50\%$ $5d_{\pi}$) contributions. The $\Pi_{1/2}$ state is also dominated by a single configuration, now corresponding to a HOMO($\sigma_{1/2}$) \rightarrow LUMO transition, and shows a rather similar picture in terms of the relative weights of the $6p_{\pi}$ and $5d_{\pi}$ orbitals, with very small contributions from the ground-state mixing due to spin-orbit coupling. The splitting in Yb^+ of $^2D_{3/2}$ and $^2D_{5/2}$ is 1372 cm^{-1} , for $^2P_{1/2}$ and $^2P_{3/2}$ it is 3330 cm^{-1} . The separation between the lowest excited $\Omega = 1/2$ states in the closed shell computation is 3779 and 3126 cm^{-1} for EOM-CCSD and IHFS-CCSD, respectively. This is an indication that the two states must be regarded as a fairly strong admixture of $6p$ and $5d$ orbitals of $j = 1/2$ or $3/2$, as one can expect a much smaller spin-orbit splitting in the axial field of the molecule (about 1/3 of the atomic spin-orbit splitting for

the P state). This picture also finds experimental support in recent measurements of hyperfine constants (d and eq_{0Q}) for the ground and $\Pi_{1/2}$ excited state of YbF^9 , where a simple ligand-field model disregarding the contributions from $5d_{\pi}$ orbitals predicted values of d a factor of 2 larger than the measurements. For bond distances much larger than the equilibrium one the system gets closer to the configurations in Yb^+ with a dominating $5d$ contribution.

As already mentioned this energy range above 18 000 cm^{-1} is dense with a large number of excited states that can mix with each other and result in new mixed states, like the [557] and [561] ones.²⁴ These will be addressed in Section 3.6.

Uttam *et al.*²⁷ reported three unidentified states with energies above 22 000 cm^{-1} which are listed in Table 7 and cannot uniquely be identified with the current results. The one at 26014.8 cm^{-1} has a larger vibrational constant indicating a more strongly bound state, possibly of the $4f^{13}\sigma_{6s}^2$ configuration. The vibrational spacing of the two states at 23 000 cm^{-1} rather points to states with a closed f shell.

3.6 Perturbation of the $3_{1/2}$ excited state

Due to the use of different sectors of Fock space to obtain the $4f^{14}$ and $4f^{13}$ configurations, the excited states with the same Ω values cannot interact among themselves, as is the case within each sector. However, from the discussion above, it is clear that dealing with states which are artificially prevented from interaction makes it difficult to establish a comparison to experiment, for states from about 18 000 cm^{-1} to about 26 000 cm^{-1} , which is where these configurations should be the most entangled. In order to remedy that, in the following we introduce a simple adiabaticization model (eqn (2)) that allows us to investigate how coupling such states would affect the overall spectra in the aforementioned energy region.

In the following we only consider the IHFS-CCSD potential energy curves, as the spectroscopic parameters are more reliable for CCSD than for KRCI. The coupled cluster results for the two methods are quite similar, and FS-CCSD was selected (because it does not include the (2h,1p) and (1h,2p) transitions with rather large uncertainties). Fig. 3 contains the original FS-CCSD curves as well as the ones obtained after adiabaticization with three different coupling constants. Looking at the potential energy curves for this energetic region, there are two $\Omega = 1/2$ and two $\Omega = 3/2$ states of $\text{Yb}(4f^{14}6p^1)\text{F}$ and $\text{Yb}(4f^{14}5d^1)\text{F}$ configurations originating from the (0h,1p) sector. For both Ω values there is an additional state with a $\text{Yb}(4f^{13}[F_{3/2}^{\circ}]\sigma_{6s}^2)\text{F}$ configuration stemming from the (1h,0p) sector. By looking at the KRCI results one expects additional states belonging to the $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]\sigma_{6s}^1)\text{F}$ configuration for this energy range, which will not be included in the current considerations.

As already mentioned earlier the lowest $\Omega = 1/2$ and $\Omega = 3/2$ states can be identified clearly and assigned to experimental observations. There are several experimental states in this energy region attributed to the mixing of states. The [557] and [561] ones²⁴ are assumed to arise from a mixing of the $3_{1/2}$ and $4_{1/2}$. The vibrational constant of the perturbing state



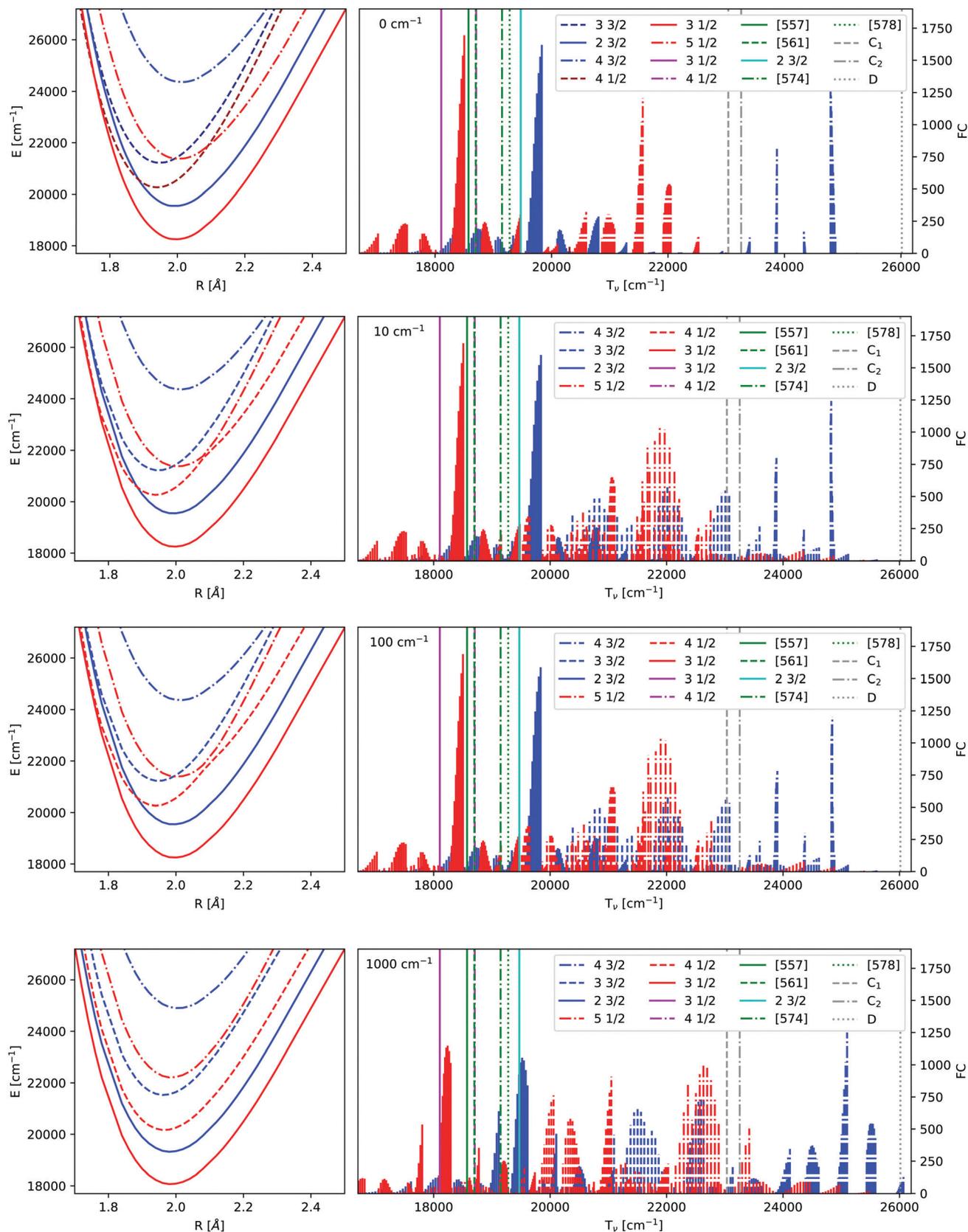


Fig. 3 Frank-Condon factors before and after adiabaticization for the IHFS-CCSD potential energy curves. C is the coupling strength in Hartree. The lowest 10 vibrational levels of the ground state as well as the lowest 60 vibrational levels of the excited state were computed using the LEVEL program.⁸² The experimental values^{18,24,26} have been added as straight lines, the labels are defined in the introduction and Table 7.



Table 8 Spectroscopic data obtained by fitting Morse potentials to the lowest points of the potential energy curves obtained with FSCC for the extrapolated basis set (CBS). This table combines results from both sectors starting either with a closed (f^{14}) or open (f^{13}) f -shell. Additionally, the table contains spectroscopic parameters after adiabaticization with a specific coupling constant (C). The transition energy (T_e), vibrational constant (ω_e), and anharmonicity constant ($\omega_e x_e$) are given in cm^{-1} , the equilibrium bond distance (r_e) in Å

Ω	CBS					$C = 100 \text{ cm}^{-1}$				
	State	r_e	ω_e	$\omega_e x_e$	T_e	State	r_e	ω_e	$\omega_e x_e$	T_e
1/2	$f^{14} - 1$	2.018	515	2.9	0	1	2.018	515	2.8	0
	$f^{13} - 2$	1.940	599	2.8	9627	2	1.940	599	2.8	9617
	$f^{14} - 2$	1.995	539	2.6	18249	3	1.995	538	2.6	18247
	$f^{13} - 3$	1.940	597	2.8	20267	4	1.935	603	8.6	20258
	$f^{14} - 3$	2.003	533	2.7	21375	5	2.002	586	0.4	21359
	$f^{14} - 4$	1.964	581	1.8	31416	6	1.964	581	1.8	31419
3/2	$f^{13} - 1$	1.944	595	2.8	10180	1	1.944	594	2.8	10170
	$f^{14} - 1$	1.992	542	2.6	19543	2	1.992	542	2.6	19540
	$f^{13} - 2$	1.948	591	2.8	21222	3	1.948	591	2.8	21217
	$f^{14} - 2$	2.012	512	2.7	24363	4	2.012	512	2.7	24369
	$f^{13} - 1$	1.949	589	2.8	10967	1	1.949	589	2.8	10960
5/2	$f^{13} - 2$	1.950	584	2.8	22127	2	1.950	583	2.8	22117
	$f^{14} - 1$	2.009	515	2.8	24919	3	2.009	516	2.7	24926

($4_{1/2}$) was estimated to be 605 cm^{-1} in ref. 18. This agrees with the $4_{1/2}$ state in Fig. 3 with a $\text{Yb}(4f^{13}[F_{5/2}^{\circ}]\sigma_{6s}^2)F$ configuration, see also Table 7. [574] and [578]²⁶ have not been identified and since their Ω value is unknown, we were not able to assign them to a configuration.

Next we take a look at the changes introduced by adiabaticization. For small and intermediate coupling strengths there are no major differences in the potential energy curves, although close to the crossing points the potentials are deformed. Intermediate coupling strengths with slightly deformed potentials close to the avoided crossings will be the most realistic description. For very large coupling strengths one obtains parallel potential energy curves due to the strong repulsion. This also results in a major change of the spectra above 19000 cm^{-1} . One of the differences between the adiabatic spectrum and the uppermost one in Fig. 3 is that the Frank-Condon factors of the $4_{1/2}$ state, which is of the $\text{Yb}(4f^{13}\sigma_{6s}^2)F$ configuration, are now noticeable and the spacing of the energy levels of the $5_{1/2}$ is changed. Similarly, transitions belonging to the $5_{1/2}$ appear.

The influence of adiabaticization on spectroscopic parameters can be investigated by comparing spectroscopic constants calculated for the IHFS-CCSD curves without and with a coupling of 100 cm^{-1} (Table 8). We observe that for this coupling strength, there are small but non-negligible changes for the excitation energies, harmonic frequencies and anharmonicity constants, for all but the fourth and fifth $\Omega = 1/2$ states; there, the coupling does seem to significantly change the anharmonicity constants. Equilibrium distances, on the other hand, are largely unperturbed in all cases. Furthermore, as expected from the preceding discussion, no changes are observed for the ground-state, since it is too separated in energy from the other electronic states.

4 Conclusion

In this manuscript we have presented a study of the ground and excited states of the YbF molecule, with 2-component multi-reference CI, equation-of-motion and Fock space coupled cluster approaches (in all cases, performing extrapolations to the complete basis set limit). In particular, we have focused on obtaining electronic states up to around 24000 cm^{-1} arising from configurations which differ in the occupation of the $4f$ shell ($4f^{14}$ and $4f^{13}$), which are very difficult to treat on the same footing due to a number of subtle correlation and relaxation effects.

In order to achieve such a balanced description, our strategy consisted of starting from YbF^+ and YbF^- , in order to arrive at the wavefunctions for YbF through the (1h,0p) and (0h,1p) sectors of Fock space. Once obtained, electronic states with same Ω values coming from these different sectors are further coupled through a simple adiabaticization model in which the coupling strength is taken as a constant.

As a general rule we find that the CI calculations do capture the essential physics of the system, though they are not as reliable as the coupled cluster approaches for excitation energies, bond lengths, harmonic vibrational frequencies and anharmonic constants. In effect, the coupled cluster calculations for the (1h,0p) and (0h,1p) sectors yield the same potential energy curves for the ground state, for internuclear distances up to around 2.8 \AA , which is sufficient to capture the bound regions of all states under consideration,

We have determined that the lowest lying excited states arise from the $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]\sigma_{6s}^2)F$ configuration, with transition energies of around 10000 cm^{-1} , and a splitting about 2000 cm^{-1} . These states are, however, not generally accessible in experiment due to their low dipolar intensity and significantly shifted minima of the potential energy curve resulting in small Frank-Condon factors.

The next set of states, coming above 18000 cm^{-1} , arise from the $\text{Yb}(4f^{14}6p^1)F$, $\text{Yb}(4f^{14}5d^1)F$, $\text{Yb}(4f^{13}[F_{5/2}^{\circ}]\sigma_{6s}^2)F$, and $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]5d^1\sigma_{6s}^1)F$ configurations. Among these, the $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]\sigma_{6s}^2)F$ configurations generally display the shortest equilibrium distances and deepest potential well, while the $\text{Yb}(4f^{14}5d^1)F$ and $\text{Yb}(4f^{14}6p^1)F$ configurations exhibit the largest bond distances and smallest harmonic frequencies, with the other configurations falling somewhere in between. The lowest $\Omega = 1/2$ and $\Omega = 3/2$ states of this group show a $\text{Yb}(4f^{14}6p^1)F$ orbital composition around the ground-state equilibrium structure, though for longer bond lengths they asymptotically approach the $\text{Yb}(4f^{14}5d^1)F$ configuration.

We note that configurations with three unpaired electrons, such as $\text{Yb}(4f^{13}[F_{7/2}^{\circ}]5d^1\sigma_{6s}^1)F$, were only considered with the KRCI method, which has larger uncertainties. This only allows us to make some qualitative statements, e.g. that their bond distances and vibrational constant should be between the values for the other configurations and that they should be higher in energy than the lowest excited $\text{Yb}(4f^{14}6p^1)F$ and $\text{Yb}(4f^{13}[F_{5/2}^{\circ}]\sigma_{6s}^2)F$ states.



A simple method was applied in order to adiabatically the curves obtained for different sectors and reference wave functions. It was applied to potential energy curves between 18000 and 26000 cm⁻¹ and small changes of the Franck-Condon factors were observed. The influence on spectroscopic constant was minor, with the exception of the asymmetry constant for two states. However, the approximation introduced (same coupling strength for all states and all geometries) is perhaps not flexible enough, and more sophisticated models should be investigated.

Conflicts of interest

There are no conflicts to declare.

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