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Structure transition of a C₆₀ monolayer on the Bi(111) surface

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The interfacial structures of C₆₀ molecules adsorbed on solid surfaces are essential for a wide range of scientific and technological processes in carbon-based nanodevices. Here, we report structural transitions of the C₆₀ monolayer on the Bi(111) surface studied *via* low-temperature scanning tunneling microscopy (STM). With an increase in temperature, the structure of the C₆₀ monolayer transforms from local-order structures to a ($\sqrt{93} \times \sqrt{93}$) R20° superstructure, and then to a (11 × 11) R0° superstructure. Moreover, the individual C₆₀ molecules in different superstructures have different orientations. C₆₀ molecules adopt the 6 : 6 C–C bond and 5 : 6 C–C bond facing-up, mixed orientations, and hexagon facing-up in the local-order structure, ($\sqrt{93} \times \sqrt{93}$) R20°, and (11 × 11) R0° superstructure, respectively. These results shed important light on the growth mechanism of C₆₀ molecules on solid surfaces.

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Introduction

C₆₀ molecule, as a prototypical fullerene molecule, has attracted widespread attention due to its potential in endohedral fullerenes,¹ photovoltaic devices,² peapod nanotubes,³ and single-molecule transistors.⁴ A C₆₀ monolayer grown on solid surfaces is critical for understanding and controlling the interfacial properties of fullerene-derived electronic and photovoltaic devices.^{5,6} STM studies demonstrated that the C₆₀ monolayer on the solid surface exhibit a variety of lattice orientations such as the “in phase” ($2\sqrt{3} \times 2\sqrt{3}$) R30°, ^{7–13} (7×7) R0°¹³ and ($\sqrt{589} \times \sqrt{589}$) R14.5°. ^{13–15} The individual molecules of fullerene and fulleride within a single domain display different orientations. In the complex orientational ordering (7×7) R0° structure, a 7-molecule C₆₀ cluster consists of a central molecule sitting atop of a gold atom and six tilted surrounding molecules.¹⁰ In the unit cell of the ($\sqrt{589} \times \sqrt{589}$) R14.5° structure, 49 C₆₀ molecules adopt 11 different orientations.¹⁴ In the ($2\sqrt{3} \times 2\sqrt{3}$) R30° structure, all C₆₀ molecules are in the same orientation.^{12,16} The complex chiral motifs have been observed.¹⁷ In C_{sn}C₆₀ fulleride films, orientational ordering appears.¹⁸ Moreover, “bright” and “dim” molecules have been widely found in the C₆₀ monolayer.^{9–17} However, the “dim” molecules in superstructures reported so far arrange irregularly.

The structure of C₆₀ monolayers grown on the solid surface is not only related to C₆₀ molecules themselves but also the substrate. In the past reports, there have been a large number of investigation on the C₆₀ monolayer structures grown on numerous metals or semiconducting substrates, such as

Ag,^{7–9,19,20} Au,^{10–16,21,22} Cu,^{23–25} graphene,^{26,27} Si,^{28,29} Ge,³⁰ C₆₀,²⁹ or NaCl.³¹ However, few reports address the superstructure of C₆₀ molecules adsorbed on semi-metal substrates. It is found that thin films of organic molecules grown on a semi-metallic Bi(111) surface shows a lot of interesting phenomena, such as the ordered crystalline layer with the standing-up orientation of pentacene molecules,³² the chiral self-assembly of rubrene molecules,³³ structural transitions in different monolayers of cobalt phthalocyanine films,³⁴ and the Moiré’ pattern in C₆₀ thin films.³⁵

In this study, we use Bi(111) as the substrate and studied the structure transition of the C₆₀ monolayer. C₆₀ molecules were deposited at 100 K form local-order structures. When the deposition temperature increased to room temperature, the local-order structures turn into a long-range ordered ($\sqrt{93} \times \sqrt{93}$) R20° superstructure. After annealing at 400 K, the ordered superstructure transforms into the (11 × 11) R0° superstructure. These superstructures are different from the structures of the C₆₀ monolayer reported so far. Furthermore, the individual C₆₀ molecules in the local-order structure, ($\sqrt{93} \times \sqrt{93}$) R20° and (11 × 11) R0° superstructure, show the 6 : 6 C–C bond and 5 : 6 C–C bond facing-up, mixed orientations, and hexagon facing-up, respectively. The 6 : 6 (5 : 6) C–C bond indicates the common side of two adjacent hexagons (pentagon and hexagon) in C₆₀ molecules.

Experimental

The experiments were conducted in an ultra-high vacuum low-temperature scanning tunneling microscope produced by Unisoku. The base pressure was kept at $\sim 1.2 \times 10^{-10}$ Torr. An Si(111) substrate was continuously degassed at ~ 870 K for 8 h

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with subsequent flashing to 1400 K for several seconds. The Bi(111) film was prepared by depositing 20 monolayers of bismuth atoms on a Si(111)- 7×7 surface at room temperature with subsequent annealing at 400 K.³⁶ C₆₀ molecules were deposited onto the Bi(111) surface by heating the tantalum cell to 700 K. The growth rate of C₆₀ molecules was about 0.4 monolayers per minute. All STM images were acquired with a tungsten tip in constant-current mode at liquid nitrogen temperature (78 K).

Results and discussion

First, a small number of C₆₀ molecules were deposited onto the Bi(111) surface when the substrate was maintained at 100 K. Fig. 1(a) shows the atomic-resolution image of the hexagonal lattices of the Bi(111) thin film. The lattice constants of the Bi(111) surface are measured to be $a_1 = a_2 = 0.45 \pm 0.02$ nm, very close to the bulk value ($a = 0.454$ nm) in Bi crystals.³⁶ Fig. 1(b) shows the isolated C₆₀ molecules on the Bi(111) surface presenting round protrusions. When reducing the bias voltage, the round protrusions are separated into two asymmetrical [Fig. 1(c)] or symmetrical [Fig. 1(d)] lobes, corresponding to the two different adsorption configurations, 5 : 6 C–C bond facing-up and 6 : 6 C–C bond facing-up, similar to C₆₀ molecules on Au(111).¹² This indicates that there are two stable adsorption orientations of isolated C₆₀ molecules on the Bi(111) substrate, 6 : 6 C–C bond, and 5 : 6 C–C bond facing-up.

When the coverage increases, C₆₀ molecules form the close-packed hexagonal structure, as shown in Fig. 2. We noticed that all the C₆₀ molecules present a uniform height, except a few dim molecules (marked by green dotted circles). The brightness contrast in images stems from the different adsorption sites of C₆₀ molecules. It is well known that metal surfaces do not behave as rigid templates for the chemisorption of C₆₀ molecules, but may reconstruct substantially to accommodate the molecules.³⁷ We speculate that Dim C₆₀ molecules are located at

the vacancies of the Bi(111) substrate, originating from the reconstruction of the Bi(111) surface, similar to C₆₀ molecules on Au(111)¹⁶ and Cu(111).³⁸

According to the arrangement of bright and dim molecules, we can see some local-order structures, though there is a lack of long-range ordering. In Fig. 2(a), there is an (11×8) R0° local-order structure (marked by red parallelogram). The lattice directions of (11×8) R0° are along with the directions of Bi(111), and the measured lattice constants are 5.00 ± 0.02 nm and 3.64 ± 0.02 nm, corresponding to 11 and 8 times of the lattice constant of the Bi(111) surface. The lattice directions of Bi(111) were obtained on the surface, which was not covered with C₆₀ molecules. In another domain, shown in Fig. 2(b), the local-order structure is mixed with three types of structures, namely (11×8) R0° (red quadrilateral), (11×11) R0° (white quadrilateral), and (10×8) R10° (blue quadrilateral). In particular, we noticed that C₆₀ molecules exhibit almost the same orientation in a single domain, and most of the individual C₆₀ molecules in the local-order structure adopt two favorite orientations (6 : 6 C–C bond and 5 : 6 C–C bond facing up) as the isolated molecules on Bi(111). For example, most of the molecules shown in Fig. 2(a) present two symmetrical lobes, corresponding to C₆₀ molecules with a 6 : 6 C–C bond facing up. However, in Fig. 2(b), the molecules present two asymmetric lobes, corresponding to the 5 : 6 C–C bond facing up. We suggest that the formation of a local-order structure is due to the low-temperature growth. Because of the low kinetic energy of C₆₀ molecules at 100 K, molecular mobility is not high enough to form a long-range ordered superstructure. The C₆₀ molecules adsorbed on Bi(111) adopt their preferred orientations (6 : 6 C–C bond and 5 : 6 C–C bond facing up), similar to the isolated molecules adsorbed on the substrate. This proves the strong molecule–substrate interaction in the local-order structure.

To investigate the influence of temperature on the structure, we deposited C₆₀ molecules on Bi(111) at room temperature. It

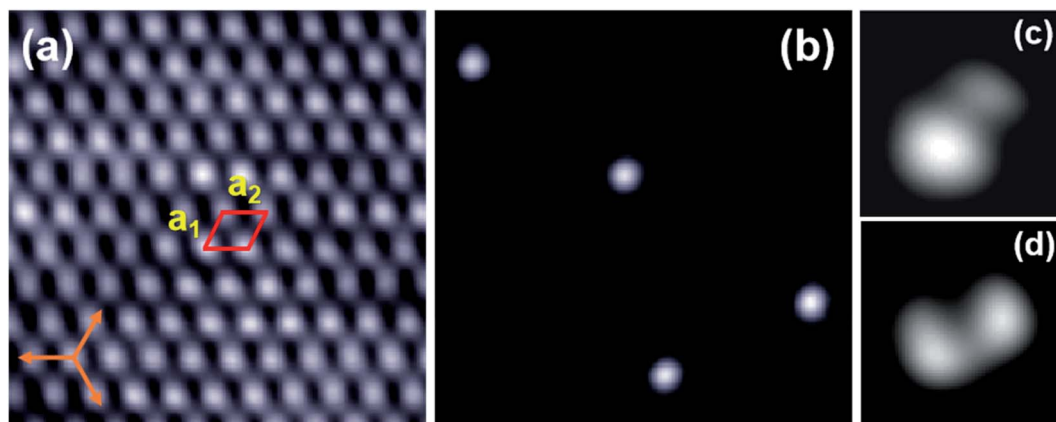


Fig. 1 The initial stage of C₆₀ molecules adsorbed on the Bi(111) surface. (a) Hexagonal lattices of the Bi(111) surface, 5 nm \times 5 nm, -0.1 V. The unit cell is marked with a red box and the orange arrows indicate the directions of the Si(111) substrate. (b) Isolated C₆₀ molecules adsorbed on Bi(111), 20 nm \times 20 nm, 2.2 V. (c) STM image of an isolated C₆₀ molecule with two asymmetrical lobes corresponding to the 5 : 6 C–C bond facing up, 1.3 nm \times 1.3 nm, 400 mV. (d) STM image of an isolated C₆₀ molecule with two symmetrical lobes corresponding to the 6 : 6 C–C bond facing up, 1.3 nm \times 1.3 nm, 200 mV.

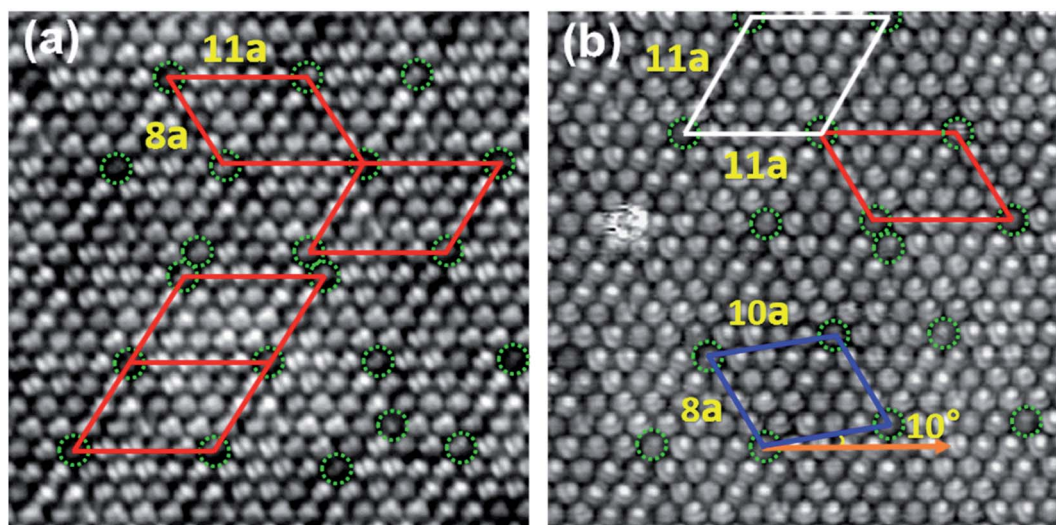


Fig. 2 Local-order structure in the monolayer C_{60} grown at a low temperature (~ 100 K). (a) Four unit cells of the (11×8) $R0^\circ$ superstructure appeared in the C_{60} monolayer, $20 \text{ nm} \times 20 \text{ nm}$, -1.2 V . The dim C_{60} molecules, located at the hollow position of Bi(111), are marked by the green dotted circles. (b) The mixture of three types of superstructures, $20 \text{ nm} \times 20 \text{ nm}$, -0.9 V . The red, white, and blue unit cells correspond to the superstructure (11×8) $R0^\circ$, (11×11) $R0^\circ$, and (10×8) $R10^\circ$.

is found that C_{60} molecules aggregate into a hexagonal structure, the same as C_{60} molecules in the local-order structure. However, the local-order structures, originating from the dim and bright molecules, turn into a long-range ordered $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure [Fig. 3(a)]. This superstructure is different from the structures of the C_{60} monolayer reported so far. There is a misorientation angle of 20° between the lattice directions of the C_{60} monolayer and the Bi(111) surface. The measured lattice constants of $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ are $b_1 = b_2 = 4.38 \pm 0.02 \text{ nm}$, agreeing well with $\sqrt{93}$ times the lattice constant of Bi(111) (0.45 nm). Fig. 3(b) shows the schematic of the $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure. Based on the lattice constant of the Bi(111) substrate, the lattice vectors of the $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure can be expressed as following matrices:

$$\begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = \begin{pmatrix} 11 & -4 \\ 4 & 7 \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix}$$

This ordered superstructure implies two things: first, the intermolecular interaction is getting stronger than that in the local-order structure prepared at low temperature (100 K). Second, the molecule–substrate interaction is also strong since the orientations of the C_{60} superstructure are commensurate with those of the substrate. Furthermore, we can clearly see that individual C_{60} molecules adopt various orientations, rather than the favorite orientations as C_{60} molecules in the local-order structure. As shown in the high-resolution STM image [Fig. 3(c)], C_{60} molecules in $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ present various shapes, such as two asymmetric lobes (white circle), two symmetrical lobes (yellow circle), and three lobes (blue circle),

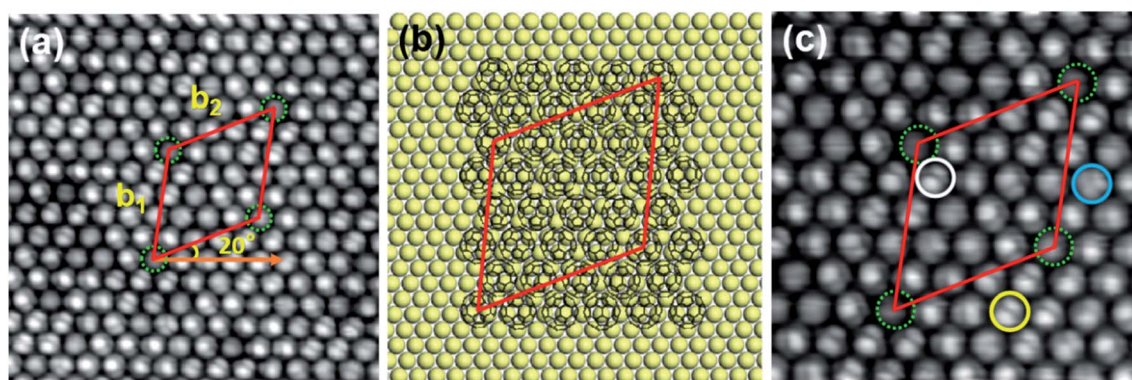


Fig. 3 (a) The STM image of the $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure, $15 \text{ nm} \times 15 \text{ nm}$, -1.2 V . (b) Schematic model of the $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure. The yellow balls and black hollow balls represent Bi atoms and C_{60} molecules. (c) High-resolution STM image of the $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure, $10 \text{ nm} \times 10 \text{ nm}$, -1.0 V . The individual molecules exhibit different orientations, such as 5 : 6 C–C bond, 6 : 6 C–C bond, and hexagon facing up, marked by white, yellow, and blue solid circles, respectively.



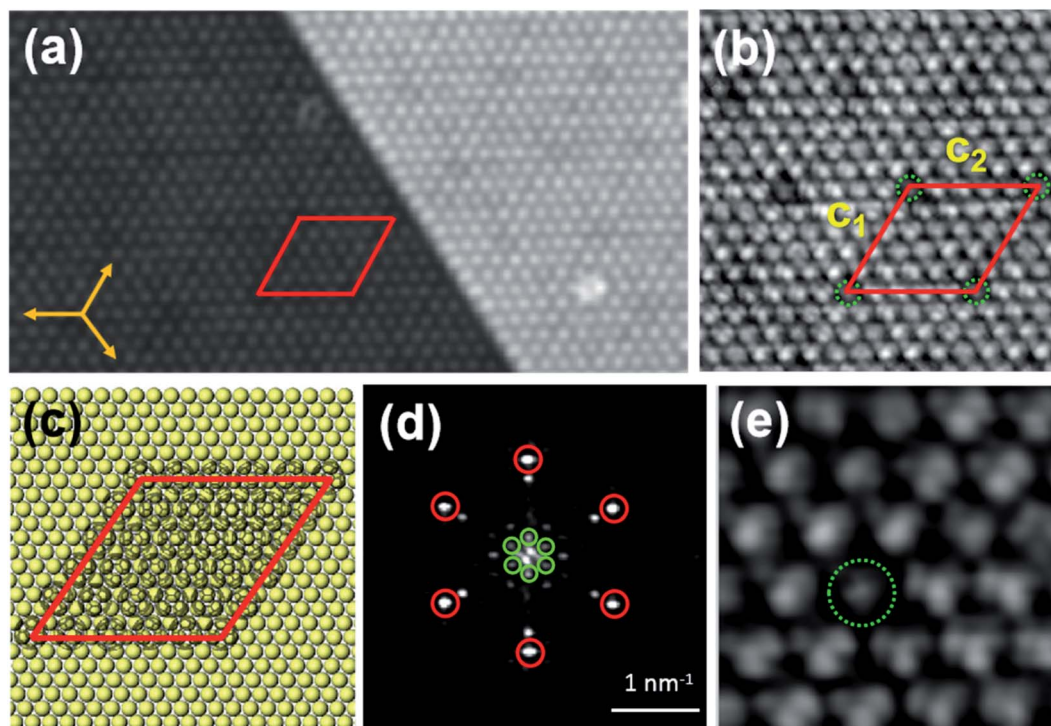


Fig. 4 (a) The STM image of the (11×11) $R0^\circ$ superstructure corresponding to the Bi(111), $39 \text{ nm} \times 21 \text{ nm}$, -1.5 V . (b) Close-up view of the (11×11) $R0^\circ$ superstructure, $14 \text{ nm} \times 14 \text{ nm}$, -0.7 V . (c) Schematic model of the (11×11) $R0^\circ$ superstructure with respect to the Bi(111) lattices. (d) FFT of the image (a). The spots marked by red circles correspond to the C_{60} hexagonal lattices, while the spots marked by the green circles represent the (11×11) $R0^\circ$ superstructure. (e) STM image with a sub-molecular resolution of the superstructure, $5 \text{ nm} \times 5 \text{ nm}$, -0.7 V .

corresponding to the $5:6$ C–C bond, $6:6$ C–C bond, and hexagon facing up. The diversity of C_{60} molecular orientations is due to the enhancement of intermolecular interaction in the $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure. The intermolecular interaction enables C_{60} molecules to overcome the molecule–substrate interaction and adopt other orientations, and then make the $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure stable.

When annealed at 400 K for about 20 min , C_{60} molecules still revealed a hexagonal lattice, while the superstructure transformed from $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ into (11×11) $R0^\circ$ superstructure [Fig. 4(a)], indicating that (11×11) $R0^\circ$ is more stable than $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$. The lattice directions of (11×11) $R0^\circ$ are along the directions of the Bi(111) substrate, and the lattice constants are $c_1 = c_2 = 5.00 \pm 0.02 \text{ nm}$ [Fig. 4(b)], corresponding to 11 times of the lattice constant of Bi(111). Fig. 4(d) is the fast Fourier transform (FFT) image of the (11×11) $R0^\circ$ superstructure, where the spots marked by red and green circles correspond to C_{60} hexagonal lattices and the (11×11) $R0^\circ$ superstructure. In the FFT image, the spots of the superstructure are clearly visible, though they are dimmer than the spots of C_{60} hexagonal lattices, implying that the (11×11) $R0^\circ$ superstructure has long-range order. The schematic model of (11×11) $R0^\circ$ is shown in Fig. 4(c). From STM images, the (11×11) $R0^\circ$ superstructure seems to have the same structure as the reported structure attributed to a Moiré' pattern in ref. 36. However, in our experiment, the (11×11) $R0^\circ$ superstructure is transformed from the $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure and

have no relationship with the Moiré' pattern. From the close-up view of the (11×11) $R0^\circ$ superstructure in Fig. 4(e), it is found that all the C_{60} molecules reveal a unified three-lobe structure, corresponding to the hexagon facing up, different from favorite orientations in the local-order structure and mixed orientations in $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$. With an increase in temperature, the superstructure of the C_{60} monolayer changes from local order to long-range order and C_{60} molecules are re-orientated. This is because the thermal diffusivities of C_{60} molecules and Bi atoms increase with the increase in temperature, which is conducive to the formation of a more orderly and stable superstructure.

Conclusions

In summary, the structure of C_{60} molecules on Bi(111) changes with temperature variation. When deposited on the Bi(111) surface at 100 K , C_{60} molecules form local-order structures, and the molecules in local-order structures adopt their favorite orientations. As the deposition temperature increases to room temperature, the local-order structures turn into a long-range ordered $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructure. The orientations of C_{60} molecules in $(\sqrt{93} \times \sqrt{93})$ $R20^\circ$ superstructures are diverse. After annealing at 400 K for about 20 min , the C_{60} film exhibits a (11×11) $R0^\circ$ superstructure, and all C_{60} molecules in this superstructure take the unified orientation, hexagon facing-up. The appearance of numerous superstructures and the molecular orientations in superstructures is due to the change in the



thermal diffusivity of C₆₀ molecules and Bi atoms at different temperatures.

Conflicts of interest

There are no conflicts to declare.

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