



Cite this: *Phys. Chem. Chem. Phys.*,
2022, 24, 8311

Received 6th October 2021,
Accepted 25th January 2022

DOI: 10.1039/d1cp04581d

rsc.li/pccp

A purely ionic voltage effect soft triode†

Elalyaa Mohamed,^a Sabine Josten^a and Frank Marlow^{a,b}

We report on the construction and characterization of an ionic soft triode intended to be based on interfacial ion adsorption and redox oxidizer depletion. The soft triode was built in a simple manner with no need for sophisticated or expensive materials. It does not utilize the control of a semiconducting channel, but an electrolyte. In different electrical circuit configurations, it can show amplification or memory effects. The device had an electrical current amplification reaching 52 and memory effects in the electrical resistance lasting for up to 6 h. These values were achieved by tuning the electrode interface, the electrolyte and diffusion properties. They are promising for neuromorphic applications.

1 Introduction

Iontronics is a wide concept of controlling electronic devices using ions. The word became widely known in 2010 when a book with this word in its title was published.¹ However, there have been several attempts in the direction of controlling electronics with electrolytes already in the past.^{2–4}

As ions move much slower than electrons, iontronic devices cannot compete with classical electronics in speed. On the other hand, iontronic devices are expected to open a new door for many applications due to their low cost, their potential coupling to biological systems, and their ability to carry chemical information that allows regulating selectively different functions and processes. Despite the outstanding performance of electronic devices, they also have limitations in using them as building blocks in so-called neuromorphic devices and systems. Iontronic devices may fill gaps here. Many of them were invented in the last few decades like ionic transistors^{5,6} ionic capacitors⁷ and ionic memory devices.

Ionic amplification was first recognized in a microfluidic ionic bipolar junction transistor.⁸ With an ionic bipolar transistor, the current gain reached 10, which was achieved by using an emitter gain voltage of 4 V. Ionic field effect transistors represent a class of ionic transistors where ionic or molecular species in micro or nanofluidic systems are controlled by an electric field on the gate electrode of the transistor.^{6,9} Most of these devices can be used as amplifiers. However, the channels in most of the mentioned transistors are semiconductors or

semiconducting polymers that need sophisticated materials and construction methods. Although many ionic devices were developed, there is not much research on signal amplification in ionic systems yet. To improve their biological and chemical sensing, more ionic devices with amplification function would offer new solutions for low-cost devices.

Another important element in electronics and iontronics is the memory. There are many different types of new memory devices but the most interesting type seems to be the resistive switching memory, also called a memristor. A resistive switching memory is a device with low energy consumption, mostly fast operation and a simple structure.^{10–12} These memory devices are expected to open a new door for computing systems and biological applications.¹³ There are different working principles for this type of memory like electrochemical reactions,¹¹ vacancy-based resistance changes/charge trapping,¹⁴ and phase-change materials.¹⁵ Most of the known resistive switching memory devices are built with a simple structure of an insulator between two metals (MIM).¹⁶ They show interesting non-linear behavior in circuits, but also some problems in stability and device integration.^{17,18}

In this work, we introduce a new, purely ionic device with amplification behavior and memory functions which is called an ionic voltage effect soft triode (IVEST). This device is a simple electrochemical micro-cell which consists of a top electrode and two bottom electrodes. The basic idea of this device is to control the concentration and diffusion of ions by the voltage applied on the top electrode. This electrode is made of porous titania, a material which can adsorb ions depending on its potential. The advantages of this device are the ability to be used as an amplifier or a memory, the ease of manufacturing, the low power consumption and the avoidance of any ultra-pure semiconducting material. The need for eco-friendly devices is increasing and using abundant materials and less sophisticated manufacturing methods is a challenge nowadays.

^a MPI für Kohlenforschung, Kaiser-Wilhelm-Platz 1, Mülheim an der Ruhr 45470, Germany. E-mail: marlow@mpi-muelheim.mpg.de

^b Center for Nanointegration Duisburg-Essen (CENIDE), University of Duisburg-Essen, Duisburg 47057, Germany

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d1cp04581d



A drawback of the device seems to be the imperfect memory contrast; however, as in biological information processing, a very high contrast is not needed for many applications. We especially believe that this device can be used in neuromorphic applications.

2 Experimental details

2.1 Device construction

The general layout of the device is shown in Fig. 1. Single parts of the device as well as their fabrication methods were adapted from the elaborated and optimized construction of dye-sensitized solar cells. In solar cells, all parts are adjusted to each other. This explains the special selection of the electrolyte as well as of sealing materials. Essentially, the device consists of three electrodes, all based on stainless steel. Variations from the basic procedure are described in the ESI.†

2.1.1 Top electrode. A metal mesh of $9 \times 5 \text{ mm}^2$ was cleaned and dip-coated with a TiO_2 precursor solution. After solidifying at 120°C , a suspension of TiO_2 nanoparticles was dip-coated and sintered at 450°C . The detailed processes have been adapted from solar cell fabrication^{19,20} and modified for the metal meshes. The resulting thickness of the nanoporous TiO_2 is about $20 \mu\text{m}$ for the non-diluted coating suspension (see below) and is reduced with suspension dilution.

2.1.2 Bottom electrodes. After wet-chemical cleaning, 1 nm of Pt was deposited on both sides of the wires by sputtering (Q150T S, Quorum GmbH).

2.1.3 Device fabrication. A polypropylene film with a thickness of $30 \mu\text{m}$ is used as the base for the device. The two bottom electrodes were put on it. A polypropylene mesh (Franz Eckert GmbH) with a typical thickness of $200 \mu\text{m}$ or a membrane (described below) was put on the top of the bottom electrodes as a spacer and separator. The top electrode was placed on the top of the spacer as shown in Fig. 1A. In addition, a polypropylene top layer was used for encapsulation and the system was glued together using a hot press (Lotus Transpresse LTS 12S) at 110°C . The electrolyte was injected inside the device using a syringe. Finally the device opening was sealed using a metal hot spot (Wetekom soldering station) at 150°C .²⁰

2.1.4 Coating solution preparation. The used solutions were prepared similarly to that described in our previous studies on DSSCs.^{20,22} Here, the following specific compositions have been used. Blocking layer (BL) precursor solution: 0.34 mL titanium isopropoxide and 0.23 mL 2 M HCl were dissolved in 98 mL isopropanol under Ar and stirred for 24 h . TiO_2 particle (TP) suspension: first, an α -terpineol-ethylcellulose mixture was prepared by melting 47.5 g α -terpineol in a round flask in an oil bath with stirring. On reaching 70°C , the heating rate was slowed down, and then 3.7 g ethylcellulose was added step by step. The mixture was kept stirred overnight. To 35.36 g of terpineol-ethylcellulose mixture, 9.02 g nano anatase powder, 1.5 mL of concentrated acetic acid and 14.29 mL isopropanol were mixed in a TopiTec mixer for 30 minutes at 2000 rpm . Finally, the resulting viscous paste was diluted by appropriate amounts of the terpineol-ethylcellulose mixture and isopropanol. Further details are given in the ESI.†

2.1.5 Electrolyte. An electrolyte adapted from ref. 23 was used as in our previous studies on DSSCs.^{20,22} To 20 mL of 3-methoxypropionitrile, 0.05 M I_2 , 0.8 M 1-methyl-3-propylimidazolium iodide, 0.3 M benzimidazole and 0.05 M guanidinium thiocyanate were dissolved and left on the stirrer for 2 h . The solution was put into a brown glass bottle and stored in the dark. In some experiments, the electrolyte was diluted with appropriate amounts of the solvent.

2.2 Characterization methods

Different measurement configurations were used to characterize the device. These configurations are described in the ESI† (S2). The most important measurement to see both the memory function and the amplification behavior turned out to be the one shown in Fig. 2.

Here the resistance between the two bottom electrodes R_{BB} is measured (IDM93N, RS PRO GmbH). Typical test currents were $40 \mu\text{A}$. Continuous as well as pulsed (typical time 1 s) measurements have been carried out.

For additional characterization, the current on the top electrode (J_{top}) was measured by applying a fixed voltage between the two bottom electrodes (V_{BB}) and using a J_V analyzer. To characterize the output current (J_{BB}), the top

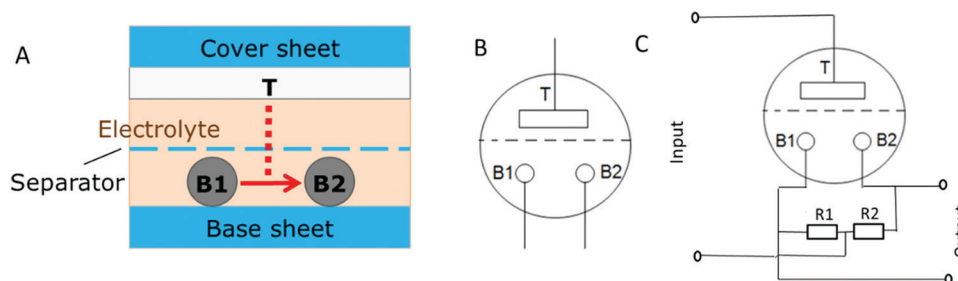


Fig. 1 (A) Device construction, (B) device symbol, and (C) two-gate configuration. T, B1, and B2 are the names of the three electrodes as explained in the text. The idea of the device is the control of the current J_{BB} (red arrow) by the potential V_T . The top electrode (T) is made of porous TiO_2 , which can adsorb ions depending on its potential as it is used in the so-called dye-sensitized solar cells.^{20,23} For device characterization, a BB symmetric configuration with $R_1 = R_2 = 2 \text{ k}\Omega$ was selected.



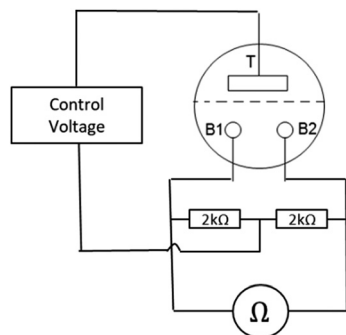


Fig. 2 The Sym-BB- Ω measurement configuration. The resistance between the bottom electrodes is measured depending on the top voltage. The B-electrodes are symmetric with respect to the top voltage.

electrode was connected to the power supply and the bottom electrodes to the JV analyzer as shown in Fig. S2 (ESI[†]).

The thickness of the TiO_2 layer of the electrodes was investigated with a SEM (TM3030, Hitachi).

3 Results

3.1 Basic characterization: the JVV field

The device can be regarded as a 2-gate with input and output. Therefore, the device was characterized with different methods adapted from classical transistor measurements.

The results in Fig. 3A reveal a strongly nonlinear dependence of the input current from the input voltage V_{top} and a weak dependence of the input current from V_{BB} . This difference of sensitivity supports our selection of the input and output gate. The weak dependence of the input current on the output voltage means a low back-action. Fig. 3A shows the result at positive values of V_{BB} only. The results for negative values can be found in the ESI[†]. They are nearly identical with only small indications for asymmetries. There were some small unavoidable asymmetries in the construction in the two bottom electrodes which may have led to small asymmetry in the results. Additionally, it is important to note that the input current is very low at $V_{\text{top}} = -0.37$ V to 0.5 V. We call this range the low-input current (LIC) range.

The output current J_{BB} was measured. Generally, the output curves are nearly linear and resemble a resistor. However, this resistor of the output gate depends on the input. In Fig. 3B it is visible that the resistance increased by increasing V_{top} . The voltage applied on the top electrode was kept lower than 1 V to avoid chemical reactions that may happen inside the system.

The magnitudes of the output current as well as of the output resistance represent a major difference of the IVEST in comparison to the other proposed ionic transistors.²⁴ Varying the input voltage from 0 to -0.6 V causes a change in the output current by 0.1 mA at $V_{\text{BB}} = 0.2$ V. This change in the output current is much higher than that in some known ionic transistors where the same change of input voltage causes a change in the output current of less than 0.002 mA at an output voltage of $V_{\text{D}} = 0.2$ V.⁶

The output shows a visible hysteresis. This hysteresis can be attributed to several reasons, but the most important ones seem to be the delay in ion motion and ion redistribution. From the first look the hysteresis seems to be a drawback of the device intending amplification and switching applications, but at further investigations it will turn to be an interesting option of the device that will be discussed later in Section 3.4. There, we discuss useful memory effects which have been achieved by systematically increasing the hysteresis.

A transfer curve usually describes the interconnections of the input and output.²¹ In Fig. 3C, the output current was plotted at different input voltages V_{top} (Fig. 3C). The output current started to change significantly at higher absolute values of V_{top} . To evaluate the performance of the device in this range, the differential transconductance g_{m} was calculated at different working points (*i.e.* at different V_{top} and V_{BB} values). The differential transconductance g_{m} is defined by the following equation which deviates slightly from the frequently used²¹ non-differential value $g'_{\text{m}} = I_{\text{out}}/V_{\text{in}}$:

$$g_{\text{m}} = \frac{\partial J_{\text{BB}}}{\partial V_{\text{top}}} \quad (1)$$

This enables the direct determination of g_{m} from transfer curves. The resulting g_{m} is shown in Fig. 3D. The values are significantly higher than the g_{m} of known EDLTs with values of 0.04 and 0.17^{25,26} but lower than values for commercial electronic devices. We will discuss this issue below again. Here we also remark an interesting V_{BB} dependence of the transconductance:

$$g_{\text{m}} = \gamma_{\text{BB}} \times V_{\text{BB}} \quad (2)$$

This allows the introduction of the “specific differential transconductance” γ_{BB} . It reflects the nearly Ohmic output and represents a big difference to the conventional transistor and the MOSFET.

In the LIC range the input curve has revealed a high resistance in comparison to the output resistance R_{BB} . When a slight change in the input current J_{top} happened, a significant change in the output current J_{BB} took place. This interconnection will be interpreted in the following section in terms of amplification.

3.2 Resistance control

Because the device has nearly Ohmic output characteristics, the characterization measurements can be simplified by using a standard Ohm-meter instead of a JV analyzer. The use of a stand-alone Ohm-meter not only shortens the measurements, but it also avoids the problem of possible ground interconnection effects between the different parts of the setup.

When a positive voltage was applied on the top electrode, the resistance between the two bottom electrodes increased and when a negative voltage was applied on the top electrode, the resistance between the two bottom electrodes decreased as shown in Fig. 4A. Similar to Fig. 3A, Fig. 4A can be split into two parts, the middle part where $V_{\text{top}} = -0.37$ V to 0.5 V and the corresponding outer part. In the middle part, the change in



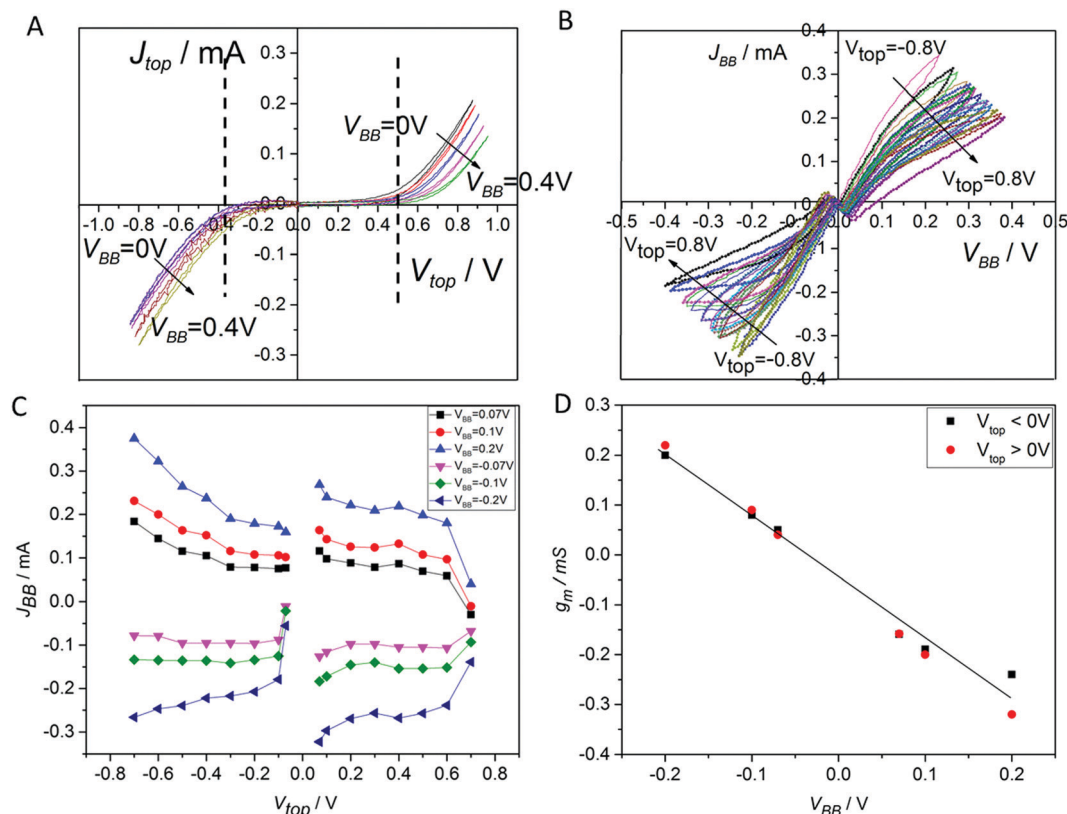


Fig. 3 Basic characterization of a non-optimized IVEST (A) Sym-BB input characteristics. At the output constant V_{BB} values were applied from 0 V to 0.4 V. The LIC range is indicated in the figure. (B) Sym-BB output characteristics. V_{BB} was scanned by a JV measurement device whereas at the input fixed values of V_{top} were applied from -0.8 V to 0.8 V. (C) Transfer curve. (D) Differential transconductance calculated from eqn (1) at small values of V_{BB} . Investigated sample: 156. As shown in the ESI†, five similar samples have shown nearly the same characteristics.

resistance is low. To describe the relationship between the output and the input, the ratio between J_{top} and J_{BB} was investigated based on the amplification definition described in eqn (3). We call it current amplification or J -amplification. Eqn (4) represents the relationship of the amplification to the system parameters at a certain working point (WP) (see the ESI†).

$$\beta_{d,WP} = \frac{\partial J_{out}}{\partial J_{in}} \quad (3)$$

$$\beta_{d,WP} = \lambda_{BB}(V_{top}) \frac{V_{BB}}{G_{d,top}(V_{top})} \quad (4)$$

Interestingly, γ_{BB} can also be determined from the output resistance measurement namely as the differential quotient between the output conductivity and the top voltage as used in the ESI† (eqn S4). Therefore, one has two possibilities for γ_{BB} determination, either from the transfer curve or from the transfer conductance which is measured utilizing a different setup (Fig. 2). The calculations and measurements of γ_{BB} and $G_{d,top}$ are described in the ESI†.

The best J -amplification after optimization was -52 at the working point WP1 $(-0.37$ V, 0.25 V). Here, the first value of

this pair WP $= (V_{top}, V_{BB})$ describes the input voltage and the second one the output voltage.

The outer part of the graph (*i.e.* the non-LIC range) is defined by high $|V_{top}|$ values when either 0.5 V $< V_{top}$ or $V_{top} < -0.37$ V. In these two ranges, the output resistance showed a rapid change when changing V_{top} . In addition, the hysteresis becomes visible in most curves when this range of voltage is touched. Likely, these effects are caused by a relatively low input current. A high input current may cause ion redistribution and non-equilibrium states. These states can be responsible for the hysteresis which could possibly lead to a useful memory effect.

To obtain a deeper insight into the hysteresis and memory behaviour, a special experiment was developed using a control voltage protocol as shown in Fig. 4B. In this experiment, no voltage was applied on the top electrode ($V_{top} = \text{open}$) during a certain “readout time”. The resistance was measured in a pulsed manner, *i.e.* with minimal currents between the B electrodes. The readings were taken each 10 s for 2 min (read-out time). Then, a positive voltage was applied on the top electrode ($V_{top} = 0.8$ V) for 1 min (writing period). After this minute, the voltage applied on the top electrode was switched off ($V_{top} = \text{open}$) and, again, the R_{BB} readings were taken each 10 s for 2 min. Then, in another writing period, a negative

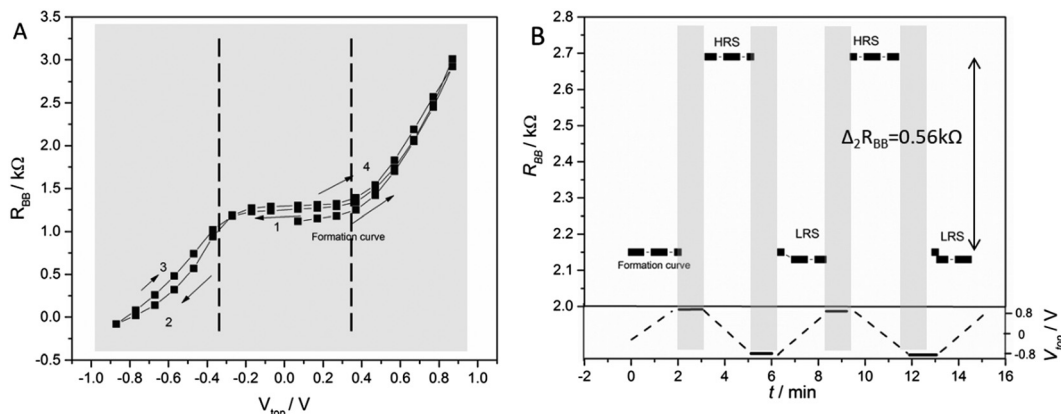


Fig. 4 (A) The BB resistance at permanent control voltage measured in a pulsed manner with about 10 s waiting time after every V_{top} change (sample 274). When a positive voltage is applied on the top electrode, R_{BB} increases and when a negative V_{top} is applied it decreases. The dashed lines indicate the range of low top-electrode current (see the ESI†). (B) Switching measurements with a voltage protocol shown in the lower part of the diagram (sample 285). A positive voltage was applied on the top electrode for 1 min. Then, the R_{BB} under open top circuit conditions was measured for 2 min. This process was repeated with negative V_{top} . A difference of resistances of $\Delta R_{\text{BB}} = 0.56 \text{ k}\Omega$ was found. A high and low resistance state (HRS and LRS, respectively) is clearly visible.

voltage was applied on the top electrode for 1 min ($V_{\text{top}} = -0.8 \text{ V}$). After that, the resistance readings were taken again for 2 min.

Obviously, the IVEST can remember the “writing” during the readout periods. The device is either in a high or in a low resistance state (HRS or LRS, respectively). Within these states, there are small changes as well, but the memory character is not destroyed. In detail, one finds that the resistance readings showed a slight increase in the LRS with time and a slight decrease in the HRS.

As visible, the two states shown in Fig. 3B are fairly stable and, therefore, represent a useful memory behavior of the device. The difference between the two states reached values up to $\Delta R_{\text{BB}} = 0.56 \text{ k}\Omega$ in this experiment. Another measure of the memory contrast is the ratio between the resistance of the HRS and the LRS which can be expressed by

$$\alpha = \frac{H}{L} = \frac{R_{\text{BB}}^{\text{H}}}{R_{\text{BB}}^{\text{L}}} \quad (5)$$

A value of about 1.3 was found for some devices as described below.

3.3 Amplification tuning

In order to investigate the nature of the amplification process and to explore the possibilities of an application, the variation of the amplification values by device variation was investigated. Two factors that were expected to have a big effect on the amplification were investigated. In addition, a working point WP1 at $V_{\text{top}} = -0.37 \text{ V}$ and $V_{\text{BB}} = 0.25 \text{ V}$ was selected for the following discussion. At this point one has the maximum amplification effects among our measured values.

The first factor important for the amplification is the electrolyte concentration X_{e} (Fig. 5A and B). A series of diluted electrolytes were tested as shown in Fig. 5A. The results showed an increase in the resistance as the electrolyte was diluted. This can be attributed to fewer amounts of ions in the electrolyte.

The amplification increased by increasing the concentration of the electrolyte until it reached the highest value at $X_{\text{e}} = 0.6$ and then it decreased as shown in Fig. 5B.

Another factor that affects the amplification is the thickness of the TiO_2 layer on the top electrode tuned by X_{TiO_2} . As shown in Fig. 5C the tuning curve exhibits a maximum at the thinnest TiO_2 layers. The negative current amplification reaches 52 at $X_{\text{TiO}_2} = 0.5$. It seems that high amounts of the semiconductor at the top electrode are not beneficial for the amplification. One may correlate this fact with the amount of dead material. When the thickness of the TiO_2 layer on the top electrode increases, more and more parts of the material are not in electrical contact with the metal mesh (“dead material”). The dead material can however adsorb electrolyte ions. Because they have no electrical contact with the electrode, they have no beneficial effect on the amplification. These ions are lost for the other processes. Thinner TiO_2 layers are technically not possible at the moment. However, one can expect a further increase of the amplification at thinner layers.

3.4 Memory tuning

In Fig. 4A, a hysteresis has appeared when the top voltage exceeded the LIC region. This hysteresis indicates a possible memory effect. Enlarging this hysteresis was, therefore, the way for optimizing this potential memory. Finally two properties, the resistance contrast $\Delta_2 R_{\text{BB}}$ and the stability time τ_2 , were used to characterize the efficiency of the memory. The stability time is the time that the two metastable states require to reach an equilibrium state. The difference of resistance between the HRS and the LRS after two minutes, $\Delta_2 R_{\text{BB}}$, was used to characterize the memory contrast.

As the adsorption of the ions on the TiO_2 layer likely plays a role in the mechanism of the device (see Section 4.2), the behavior of the IVEST was investigated at different electrode thicknesses realized by different concentrations of TiO_2 in the dip-coating suspension. In Fig. 6A one can see that $\Delta_2 R_{\text{BB}}$



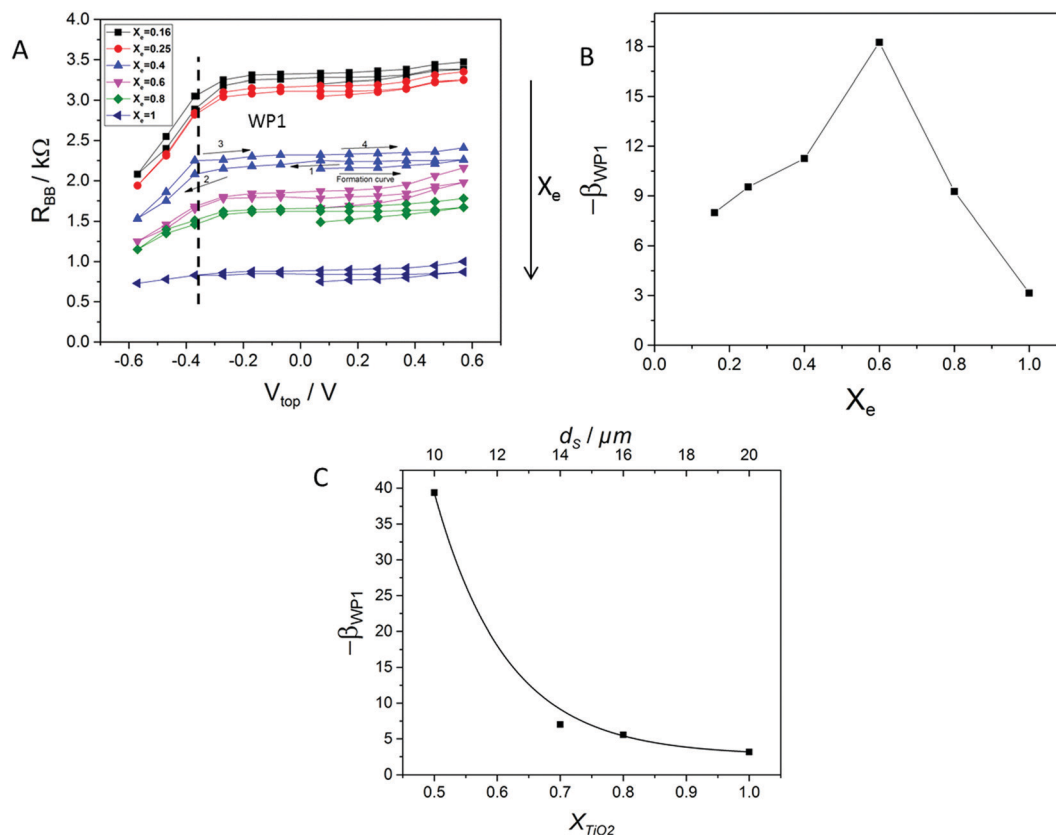


Fig. 5 Tuning of device parameters. (A) Effect of electrolyte dilution on R_{BB} measured using a pulsed DC resistance measurement. (B) Calculated amplification at WP1 based on eqn S4 (ESI†), at different electrolyte concentrations. (C) Calculated amplification at different TiO_2 thicknesses, determined analogously as in (A) and (B). The thickness scale was constructed from two SEM investigations of the electrodes.

moderately increased by decreasing the concentration of TiO_2 . This increase in $\Delta_2 R_{BB}$ could likely be attributed to the decrease of the dead material or the rim on the top electrode.

The second factor that affects $\Delta_2 R_{BB}$ is the electrolyte. The difference between the two switching states $\Delta_2 R_{BB}$ increased when the electrolyte was diluted (Fig. 6B).

The pore size of the membrane separating the device parts is also expected to have an influence on the memory behavior. This was indeed found. As the pore size decreased, $\Delta_2 R_{BB}$ showed an increase. The factor that affects the stability time the most was also found to be the pore size of the membrane. By decreasing the pore size, τ_2 increased. There was a big difference in τ_2 when using a membrane with 150 μm pore size and membranes with pore sizes $< 1.5 \mu m$.

Based on the previous results, an optimized cell for memory effects was built with $X_e = 0.16$, $X_{TiO_2} = 0.5$, and a spacer of 1.2 μm pore size. In the case of an optimized cell, values of $\Delta_2 R_{BB} = 0.56 k\Omega$ and $\tau_2 = 5.8$ h were found, leading to a contrast factor of $\alpha = 1.28$. This cell is also shown in Fig. 4B.

4 Discussion

4.1 Working points and amplification mechanism

The results show that the amplification is strongly WP-dependent and that a working point WP1 with

$V_{top} = -0.37$ V seems to be optimal to have a maximal amplification for the investigated voltages. How can this be understood?

When the semiconductor material comes in contact with the electrolyte, an electronic depletion zone is likely created in the semiconductor as shown in Fig. 7. This picture reminds on the Schottky contact, but here the Fermi energy of the metal is replaced by the redox potential of the electrolyte.²⁷ This depletion zone near the electrolyte represents a positive charge. The electrolyte compensates this charge by forming different ionic layers. These layers must contain just the opposite ionic excess charge as the positive charge of the semiconductor. In addition, the ionic layers must contain a lot of additional electrolyte ions for screening all these charges. This enhanced concentration of ions near the electrode surface can be interpreted as variable storage. The amount of ions stored in this layer will be missing in the B-part of the IVEST which can explain the influence on the BB-resistance.

A change in potential would change the depletion layer inducing a change in the ionic adsorption layer. This, in turn, would change the ion concentration in the B-part of the device leading to the resistance change. Since the adsorption/desorption phenomena are nearly current-free, the input current is low whereas the output current mainly depends on the output voltage. So, the output current is normally not low. This



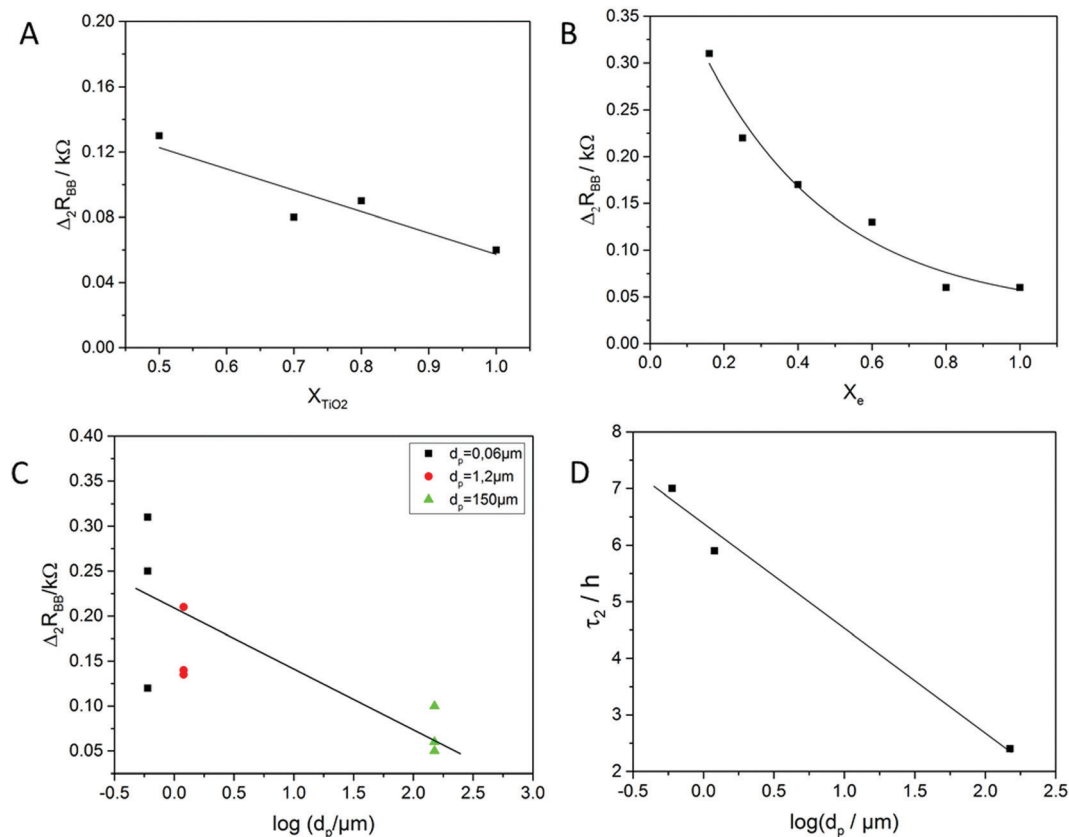


Fig. 6 (A) Effect of TiO_2 thickness on $\Delta_2 R_{BB}$. (B) Effect of electrolyte concentration on $\Delta_2 R_{BB}$. (C) Effect of membrane pore size on $\Delta_2 R_{BB}$. (D) Effect of membrane pore size on the stability. There is a big increase in τ_2 when a membrane with a small pore size is used for the separation the device parts. The solid lines are guides to the eye.

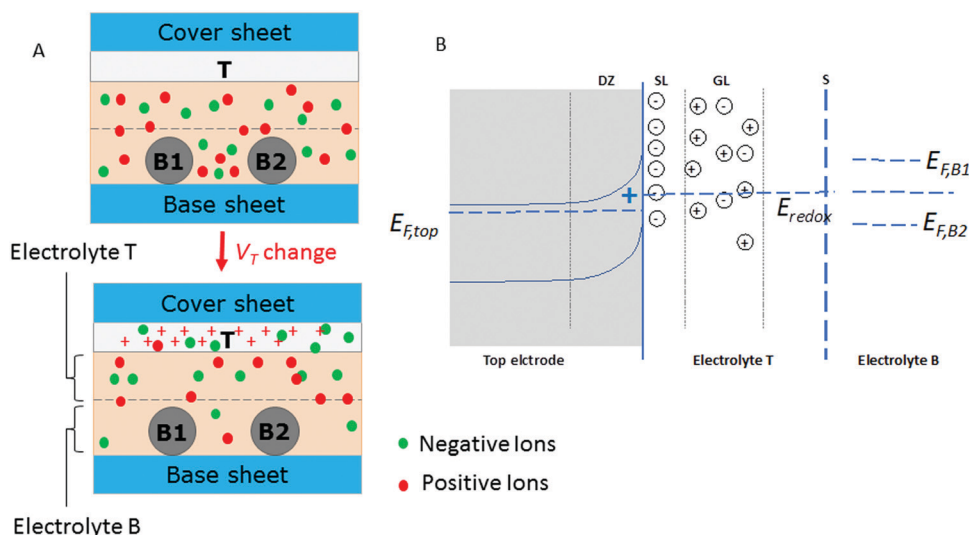


Fig. 7 (A) Schematic description of the mechanism of the device. (B) The IVEST energy levels and relevant charge layers. The image shows the electronic depletion zone (DZ), the Stern layer (SL), the Gouy layer (GL), and the separator (S). The left half of Fig. 7B shows the classical description of a semiconductor electrolyte interface as e.g. in ref. 27. The right half shows the proposed coupling to the bottom part of the electrolyte. The Fermi energy offsets are generated by external voltages.

fundamental difference of the input and output current is the root of the amplification of the device.

For a quantitative estimation, we use the Mott-Schottky theory.²⁸ In addition, the basic equation for the coupling of the semiconductor with the electrolyte is assumed to be:

$$N_E = en_d A w \gamma \quad (6)$$

where N_E represents the total number of additional ions bound near the electrode, γ is the enhancement factor between the number of depleted electrons and the number of ions involved in the charge compensation, and w is the width of the (electronic) depletion zone. This equation reflects the fact that the electronic depletion is counter-balanced by ions. These ions are bound to the top electrode and cannot contribute to the conductivity of the electrolyte in parts a little away from the top electrode. For an open separator, they directly reduce the concentration of ions in all parts of the electrolyte simultaneously and, therefore, they reduce the conductance in the B-part. This results in the equation:

$$G_{BB} = G_0 - \alpha_1 (V_{top} - V_{fb})^{1/2} \text{ for } V_{top} > V_{fb} \quad (7)$$

with the flat-band potential V_{fb} and an appropriate constant α_1 . Here, an explicit expression for w was used which has been adapted from the Mott-Schottky theory. Finally one gets for the specific differential transconductance:

$$\gamma_{BB} = \frac{\partial G_{BB}}{\partial V_{top}} = \alpha_2 (V_{top} - V_{fb})^{-1/2} \quad (8)$$

with another constant α_2 .

The physical meaning of this formula is very similar to the Mott-Schottky formula. The depletion layer disappears near the flat-band potential; however, it also shows the strongest modifications there. Therefore, the selection of the working point near the flat-band potential seems to be most appropriate. On the other hand, in this region also the conductivity of the interface increases. Therefore, the optimization of the amplification is a delicate balance between increasing the transconductance and keeping the resistance high.

4.2 The potentiometric effect

The basic action of the device can be regarded as a potentiometer with external control voltage. This control voltage has a very nonlinear effect as can be seen from Fig. 4. A part of these effects is instantaneous, and another part can be modified by a waiting time. Therefore, we think that two mechanisms contribute to the effect. The first is the voltage-controlled adsorption of ions in the TiO_2 electrode and the second is the ion transport through the membrane. Both mechanisms can change the concentration of ions. The instantaneous effect can be good to construct an amplifier. Using a relatively long writing period, the transport phenomena become dominating.

The control current created by the voltage can be seen from Fig. S2 in the ESI.† This Figure also clarifies why a relatively high control voltage was used for the writing phase. Here a significant current is flowing.

The nonlinearity of the control current dependent on the control voltage might be regarded as a useful feature. It creates an in-built threshold switch or, more generally, an activation function in neural networks. The basic potentiometer characteristics in the output are well suited for perceptrons or neuromorphic networks for realizing signal weighting.

4.3 Memory effects and the role of membranes

While *trans*-conductance phenomena only need instantaneous interactions and nonlinearities for explanation, memory effects are based on multi-stability or non-equilibrium states. For our device, non-equilibrium phenomena seem to be most important. A distinction should be made between the “main memory effect”, *i.e.* the memory effect which is intended to be used in a device, and other memory effects which only disturb our measurements.

The main memory effect found in the IVESTs can likely be explained by the different concentrations of ions between the two parts of the device. The T-part of IVEST is the top electrode and a part of the electrolyte, while the second part of the IVEST consists of two bottom electrodes and the electrolyte surrounding them (“B-part”). When a voltage is applied, gradients of the electrolyte ions are formed in the device creating non-equilibrium states between the two parts of the device. In the readout phase (*i.e.* without control voltage), the ions need time to redistribute again and to reach an equilibrium state. This time (stability time) is usable for the memory.

Fig. 8 shows a brownish color between and near the two bottom electrodes. We observed this coloration when applying higher voltages ($|V_{top}| > 0.6$ V) on the top electrode. This coloration is a clear indication of I_3^- or polyiodides accumulated near the two bottom electrodes creating a non-equilibrium state. Very likely, also I_3^- depletions occur and also the iodide ions form such gradients. The top electrode is expected to play a significant role in the enhancement of the non-equilibrium concentration between the two parts of the device. Fig. 6A shows that by increasing the thickness of TiO_2 on the top electrode, $\Delta_2 R_{BB}$ decreases. This effect is a bit counter-intuitive. But one needs to consider that the TiO_2

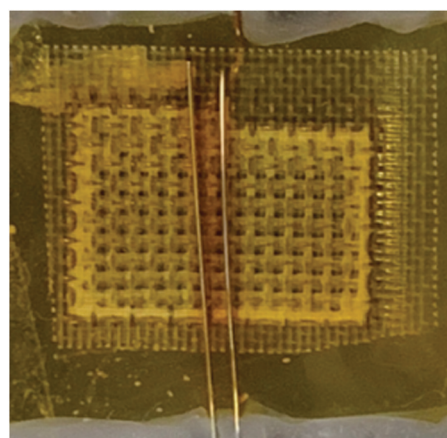


Fig. 8 Photo of the coloration of the B-electrodes.



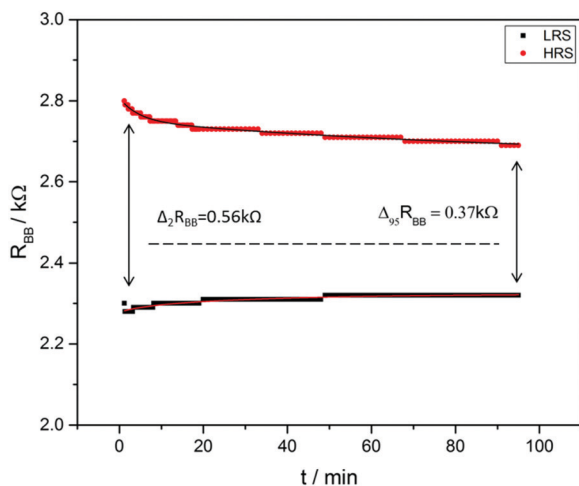


Fig. 9 The behavior of the two resistance states over 95 min. The dashed line represents the equilibrium state of the device. The measurement was carried out at an optimized cell (sample 285) with a separator pore size of 1.2 μm , an electrolyte with $X_e = 0.16$ and a titania layer with $X_{\text{TiO}_2} = 0.5$.

electrodes are not ideal and have an active part and some dead material which is not interconnected to the metal support. This material would explain the negative effect of a thick TiO_2 layer.

The separator typically used for the construction of the device can be varied in a wide range of pore sizes. For small pore sizes, it acts as a membrane which can slow down the redistribution of ions between the T-part and the B-part of the device. It increases the lifetime of the non-equilibrium state and may, therefore, enhance the memory effect. Such an enhancement can actually be seen in Fig. 6D. Here, decreasing the pore size led to an increase in the stability time τ_2 . The membrane pore size influenced $\Delta_2 R_{\text{BB}}$ and is shown in Fig. 6B. Using a membrane with a smaller pore size increased $\Delta_2 R_{\text{BB}}$ by a factor of 5.

Another memory effect becomes visible during variations of the read-out process (*i.e.* variation of the readout time and used current). There one can find different memory curves. In addition, the coloration of one of the B-electrodes becomes visible. Obviously, redistribution of ions within the B-part occurs. This second memory effect might be useful for special circuit layouts. However in the frame of this work, it is only a disturbance of the characterization measurements.

4.4 Adaptation, stability time of memory states and physical effects

The device has two metastable states HRS and LRS in the readout phase. The meta-stability of each state becomes visible in two kinetic processes with different time constants (Fig. 9). The first one can be called adaptation and is visible within a time of < 2 min. It is characterized by the time constant τ_1 . Since it occurs directly after the writing process, the gradients created during writing might be responsible for this relatively fast decay.

The second visible process during readout has a time constant of $\tau_2 > 2$ min and can durate over times up to hours.

The constant is called the stability time and determines how long the state can be used for storing information. After this time, the two metastable states are assumed to reach the equilibrium state of the device.

Mathematically, the sum and the difference of the two resistances were fitted to obtain τ_2 and ΔR_{BB} in a reliable manner (see the ESI†). Values of $\Delta_2 R_{\text{BB}} = 0.56$ k Ω and $\tau_2 = 5.8$ h were obtained indicating a very useful kind of memory. In particular, the τ_2 -time limits the usability of the memory since after this time a refreshment process is needed to store information for a longer time.

For the interpretation of the dynamical phenomena we use the same physical picture as in Section 4.2 with focus on the redistribution of I_2 . In particular, the low concentrations will lead to a (variable) depletion of I_3^- in the B-part. Now one has to consider that the resistance between two electrodes includes different contributions: the two electrode surfaces including the ion layers in the electrolyte and the volume contribution of the electrolyte. Whereas the volume contribution is not strongly dependent on the kind of ion, the electrode contributions are very sensitive. The exchange of I_3^- with I^- has a dramatic effect on the redox reactions and, therefore, on the electrode resistance. For a quantitative estimation of such effects, one can use models as *e.g.* described in ref. 29.

As said above, the decisive influences can be expected from the extinction of I_3^- which is the oxidizer part of the redox couple. Therefore, we would like to call the proposed mechanism “redox oxidizer depletion”. In the writing process this depletion is generated or removed mainly between the two cell parts. However in addition, intra-part gradients are formed. In the beginning of the readout process these gradients will decay. We ascribe this to the typical time τ_1 . Later in the readout phase, the inter-part gradients are of importance only. They lead to the typical decay time τ_2 . One could now make an estimation on the $\tau_2(d_p)$ -dependence. However, our measured dependence includes not a variation of d_p alone, but also of the thickness and kind of material. Therefore, quantitative agreement cannot be expected. In addition to this open point, the investigation of other ionic gradients is still in progress.

5 Conclusions

A new purely ionic soft triode was constructed. It can be used as a memory or an amplifier. Two physico-chemical principles are likely to be realized in the device, ion adsorption in a porous electrode and redox oxidizer depletion. Both principles allow a wide range of chemical tuning. A non-equilibrium of ions occurs between the two parts of the device, which can be used to control the current in the device.

Different simple tuning factors were studied. The amplification of the device was optimized and reached a value of $\beta = -52$. The memory effect was optimized and reached $\Delta_2 R_{\text{BB}} = 0.56$ k Ω and $\tau_2 = 5.8$ h by using a membrane with a pore size of 1.2 μm , an diluted electrolyte (16% of our standard) and a thinner titania layer (of about 10 μm). The high stability time of the



device allows it to be used as a memory element in different applications such as sensor systems,³⁰ health care applications³¹ and neuromorphic memories. This device could be a building block for neuromorphic information processing systems and could, therefore, be a step towards a new computational technology.

The device has several advantages in comparison with known devices such as the possibility of a green fabrication involving no high-temperature step, the possibility of incorporation into printable electronics, and the low power consumption of such devices. It should be underlined that there is no need to use a semi-conductor material to create a channel between the bottom electrodes. This would always increase the production cost. Liquid electrolytes seem to be a drawback, but existing widely used devices as batteries and electrolytic capacitors show that one can well manage the technical requirements. In addition, solid or gel electrolytes can be used for such devices as well.

From the device point of view, it is interesting that the IVESTs have a controllable linear Ohmic output. This is a difference to the known devices, which have strongly non-linear output characteristics. However, this fact enables and enforces modified circuit layouts.

Conflicts of interest

The authors declare no conflict of interest.

Acknowledgements

The authors would like to thank U. Petrat and H. Köster for providing support in carrying out the experiments. We thank W. Schmidt for a critical reading of the manuscript. E.M. acknowledges the Erasmus+ program (2016-KA1-Student mobility) of the European Commission for support.

References

- 1 J. Leger, M. Berggren and S. Carter, *Iontronics: Ionic carriers in organic electronic materials and devices*, CRC Press, 2010.
- 2 L. Hegedus, N. Kirschner, M. Wittmann and Z. Noszticzius, *J. Phys. Chem. A*, 1998, **102**, 6491–6497.
- 3 W. H. Brattain and R. B. Gibney, *Three-electrode circuit element utilizing semiconductor materials*, US Patent, US2524034A, 1948.
- 4 R. E. Hetrick, *Electrochemical transistor structure with two spaced electrochemical cells*, US Patent, 4,332,003, 1984.
- 5 G. Bruno, G. Canavese, X. Liu, C. S. Filgueira, A. Sacco, D. Demarchi, M. Ferrara and A. Grattoni, *Nanoscale*, 2016, **8**, 18718–18725.
- 6 S.-W. Nam, M. J. Rooks, K.-B. Kim and S. M. Rossnagel, *Nano Lett.*, 2009, **9**, 2044–2048.
- 7 P. Janson, E. O. Gabrielsson, K. J. Lee, M. Berggren and D. T. Simon, *Adv. Mater. Technol.*, 2019, **4**, 1800494.
- 8 K. Tybrandt, K. C. Larsson, A. Richter-Dahlfors and M. Berggren, *Proc. Natl. Acad. Sci. U. S. A.*, 2010, **107**, 9929–9932.
- 9 F. Wang, P. Stepanov, M. Gray, C. N. Lau, M. E. Itkis and R. C. Haddon, *Nano Lett.*, 2015, **15**, 5284–5288.
- 10 M. T. Ghoneim and M. M. Hussain, *Electronics*, 2015, **4**, 424–479.
- 11 R. Waser, R. Dittmann, G. Staikov and K. Szot, *Adv. Mater.*, 2009, **21**, 2632–2663.
- 12 E. Madai, B. Matejczyk, A. Dallos, M. Valisko and D. Boda, *Phys. Chem. Chem. Phys.*, 2018, **20**, 24156–24167.
- 13 H. Zhang, C. Cheng, H. Zhang, R. Chen, B. Huang, H. Chen and W. Pei, *Phys. Chem. Chem. Phys.*, 2019, **21**, 23758–23763.
- 14 Y. Shan, Z. Lyu, X. Guan, A. Younis, G. Yuan, J. Wang, S. Lia and T. Wu, *Phys. Chem. Chem. Phys.*, 2018, **20**, 23837–23846.
- 15 W. Zhang, R. Mazzarello, M. Wuttig and E. Ma, *Nature*, 2019, **4**, 150–168.
- 16 S. Menzel, S. Tappertzhofen, R. Waser and I. Valov, *Phys. Chem. Chem. Phys.*, 2013, **15**, 6945–6952.
- 17 Z. Wang, K. L. Wang, H. Wu, Q. Xia, G. W. Burr, J. J. Yang and C. S. Hwang, *Nat. Rev. Mater.*, 2020, **5**, 173–195.
- 18 A. Chen, *Solid state electrochemistry II: electrodes, interfaces and ceramic membranes*, John Wiley & Sons, 2012.
- 19 S. O. Abdellatif, S. Josten, A. S. G. Khalil, D. Erni and F. Marlow, *IEEE J. Photovolt.*, 2020, **10**, 522–530.
- 20 F. Marlow, A. Hullermann and L. Messmer, *Adv. Mater.*, 2015, **27**, 2447–2452.
- 21 P. Horowitz and W. Hill, *The Art of Electronics*, Cambridge Univ. Press, 1989.
- 22 S. Abdellatif, P. Sharifi, K. Kirah, R. Ghannam, A. S. G. Khalil, D. Erni and F. Marlow, *Microporous Mesoporous Mater.*, 2018, **264**, 84–91.
- 23 P. R. F. Barnes, K. Miettunen, X. Li, A. Y. Anderson, T. Bessho, M. Grätzel and B. C. O'Regan, *Adv. Mater.*, 2013, **25**, 1881–1922.
- 24 K. Tybrandt, E. O. Gabrielsson and M. Berggren, *J. Am. Chem. Soc.*, 2011, **133**, 10141–10145.
- 25 H. Yuan, H. Shimotani, A. Tsukazaki, A. Ohtomo, M. Kawasaki and Y. Iwasa, *Adv. Funct. Mater.*, 2009, **19**, 1046–1053.
- 26 L. Herlogsson, X. Crispin, N. D. Robinson, M. Sandberg, O.-J. Hagel, G. Gustafsson and M. Berggren, *Adv. Mater.*, 2007, **19**, 97–101.
- 27 X. G. Zhang, *Electrochemistry of Silicon and its Oxide*, Springer, 2007.
- 28 A. J. Bard and L. R. Faulkner, *Electrochemical Methods: Fundamentals and Applications*, Wiley, New York, 1980.
- 29 J. Akilavasan and F. Marlow, *J. Phys. Chem. C*, 2020, **124**, 4101–4108.
- 30 A. Chhetry, J. Kim, H. Yoon and J. Y. Park, *ACS Appl. Mater. Interfaces*, 2018, **11**, 3438–3449.
- 31 B. Nie, R. Li, J. Cao, J. D. Brandt and T. Pan, *Adv. Mater.*, 2015, **27**, 6055–6062.

