


Cite this: *RSC Adv.*, 2022, 12, 25048

Growth of turbostratic stacked graphene using waste ferric chloride solution as a feedstock†

Piyaporn Surinlert,^{a,c} Phurida Kokmat^b and Akkawat Ruammaitree  ^{*bc}

Monolayer graphene has excellent electrical properties especially a linear dispersion in the band structure at the *K*-point in the Brillouin zone. However, its electronic transport properties can be degraded by surface roughness and attachment of charge impurities. Although multilayer graphene can reduce the surface roughness and attachment of charge impurities, the increase in the number of graphene layers can degrade the electronic transport properties due to interlayer interactions. Turbostratic graphene can significantly reduce the effect of interlayer interaction of multilayer graphene resulting in electrical properties similar to those of monolayer graphene. In this report, we have demonstrated the growth of turbostratic stacked graphene using waste ferric chloride solution as a feedstock by vaporization and calcination at 700 °C for 6 hours under an argon atmosphere. SEM images and EDX elemental distribution maps showed graphene can be grown on iron and nickel catalysts. XRD results and Raman spectra confirmed the presence of turbostratic stacked graphene with the interlayer spacing in the range of 3.41 Å to 3.44 Å. The Raman spectra in all samples also displayed a weak intensity peak of ITALO^- and a well-fitted 2D band by a single Lorentzian peak indicating the presence of turbostratic stacked graphene. In addition, XPS spectra reveal the growth mechanism of the turbostratic stacked graphene. This synthesis process of turbostratic stacked graphene is not only simple, low-cost, and suitable for large-scale production but also decreases the environmental issues from releasing waste ferric chloride solution with improper disposal.

Received 27th April 2022
Accepted 27th August 2022

DOI: 10.1039/d2ra02686d

rsc.li/rsc-advances

Introduction

Graphene is a monolayer of carbon atoms arranged in a honeycomb crystal lattice. Graphene has attracted attention from many researchers due to its exotic properties such as high electron mobility,¹ linear dispersion (Dirac cone) at the *K*-point in the Brillouin zone,² high intrinsic strength³ and superior thermal conductivity.⁴ However, these desirable properties, especially electrical properties, can be degraded with an increase in the number of graphene layers due to the interlayer interactions between adjacent graphene layers. In theory, monolayer graphene should have the best electronic transport properties because of no interlayer interaction, but realistically its electronic transport properties can be degraded by surface roughness⁵ and attachment of charge impurities.⁶ High levels of the surface roughness and charge impurities can often appear

on the monolayer graphene because it has only single carbon atom thick.

Turbostratic stacked graphene is multilayer graphene which contains the relative rotations between adjacent graphene layers. The interlayer rotations of turbostratic graphene can reduce the effect of interlayer interaction leading to the linear dispersion and carrier mobilities similar to those of monolayer graphene.⁷ In addition, the turbostratic stacked multilayer graphene has lower level of surface roughness and charge impurities than monolayer graphene resulting in its conductivity and carrier mobility are higher than those of the monolayer graphene.⁸ Therefore, the turbostratic multilayer graphene is a promising candidate material for the application of high-performance electronic devices.

Chemical vapor deposition (CVD)⁹ is generally utilized for the growth of graphene sheet because it can provide large-area graphene at low cost.¹⁰ Nickel is commonly utilized as metal catalyst for the growth of multilayer graphene by CVD.¹¹ After the CVD growth procedure, the nickel catalyst must be etched by immersion in the ferric chloride solution.¹² The waste ferric chloride solution can cause water pollution and harmful chronic effects on the aquatic biota if it is thrown overboard without proper disposal.¹³

Various methods for synthesis of graphene powder have been developed to generate high-quality graphene powder with

^aChulabhorn International College of Medicine, Thammasat University, Pathum Thani 12120, Thailand

^bDepartment of Physics, Faculty of Science and Technology, Thammasat University, Pathum Thani 12120, Thailand. E-mail: u4605070@hotmail.com

^cThammasat University Research Unit in Synthesis and Applications of Graphene, Thammasat University, Pathum Thani 12120, Thailand

† Electronic supplementary information (ESI) available. See <https://doi.org/10.1039/d2ra02686d>


low cost and large-scale production. Micromechanical exfoliation of graphite can produce high-quality graphene monolayer, but the scale production is small.¹⁴ The oxidation–reduction method is commonly used to fabricate graphene powder because it can achieve the large-scale production of reduced graphene oxide at low cost. However, the quality of the resultant graphene powder is low due to the presence of high level of defects and oxygen content on the graphene surface.¹⁵ In 2014, Binbin Zhang *et al.* propose the different method which can achieve mass production of high-quality graphene with low cost by calcinating glucose and ferric chloride solution under argon atmosphere at 700 °C for 6 hours.¹⁶ For this method, the iron is utilized as metal catalyst for synthesis of graphene. Nevertheless, nickel is also the metal catalyst which is commonly utilized to synthesis graphene. Therefore, the mixture of iron and nickel can also be utilized as metal catalyst for graphene growth.

In this study, we demonstrated the growth of turbostratic stacked graphene using waste ferric chloride solution as a feedstock. The utilization of the waste ferric chloride solution as a raw material for the synthesis process of turbostratic stacked graphene can not only reduce the cost production of turbostratic stacked graphene but also decrease the environmental issues.

Experimental

Preparation of waste ferric chloride solution

The ferric chloride solution was prepared by dissolving 80 g ferric chloride in 100 ml deionized water. A graphene on nickel foam sample (ESI†) was immersed in the ferric chloride solution for 24 hours to etch nickel. The mass of nickel in the waste solution was calculated by the mass difference of the sample before and after immersion in the ferric chloride solution. In this study, the nickel masses of 0 g, 0.6719 g, 1.6407 g and

2.5981 g was dissolved in the waste solution for studying the growth of turbostratic stacked graphene. The sample names were designated by the mass of nickel dissolved in the waste solution *i.e.*, the GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981 were prepared by the mixture of 80 g ferric chloride, 100 ml deionized water and the dissolved nickel of 0 g, 0.6719 g, 1.6407 g and 2.5981 g, respectively.

Preparation of turbostratic stacked graphene

The 5 ml waste ferric chloride solution and 2 g sucrose were mixed and poured into an alumina crucible boat. Subsequently, the sample was vaporized at 90 °C for 24 hours. After that, the sample was transferred into the quartz tube furnace. Before annealing the sample, the argon gas was introduced for 30 minutes to evacuate air in the quartz tube. Thereafter, the sample was calcined in the quartz tube under argon atmosphere at 700 °C for 6 hours. Then the sample was fast cooled down to room temperature under argon atmosphere. The graphene-wrapped metal (iron and nickel) was obtained. The metal in the sample can be eliminated by immersing the sample in 6 M HCl for 6 hours to obtain turbostratic stacked graphene powder.

Characterization

Field emission scanning electron microscope (Fe-SEM) and energy-dispersive X-ray spectroscopy (EDX) were carried out by Hitachi UHR Fe-SEM SU8010 with incident beam of 20 kV. X-ray diffraction (XRD) measurement was performed using benchtop X-ray powder diffractometer (Bruker) with Cu-K α radiation ($\lambda = 0.154184$ nm). Micro-Raman measurement was conducted at room temperature using a 532 nm laser. The laser beam size is about 1 μ m in diameter. X-ray photoelectron spectroscopy (XPS) was measured by Axis Supra, Kratos using 225 W Al-K α monochromator. The base pressure was about 2.7×10^{-9} torr.

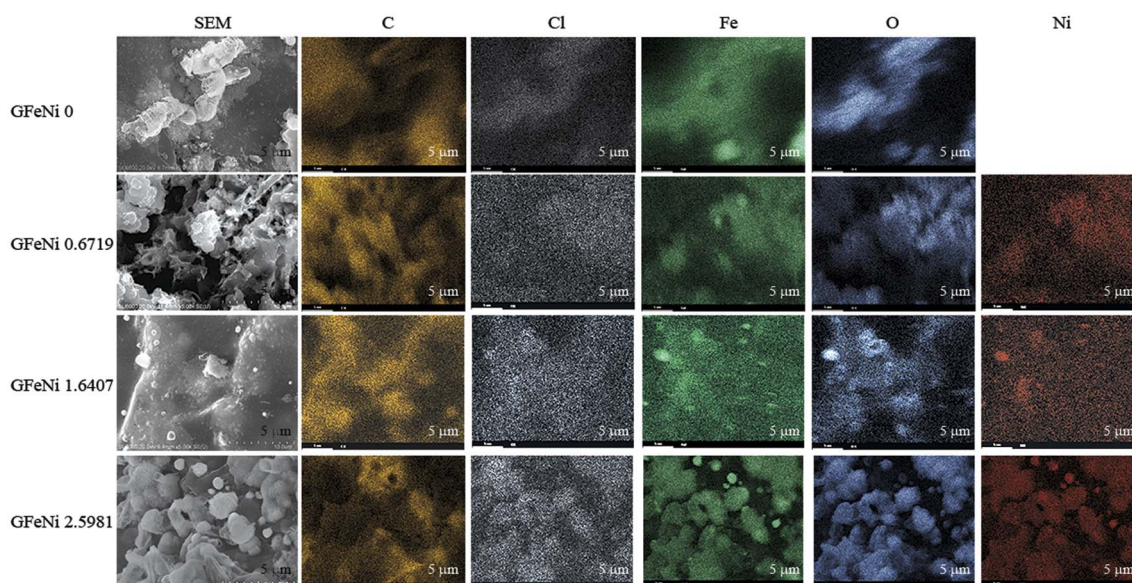


Fig. 1 Comparison of SEM images (left column) and EDX elemental distribution maps for carbon (C, orange), chlorine (Cl, white), iron (Fe, green), oxygen (O, blue) and nickel (Ni, red) of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981.

Results and discussion

Fig. 1 shows SEM images and EDX maps of the GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981. EDX maps display carbon element distributes on all sample surfaces meanwhile nickel element distributes on the sample surfaces of GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981 revealing that graphene can be grown on nickel and iron. In addition, the non-uniform carbon signal implies the non-uniform thickness of graphene.

XRD results of the GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981 are displayed in Fig. 2(a). The XRD patterns of all samples show the characteristic of graphene oxide and graphene peak at 2θ of $\sim 10^\circ$ and $\sim 26^\circ$, respectively, indicating the presence of graphene oxide and graphene on all samples. The interlayer spacing of graphene oxide is much larger than that of graphene due to the presence of oxygen-containing groups on the edge of each layer.¹⁷ In addition, the XRD pattern of GFeNi2.5981 displays the XRD peaks at 2θ of 35.39° , 42.84° , 43.81° , 44.59° , 51.05° and 74.92° which correspond to NiCl_2 ,¹⁸ FeCl_3 ,¹⁹ Fe_2O_3 (ref. 20) and NiO ,²¹ Ni^{2+} and Fe ,²⁰ Ni ,²¹ NiO ,²¹ respectively, implying that graphene can be grown on the surfaces of nickel and iron. It is in good agreement with EDX results which show Cl, Fe, O and Ni elements distribute on whole sample surface. Fig. 2(b) shows XRD experimental pattern and fitting curve around the graphene peak of GFeNi2.5981. The curve fitting is calculated using the following equation:²²

$$|F|^2 \propto |f(\theta)|^2 \left| \sum_{j=0}^{N-1} \beta_j e^{ikaj} \right|^2$$

where, F is a structure factor, $f(\theta)$ is an atomic scattering of carbon.²³ N is the number of graphene layer, β_j is an occupancy of $(j+1)^{\text{th}}$ graphene layer. $ka_j = (4\pi d_j \sin \theta)/\lambda$ where d_j is a graphene layer spacing. θ is an angle between the incident beam and the scattering planes. λ is wavelength of incident X-ray beam. For the graphene peak of GFeNi2.5981, the fitting curve is calculated using the parameters in Table S1.† The calculation curve reveals that the graphene layer spacing is 3.42 \AA and the thickness of graphene is non-uniform. The information about the graphene thickness and layer spacing of GFeNi2.5981 is shown in Fig. 2(c) and Table S2.† In addition, the XRD patterns reveal the graphene layer spacings of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981 are 3.43 \AA , 3.44 \AA , 3.44 \AA and 3.42 \AA , respectively. These graphene layer spacings are much larger than the interlayer spacing of AB-stacked graphene (3.35 \AA) implying all samples contain turbostratic stacked graphene.²⁴

Fig. 3(a) presents Raman spectra of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981. The results show the characteristic graphene peaks at $\sim 1570 \text{ cm}^{-1}$ (G band) and $\sim 2680 \text{ cm}^{-1}$ (2D band) indicating the presence of graphene in all samples. The intensity of G peak is higher than that of 2D peak. It is corresponding to the turbostratic stacked graphene prepared by laser-assisted process²⁵ and direct carbon ions implantation.²⁶ Whereas it is different from the turbostratic

stacked graphene prepared by physical vapor deposition (PVD)⁷ which shows the Raman intensity of 2D peak is significantly greater than that of G peak. In 2009, Casiraghi proved that the intensity of 2D peak drastically decreases when doping increases while the intensity of G peak is independent on the doping.²⁷ Therefore, the synthesis of graphene by different methods may produce turbostratic stacked graphene with different doping. The Raman spectra also show graphene D-band ($\sim 1340 \text{ cm}^{-1}$) which originates from the sp^2 -hybridized disordered carbon materials.²⁸ Fig. 3(b) displays the intensity ratio of D peak to G peak (I_D/I_G) and the graphene grain size of all samples. The graphene grain size is calculated by the following equation:²⁹

$$L_a(\text{nm}) = \frac{560}{E_{\text{laser}}^4(\text{eV})} \left(\frac{I_D}{I_G} \right)^{-1}$$

where L_a is the graphene grain size. E_{laser} is the energy of incident laser in the Raman measurements. I_D/I_G is the intensity ratio of graphene D peak to G peak. The calculation results show

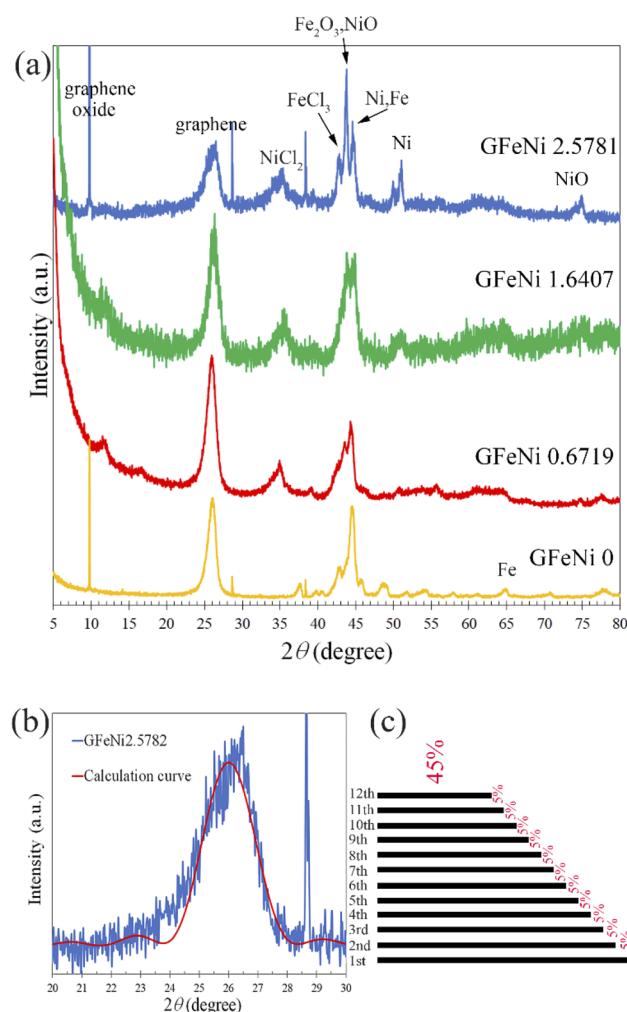


Fig. 2 (a) XRD patterns of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981 before etching the metal. (b) XRD experimental results (blue curve) and XRD calculation curve (red curve) of GFeNi2.5981 at graphene peak. (c) Proportion of graphene thickness in the GFeNi2.5981.



the graphene grain size decreases when the nickel mass increases. Moreover, all samples contain the weak intensity peak of iTA^- at $\sim 1850\text{ cm}^{-1}$ as shown in Fig. 3(c). The iTA^- Mode is originated from a combination of in-plane transverse acoustic (iTA) and the longitudinal optic (LO) phonons. The intensity of the iTA^- peak is high for monolayer graphene and turbostratic graphene but it reduces with increasing number of graphene layers. The iTA^- peak almost disappears in the case of AB-stacked few-layer graphene and it becomes absent for AB-stacked graphite.³⁰ Therefore, the appearance of iTA^- indicates the presence of turbostratic stacked graphene on all samples.

M band ($\sim 1750\text{ cm}^{-1}$) is an overtone of the oTO phonon. M band has been observed in bilayer and few-layer AB stacked graphene and highly ordered pyrolytic graphite (HOPG) because the M band is activated by strong interlayer interactions between adjacent graphene layers.³⁰ The M band is absent for single layer graphene and turbostratic stacked graphene. Therefore the absence of M band from the samples implies that the stacking pattern of graphene on all samples is turbostratic. It is in good agreement with the XRD results. Fig. 3(d)–(g) display the curve fitting on Raman spectra of 2D band of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981 that can well fit using a single Lorentzian peak (red) with the full width at half maximum (FWHM) of 85 cm^{-1} , 92 cm^{-1} , 96 cm^{-1} and 100 , respectively. These FWHM values are corresponding to the turbostratic stacked graphene prepared by laser-assisted process²⁵ and direct carbon ions implantation.²⁶ Whereas, these FWHM values are much higher than that of monolayer graphene ($\sim 24\text{ cm}^{-1}$)²⁴ and the turbostratic stacked graphene prepared by PVD ($27\text{--}65\text{ cm}^{-1}$). Although the FWHM of these samples is wide and similar to AB stacked multilayer graphene ($\sim 100\text{ cm}^{-1}$),⁷ the well-fitted 2D band by a single Lorentzian peak in all samples can confirm the structure of graphene is turbostratic stacking. In addition, the single Lorentzian fit can

reveal the value of the c -axis lattice constant using the following empirical formula³¹

$$c(\text{nm}) = 0.682 - 0.11 \left| \frac{I_{G'_{3DB}}}{I_{G'_{3DB}} + I_{G'_{2D}}} \right|$$

where $I_{G'_{2D}}$ and $I_{G'_{3DB}}$ are the intensities of the single Lorentzian peak and its coexisting Lorentzian, respectively, at 2D band. Since all samples can be well fit by a single Lorentzian peak, the $I_{G'_{3DB}}$ vanishes resulting in the c -axis lattice constant and interlayer spacing of graphene is 0.682 nm and 0.341 nm , respectively (the interlayer spacing of graphene is half of the value of c -axis lattice constant³²). The interlayer spacing of 3.41 \AA is near the XRD results that reveal the graphene layer spacings of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981 are 3.43 \AA , 3.44 \AA , 3.44 \AA and 3.42 \AA , respectively.

The vaporization is an important process to arrange carbon atoms around the metal before the growth of graphene by calcination at $700\text{ }^\circ\text{C}$. Fig. 4 shows XPS spectra of the C 1s of the GFeNi0 before and after calcination at $700\text{ }^\circ\text{C}$ for 6 hours. The XPS spectrum of GFeNi0 before calcination process is deconvoluted into 5 peaks at 283.0 eV , 284.6 eV , 286.3 eV , 287.5 eV and 289.1 eV , which correspond to the bonds of C-Fe, C-C (sp^3), C-O, C=O and O-C=O, respectively. The high intensity peak of C-Fe implies that after the sample underwent the vaporization process, the metal catalyst is surrounded by amorphous carbon and the C-Fe bonds occur between the metal catalyst and its nearest carbon atoms. After calcination at $700\text{ }^\circ\text{C}$ for 6 hours, the XPS spectrum of the GFeNi0 is deconvoluted into 5 peaks at 284.1 eV , 285.2 eV , 286.2 eV , 288.3 eV and 290.5 eV , which correspond to the bonds of C=C (sp^2), C-C (sp^3), C-O, C=O and O-C=O, respectively. The peak areas of C-O bond, C=O bond and O-C=O bond decrease significantly due to the carbon oxidation.³³ The dominant C=C (sp^2) corresponds to the π -bonded carbon atoms of the graphene network.³⁴

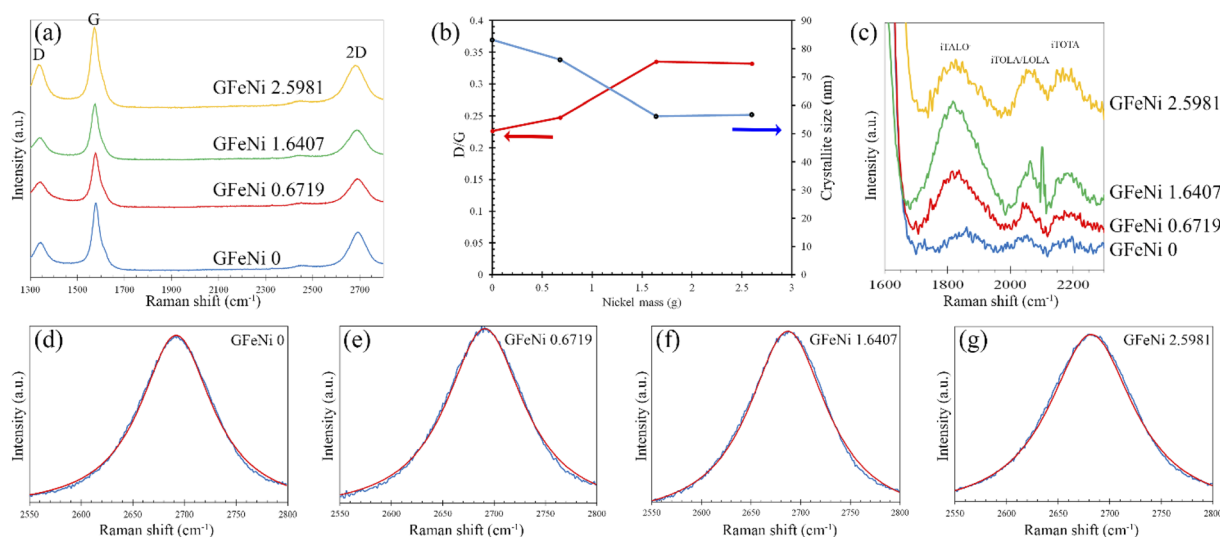


Fig. 3 (a) Raman spectra of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981. (b) D/G (left axis) and graphene grain size (right axis) of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981. (c) Magnified Raman spectra (after baseline subtraction) between 1600 cm^{-1} and 2300 cm^{-1} of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981. (d)–(g) Experimental Raman spectra (blue) and single Lorentzian fit (red) of the 2D band of GFeNi0, GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981, respectively.



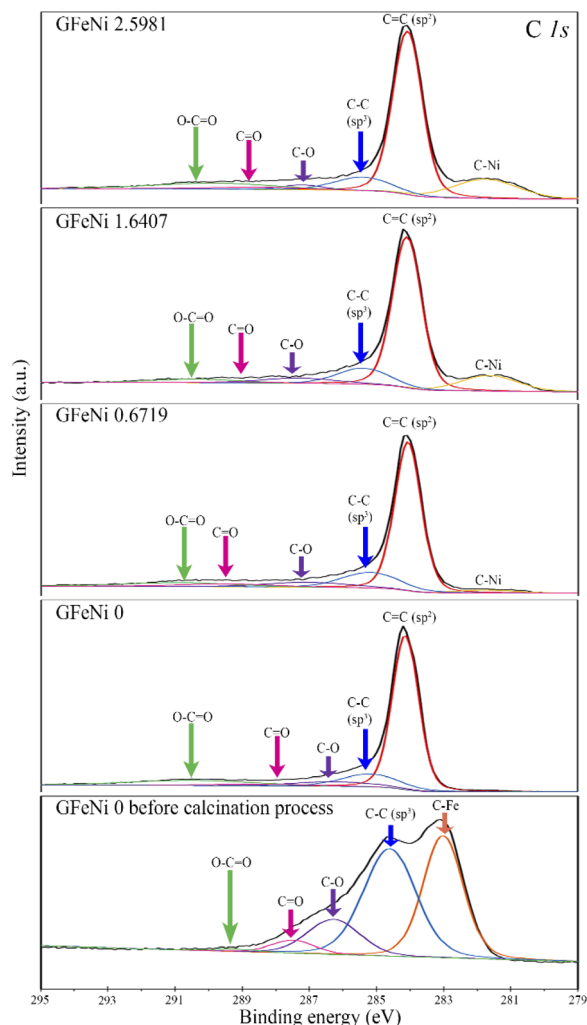


Fig. 4 XPS spectra of C 1s of the GFeNi0 before and after calcination process, GFeNi0.6719, GFeNi1.6407, GFeNi2.5981.

However, the C-Fe peak is absent from the XPS spectrum indicating there is not C-Fe bonding between iron and the first carbon layer (innermost carbon layer). Therefore, the first carbon layer is graphene layer. Fig. 4 also displays the XPS spectra of the C 1s of GFeNi0.6719, GFeNi1.6407 and GFeNi2.5981. A peak at ~ 281.7 eV corresponds to C-Ni bonds. The peak area of C-Ni bond increases when nickel content increases. The binding energy of this C-Ni peak is lower than that of nickel carbide thin film (283.3 eV).³⁵ The presence of C-Ni bonds implies that there are C-Ni bonds between nickel and the first carbon layer (innermost carbon layer). Therefore, the first carbon layer is buffer layer while the second carbon layer is graphene layer. It is similar to the growth of epitaxial graphene on Si-terminated SiC that the first carbon layer is buffer layer which contains Si-C bonds between the first carbon layer and the silicon layer underneath.^{36,37}

Fig. S3(a)† shows XRD patterns of the GFeNi0 prepared with and without vaporization process. The intensity ratio of graphene peak ($\sim 26^\circ$) to Fe peak ($\sim 44^\circ$) ($I_{\text{graphene}}/I_{\text{Fe}}$) for the GFeNi0 prepared with and without vaporization process are 0.72 and 0.33, respectively, revealing that the GFeNi0 prepared

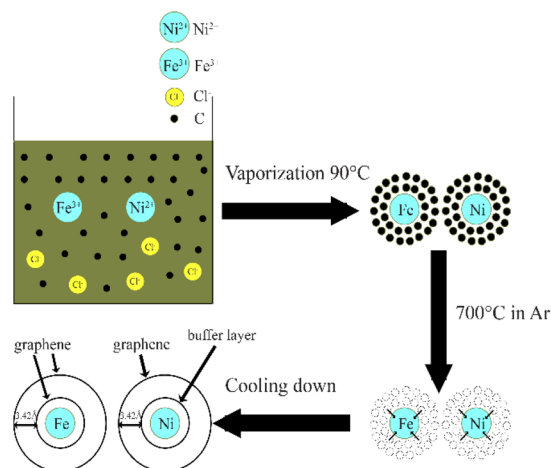


Fig. 5 Schematic diagram of growth mechanism of turbostratic graphene using waste ferric chloride solution as a feedstock.

with vaporization process contains a higher quantity of graphene since only the carbon atoms near the metal catalyst can be dissolved and formed graphene after the calcination at 700°C for 6 hours. After the graphene growth process, all samples can be attracted by a magnet (Fig. S3(b)†) due to the presence of metal inside. However, the magnetic property can be removed (Fig. S3(c)†) after etching the metal catalyst inside by immersion in HCl.

The growth mechanism of turbostratic stacked graphene using waste ferric chloride solution as a feedstock is illustrated schematically in Fig. 5. First, the sucrose is dissolved in waste ferric chloride solution. Subsequently, the solution was vaporization at 90°C for 24 hours resulting in the amorphous carbon surrounds the nickel and iron. The C-metal bonds occur between the metal and the nearest carbon atoms. After that the sample is calcined in the quartz tube under argon atmosphere at 700°C for 6 hours. At this annealing temperature, the surrounding carbon atoms are dissolved into the nickel and iron. Thereafter the sample is fast cooled down to room temperature, the carbon atoms in the nickel and iron precipitate and form turbostratic stacked graphene encloses the nickel and iron. The formation of turbostratic stacked graphene may arises from the fast heating and cooling rate.^{25,38,39} However, in the case of nickel metal catalyst, the first carbon layer is buffer layer while the second carbon layer is the first layer of graphene. The layer spacing of the turbostratic stacked graphene is ~ 3.42 Å.

Conclusion

We have demonstrated the growth of turbostratic stacked graphene using waste ferric chloride solution as a feedstock by vaporization at 90°C for 24 hours and calcination at 700°C for 6 hours. SEM images and EDX elemental distribution maps showed graphene can be grown on iron and nickel catalysts meanwhile XRD patterns confirm the presence of turbostratic stacked graphene with the interlayer spacing in the range of 3.42 Å to 3.44 Å. Moreover, the Raman spectra display the weak intensity peak of iTALO^- and the well-fitted 2D band by a single



Lorentzian peak indicating the presence of turbostratic stacked graphene. In addition, XPS spectra reveal the growth mechanism of the turbostratic stacked graphene. The first carbon layer (innermost carbon layer) which surrounds iron is graphene layer meanwhile the first carbon layer which surrounds nickel is buffer layer. This synthesis process of turbostratic stacked graphene is not only simple, low-cost, and large-scale production but also decreases the environmental issues from releasing the waste ferric chloride solution with improper disposal.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was supported by the Thailand Science Research and Innovation Fundamental Fund, and Thammasat University Research Fund, Contract No. TUFT-FF 8/2565, and Thammasat University Research Unit in Synthesis and Applications of Graphene.

Notes and references

- 1 A. K. Geim and K. S. Novoselov, *Nat. Mater.*, 2007, **6**, 183.
- 2 T. Ohta, A. Bostwick, J. L. McChesney, T. Seyller, K. Horn and E. Rotenberg, *Phys. Rev. Lett.*, 2007, **98**, 206802; K. V. Emtsev, F. Speck, Th. Seyller and L. Ley, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2008, **77**, 155303.
- 3 I. A. Ovid'ko, *Rev. Adv. Mater. Sci.*, 2013, **34**, 1.
- 4 A. A. Balandin, S. Ghosh, W. Bao, I. Calizo, D. Teweldebrhan, F. Miao and C. N. Lau, *Nano Lett.*, 2008, **8**, 902.
- 5 S. B. Touski and M. Hosseini, *Phys. E*, 2020, **116**, 113763.
- 6 J. Yang, Y. He, X. Zhang, W. Yang, Y. Li, X. Li, Q. Chen, X. Chen, K. Du and Y. Yan, *J. Mater. Res. Technol.*, 2021, **15**, 3005.
- 7 J. A. Garlow, L. K. Barrett, L. Wu, K. Kisslinger, Y. Zhu and J. F. Pulecio, *Sci. Rep.*, 2016, **6**, 19804.
- 8 C. Wei, R. Negishi, Y. Ogawa, M. Akabori, Y. Taniyasu and Y. Kobayashi, *Jpn. J. Appl. Phys.*, 2019, **58**, SIIB04.
- 9 A. Ruammaitree, *Surf. Rev. Lett.*, 2018, 1840003.
- 10 S. Bae, H. Kim, Y. Lee, X. Xu, J. S. Park, Y. Zheng, J. Balakrishnan, T. Lei, H. R. Kim, Y. Song, Y. J. Kim, K. S. Kim, B. Özyilmaz, J. H. Ahn, B. H. Hong and S. Iijima, *Nat. Nanotechnol.*, 2010, **5**, 574.
- 11 A. Reina, X. Jia, J. Ho, D. Nezich, H. Son, V. Bulovic, M. S. Dresselhaus and J. Kong, *Nano Lett.*, 2009, **9**, 30.
- 12 L. Huang, Q. H. Chang, G. L. Guo, Y. Liu, Y. Q. Xie, T. Wang, B. Ling and H. F. Yang, *Carbon*, 2012, **50**, 551.
- 13 R. B. S. Santos, O. Rocha and J. Povinelli, *Chemosphere*, 2007, **68**, 628.
- 14 K. S. Novoselov, A. K. Geim, S. V. Morozov, D. Jiang, Y. Zhang, S. V. Dubonos, I. V. Grigorieva and A. A. Firsov, *Science*, 2004, **306**, 666.
- 15 R. Mangaiyarkarasi, N. Santhiya and S. Umadevi, *Colloids Surf., A*, 2022, **642**, 128673.
- 16 B. Zhang, J. Song, G. Yang and B. Han, *Chem. Sci.*, 2014, **5**, 4656.
- 17 F. T. Johra, J. Lee and W. Jung, *J. Ind. Eng. Chem.*, 2014, **20**, 2883.
- 18 Y. Chen, S. Shen, J. Gu, Z. Zhang and N. Li, *ACS Omega*, 2020, **5**, 27278.
- 19 Sh. M. Nasef, N. A. Badawy, F. H. Kamal, S. F. Sherbiny and E. M. El-Nesr, *Arab J. Nucl. Sci. Appl.*, 2019, **52**, 209.
- 20 S. H. Khezri, A. Yazdani and R. Khordad, *Eur. Phys. J.: Appl. Phys.*, 2012, **59**, 30401.
- 21 J. T. Richardson, R. Scates and M. V. Twigg, *Appl. Catal., A*, 2003, **246**, 137.
- 22 A. Ruammaitree, H. Nakahara, K. Akimoto, K. Soda and Y. Saito, *Appl. Surf. Sci.*, 2013, **282**, 297.
- 23 P. Brown, A. Fox, E. Maslen, M. O'Keefe and B. Willis, *International Tables for Crystallography C*, 2004, vol. 6, p. 554.
- 24 L. M. Malard, M. A. Pimenta, G. Dresselhaus and M. S. Dresselhaus, *Phys. Rep.*, 2009, **473**, 51.
- 25 M. Athanasiou, N. Samartzis, L. Sygellou, V. Dracopoulos, T. Loannides and S. N. Yannopoulos, *Carbon*, 2021, **172**, 750.
- 26 K. Liu, F. Lu, K. Li, Y. Xu and C. Ma, *Appl. Surf. Sci.*, 2019, **493**, 1255.
- 27 C. Casiraghi, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2009, **80**, 233407.
- 28 M. J. Matthews, M. A. Pimenta, G. Dresselhaus, M. S. Dresselhaus and M. Endo, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, R6585.
- 29 P. Machac and T. Hrebicek, *J. Mater. Sci.: Mater. Electron.*, 2017, **28**, 12425.
- 30 R. Rao, R. Podila, R. Tsuchikawa, J. Katoch, D. Tishler, A. M. Rao and M. Ishigami, *ACS Nano*, 2011, **5**, 1594.
- 31 L. G. Cancado, K. Takai, T. Enoki, M. Endo, Y. A. Kim, H. Mizusaki, N. L. Speziali, A. Jorio and M. A. Pimenta, *Carbon*, 2008, **46**, 272.
- 32 J. Hass, W. A. Heer and E. H. Conrad, *J. Phys.: Condens. Matter*, 2008, **20**, 323202.
- 33 L. Zhou, Chapter 3 – Fundamentals of Combustion Theory, in *Theory and Modeling of Dispersed Multiphase Turbulent Reacting Flows*, 2018, vol. 15.
- 34 M. K. Rabchinskii, S. D. Saveliev, D. Y. Stolyarova, M. Brzhezinskaya, D. A. Kirilenko, M. V. Baidakova, S. A. Ryzhkov, V. V. Shnitov, V. V. Sysoev and P. N. Brunkov, *Carbon*, 2021, **182**, 593.
- 35 A. Furlan, J. Lu, L. Hultman, U. Jansson and M. Magnuson, *J. Phys.: Condens. Matter*, 2014, **26**, 415501.
- 36 A. Ruammaitree, H. Nakahara and Y. Saito, *Appl. Surf. Sci.*, 2014, **307**, 136.
- 37 W. Strupinski, K. Grodecki, P. Caban, P. Ciepielewski, I. Jozwik-Biala and J. M. Baranowski, *Carbon*, 2015, **81**, 63.
- 38 S. N. Yannopoulos, A. Siokou, N. K. Nasikas, V. Dracopoulos, F. Ravani and G. N. Papatheodorou, *Adv. Funct. Mater.*, 2012, **22**, 113.
- 39 D. X. Luong, K. V. Bets, W. A. Algozeeb, M. G. Stanford, C. Kittrell, W. Chen, R. V. Salvatierra, M. Ren, E. A. McHugh, P. A. Advincula, Z. Wang, M. Bhatt, H. Guo, V. Mancevski, R. Shahsavari, B. I. Yakobson and J. M. Tour, *Nature*, 2020, **577**, 647.

