# RSC Advances



# PAPER

Cite this: RSC Adv., 2022, 12, 35418

Received 14th September 2022 Accepted 25th November 2022

DOI: 10.1039/d2ra05805g

rsc.li/rsc-advances

## 1. Introduction

The burning of fossil fuels for production of electricity and heat is the most important source of greenhouse gas emissions, especially carbon dioxide  $(CO<sub>2</sub>)$ , which causes climate changes and global warming. Consequently, various scientific efforts have been made to capture  $CO<sub>2</sub>$  and protect the environment from its dangerous effects. In the last decade, several methods have been developed for  $CO<sub>2</sub>$  capture, for example adsorption,<sup>1</sup> absorption,<sup>2</sup> membrane separation<sup>3</sup> and electrochemical separation.<sup>4</sup> At present, industries use aqueous amine solutions

# DFT and COSMO-RS studies on dicationic ionic liquids (DILs) as potential candidates for  $CO<sub>2</sub>$ capture: the effects of alkyl side chain length and symmetry in cations†

Mehrangiz Torkzadeh and Majid Moosavi<sup>D\*</sup>

A few studies on CO<sub>2</sub> capture using dicationic ionic liquids (DILs) show that they are more promising absorbents for  $CO<sub>2</sub>$  capture than monocationic ILs (MILs). Ion–ion, ion– $CO<sub>2</sub>$  and DIL molecule– $CO<sub>2</sub>$ interactions are important for understanding the performance–structure–property relationships for the rational design of DILs for  $CO<sub>2</sub>$  capture applications. However, the role of these interactions in determining  $CO<sub>2</sub>$  solubility in DILs is unclear. In this study, we used DFT methods to understand these interactions in three selected DILs)considering the effects of alkyl side chain length and symmetry in cations (by exploring different aspects, such as the electronic and geometrical structures, topological properties and the strength and nature of interactions, charge transfer, etc. The results showed that the most suitable solvent for CO<sub>2</sub> is the symmetric DIL with a longer side chain length, *i.e.* [Bis(mim)C<sub>5</sub>-(C<sub>4</sub>)<sub>2</sub>] [NTf<sub>2</sub>]<sub>2</sub>. In addition, we used the COSMO-RS calculations to obtain the macroscopic solubility of CO<sub>2</sub> in the studied DILs, which was in good agreement with the DFT results. Gas selectivity results calculated using COSMO-RS theory indicated that the selectivity of  $CO<sub>2</sub>$  from H<sub>2</sub>, CO and CH<sub>4</sub> gases decreases slightly with increasing the length of side alkyl chains. PAPER<br> **(A)** Check for updates<br>
Case this scale, 2007, 2007 and COSMO-RS studies on dicationic ionic<br>
Case this scale, 2007, 2008 **Case in the Case of the Case o** 

because of their high tendency to adsorb gas, specifically for  $CO<sub>2</sub>$  capture. However, these solvents have disadvantages such as high volatility, poor stability, high cost, high energy demand for regeneration, corrosion of equipment, environmental damage associated with amines,  $etc.56$  During the past years, ionic liquids (ILs), composed of inorganic/organic anions and organic cations, have been receiving more attention as potential alternatives to conventional amine solutions due to their desirable properties such as negligible vapor pressure, high thermal stability, fine gas solubility and so on.<sup>7,8</sup> They can be designed and assembled by altering the structure of their cations and anions to develop functional ILs for selective and effective  $CO<sub>2</sub>$  capture.<sup>9</sup> Dicationic ionic liquids (DILs) are a special class of ILs composed of a doubly charged cation which have two singly charged constituents linked by an alkyl chain paired with two anions.<sup>10</sup> It is found that DILs possess wider liquid range, higher melting point, and better thermal stability compared to conventional monocationic ILs (MILs). The physicochemical properties of DILs are more tunable than the corresponding MILs due to the higher combinations of anions and cations.<sup>11</sup> From the structural point of view, additional variability of DILs arise from the nature of the cation or from the type and length of their linkage chain. Therefore, DILs present several advantages over traditional MILs in their applications as solvents,<sup>12</sup> lubricants,<sup>13</sup> and separation media,<sup>14</sup> etc. Although, a noticeable quantity of experimental and

Department of Physical Chemistry, Faculty of Chemistry, University of Isfahan, Isfahan 81746-73441, Iran. E-mail: m.mousavi@sci.ui.ac.ir; Fax: +98-313-668-9732; Tel: +98- 313-7934942

<sup>†</sup> Electronic supplementary information (ESI) available: Fig. S1. Calculated spatial distribution of HOMO and LUMO for isolated ions; Fig. S2. Topological graphics of studied DILs obtained from QTAIM calculations; Fig. S3 and S4. Optimized structures of isolated ion-CO<sub>2</sub> and DIL-CO<sub>2</sub> complexes; Fig. S5. Calculated spatial distribution of HOMO and LUMO for DIL-CO<sub>2</sub> complexes; Fig. S6. Color-filled RDG isosurfaces and scatter graphs of RDG for DIL-CO<sub>2</sub> complexes; Fig. S7 and S8. Topological graphics of isolated ion– $CO<sub>2</sub>$  and DIL– $CO<sub>2</sub>$ complexes obtained from QTAIM calculations; Table S1. Obtained parameters from QTAIM calculations in the studied DILs. Table S2. Charge transfer calculated in the studied DILs using CHELPG and NBO methods. Table S3. NBO analysis of the studied DILs; Table S4. Interaction energies, empirical dispersion and basis set superposition errors for studied ion– $CO<sub>2</sub>$  and DIL– $CO<sub>2</sub>$ complexes. See DOI: <https://doi.org/10.1039/d2ra05805g>

theoretical works has been invested in studying  $CO<sub>2</sub>$  capture using ILs, most of them have focused on conventional imidazolium-,<sup>15,16</sup> pyridium-<sup>17</sup> and phosphonium-based MILs<sup>18</sup> and DILs have received very little attention of researchers. Few studies on  $CO<sub>2</sub>$  capture using DILs show that they are more promising absorbents for  $CO<sub>2</sub>$  capture than MILs.<sup>19–22</sup> Zhang et al.<sup>19</sup> showed that the  $CO<sub>2</sub>$  absorption capacities of six DILs, including four amino acid-functionalized DILs (AADILs) and two traditional DILs (T-DILs), were enhanced heavily compared to MILs. Hojniak et  $al.^{20}$  reported that DILs-incorporated membranes showed the increased  $CO<sub>2</sub>/N<sub>2</sub>$  selectivity and adsorption capability compared to MILs-incorporated membrane, which was attributed to the increased  $CO<sub>2</sub>$  interaction sites in DILs. Similarly, Zhang et  $al.^{22}$  observed the enhanced  $CO<sub>2</sub>$  capture by imidazolate-based DILs, ([Bis(mim)  $C_2$ [Im]<sub>2</sub>) and ([Bis(mim)C<sub>4</sub>]-[Im]<sub>2</sub>) in which the CO<sub>2</sub> capture by DILs is twice that of their MIL analogues. Therefore, the broader and deeper knowledge of structural and electronic properties of ions of DILs is important for designing new DILs for  $CO<sub>2</sub>$ capture.

Theoretical studies are a powerful tool in the insight of nanoscopic features in controlling  $CO<sub>2</sub>$  absorption by ILs.<sup>23-27</sup> The computational studies of gas phase structure of  $ILs-CO<sub>2</sub>$ systems are of great importance, since the ionic interactions in ILs are often strong to maintain gas-phase structure in liquid phase as well.<sup>28</sup> On the other hand, charge distributions are important in describing the structural and dynamical properties in liquid phase, and the partial atomic charges in the force fields of molecular dynamics (MD) simulations of liquid phase are obtained using electronic density and molecular electrostatic potential in gas phase.<sup>29</sup> Quantum mechanical methods have proven capable in characterizing  $IL-CO<sub>2</sub>$  systems, providing valuable information on various aspects such as structural changes, charge transfer, intermolecular hydrogen bonding, etc.<sup>30–34</sup> The quantum studies on large systems containing large numbers of ions and  $CO<sub>2</sub>$  molecules are not possible due to computational constraints. However, studies on anion–cation and ion– $CO<sub>2</sub>$  pairs and the analysis of the shortrange interactions in these systems are useful and provide valuable information about these systems. In general, in previous quantum studies on MILs– $CO<sub>2</sub>$  systems, the interactions have been calculated from placing  $CO<sub>2</sub>$  molecules around individual ions or IL molecules.<sup>35-37</sup> In summary, for common MILs, the  $CO<sub>2</sub>$  solubility depends more on the cation-anion interactions rather than the  $CO_2$ -anion or  $CO_2$ -cation interactions. However, the correlations are more complicated for DILs, which deserve much further analyses.

Although, several studies have been investigated the properties of pure DILs by quantum methods,<sup>38-48</sup> to the best of our knowledge, no research has been performed on the structural and electronic properties of DILs in the presence of  $CO<sub>2</sub>$  gas using quantum mechanical methods. Therefore, in this work, in order to identify the main sites effective in  $CO<sub>2</sub>$  capture, the structures of three selected DILs (considering the effects of alkyl side chain length and symmetry in cations) in the absence and presence of  $CO<sub>2</sub>$  are investigated. For each system, the interaction sites, the electronic and geometrical structures, topological properties and the strength and nature of intercations, charge transfer, etc. have been explored. Furthermore, some previous studies, have been also shown that COSMOthermX<sup>49</sup> is able to predict the solubility of a gas in the ILs.<sup>50-52</sup> Therefore, at the end of this work, the predicted gas absorption in the studied DILs is determined using COSMOthermX based on the absorption isotherms and the results compared with the corresponding quantum mechanical findings.

### 2. Computational details

### 2.1 Quantum mechanical calculations

At first, in order to investigate the molecular structure and cation–anion interactions in the selected DILs, one chargeneutral DIL molecule consisting of one cation and two anions (ion triplet) was further investigated by density functional theory (DFT) calculations. The structures of the selected cations and anions along with the labeling of the atoms are shown in Fig. 1. Since the size of a cation in a DIL is large, one DIL molecule (one cation and two anions) may has hundreds of relatively stable configurations. As a result, to find the most stable conformation for each DIL, each DIL was simulated in a big simulation box for 2 ns using NVT ensemble at 298 K. The conformations obtained from simulations were optimized using M06-2X functional and the cc-pVDZ basis set $53,54$  with Gaussian 09 package.<sup>55</sup> Also, the presence of true minima was confirmed by the absence of imaginary frequencies in the calculated vibrational spectra. Recently, researchers have used the M06-2X functional to investigate the physicochemical properties of ILs. M06-2X has been shown remarkable success in predicting reliable structures of the complexes<sup>56</sup> and describes the non-covalent interactions better than density functional which are commonly used. Paper<br>
Recording the state is the most of the most one of the most one of the most one of the most one of the most of the most one of the most of the

After exploring DILs in pure state, their interaction with  $CO<sub>2</sub>$ was studied by considering isolated ions– $CO<sub>2</sub>$  and DIL– $CO<sub>2</sub>$ systems. For this purpose, in the second step, after the separate optimization of ions and  $CO<sub>2</sub>$  molecule, the cation– $CO<sub>2</sub>$ , anion–  $CO<sub>2</sub>$  and DIL– $CO<sub>2</sub>$  complexes were optimized at the same theoretical level of the optimization of pure DILs  $(i.e.$  M06-2X/ cc-pVDZ). For each complex, different orientations of  $CO<sub>2</sub>$  in interaction with isolated ions and DILs were considered and the most stable structures have been chosen to carry out further calculations. The total electronic energies for all the studied systems were corrected using Grimme's DFT-gCP-D3 correction method in which empirical dispersion correction (D3) has been developed by Grimme<sup>57</sup> and also the basis set superposition error (BSSE)<sup>58</sup> corrections were applied. Charge transfer in all systems was calculated by charges from the electrostatic potential using a grid-based method (CHELPG)<sup>59</sup> and natural bonding orbital (NBO) methods at the same theoretical level. Furthermore, to gain further insight into the nature of interactions within the studied systems, NBO analysis has been carried out.<sup>60</sup> In order to analyze bond critical point (BCP) data in the studied systems, the quantum theory of atoms in molecules  $(QTAIM)^{61}$  was done using AIM2000 package.<sup>62</sup> HOMO– LUMO band gap of all systems were calculated using MultiWFN program.<sup>63</sup> Moreover, for identification of non-covalent





 $[NTf_2]$ [bis(trifluoromethanesulfonyl)imide]

Fig. 1 Structures of [Bis(mim)C<sub>5</sub>]<sup>2+</sup>, [N<sub>111</sub>-C<sub>5</sub>-mim]<sup>2+</sup>, [Bis(mim)C<sub>5</sub>-(C<sub>4</sub>)<sub>2</sub>]<sup>2+</sup> cations and [NTf<sub>2</sub>]<sup>–</sup> anion along with their atomic labeling schemes.

interactions (NCIs), the reduced density gradient (RDG) was calculated using NCIPLOT program<sup>64</sup> and RDG plots were generated through the visual molecular dynamics program  $(VMD).<sup>65</sup>$ 

### 2.2 COSMO-RS theory calculations

It has been verified experimentally that the conductor-like screening model for real solvents (COSMO-RS) theory is a reliable solubility prediction method for ILs systems that combines statistical thermodynamics and quantum chemical calculations.<sup>49</sup> Therefore, to obtain the solubility of  $CO<sub>2</sub>$  in DILs

macroscopically, COSMO-RS theory of molecular thermodynamics was used to get macroscopic gas solubility data. In this work, BP functional and TZVP basis set from Turbomole program<sup>66</sup> were used for COSMO calculations. Then, the generated COSMO files were imported into COSMOtherm program<sup>67</sup> to compute  $CO<sub>2</sub>$  solubility in the studied DILs.

### 3. Results and discussion

### 3.1 DFT calculations

3.1.1 Pure DILs (ion triplet). The possible interaction sites around cations and anions can be characterized from charge



Fig. 2 Calculated ESP maps (isovalue = 0.0004) of (a) isolated ions and (b) the studied DILs at M06-2X/ccpVDZ level of theory.

distribution obtained by electrostatic potential (ESP) maps.<sup>68</sup> The ESP maps help to better visualize the charged regions in species. Fig. 2(a) shows the ESP maps calculated for isolated cations and anions at M06-2X/ccpVDZ level of theory, where the blue-colored surfaces signify positive electrostatic potential and red-colored surfaces represent negative ones. The color codes of

these maps for the ions are in the range of −0.200 to 0.200 a.u for  $[NTf_2]^-$  anion,  $-0.264$  to 0.264 a.u,  $-0.268$  to 0.268 a.u and  $-0.262$  to 0.262 a.u for [Bis(mim)C<sub>5</sub>]<sup>2+</sup>, [N<sub>111</sub>-C<sub>5</sub>-mim]<sup>2+</sup> and [Bis(mim) $C_5$ - $(C_4)_2$ ]<sup>2+</sup> cations, respectively. From ESP values, it can be found that the electron accepting property of  $[N_{111}$ -C<sub>5</sub>- $\text{min}^{2+}$  cation is greater than other ones. As it can be seen, all

regions around cations have positive electrostatic potential. The intensity of colors indicates that due to the high potential (dark blue), imidazolium ring and ammonium head group of cations are favorable for interactions with the oxygen and nitrogen atoms of anions (dark red). On the other hand, positive charges in free cation are localized on the H atoms of imidazolium rings, and methyl groups of ammonium head. Fig. 2(b) shows the calculated ESP maps of the studied DILs with a combination of red, green, blue, and yellow colors. The color codes of these maps for DILs are in the range between −0.079 to 0.079 a.u,  $-0.090$  to 0.090 a.u and  $-0.074$  to 0.074 a.u for [Bis(mim)C<sub>5</sub>]  $[NTf_2]_2$ ,  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> and  $[Bis(min)C_5$ - $(C_4)_2$ ][NTf<sub>2</sub>]<sub>2</sub>, respectively. From ESP values, it seems that the configuration of  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> is more stable than the other two symmetric DILs.

The possible configurations of anions around cation have many effects on hydrogen bondings, interaction energies and physicochemical properties of DILs.<sup>45</sup> The most stable structures obtained for the studied DILs have been represented in Fig. 3. It can be found that in all these structures, there exist multiple hydrogen bonds between the cations and the anions, which play an important role for stabilizing the DILs. As it can be seen, in all three structures, the linkage chain entangles in order to maximize the interactions of the two heads of DILs with two anions. Compared with asymmetric  $[N_{111} - C_5 - mim][NTf_2]_2$ , in two symmetric DILs, the middle chain folding is more and two imidazolium rings are almost parallel to each other. Bodo et al.<sup>39</sup> and Sun et al.<sup>38</sup> showed that the stability of ion pairs increase as two imidazolium rings become more compact for  $[C_n]$ dmim][NTf<sub>2</sub>]<sub>2</sub> with  $n = 3$ , 9 and for  $[C_3]$ (mim)<sub>2</sub>]Br<sub>2</sub> DILs, respectively. Also, Soltanabadi and Bahrami<sup>47</sup> showed that the two rings bend towards each other and two halides are placed in the middle of two rings for imidazolium-based DILs. This situation allows the halide anions to have better interaction with special sites of cation. As it can be seen in Fig. 3, the minimum observed distance in the studied DILs is between the anion atoms and the HR (HR′ ) atoms of the rings which is in agreement with the observed results for imidazolium based  $MILs$ <sup>69</sup> DILs<sup>70</sup> and tricationic ILs (TILs).<sup>71</sup> The minimum distances between the oxygen atoms of anions and HR (HR′ ) atoms are 2.19, 2.18 and 2.23 Å for [Bis(mim)C<sub>5</sub>][NTf<sub>2</sub>]<sub>2</sub>, [N<sub>111</sub>- $C_5$ -mim][NTf<sub>2</sub>]<sub>2</sub> and [Bis(mim) $C_5$ - $(C_4)_2$ ][NTf<sub>2</sub>]<sub>2</sub>, respectively.

Interaction energy  $(\Delta E_{int})$  is one of the most powerful tools for estimating the strength of non-covalent interactions, which is defined as difference between electron energy of DIL molecule and sum of the energies of cation and anion species:

$$
\Delta E_{\rm int} = E_{\rm DIL} - (E_{\rm C} + 2E_{\rm A})\tag{1}
$$

The values of interaction energies, and also BSSE and D3 corrections for the most stable structures observed in studied DILs (Fig. 3) are reported in Table 1. According to this table, the trend of interaction energy values (with corrections) is as follows:  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> > [Bis(mim)C<sub>5</sub>][NTf<sub>2</sub>]<sub>2</sub> > [Bis(mim)  $C_5$ - $(C_4)_2$ [NTf<sub>2</sub>]<sub>2</sub>. It seems that in asymmetric DIL [N<sub>111</sub>-C<sub>5</sub>-mim]  $[NTf<sub>2</sub>]<sub>2</sub>$ , by replacing a imidazolium group with the ammonium group, the interaction energy between the cations and anions



Fig. 3 Optimized structures of (a)  $[Bis(min)C_5][NTf_2]_2$ , (b)  $[N_{111} - C_5$  $mim][NTf_2]_2$  and (c)  $[Bis(min)C_5-(C_4)_2][NTf_2]_2$  at M06-2x/cc-pVDZ level of theory. Broken lines show hydrogen bonding interactions (distances are in Å)

increases due to the more localized charge in the ammonium group. Similarly, Fernandes et al.<sup>72</sup> showed that aromatic-based ILs (pyridinium- and imidazolium-based) due to charge delocalization in aromatic rings have lower relative interaction strength values compared with their saturated counterparts (pyrrolidinium- and piperidinium-based). Also, by increasing "side" chain length in [Bis(mim) $C_5$ - $(C_4)_2$ ][NTf<sub>2</sub>]<sub>2</sub>, the interaction energy decreases. A similar trend of decrease of interaction

Table 1  $\;$  Interaction energies ( $\Delta E_{\rm int}$  in kcal mol $^{-1}$ ), empirical dispersion  $(D_3)$ , and basis set superposition errors (BSSEs in kcal mol<sup>-1</sup>) for the studied DILs.  $\Delta E_c$  is the interaction energy corrected by  $D_3$  and BSSE

System	$\Delta E_{\rm int}$	$D_3$	$\Delta E_{\rm{BSSE}}$	$\Delta E_c$
"Entangled" configuration				
$[\text{Bis}(\text{mim})\text{C}_5]\text{[NTf}_2]_2$	$-253.22$	$-6.67$	29.08	$-230.81$
$[N_{111}$ -C <sub>5</sub> -mim][NTf <sub>2</sub> ] <sub>2</sub>	$-254.47$	$-6.40$	26.40	$-234.47$
$[\text{Bis}(mim)C_5-(C_4)_2][\text{NTf}_2]_2$	$-248.41$	$-9.17$	33.47	$-224.11$
"Stretched" configuration				
$[\text{Bis}(\text{mim})\text{C}_5]\text{[NTf}_2]_2$	$-224.10$	$-5.87$	43.10	$-186.87$
$[N_{111}$ -C <sub>5</sub> -mim][NTf <sub>2</sub> ] <sub>2</sub>	$-229.22$	$-5.43$	39.10	$-195.55$
$[\text{Bis}(mim)C_5-(C_4)_2][\text{NTf}_2]_2$	$-208.03$	$-5.86$	36.40	$-177.49$

energies with increasing the length of "linkage" alkyl chain between rings in DILs has been reported.<sup>70</sup>

In Table 1, the interaction energies of the most stable stretched structures have been also tabulated. As it can be seen, the structures with a "stretched" configuration of the linkage chain have higher energies than the "entangled" structures. The "stretched" configurations of  $[\text{Bis}(\text{min})\text{C}_5][\text{NTf}_2]_2$ ,  $[\text{N}_{111}\text{-C}_5]$ mim][NTf<sub>2</sub>]<sub>2</sub> and [Bis(mim)C<sub>5</sub>-(C<sub>4</sub>)<sub>2</sub>][NTf<sub>2</sub>]<sub>2</sub> have about 43.94, 38.92, 46.62 kcal mol−<sup>1</sup> higher than the entangled ones, respectively. Similarly, Bodo et al.<sup>39</sup> reported that the energy of "stretched" chain structure is about 40 kcal mol<sup>-1</sup> higher than the "entangled" one for  $[C_9]$ dmim $[[NTf_2]_2]$ . In general, the calculated interaction energies between cation and two anions in the studied DILs are much higher than reported interaction energies between ion pairs in MILs.73,74 This can be due to the stronger coulombic anion–cation attractions in DILs compared to the interaction between ions in MILs which is responsible for higher thermal stability and melting point in DILs. This result is in agreement with those obtained by Soltanabadi and Bahrami for TILs.<sup>47</sup> They reported that the interaction energies of TILs are more than three times of their corresponding MILs. Also, Farmanzadeh et al.<sup>40</sup> reported the following trend for the interaction energies in comparison among several MILs and DILs with halide anions:  $C_6$ (mim)<sub>2</sub>Cl<sub>2</sub> >  $C_6$ (mim)<sub>2</sub>Br<sub>2</sub> >  $C_6(min)_2I_2 > C_6(min)Cl > C_6(min)Br > C_6(min)I.$ Puper<br>
The 1 interaction energy Articles. Published on 09 December 2022. Downloaded on 1/25/2025 1:49:46 PM. This article is licensed under the main of the creation of the creation-non-<br>
The article is the meaning proposi

The frontier molecular orbitals including the lowest unoccupied molecular orbital (LUMO) and the highest occupied molecular orbital (HOMO) are key factors in chemical stability of a system.<sup>75</sup> Generally, LUMO as an electron acceptor

represents the ability of molecule to capture an electron; whereas HOMO represents the ability of molecules to donate an electron. The electronic properties of a chemical species may be described in terms of global quantum molecular descriptors such as HOMO–LUMO gap ( $\Delta E_{\text{gap}}$ ), chemical hardness ( $\eta$ ) and chemical potential  $(\mu)$  which are calculated as:<sup>76</sup>

$$
\eta = \frac{(E_{\text{LUMO}} - E_{\text{HOMO}})}{2} \tag{2}
$$

$$
\mu = \frac{(E_{\text{HOMO}} + E_{\text{LUMO}})}{2} \tag{3}
$$

The calculated values of the HOMO, LUMO energies,  $\Delta E_{\text{can}}$ ,  $\eta$ and  $\mu$  for the isolated cations, anions and the studied DILs have been reported in Table 2. The magnitude of band gaps ( $\Delta E_{\text{gap}}$ ) implies the strength of interactions in an IL, that is, stronger interactions between cations and anions lead to a further reduction in energy gap. As it can be seen, the HOMO–LUMO energy gap follows the order of  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> (−5.91 eV)  $\langle$  [Bis(mim)C<sub>5</sub>][NTf<sub>2</sub>]<sub>2</sub> (−5.99 eV)  $\langle$  [Bis(mim)C<sub>5</sub>-(C<sub>4</sub>)<sub>2</sub>][NTf<sub>2</sub>]<sub>2</sub> (−6.33 eV) which is in agreement with trend of interaction energies. Chemical hardness is a measure of inhibition of charge transfer to the molecule, and species with less polarizability have higher hardness values. Considering chemical hardness, if a molecule has a large  $\Delta E_{\text{gap}}$ , it is a hard molecule, and if it has a small  $\Delta E_{\text{gap}}$ , it is a soft molecule. Also, the stability of a molecule can be related to its hardness, as a small  $\Delta E_{\text{gap}}$ leads to high chemical reactivity.

According to the values in Table 2, the calculated chemical hardness has the trend  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> < [Bis(mim)C<sub>5</sub>]  $[NTf_2]_2$  <  $[Bis(min)C_5-(C_4)_2][NTf_2]_2$  for the studied DILs. Furthermore, electronic chemical potential describes tendency of the electron to escape from a stable system. The observed trend for the chemical potential of DILs is the reverse of the chemical hardness trend. Therefore, it is plausible to say that with increasing symmetry and increasing side chain length of cations, the stability of DILs increases. The spatial distribution of frontier molecular orbitals is important in describing electrical properties of ILs. Thus, spatial distribution of HOMO and LUMO for isolated cation, anions and the studied DILs are visualized in Fig. S1<sup>†</sup> and 4, respectively. From these figures, it can be seen that HOMOs are mainly concentrated on two anions, while LUMOs are mainly located on the imidazolium rings of cations. Thus, the charge transfer can be occurred from

Table 2 Energies of HOMO and LUMO, HOMO–LUMO energy gap ( $E_{gap}$ ), chemical hardness ( $\eta$ ) and chemical potential ( $\mu$ ) of isolated ions and DILs calculated at M06-2X/ccpVDZ level

System	$E_{\text{HOMO}}$ /eV	$E_{\text{LUMO}}/\text{eV}$	$E_{\rm gap}$ /eV	$\eta$ /eV	$\mu$ /eV	
$[\text{Bis}(mim)C_5]^{2+}$	$-13.24$	$-6.92$	$-6.31$	3.16	$-10.08$	
$\rm [N_{111}\text{-}C_5\text{-}min]^{2+}$	$-13.54$	$-6.92$	$-6.62$	3.31	$-10.23$	
$[\text{Bis}(mim)C_5-(C_4)_2]^{2+}$	$-12.93$	$-6.67$	$-6.25$	3.13	$-9.80$	
$[NTf_2]$	$-3.90$	3.54	$-7.44$	3.72	$-0.18$	
$[\text{Bis}(mim)C_5]\text{[NTf}_2]_2$	$-7.16$	$-1.17$	$-5.99$	2.99	$-4.16$	
$[N_{111}$ -C <sub>5</sub> -mim][NTf <sub>2</sub> ] <sub>2</sub>	$-7.21$	$-1.29$	$-5.91$	2.96	$-4.25$	
$[Bis(min)C_5-(C_4)_2][NTf_2]_2$	$-7.28$	$-0.94$	$-6.33$	3.17	$-4.11$	



Fig. 4 Calculated spatial distribution of HOMO and LUMO for DILs at M06-2X/ccpVDZ level.

HOMO of  $[NTF_2]$  anion to the LUMO of the imidazolium rings in all DILs.

As seen in Fig. 3, there are several intermolecular interactions between cations and anions. The nature of these interactions in DILs may be analyzed by Bader's topological QTAIM analysis using topological parameters such as electron density  $(\rho(r))$ , Laplacian of electron density  $(\nabla^2 \rho(r))$ , kinetic energy density  $(G(r))$ , potential energy density  $(V(r))$ , the total energy density  $(H(r))$  (where  $H(r) = G(r) + V(r)$ ) and  $-(G(r)/V(r))$  ratio at bond critical point (BCP) between the cation and anions of DILs. The topological graphics of the lowest energy structures of studied GDILs have been shown in Fig. S2.† Also, the values of mentioned parameters at M06-2X/ccpvDZ level have been summarized in Table S1.† The hydrogen bonding is characterized from  $\rho(r)$  values in the range of 0.002–0.035 a.u and with

corresponding  $\nabla^2 \rho(r)$  in the range of 0.024–0.139 a.u. For all of the studied DILs, values of  $\rho(r)$  and  $\nabla^2 \rho(r)$  are in the mentioned ranges.  $\rho(r)$  is used to describe the strength of a bond, the stronger the bond, the larger the value of  $\rho(r)$ . As it can be seen, interaction of oxygens of anions with HR (or HR′ ) atoms of the rings is stronger than interaction with the rest of hydrogens and there are at least two interactions between the HR (or HR′ ) atoms and the anions in DILs. The hydrogen bonds in ILs are classified as medium and weak ( $\nabla^2 \rho(r) > 0$ , and  $H(r) > 0$ ) and strong ( $\nabla^2 \rho(r) > 0$ , and  $H(r) < 0$ ) in nature. Therefore, according to the values in the table, hydrogen interactions in studied DILs are often of the type of weak and medium interactions. In asymmetric  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> DIL, total electron density  $(\sum \rho(r))$  at all critical points of the bond is higher than the other two symmetric DILs. Also, increasing the side chain length in

 $[\text{Bis}(mim)C_5-(C_4)_2][\text{NTf}_2]_2$  DIL reduces the electron density in the system which is in agreement with the results of interaction energies. The  $-(G(r)/V(r))$  ratio serves as a useful descriptor for separation of various interactions in ILs.<sup>77</sup> As it can be seen, the  $-(G(r)/V(r))$  ratio turns out to be between 0.99 a.u and 1.27 a.u, which suggests that the studied DILs are stabilized by van der Waals (vdW) and weak hydrogen bonding interactions. It has been shown that hydrogen bond interactions may not reveal its signature with the emergence of BCP. Extension of the QTAIM theory introduces the  $G(r)$  component corresponding to  $S_{\text{min}}$ (equivalent to BCP in QTAIM theory) that is directly proportional to the electron density mobility in BCP. Thus,  $G(r)$ component offers certain advantages over conventional method used to probe the electron density properties and is useful to characterize the hydrogen bondings. Saleh et al.<sup>78</sup> demonstrated that there is a correlation between OH stretching frequencies and the  $G(r)$  in substituted amino alcohols that have intramolecular hydrogen bonds. Lane et al.<sup>79</sup> showed that the strength of intermolecular hydrogen bonds which reflect in frequency down (red) shift of the OH-stretching vibration can be estimate with  $G(r)$  component. A plot of  $G(r)$  as a function of Paper<br>
Distantion-C<sub>(s)</sub>hyrri<sub>d</sub>, DL reduces the electron density in interaction energies has been abread to the space which is licensed under the energy of the space of



Fig. 5 (a) Kinetic energy density, G(r), as a function of interaction energies and (b)  $E_{HB}$  as a function of  $G(r)$ .

interaction energies has been shown in Fig. 5(a) turns out to be almost linear. Espinosa et  $al.^{80}$  showed that hydrogen bond energy  $(E_{HB})$  can be estimated from  $V(r)$  at the corresponding BCP, through  $E_{\text{HB}} = -0.5$  V(r) which a larger energy indicates a stronger bond. A plot of  $E_{HB}$  as a function of  $G(r)$  displayed in Fig. 5(b) turns out to be linear. As it can be seen, the asymmetric  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> possesses stronger cation–NTf<sub>2</sub> binding.

Charge transfer is the process in which a fraction of electron charge is transferred within a molecule or between different molecules. Generally, the charge is transferred from the HOMO to LUMO orbitals. In this section, NBO and CHELPG methods have been used to study charge transfer in the studied DILs. The results of both methods have been reported in Table S2.† According to this table, in both methods, the charge transfer for DILs has the following trend:  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> > [Bis(mim)  $C_5$ [NTf<sub>2</sub>]<sub>2</sub> > [Bis(mim) $C_5$ - $(C_4)$ <sub>2</sub>][NTf<sub>2</sub>] which is in agreement with decreasing the HOMO–LUMO energy gap from [Bis(mim)  $C_5$ - $(C_4)_2$ [NTf<sub>2</sub>]<sub>2</sub> to [N<sub>111</sub>- $C_5$ -mim][NTf<sub>2</sub>]<sub>2</sub> reported in Table 2.

In DILs, multiple hydrogen bonds are able to form. For example, an anion can interact simultaneously with two hydrogens which both of them can belong to alkyl chains or one hydrogen belongs to the alkyl chains and one hydrogen belongs to the ring. The NBO analysis (according to second order perturbation theory) has been extensively employed in studying hydrogen bonding.<sup>77</sup> The NBO analysis of DILs considering second-order perturbation energy  $(E(2))$ , the energy difference between donor and acceptor  $(\Delta E_{ij})$ , and Fock matrix element between donor and acceptor  $(F_{ii})$  has been reported in Table S3.†  $E(2)$  describes the strength of interactions between antibonding orbitals (BD\*) of C–H bonds of cations and lone pair electrons (LP) of N, F and O atoms of anion. Strong hydrogen bonding in the studied DILs, are characterized by large  $E(2)$ values which are produced by low  $\Delta E_{ij}$  and large  $F_{ij}$ . The  $E(2)$ values in Table S3† show that there is a high stabilization energy corresponding to charge transfers between oxygen atoms of anions and HR atoms of rings. The trend for these strong hydrogen bonds follows the order of  $[N_{111} - C_5 - m_1m][NT_2]_2$  $[\text{Bis}(mim)C_5][\text{NTf}_2]_2 > [\text{Bis}(mim)C_5-(C_4)_2][\text{NTf}_2]_2$ , which is in agreement with the trend of AIM and interaction energies. In the studied DILs, alkyl chains bend and more atoms of alkyl chains are able to form hydrogen bonds (see Fig. 3). Besides the strong hydrogen bonds forming between HR atoms of the rings and anions, some of the weaker hydrogen bonds with lower  $E(2)$ , which form between the hydrogens of alkyl chains and anions, are shown in Table S3.†

3.1.2 Isolated ion– $CO<sub>2</sub>$  and ion triplet (DIL molecule)– $CO<sub>2</sub>$ interactions. To understand the relation between the change of structure and the performance of  $CO<sub>2</sub>$  absorption, the stable absorption sites of ions and DILs must be clearly identified. In order to find the most stable structures of isolated ion– $CO<sub>2</sub>$  and  $DIL-CO<sub>2</sub>$  complexes, the  $CO<sub>2</sub>$  molecule was placed in several different positions around ions and DILs molecules. Then, the mentioned complexes were optimized at the same theoretical level with the pure state (M06-2X/cc-pVDZ). Fig. S3 and S4† show some more stable structures of the ion– $CO<sub>2</sub>$  and DIL– $CO<sub>2</sub>$ complexes, respectively. The interaction energy  $(\Delta E_{\text{int}})$  is defined as difference between total energy of the relaxed



Fig. 6 Most stable structures of isolated ion–CO<sub>2</sub> complexes at M06-2X/ccpVDZ level. Broken lines show the most important intermolecular interactions (distances are in Å)

complex  $(E_{\text{complex}})$  and sum of the energies of the isolated ion or DIL molecule and  $CO_2$  molecule ( $E_{\text{ion or DIL}}$  and  $E_{\text{CO}_2}$ , respectively), were calculated using the following equation:

$$
\Delta E_{\text{int}} = E_{\text{complexes}} - (E_{\text{ion or DIL}} + E_{\text{CO}_2})
$$
(4)

The values of interaction energies, BSSE and D3 corrections for the structures shown in Fig. S3 and S4† have been reported in Table 5. Fig. 6 and 7 show the most stable structures for ion–  $CO<sub>2</sub>$  and DIL– $CO<sub>2</sub>$  complexes, respectively, which were selected according to the values of interaction energies for further analyses. Due to the presence of two different heads in the asymmetric cation, for  $[N_{111}$ -C<sub>5</sub>-mim]<sup>2+</sup>-CO<sub>2</sub> and  $[N_{111}$ -C<sub>5</sub>-mim] [NTf<sub>2</sub>]<sub>2</sub>-CO<sub>2</sub> complexes, two stable structures were considered to investigate the differences between two heads in the performance of  $CO<sub>2</sub>$  absorption.

As it can be seen in Fig. 6, based on intermolecular cation–  $CO<sub>2</sub>$  distances, the most stable sites for  $CO<sub>2</sub>$  absorption in the symmetric cations are mainly the co-absorbing position by HR hydrogen in imidazolium ring and methylene group attached to

the ring. This result is consistent with previous studies on the interaction of imidazolium-based MILs with  $CO<sub>2</sub>$  molecules.<sup>81</sup> However, in the asymmetric  $[N_{111}$ -C<sub>5</sub>-mim<sup>2+</sup> cation, the most stable structure is related to the placement of the  $CO<sub>2</sub>$  molecule in front of the ammonium head (structure no. 3 in Fig. S3†) probably due to the more localized charge in the ammonium group. As it can be seen, for  $[NTf_2]$  anion, the most stable sites for CO<sub>2</sub> absorption are nitrogen and oxygen atoms of anion. Comparing the interactions of anions and cations with the  $CO<sub>2</sub>$ molecule shows that, with the exception of the structure in which the  $CO<sub>2</sub>$  molecule interacts with the asymmetric cation from the ammonium head, the interaction energy of the anion–  $CO<sub>2</sub>$  is higher than that for the cation– $CO<sub>2</sub>$ . It can be said that the main intermolecular interactions of anion– $CO<sub>2</sub>$  are due to the strong electrostatic interactions, while for cation– $CO<sub>2</sub>$  are the weak hydrogen bonds.

According to the values of energies reported in Table S4,† in symmetric  $DIL-CO<sub>2</sub>$  complexes, the most stable structure is obtained from the orientation of the  $CO<sub>2</sub>$  molecule on top of the



Fig. 7 Most stable structures of DIL–CO<sub>2</sub> complexes at M06-2X/ccpVDZ level. Broken lines show the most important intermolecular interactions (distances are in Å).

HR hydrogen of ring, and the interaction energy decreases as the  $CO<sub>2</sub>$  molecule moves away from the ring. In the  $[N<sub>111</sub>-C<sub>5</sub>]$  $\text{min}[\text{NTf}_2]_2$ –CO<sub>2</sub> complexes, the most stable structure is related to the placement of the  $CO<sub>2</sub>$  molecule in the vicinity of the ammonium head. In the  $[N_{111} - C_5 - \text{min}]^{2+} - \text{CO}_2$  complexes,

we saw that the most stable structure is formed by the orientation of one of the oxygen atoms of  $CO<sub>2</sub>$  in front of the ammonium head. However, in  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub>-CO<sub>2</sub> complex due to the presence of anion,  $CO<sub>2</sub>$  molecule is oriented in the ammonium head in such a way that it has the most

Table 3 Charge transfer calculated in the studied ion–CO<sub>2</sub> and DIL–CO<sub>2</sub> complexes using CHELPG and NBO methods.  $\Delta \phi = 180 - \phi$  angle deviation ( $\phi$  is O–C–O angle in CO<sub>2</sub> molecule)



interaction with the anion and cation at the same time and presents a more accurate picture of reality.

The trend of interaction energies for  $DIL-CO<sub>2</sub>$  complexes is as follows:  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub>-CO<sub>2</sub> < [Bis(mim)C<sub>5</sub>][NTf<sub>2</sub>]<sub>2</sub>- $CO_2$  < [Bis(mim) $C_5$ - $(C_4)_2$ ][NTf<sub>2</sub>]<sub>2</sub>- $CO_2$  which is the reverse of the trend of interaction energies for pure DILs. It seems that with decreasing interaction energy between anions and cations, the  $CO<sub>2</sub>$  capture in DILs increases. This result is in agreement with the studies of Babarao et al.<sup>82</sup> and Gupta et al.<sup>83</sup> on MILs–CO<sub>2</sub> systems. They showed that  $CO<sub>2</sub>$  solubility is governed by the binding energy of cation–anion rather than anion– $CO<sub>2</sub>$ .

In Table 3, the values of charge transfer in ion- $CO<sub>2</sub>$  and DIL- $CO<sub>2</sub>$  complexes with NBO and CHELPG methods have been reported. As it can be seen, the CHELPG method predicts higher values of charge transfer, and as expected in cation– $CO<sub>2</sub>$ complexes the charge transfer is from the  $CO<sub>2</sub>$  molecule to cations and in the anion– $CO<sub>2</sub>$  complex the charge transfer is from anion to  $CO<sub>2</sub>$  molecule. Charge transfer calculated by both methods indicates higher charge transfer in the anion– $CO<sub>2</sub>$ complex compared to cation– $CO<sub>2</sub>$ , which shows stronger electrostatic interactions in this complex. Also, in both methods, the asymmetric cation– $CO<sub>2</sub>$  complex shows higher charge transfer compared to the two symmetric cations. Another interesting feature presented in Table 3 is the bending of O–C–O angle in the  $CO<sub>2</sub>$  molecule for anion– $CO<sub>2</sub>$  system, which is significant compared to cation– $CO<sub>2</sub>$  complexes.

In DIL-CO<sub>2</sub> complexes, the negative charge of  $CO<sub>2</sub>$  molecule shows that more charge from the anion has been transferred to  $CO<sub>2</sub>$  molecule. According to the results presented in the table, charge transfer by both CHELPG and NBO methods for DIL–  $CO<sub>2</sub>$  complexes obeys the following order:  $[N<sub>111</sub>-C<sub>5</sub>-min]$  $[NTf_2]_2$ –CO<sub>2</sub> <  $[Bis(min)C_5][NTf_2]_2$ –CO<sub>2</sub> <  $[Bis(min)C_5-(C_4)_2]$  $[\text{NTf}_2]_2$ –CO<sub>2</sub>

This trend is in agreement with interaction energies reported in Table S4.† In some studies, the bending of O–C–O angle is considered as a criterion of the tendency to adsorb  $CO<sub>2</sub>$  in the ILs.<sup>27,32</sup> By comparing the values of the angle deviation for  $CO<sub>2</sub>$ molecules in different  $DIL-CO<sub>2</sub>$  complexes, it seems that the interactions between  $CO<sub>2</sub>$  molecule and the studied DILs have an enhanced tendency as side chains extend.

The values of the HOMO, LUMO energies,  $\Delta E_{\text{gap}}$ ,  $\eta$  and  $\mu$ calculated for systems containing  $CO<sub>2</sub>$  have been listed in Table 4. Comparing the values of  $\Delta E_{\text{gap}}$  in Table 2 and its values in this table (after interaction with  $CO<sub>2</sub>$ ), it is observed that the  $\Delta E_{\text{gap}}$  of anion–CO<sub>2</sub> complexes decreases from 7.44 eV to 6.57 eV, but no significant change is observed in the cation– $CO<sub>2</sub>$ complexes. It can be said that the magnitude of the decrease in the  $\Delta E_{\text{gap}}$  is due to the increase in the strength of the interactions between the anion and  $CO<sub>2</sub>$ , and stronger interactions between the orbitals lead to a more stable anion– $CO<sub>2</sub>$  complex. For DIL–CO<sub>2</sub> complexes, the  $[Bis(min)C_5-(C_4)_2][NTf_2]_2-CO_2$ system has the smallest energy gap compared to other systems which is in agreement with interaction energy and charge transfer results. Fig. S5† shows the spatial distribution of HOMO and LUMO orbitals for studied complexes. As it can be seen, similar to pure DILs, the HOMO orbitals are mainly located on two anions, while the LUMO orbitals for all systems are mainly located on the imidazolium rings. RSC Advances<br>
Interaction with the anion and casion at the same time and<br>
Interaction systems complexes are free for the proposed are for the systems are commonly CG, pre-<br>
interaction and the systems CG, commonly ( $\alpha$ ),

In Table 5, the NBO analysis for studied complexes containing  $CO<sub>2</sub>$  molecules has been reported. In the case of cation–  $CO<sub>2</sub>$  complexes, interactions are characterized with charge transfer from one of the lone pairs of oxygen atoms in  $CO<sub>2</sub>$ molecule to anti-bonding orbitals in the C–H bonds of cations and in the case of anion– $CO<sub>2</sub>$  complexes, interactions are characterized with charge transfer from nitrogen and oxygen atoms to the anti-bonding orbital in  $CO<sub>2</sub>$  molecules. As mentioned before, the larger value of  $E(2)$  indicates the stronger interaction, which may result from a lower  $\Delta E_{ii}$  or a larger  $F_{ii}$  (*i.e.* better symmetry between the electron donor and electron acceptor). In agreement with previous results, the largest value of  $E(2)$  in cations arises from the charge transfer of  $CO<sub>2</sub>$  molecule to anti-bonding orbitals in C–HR bond of the rings. The trend of  $E(2)$  for this charge transfer in cation–CO<sub>2</sub> complexes is as follows:  $>$   $[ Bis(min]C_5]^{2+}$  -CO<sub>2</sub>  $[N_{111}$  -C<sub>5</sub>-mim]<sup>2+</sup> -CO<sub>2</sub>  $>$   $[ Bis( \text{min}$  $\text{C}_5$ - $\text{C}_4$ <sub>2</sub> $\text{C}_2$ <sup>2+</sup>-CO<sub>2</sub>

The larger  $E(2)$  value in the  $[N_{111}$ -C<sub>5</sub>-mim<sup> $]$ 2+</sup>-CO<sub>2</sub> arises from larger  $F_{ij}$  value, which shows strong donor-acceptor symmetry in this complex. Also, the stronger interaction energy and more charge transfer in anion– $CO<sub>2</sub>$  complexes with respect to cation–  $CO<sub>2</sub>$  complexes (see Tables S4 $\dagger$  and 3, respectively) can be explained based on larger  $E(2)$  values which in turn arises from

Table 4 Energies of LUMO and HOMO, HOMO–LUMO energy gap ( $E_{\text{gap}}$ ), chemical hardness ( $\eta$ ) and chemical potential ( $\mu$ ) of the ion–CO<sub>2</sub> and DIL–CO2 complexes calculated at M06-2X/ccpVDZ level



### Table 5 NBO analysis of studied ion– $CO<sub>2</sub>$  and DIL– $CO<sub>2</sub>$  complexes calculated at M06-2x/cc-pVDZ level



low  $\Delta E_{ij}$  coupled with suitable donor-acceptor symmetry and lead to more effective interactions for anions than cations.

For DIL–CO<sub>2</sub> complexes, the largest  $E(2)$  in all systems arises from the charge transfer of the anions in DILs to  $CO<sub>2</sub>$  molecules. Compared to cations, anions have a stronger interaction with  $CO<sub>2</sub>$  molecules which is in agreement with previous results. Furthermore, according to  $E(2)$  values, the strongest interactions in the cations of all DILs (including imidazolium head for asymmetric DIL) are due to the charge transfer from  $CO<sub>2</sub>$ molecule to anti-bonding orbitals in C–HR bond of the rings. The trend of the  $E(2)$  (charge transfer from  $CO<sub>2</sub>$  molecule to C– HR bond) for cations of DILs in the presence of anions is as follows:

 $[N_{111}$ -C<sub>5</sub>-mim]<sup>2+</sup>-CO<sub>2</sub> < [Bis(mim)C<sub>5</sub><sup>2+</sup>-CO<sub>2</sub> < [Bis(mim)C<sub>5</sub>- $(\mathrm{C}_4)_2]^{2^+}\text{--CO}_2.$  The different trend of  $\mathit{E}(2)$  for isolated cations and for cations in the presence of anions can be attributed to the effect of the interaction energy between cations and anions on  $CO<sub>2</sub>$  capture. The weaker interaction energy between the anions and cations, the stronger interaction of  $CO<sub>2</sub>$  molecule with cation. It should be noted that since the interaction energy of cation–anion in MILs is lower than that of DILs, it is expected that MILs have higher gas solubility. While, experimental studies show the opposite.<sup>19,20,22</sup> In fact, in addition to the interaction energy between ions, other factors such as the

number of interaction sites, free volume fraction, types of ions, *etc.* are effective in the amount of  $CO<sub>2</sub>$  absorption by ILs and should be taken into consideration.<sup>84,85</sup>

RDG isosurfaces and 2D scatter plots for studied complexes have been displayed in Fig. 8 and S6.† Non-covalent interactions (NCIs) such as hydrogen bonding, vdW interactions and steric effects can be evaluated using RDG analysis. RDG plots can be provided information about the type and strength of interactions in IL-CO<sub>2</sub> systems.<sup>86</sup> The sign of  $\lambda_2$  (second eigenvalues of the electron-density of Hessian matrix) and electron density  $\rho(r)$ value determine the type and strength of the interactions, respectively, where negative and positive sign  $(\lambda_2)$  values describe the attractive and repulsive interactions, respectively. In 2D scatter RDG plots, spikes in large negative values of sign  $(\lambda_2)^*\rho$  ( $\rho > 0$ ,  $\lambda_2 < 0$ ) are indicative of hydrogen bond interactions, whereas the spikes located at large positive values sign  $(\lambda_2)^*\rho$  ( $\rho$  $> 0$ ,  $\lambda_2 > 0$ ) are the steric effects. In addition, spikes situated at sign  $(\lambda_2)^*\rho$  values approach zero ( $\rho \sim 0$ ,  $\lambda_2 \sim 0$ ) indicate very weak vdW interactions. Also, the colored RDG isosurfaces can be used for qualitative understanding of the nature and the strength of NCIs. The blue, red, and green regions are used for hydrogen bondings, strong repulsion interactions, and attractive vdW interactions, respectively.

In 2D scatter plots for symmetric cation– $CO<sub>2</sub>$  complexes (Fig. 8(a) and (c)), two spikes at about  $-0.01$  a.u. and  $-0.12$  a.u. related to the interaction of oxygen of  $CO<sub>2</sub>$  molecule with  $H<sub>1</sub>$  and HR atoms of cations, respectively. For the asymmetric cation–  $CO<sub>2</sub>$  complex (Fig. 8(b)), a strong single spike at  $-0.01$  a.u. indicates the interaction of oxygen of  $CO<sub>2</sub>$  molecule with methyl



Fig. 8 Color-filled RDG isosurfaces and scatter graphs of RDG for (a) [Bis(mim)C<sub>5</sub>]<sup>2+</sup>–CO<sub>2</sub>, (b) [N<sub>111</sub>-C<sub>5</sub>-mim]<sup>2+</sup>–CO<sub>2</sub>, (c) [Bis(mim)C<sub>5</sub>-(C<sub>4</sub>)<sub>2</sub>]<sup>2+</sup>–CO<sub>2</sub> and (d)  $[NTf<sub>2</sub>]<sup>-</sup>$  CO<sub>2</sub>. Green and red colors denote vdW interactions and steric contributions, respectively.

groups at the ammonium head of cation. Furthermore, it is observed that in areas close to zero, the number of spikes and their intensity for anion– $CO<sub>2</sub>$  system increase compared to the cation– $CO<sub>2</sub>$  systems, which indicates an increase in vdW interactions in the anion– $CO<sub>2</sub>$  system. In the RDG isosurfaces, the extension of green region shows the local character of an interaction. A larger size of green region for asymmetric cation shows that  $CO<sub>2</sub>$  has stronger interactions with the ammonium head of the asymmetric cation than the imidazolium head which is in agreement with the previous analyses. In RDG isosurfaces of  $DIL-CO<sub>2</sub>$  complexes (Fig. S6†), the green areas are formed more between anion and  $CO<sub>2</sub>$ , which indicates the greater contribution of anions in interaction with  $CO<sub>2</sub>$  molecules. Also, for asymmetric DIL-CO<sub>22</sub> system (Fig. S6(b) and (c)<sup>†</sup>), the size of green region between  $CO<sub>2</sub>$  and DIL is less than the other two symmetric DILs–CO<sub>2</sub> systems (Fig. S6(a) and (d)†). Comparing the scatter plots, it can be seen that in the [Bis(mim)  $C_5$ - $(C_4)_2$ [NTf<sub>2</sub>]<sub>2</sub>– $CO_2$  system, a strong spike in −0.03 a.u. indicates the existence of stronger hydrogen bond interactions and the more tendency of this DIL for interaction with CO2 molecule. In general, the green region between DILs and  $CO<sub>2</sub>$ 

indicates that vdW interactions are driving force to  $CO<sub>2</sub>$  capture in DILs.

The AIM analysis of isolated ion– $CO<sub>2</sub>$  and DIL– $CO<sub>2</sub>$  interactions has been reported in Table 6 and Fig. S7 and S8,† respectively. As it can be seen, according to the values of  $\rho(r)$ , in all systems, the interactions of oxygens of  $CO<sub>2</sub>$  molecules with HR atoms of rings are stronger than other hydrogens. In DIL–  $CO<sub>2</sub>$  complexes, the number of interactions between the  $CO<sub>2</sub>$ molecules and the anions is more than that for cations, and  $CO<sub>2</sub>$ molecules interact with the anions through the oxygen and fluorine atoms. In accordance with the results of previous analyses, the trend of  $\sum \rho(r)$  at BCPs between DILs and CO<sub>2</sub> molecules is as follows:  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub>-CO<sub>2</sub> < [Bis(mim)  $C_5$ ][NTf<sub>2</sub>]<sub>2</sub>– $CO_2$  < [Bis(mim) $C_5$ - $(C_4)$ <sub>2</sub>][NTf<sub>2</sub>]<sub>2</sub>– $CO_2$ 

Which indicates that DILs with symmetric cation and a longer side alkyl chain are more desirable for interaction with CO<sub>2</sub> molecules. Whereas in isolated ion–CO<sub>2</sub> systems,  $\sum \rho(r)$  for  $[N_{111}$ -C<sub>5</sub>-mim<sup>2+</sup>-CO<sub>2</sub> is more and the replacement of the ammonium group has a greater effect on the interaction with  $CO<sub>2</sub>$  molecule compared to increasing the length of side alkyl chain. Therefore, the results indicate the important effect of



anion in the interaction of  $CO<sub>2</sub>$  with cations of DILs, which cannot be ignored.

### 3.2 COSMO-RS calculations

3.2.1 COSMO-RS analysis of  $CO<sub>2</sub>$  solubility in the studied DILs. In this section, from molecular thermodynamics point of view, the  $CO<sub>2</sub>$  solubility data in the studied DILs were obtained by COSMO-RS theory employing DFT calculations at the level of BP/TZVP. Fig. 9(a) shows the solubility of  $CO<sub>2</sub>$  in the studied DILs at a pressure of 1 bar and at different temperatures. As expected, the solubility of  $CO<sub>2</sub>$  decreases with increasing temperature. Also, at each temperature, the solubility of  $CO<sub>2</sub>$  in the DILs has the following trend:  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> < [Bis(- $\text{min}$ )C<sub>5</sub>][NTf<sub>2</sub>]<sub>2</sub>< [Bis(mim)C<sub>5</sub>-(C<sub>4</sub>)<sub>2</sub>][NTf<sub>2</sub>]<sub>2</sub>. As it can be seen, symmetric DILs tend to accommodate a little more  $CO<sub>2</sub>$  molecules and the solubility of  $CO<sub>2</sub>$  slightly increases as side alkyl chains of DILs extend longer, which accords with the quantum results in previous sections. Payagala et al.<sup>87</sup> reported the trend of the densities of these DILs as  $[Bis(min)C_5][NTf_2]_2$ (1.57 g cm<sup>-3</sup>) > [N<sub>111</sub>-C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> (1.54 g cm<sup>-3</sup>) > [Bis(mim)  $C_5$ - $(C_4)_2$ [[NTf<sub>2</sub>]<sub>2</sub> (1.44 g cm<sup>-3</sup>) at 303.15 K. Also, they reported that the viscosity of these systems decreases as  $[N_{111}$ -C<sub>5</sub>-mim]  $[NTf_2]_2$  (292.60 mPa s) >  $[Bis(min)C_5][NTf_2]_2$  (186.83 mPa s) >  $[\text{Bis}(mim)C_5-(C_4)_2][\text{NTf}_2]_2$  (145.44 mPa s) at 323.15 K which is in agreement with the trend of interaction energies of pure DILs. In comparison among these three DILs, the density and viscosity of  $[\text{Bis}(mim)C_5-(C_4)_2][\text{NTf}_2]_2$  is the lowest and it has the largest free volume. Indeed, a longer side alkyl chain increases the free volume in DILs, providing more available space for  $CO<sub>2</sub>$ absorption. As it can be seen, by replacing an imidazolium head of the symmetric  $[\mathrm{Bis}(\mathrm{mim})\mathrm{C}_5]^{2^+}$  cation with a trimethylammonium group, the electrostatic attraction between anion and cation increases due to the localized charged density on the nitrogen of quaternary amine and the viscosity of  $[N_{111} - C_5 - m_1m]$  $[NTf<sub>2</sub>]$ <sub>2</sub> increases. Also, by increasing the length of side alkyl chain in cations, the strength of vdW interactions increase, while the dispersion of the charge centers weakens the electrostatic attractions between anion and cation. It seems that the weakening of electrostatic interactions is dominant and the viscosity of  $[\text{Bis}(\text{min})C_5-(C_4)_2][\text{NTf}_2]_2$  decreases. Similarly, Su et al.<sup>88</sup> showed that  $CO<sub>2</sub>$  solubility increases as alkyl chains of quaternary ammonium mono cationic ILs (QAMILs) extend at the same temperature and pressure. In industry,  $CO<sub>2</sub>$  absorption at high pressure is common; therefore, the solubility of  $CO<sub>2</sub>$ at pressures up to 10 bar was also calculated. According to Fig. 9(b), it can be seen that high pressures significantly improve  $CO<sub>2</sub>$  absorption by DILs. The solubility of  $CO<sub>2</sub>$  in the studied DILs at 10 bar is approximately 7 times higher than that under ambient pressure, which shows the potential industrial application of DILs in pressure swing adsorption (PSA) of  $CO<sub>2</sub>$ .

3.2.2 Selectivity of  $CO_2$  from  $CO$ ,  $H_2$  and  $CH_4$  in the studied DILs. The solubility of a gas in a liquid at infinite dilution can be measured by Henry's law constant, which is calculated using the following equation:

$$
K_{\rm H} = \frac{RT\rho}{M} \exp\left(\frac{\Delta G_{\rm sol}}{RT}\right) \tag{5}
$$



Fig. 9 (a)  $CO<sub>2</sub>$  solubility in the studied DILs at 1 bar and at different temperatures. (b)  $CO<sub>2</sub>$  solubility in DILs at 303.15 K and at different pressures calculated by COSMO-RS theory.

where  $\Delta G_{\rm sol}$  is the free energy of solvation,  $\rho$  is the density and M is the molecular mass of pure DIL. Henry's constant is directly proportional to gas solubility. In Fig. 10(a), the changes of the Henry's constant of  $CO<sub>2</sub>$  gas with temperature in the studied DILs have been shown. As the temperature increases from 303.15 to 373.15 K, Henry's constant changes from about 60 bar to 200 bar, which shows that high temperature conditions are not suitable for the better physical absorption. The values of Henry's constant show that the most suitable solvent for  $CO<sub>2</sub>$  is the symmetric DIL [Bis(mim)C<sub>5</sub>-(C<sub>4</sub>)<sub>2</sub>][NTf<sub>2</sub>]<sub>2</sub> with a longer side chain length.

Separation of  $CO<sub>2</sub>$  from gas mixture is very common and necessary in the chemical industry. For example, in the production of natural gas, in addition to the main component  $CH<sub>4</sub>$ , in most cases  $CO<sub>2</sub>$  is also present in the mixed gas. The presence of  $CO<sub>2</sub>$  reduces the combustion efficiency of natural gas and can cause pipeline corrosion.<sup>89</sup> Therefore, separation of CO2 from some gases is very important. In this work, in order to investigate the effect of side chain length and symmetry in cations on gas separation, in Fig. 10(b), the selectivities of  $CO<sub>2</sub>/$ CO,  $CO_2/CH_4$  and  $CO_2/H_2$  mixtures were calculated using



Fig. 10 (a) Henry's constant versus temperature for the studied DILs. (b) Selectivity of  $CO<sub>2</sub>$  from CO, H<sub>2</sub> and CH<sub>4</sub> in the studied DILs at 303.15 K calculated by COSMO-RS theory.

COSMO-RS theory. The ideal gas selectivity,  $S_{m/n}$ , can be obtained by eqn  $(6)$ :<sup>90</sup>

$$
S_{m/n} = \frac{H_m}{H_n} \tag{6}
$$

where  $S_{m/n}$  is the selectivity of gas *n* to gas *m*, and  $H_m$  and  $H_n$  are constants of Henry's law of gas  $m$  and  $n$ , respectively. As it can be seen, in all DILs, the trend of selectivity is as follows:  $CO<sub>2</sub>/H<sub>2</sub>$  $>$  CO<sub>2</sub>/CO  $>$  CO<sub>2</sub>/CH<sub>4</sub>. Apparently, the selectivity of CO<sub>2</sub> from the other three gases decreases slightly with increasing the length of side alkyl chain. Therefore, if we want to achieve optimal separation of  $CO<sub>2</sub>$  from CO,  $H<sub>2</sub>$  and CH<sub>4</sub> at the same time, it is better to consider short side chains for DILs. For the asymmetric  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub>, the selectivity of CO<sub>2</sub>/CH<sub>4</sub> and CO2/CO mixtures slightly increases compared to the symmetric [Bis(mim)C<sub>5</sub>][NTf<sub>2</sub>]<sub>2</sub> and for selectivity of CO<sub>2</sub>/H<sub>2</sub> mixture decreases slightly. Su et al.<sup>88</sup> for ammonium-based MILs, reported a decrease in selectivity of  $CO<sub>2</sub>$  from CO and CH<sub>4</sub> and an increase in the selectivity of  $CO<sub>2</sub>$  from  $H<sub>2</sub>$  by increasing the length of side chains.

# 4. Conclusion

In summary, we have investigated three selected DILs (considering the effects of alkyl side chain length and symmetry in cations) in the absence and presence of  $CO<sub>2</sub>$  using quantum mechanical methods. The results showed that for all three structures, the linkage chain entangles in order to maximize interactions of two heads of DILs with two anions. The trend of interaction energy values was as follows:  $[N_{111} - C_5 - mim][NTf_2]_2$  $[Bis(min)C_5][NTf_2]_2 > [Bis(min)C_5-(C_4)_2][NTf_2]_2$ . It seems that in asymmetric DIL  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub>, by replacing a imidazolium group with the ammonium one, the interaction energy between the cations and anions increases due to the more localized charge in the ammonium group. Furthermore, AIM and NBO analyses for all three DILs, in addition to confirming the trend of interaction energies, showed that interaction of oxygens of anions with HR (or HR′ ) atoms of the rings is stronger than interaction with the rest of hydrogens.

In ion– $CO<sub>2</sub>$  complexes, it was observed that the most stable sites for  $CO<sub>2</sub>$  absorption in the symmetric cations are mainly the co-absorbing positions by HR hydrogen in the imidazolium ring and methylene group attached to the ring. However, in the asymmetric  $[N_{111}$ -C<sub>5</sub>-mim]<sup>2+</sup> cation, the most stable structure is related to the placement of  $CO<sub>2</sub>$  molecule in front of the ammonium head probably due to the more localized charge in the ammonium group. The results of charge transfer values, O-C-O angle deviations in the  $CO<sub>2</sub>$  molecule, AIM and NBO showed that, with the exception of the structure in which the  $CO<sub>2</sub>$  molecule interacts with the asymmetric cation from the ammonium head, the tendency of the anions to  $CO<sub>2</sub>$  is more than cations. It can be said that the intermolecular interactions of the anion– $CO<sub>2</sub>$  are due to the strong electrostatic interactions, while for cation– $CO<sub>2</sub>$  are the weak hydrogen bonds.

For  $DIL$ – $CO<sub>2</sub>$  complexes, the trend of interaction energies was as follows:  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub>-CO<sub>2</sub> < [Bis(mim)C<sub>5</sub>][NTf<sub>2</sub>]<sub>2</sub>-CO<sub>2</sub>  $\langle$  [Bis(mim)C<sub>5</sub>-(C<sub>4</sub>)<sub>2</sub>][NTf<sub>2</sub>]<sub>2</sub>-CO<sub>2</sub> which is the reverse of the trend of interaction energies for pure DILs. It seems that with decreasing the interaction energy between ions, the  $CO<sub>2</sub>$  capture in DILs increases. The results of NBO and AIM analyses indicated that DIL with symmetric cation and longer side chain is more favorable for  $CO<sub>2</sub>$  absorption which was in agreement with the results of interaction energies and charge transfers. Moreover, the green regions between DILs and  $CO<sub>2</sub>$  showed that the vdW interactions are driving force to  $CO<sub>2</sub>$  capture in DILs.

COSMO-RS results indicated that the solubility of  $CO<sub>2</sub>$  in the DILs has the following trend:  $[N_{111}$ -C<sub>5</sub>-mim][NTf<sub>2</sub>]<sub>2</sub> < [Bis(mim)C<sub>5</sub>]  $[NTF_2]_2$  <  $[Bis(min)C_5-(C_4)_2][NTf_2]_2$ . Indeed, both symmetric DILs tend to accommodate a little more  $CO<sub>2</sub>$  molecules and the solubility of  $CO<sub>2</sub>$  slightly increases as side alkyl chains of a DIL extends longer, which is in agreement with quantum results. Also, results showed that the selectivity of  $CO<sub>2</sub>$  from  $H<sub>2</sub>$ , CO and CH<sub>4</sub> gases decreases slightly with increasing the length of side alkyl chains.

### Author contributions

Mehrangiz Torkzadeh: software, investigation, methodology, visualization, data curation, writing - original draft. Majid Moosavi: supervision, conceptualization, methodology, validation, writing – review & editing.

# Conflicts of interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper. **PSC Articles:**<br>
Notest Articles: The mass are computed on the computed on the common C Phase Common C Pass Article is licensed under the time mass Articles.<br> **CONNECT So Finder Case Article 2022.**<br>
2.1974-2315.<br>
2.1974-2

# Acknowledgements

This research was supported financially by Iran national science Foundation (INSF) and the Research Council of University of Isfahan [grant number 4002409].

## References

- 1 C. A. Trickett, A. Helal, B. A. Al-Maythalony, Z. H. Yamani, K. E. Cordova and O. M. Yaghi, Nat. Rev. Mater., 2017, 2, 1– 16.
- 2 I. Sreedhar, T. Nahar, A. Venugopal and B. Srinivas, Renewable Sustainable Energy Rev., 2017, 76, 1080–1107.
- 3 Z. Dai, R. D. Noble, D. L. Gin, X. Zhang and L. Deng, J. Membr. Sci., 2016, 497, 1–20.
- 4 B. Gurkan, F. Simeon and T. A. Hatton, ACS Sustainable Chem. Eng., 2015, 3, 1394–1405.
- 5 E. Masiren, N. Harun, W. Ibrahim and F. Adam, Int. J. Eng. Sci. Res. Technol., 2016, 3, 43–51.
- 6 F. Moosavi, F. Abdollahi and M. Razmkhah, Int. J. Greenhouse Gas Control, 2015, 37, 158–169.
- 7 M. J. Earle and K. R. Seddon, Pure Appl. Chem., 2000, 72, 1391–1398.
- 8 X. Zhang, X. Zhang, H. Dong, Z. Zhao, S. Zhang and Y. Huang, Energy Environ. Sci., 2012, 5, 6668–6681.
- 9 G. Cui, J. Wang and S. Zhang, Chem. Soc. Rev., 2016, 45, 4307–4339.
- 10 B. Bhargava and M. L. Klein, J. Phys. Chem. B, 2011, 115, 10439–10446.
- 11 H. Shirota, T. Mandai, H. Fukazawa and T. Kato, J. Chem. Eng. Data, 2011, 56, 2453–2459.
- 12 X. Han and D. W. Armstrong, Org. Lett., 2005, 7, 4205–4208.
- 13 M. Mahrova, F. Pagano, V. Pejakovic, A. Valea, M. Kalin, A. Igartua and E. Tojo, Tribol. Int., 2015, 82, 245–254.
- 14 M. Qi and D. W. Armstrong, Anal. Bioanal. Chem., 2007, 388, 889–899.
- 15 C. Cadena, J. L. Anthony, J. K. Shah, T. I. Morrow, J. F. Brennecke and E. J. Maginn, J. Am. Chem. Soc., 2004, 126, 5300–5308.
- 16 S. Tanaka, K. Kida, H. Fujimoto, T. Makino and Y. Miyake, Langmuir, 2011, 27, 7991–7995.
- 17 S. Supasitmongkol and P. Styring, Energy Environ. Sci., 2010, 3, 1961–1972.
- 18 Y. Zhang, S. Zhang, X. Lu, Q. Zhou, W. Fan and X. Zhang, Eur. J. Chem., 2009, 15, 3003–3011.
- 19 Y. Zhang, P. Yu and Y. Luo, Chem. Eng. J., 2013, 214, 355–363.
- 20 S. D. Hojniak, A. L. Khan, O. Holloczki, B. Kirchner, I. F. Vankelecom, W. Dehaen and K. Binnemans, J. Phys. Chem. B, 2013, 117, 15131–15140.
- 21 S. Li, W. Zhao, G. Feng and P. T. Cummings, Langmuir, 2015, 31, 2447–2454.
- 22 Y. Zhang, Z. Wu, S. Chen, P. Yu and Y. Luo, Ind. Eng. Chem. Res., 2013, 52, 6069–6075.
- 23 X. Liu, J. E. Bara and C. H. Turner, J. Phys. Chem. B, 2021, 125, 8165–8174.
- 24 X. Liu, K. E. O'Harra, J. E. Bara and C. H. Turner, Phys. Chem. Chem. Phys., 2020, 22, 20618–20633.
- 25 X. Liu and C. H. Turner, Chem. Phys. Lett., 2022, 786, 139204.
- 26 J. Ma, Y. Wang, M. Zhu, X. Yang and B. Wang, J. Mol. Liq., 2020, 320, 114408.
- 27 V. Sanz, R. Alcalde, M. Atilhan and S. Aparicio, J. Mol. Model., 2014, 20, 1–14.
- 28 G. E. Logotheti, J. Ramos and I. G. Economou, J. Phys. Chem. B, 2009, 113, 7211–7224.
- 29 H. Li and M. N. Kobrak, J. Chem. Phys., 2009, 131, 194507.
- 30 B. E. Gurkan, J. C. de la Fuente, E. M. Mindrup, L. E. Ficke, B. F. Goodrich, E. A. Price, W. F. Schneider and J. F. Brennecke, J. Am. Chem. Soc., 2010, 132, 2116–2117.
- 31 M. I. Cabaço, M. Besnard, Y. Danten and J. Coutinho, J. Phys. Chem. B, 2011, 115, 3538–3550.
- 32 S. Aparicio and M. Atilhan, J. Phys. Chem. B, 2012, 116, 9171– 9185.
- 33 J.-J. Chen, W.-W. Li, X.-L. Li and H.-Q. Yu, Phys. Chem. Chem. Phys., 2012, 14, 4589–4596.
- 34 C. Wu, T. P. Senftle and W. F. Schneider, Phys. Chem. Chem. Phys., 2012, 14, 13163–13170.
- 35 A. Indarto and J. Palgunadi, Ionics, 2012, 18, 143–150.
- 36 J. A. Steckel, J. Phys. Chem. A, 2012, 116, 11643–11650.
- 37 H. Tang and C. Wu, ChemSusChem, 2013, 6, 1050–1056.
- 38 H. Sun, D. Zhang, C. Liu and C. Zhang, J. Mol. Struct.: THEOCHEM, 2009, 900, 37–43.
- 39 E. Bodo and R. Caminiti, J. Phys. Chem. A, 2010, 114, 12506– 12512.
- 40 D. Farmanzadeh, A. Soltanabadi and S. Yeganegi, J. Chin. Chem. Soc., 2013, 60, 551–558.
- 41 P. L. Verma, L. J. Bartolotti and S. P. Gejji, J. Phys. Chem. A, 2016, 120, 7732–7744.
- 42 S. M. Alavi and S. Yeganegi, J. Mol. Liq., 2018, 256, 330–343.
- 43 P. L. Verma and S. P. Gejji, J. Mol. Graphics Modell., 2018, 85, 304–315.
- 44 L. Guglielmero, L. Guazzelli, A. Toncelli, C. Chiappe, A. Tredicucci and C. S. Pomelli, RSC Adv., 2019, 9, 30269– 30276.
- 45 H. Roohi, S. F. G. Gildeh, K. Ghauri and P. Fathei, Ionics, 2020, 26, 1963–1988.
- 46 H. Roohi, S. F. G. Gildeh, M. Mehrdad and K. Ghauri, J. Mol. Liq., 2020, 310, 113060.
- 47 A. Soltanabadi and M. Bahrami, J. Mol. Graphics Modell., 2020, 96, 107529.
- 48 P. L. Verma and S. P. Gejji, J. Mol. Struct., 2020, 1201, 127112.
- 49 A. Klamt, J. Phys. Chem., 1995, 99, 2224–2235.
- 50 J. E. Sosa, R. Santiago, A. E. Redondo, J. Avila, L. F. Lepre, M. C. Gomes, J. o. M. Araújo, J. Palomar and A. B. Pereiro, Environ. Sci. Technol., 2022, 56, 5898–5909.
- 51 N. A. Manan, C. Hardacre, J. Jacquemin, D. W. Rooney and T. G. Youngs, J. Chem. Eng. Data, 2009, 54, 2005–2022.
- 52 M. K. Hadj-Kali, M. Althuluth, S. Mokraoui, I. Wazeer, E. Ali and D. Richon, Chem. Eng. Commun., 2020, 207, 1264–1277.
- 53 Y. Zhao, N. E. Schultz and D. G. Truhlar, J. Chem. Theory Comput., 2006, 2, 364–382.
- 54 Y. Zhao and D. G. Truhlar, Theor. Chem. Acc., 2008, 120, 215– 241.
- 55 M. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, B. Mennucci and G. Petersson, Gaussian 09, Revision D.01, Gaussian Inc., Wallingford CT, 2009. Published on 09 December 2022. Downloaded on 2022. Downloaded on 2022. Downloaded on 2022. Downloaded on 2022. December 2022. Decemb
	- 56 M. Walker, A. J. Harvey, A. Sen and C. E. Dessent, J. Phys. Chem. A, 2013, 117, 12590–12600.
	- 57 S. Grimme, J. Antony, S. Ehrlich and H. Krieg, J. Chem. Phys., 2010, 132, 154104.
	- 58 S. F. Boys and F. Bernardi, Mol. Phys., 1970, 19, 553–566.
	- 59 C. M. Breneman and K. B. Wiberg, J. Comput. Chem., 1990, 11, 361–373.
	- 60 A. E. Reed, L. A. Curtiss and F. Weinhold, Chem. Rev., 1988, 88, 899–926.
	- 61 R. F. Bader and T. Nguyen-Dang, Adv. Quantum Chem., 1981, 14, 63–124.
	- 62 F. Biegler-Konig, J. Schonbohm and D. Bayles, J. Comput. Chem, 2001, 22, 545–559.
	- 63 T. Lu and F. Chen, J. Comput. Chem, 2012, 33, 580–592.
	- 64 J. Contreras-García, E. R. Johnson, S. Keinan, R. Chaudret, J.-P. Piquemal, D. N. Beratan and W. Yang, J. Chem. Theory Comput, 2011, 7, 625–632.
	- 65 W. Humphrey, A. Dalke and K. Schulten, J. Mol. Graph, 1996, 14, 33–38.
	- 66 R. Ahlrichs, M. Bar and M. Ha, Chem. Phys. Lett, 1989, 162, 165–169.
	- 67 F. Eckert and A. C. Klamt, COSMOTherm, Version C3. 0, Release 12.01, COSMOlogic GmbH & Co. KG, Leverkusen, Germany, 2013.
	- 68 P. Politzer, P. R. Laurence and K. Jayasuriya, Environ. Health Perspect., 1985, 61, 191–202.
	- 69 W. Li, C. Qi, H. Rong, X. Wu and L. Gong, Chem. Phys. Lett, 2012, 542, 26–32.
- 70 M. Moosavi, F. Khashei and E. Sedghamiz, Phys. Chem. Chem. Phys, 2018, 20, 435–448.
- 71 E. Sedghamiz, F. Khashei and M. Moosavi, J. Mol. Liq., 2018, 271, 96–104.
- 72 A. M. Fernandes, M. A. Rocha, M. G. Freire, I. M. Marrucho, J. A. Coutinho and L. M. Santos, J. Phys. Chem. B, 2011, 115, 4033–4041.
- 73 K. Dong, S. Zhang, D. Wang and X. Yao, J. Phys. Chem. A, 2006, 110, 9775–9782.
- 74 P. A. Hunt, I. R. Gould and B. Kirchner, Aust. J. Chem., 2007, 60, 9–14.
- 75 K. Fukui, Frontier orbitals and reaction paths: selected papers of Kenichi Fukui, World Scientific, 1997.
- 76 S. Armaković, S. J. Armaković, M. Vraneš, A. Tot and S. Gadžurić, J. Mol. Lig., 2016, 222, 796-803.
- 77 D. Cremer and E. Kraka, Angew. Chem., Int. Ed. Engl., 1984, 23, 627–628.
- 78 G. Saleh, C. Gatti and L. L. Presti, Comput. Theor. Chem., 2015, 1053, 53–59.
- 79 J. R. Lane, S. D. Schrøder, G. C. Saunders and H. G. Kjaergaard, J. Phys. Chem. A, 2016, 120, 6371–6378.
- 80 E. Espinosa, E. Molins and C. Lecomte, Chem. Phys. lett., 1998, 285, 170–173.
- 81 B. Li, C. Wang, Y. Zhang and Y. Wang, Green Energy Environ., 2021, 6, 253–260.
- 82 R. Babarao, S. Dai and D.-e. Jiang, J. Phys. Chem. B, 2011, 115, 9789–9794.
- 83 K. M. Gupta and J. Jiang, J. Phys. Chem. C, 2014, 118, 3110– 3118.
- 84 M. S. Shannon, J. M. Tedstone, S. P. O. Danielsen, M. S. Hindman, A. C. Irvin and J. E. Bara, Ind. Eng. Chem. Res., 2012, 51, 5565–5576.
- 85 E. J. Beckman, Chem. Commun., 2004, 1885–1888.
- 86 B. A. Marekha, O. N. Kalugin and A. Idrissi, Phys. Chem. Chem. Phys., 2015, 17, 16846–16857.
- 87 T. Payagala, J. Huang, Z. S. Breitbach, P. S. Sharma and D. W. Armstrong, Chem. Mater., 2007, 19, 5848.
- 88 T. Su, Z. Tang, C. Yin, Y. Yang, H. Wang, L. Peng, Y. Su, P. Su and J. Li, J. Mol. Liq., 2021, 327, 114857.
- 89 J. Adewole, A. Ahmad, S. Ismail and C. Leo, Int. J. Greenhouse Gas Control, 2013, 17, 46–65.
- 90 Y. Zhao, R. Gani, R. M. Afzal, X. Zhang and S. Zhang, AIChE J., 2017, 63, 1353–1367.