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# Accurately predicting solvation free energy in aqueous and organic solvents beyond 298 K by combining deep learning and the 1D reference interaction site model†

Daniel J. Fowles, Rose G. McHardy, Abdullah Ahmad and David S. Palmer \*

We report a method to predict the absolute solvation free energy (SFE) of small organic and druglike molecules in water, carbon tetrachloride and chloroform solvents beyond 298 K by combining the 1 Dimensional Reference Interaction Site Model (1D-RISM) and deep learning. RISM is a statistical mechanics based method for modelling molecular solutions that is computationally inexpensive but is too inaccurate for routine SFE calculations in its common form. By replacing the 1D-RISM SFE functional with a 1D convolutional neural network (CNN) trained on RISM correlation functions, we show that predictions approaching chemical accuracy can be obtained for aqueous and non-aqueous solvents at a wide-range of temperatures. This method builds upon the previously reported RISM-MOL-INF procedure which applied RISM to accurately characterise solvation and desolvation processes through solute–solvent correlation functions [Palmer *et al.*, *Mol. Pharm.*, 2015, **12**, 3420–3432]. Unlike RISM-MOL-INF however, the newly developed pyRISM-CNN model applied here is capable of rapidly modelling these processes in several different solvents and at a wide-range of temperatures. The pyRISM-CNN functional reduces the predictive error by up to 40-fold as compared to the standard 1D-RISM theory. Prediction errors below 1 kcal mol<sup>-1</sup> are obtained for organic solutes in carbon tetrachloride or chloroform solvent systems at 298 K and water solvent systems at 273–373 K. pyRISM-CNN has been implemented in our in-house 1D-RISM solver (pyRISM), which is made freely available as open-source software.

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## 1 Introduction

Solvation thermodynamic parameters are important in modelling many industrial processes, from the behaviour of candidate drugs in the body, to the distribution of potential pollutants in the environment. State-of-the-art computational methodologies are capable of making accurate predictions of solvation thermodynamics for aqueous systems, and have commonly done so in the calculation of p*K*<sub>a</sub>,<sup>1–3</sup> protein–ligand binding affinities<sup>4</sup> and aqueous solubility.<sup>5,6</sup> However, there has been a lack of development for organic solvents or non-ambient temperatures.

Methods of estimating SFE can generally be separated into two categories, implicit and explicit solvent models. The most common implicit models treat bulk solvent as a uniform polarisable medium defined by a dielectric constant, and have found extensive use through models such as the solvation model based on solute electron density (SMD)<sup>7</sup> and the

polarisable continuum model (PCM).<sup>8,9</sup> However, implicit models fail to capture important short-ranged solute–solvent interactions and only include molecular level details intrinsic to the dielectric constant, which limits their applicability as an approach to determining solvation free energy for complex systems. Explicit solvent models, such as molecular dynamics (MD), offer a viable alternative to implicit continuum based approaches<sup>10</sup> but at far greater computational cost.

The reference interaction site model (RISM) is a third approach, capable of calculating solvation dependent thermodynamic parameters at a lower computational cost than explicit models, whilst modelling specific solute–solvent interactions. The RISM theory uses a simplified form of the high-dimensional molecular Ornstein–Zernike (MOZ) equations to model solvent density distribution around a solute molecule through a set of correlation functions, from which two distinct methods have been developed. The most commonly used of these is 3D-RISM, which approximates the MOZ equations by a set of three-dimensional integral equations. With the recent development of several semi-empirical<sup>11,12</sup> and theoretical free energy functionals,<sup>13,14</sup> 3D RISM has found frequent use as a method to predict SFE.<sup>15–19</sup> By contrast, the 1D-RISM theory, in

Department of Pure and Applied Chemistry, University of Strathclyde, Thomas Graham Building, 295 Cathedral Street, Glasgow G1 1XL, Scotland, UK. E-mail: david.palmer@strath.ac.uk

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which the MOZ equations are approximated as a set of one-dimensional integral equations, is rarely used for quantitative calculations of solvation thermodynamics because it is considered to be too inaccurate in its common form.

Within the RISM framework, solvation free energy predictions are made analytically using one of several available free energy functionals. In 1D-RISM many of these functionals fail to accurately predict the energetic parameters of the chemical system under investigation. These functionals, such as the Hyper-Netted Chain (1D-RISM/HNC) model,<sup>20</sup> are too inaccurate for routine use and typically achieve absolute prediction errors above 20 kcal mol<sup>-1</sup>. Much effort has been put into improving the predictive capabilities of 1D-RISM based functionals for SFE calculations. Some of these improved models, such as the Gaussian Fluctuations (1D-RISM/GF) and Partial Wave models (1D-RISM/PW), can more accurately predict SFE than previous methods.<sup>21,22</sup> Although reasonable qualitative agreement with experimental data has been reported, large predictive errors are still commonly observed for many chemical systems.

A novel method for improving the accuracy of 1D-RISM calculated solvation free energies was introduced by Ratkova *et al.*<sup>23</sup> This method combines RISM and cheminformatics into a hybrid approach by making empirical corrections to 1D-RISM calculated SFE. By including correction parameters determined from chemical descriptors, this structural descriptors correction (SDC) model was able to lower the prediction error of small organic molecules in aqueous solvent at 298 K to 1.2 kcal mol<sup>-1</sup>. However, the inclusion of dataset specific descriptors limits the wider applicability of this approach, with the potential need for reparameterisation when new molecules are introduced. SFE calculations are limited to 298 K and aqueous solvent by the standard SDC model.

In previous work, Palmer *et al.* reported a new method of accurately predicting physico-chemical properties of drug-like molecules from 1D-RISM characterised solvation and desolvation processes.<sup>24</sup> This method, RISM-MOL-INF, trained partial least squares (PLS) models on 1D-RISM correlation functions generated from a dataset of small organic molecules. By replacing the inaccurate 1D-RISM free energy functionals, it was shown that RISM-MOL-INF could make accurate predictions of hydration free energy and caco-2 cell permeability. However, RISM-MOL-INF is untested beyond a limited subset of small organic molecules in aqueous solvent at 298 K. As partial least squares was the only reported model, it is unknown whether a non-linear statistical method may be more suited for this approach. Further, the RISM-MOL-INF method used the RISM-MOL solver, which includes hard-coded solvent and temperature limitations, preventing simulations beyond aqueous solvent and 298 K.

Here, we present an overhaul of the RISM-MOL-INF process with our in-house 1D-RISM solver.<sup>25</sup> This solver, pyRISM, provides a more adaptable solver for the solute-solvent 1D RISM equations. Unlike its predecessors, which were limited to aqueous solutions at 298 K, pyRISM is capable of rapidly modelling solvation thermodynamics in both water and most common organic solvents, and at a wide-range of temperatures. By replacing existing machine learning models with a deep learning approach, pyRISM-CNN can make significantly more

accurate SFE predictions and has been tested on a considerably larger dataset of organic molecules. These predictions, just as with the pyRISM 1D-RISM implementation, can be expanded to organic solvents and temperatures beyond 298 K. Moving from 1D-RISM calculation to pyRISM-CNN prediction requires minimal additional computational expense as descriptors can be generated as part of the typical 1D-RISM workflow.

## 2 Theory

### 2.1 1D-RISM

The details of the general RISM theory have been discussed in depth elsewhere,<sup>26</sup> and so only the 1D-RISM theory will be explained here. 1D-RISM uses an approximated one-dimensional form of the molecular Ornstein-Zernike equation with spherically symmetric site-site correlation functions for the modelling of molecular solutions. Both solute and solvent molecules are treated as spherically symmetric sites that solely depend on the distance between these sites, with an individual atom being the simplest representation. There are three types of site-site correlation functions that are considered in RISM: intramolecular correlation functions, total correlation functions and direct correlation functions. The intramolecular correlation functions describe the structure of a given molecule. For two sites within a molecule,  $s$  and  $s'$ , the intramolecular correlation function is written as

$$\omega_{ss'}(r) = \frac{\delta(r - r_{ss'})}{4\pi r_{ss'}^2} \quad (1)$$

where  $r_{ss'}$  is the distance between sites and  $\delta(r - r_{ss'})$  is the Dirac delta function.

Intermolecular solute-solvent correlations are defined for each pair of solute and solvent sites by the total correlation functions  $h_{s\alpha}(r)$  and direct correlation functions  $c_{s\alpha}(r)$ . Here,  $s$  refers to a solute site and  $\alpha$  to a solvent site. The total correlation functions are closely related to the radial distribution function (RDF) as

$$h_{s\alpha}(r) = g_{s\alpha}(r) - 1 \quad (2)$$

where  $g_{s\alpha}(r)$  is the radial distribution function of solvent sites around a given solute site.

The total and direct correlation functions are related *via* a set of RISM equations

$$h_{s\alpha}(r) = \sum_{s'=1}^M \sum_{\xi=1}^N \int_{R^3} \int_{R^3} \omega_{ss'}(|r_1 - r'|) \times c_{s'\xi}(|r' - r''|) \chi_{\xi\alpha}(|r'' - r_2|) dr'' dr' \quad (3)$$

where  $r = |r_1 - r_2|$ ,  $\chi_{\xi\alpha}(r)$  are the bulk solvent susceptibility functions and  $M$  and  $N$  are the number of solute and solvent sites respectively. Any mutual correlations between bulk solvent sites are described by the solvent susceptibility functions  $\chi_{\xi\alpha}^{\text{solv}}(r)$ , which are determined from solvent-solvent site total correlation functions  $h_{\xi\alpha}^{\text{solv}}(r)$ , intramolecular correlation function  $\omega_{\xi\alpha}^{\text{solv}}(r)$  and the solvent bulk number density  $\rho$ .

$$\chi_{\xi\alpha}(r) = \omega_{\xi\alpha}^{\text{solv}}(r) + \rho h_{\xi\alpha}^{\text{solv}}(r) \quad (4)$$



The solvent–solvent site  $h_{\xi\alpha}^{\text{solv}}(r)$  and  $\omega_{\xi\alpha}^{\text{solv}}(r)$  are obtained from preliminary solvent–solvent 1D-RISM calculations and molecular structure. To complete the set of RISM equations, closure relations must be introduced

$$h_{s\alpha}(r) = \exp(-\beta u_{s\alpha}(r) + \gamma_{s\alpha}(r) + B_{s\alpha}(r)) - 1 \quad (5)$$

where  $u_{s\alpha}(r)$  is the atom–atom potential,  $B_{s\alpha}(r)$  is a bridge function,  $\beta = 1/k_{\text{B}}T$  and  $\gamma_{s\alpha}$  is the indirect correlation function ( $\gamma_{s\alpha}(r) = h_{s\alpha}(r) - c_{s\alpha}(r)$ ).

The exact bridge functions are typically unknown and so an approximation is needed to solve for the total correlation functions and direct correlation functions. A commonly used closure is the Kovalenko and Hirata (KH) closure<sup>27</sup>

$$h_{s\alpha}(r) = \begin{cases} \exp(\Xi_{s\alpha}(r)) - 1 & \Xi_{s\alpha}(r) \leq C \\ \exp(\Xi_{s\alpha}(r)) + \exp(C) - C - 1 & \Xi_{s\alpha}(r) > C \end{cases} \quad (6)$$

where  $\Xi_{s\alpha}(r) = -\beta u_{s\alpha}(r) + \gamma_{s\alpha}(r)$ . In some cases the argument of the exponent can grow uncontrollably, leading to a divergence of the numerical solution of RISM equations. To counteract this, a threshold constant  $C$  is introduced to linearize the exponent when its argument is larger than  $C$ .

There are multiple expressions available within RISM for determining solvation free energy once the total and direct correlation functions have been solved. The functional is usually selected to be consistent with the closure used within the 1D-RISM calculations. The Gaussian fluctuations approximation (GF),<sup>28</sup> KH<sup>29</sup> and hypernetted chain (HNC)<sup>20</sup> expressions are shown below.

$$\Delta G_{\text{GF}} = 2\pi\rho kT \sum_{s\alpha} \int_0^{\infty} [-2c_{s\alpha}(r) - h_{s\alpha}(r)c_{s\alpha}(r)]r^2 dr \quad (7)$$

$$\Delta G_{\text{KH}} = 2\pi\rho kT \sum_{s\alpha} \int_0^{\infty} [-2c_{s\alpha}(r) - h_{s\alpha}(r)c_{s\alpha}(r) + h_{s\alpha}^2(r)\Theta(-h_{s\alpha}(r))]r^2 dr \quad (8)$$

$$\Delta G_{\text{HNC}} = 2\pi\rho kT \sum_{s\alpha} \int_0^{\infty} [-2c_{s\alpha}(r) - h_{s\alpha}(r)c_{s\alpha}(r) + h_{s\alpha}^2(r)]r^2 dr \quad (9)$$

## 2.2 pyRISM

The pyRISM program<sup>25</sup> includes a general method of obtaining variables which contain solvation and desolvation relevant descriptors from the standard 1D-RISM free energy functionals. Previously this method was applied within the RISM-MOL framework and has been described in detail elsewhere,<sup>24,30</sup> so only a short summary of the process will be described here. Each of the free energy functionals described in eqn (7)–(9) can be condensed into a generalised form:

$$\Delta G_{\text{RISM}} = \int_0^{\infty} w(r)dr \quad (10)$$

where the integrand functional  $w(r)$  combines the prefactor ( $2\pi\rho kT$ ), and the total and direct correlation functions of a single solute into an individual function of  $r$  which is referred

to as the solvation free energy density (SFED). By then omitting the integration over  $r$ , this functional can be used to obtain variables that quantify the response of solvent molecules to the solute at chosen distances  $r$  from the solute site. The SFED functions derived from the GF, KH and HNC SFE functionals are given below:

$$gf\_w(r) = 2\pi\rho kT \sum_{s\alpha} [-2c_{s\alpha}(r) - h_{s\alpha}(r)c_{s\alpha}(r)] \quad (11)$$

$$kh\_w(r) = 2\pi\rho kT \sum_{s\alpha} [-2c_{s\alpha}(r) - h_{s\alpha}(r)c_{s\alpha}(r) + h_{s\alpha}^2(r)\Theta(-h_{s\alpha}(r))] \quad (12)$$

$$hnc\_w(r) = 2\pi\rho kT \sum_{s\alpha} [-2c_{s\alpha}(r) - h_{s\alpha}(r)(c_{s\alpha}(r) - h_{s\alpha}(r))] \quad (13)$$

When the 1D-RISM equations are solved, the total and direct correlation functions are represented on a fine grid. The values of the SFED functions at selected grid points provide variables that are denoted as  $m\_w\_n$ , where  $m$  is the 1D-RISM free energy functional from which the variable is based and  $n$  is the grid point at which the variable is evaluated. Machine learning algorithms are then trained on these variables and the subsequent model can be used for solvation free energy prediction.

## 3 Methods

### 3.1 Dataset preparation

Two sources with known experimental solvation free energies for small organic molecules were used to generate the datasets applied within this work: the Minnesota Solvation Database (MSD)<sup>31</sup> and those developed by Chamberlin *et al.*<sup>32,33</sup> The MSD contains experimental solvation free energies in water and several organic solvents at 298 K. Experimental data obtained from Chamberlin *et al.* includes hydration free energies in a 273–373 K temperature range.

As hydration free energies were available over a range of temperatures, two aqueous solvent datasets were compiled. The first dataset contained 521 solute molecules with known experimental hydration free energies at 298 K, 133 of which were obtained from Chamberlin *et al.*<sup>32,33</sup> and 388 from the Minnesota Solvation Database. The second dataset was exclusively taken from Chamberlin *et al.* and included free energy data in a 273–373 K temperature range for 272 solute molecules. By using free energies over a range of temperatures, a total of 3053 datapoints were available for this multiple temperature dataset. As experimental solvation free energies were only available from the MSD at 298 K, a single dataset was compiled each for chloroform and carbon tetrachloride with 109 and 79 solute molecules respectively. Three additional datasets were compiled, containing SFED descriptors from each solvent. These multi-solvent datasets were separated by temperature into a single 298 K and two 273–373 K datasets. The 298 K dataset was made up of aqueous and organic solvent data taken from the MSD. The first of the multi-temperature datasets



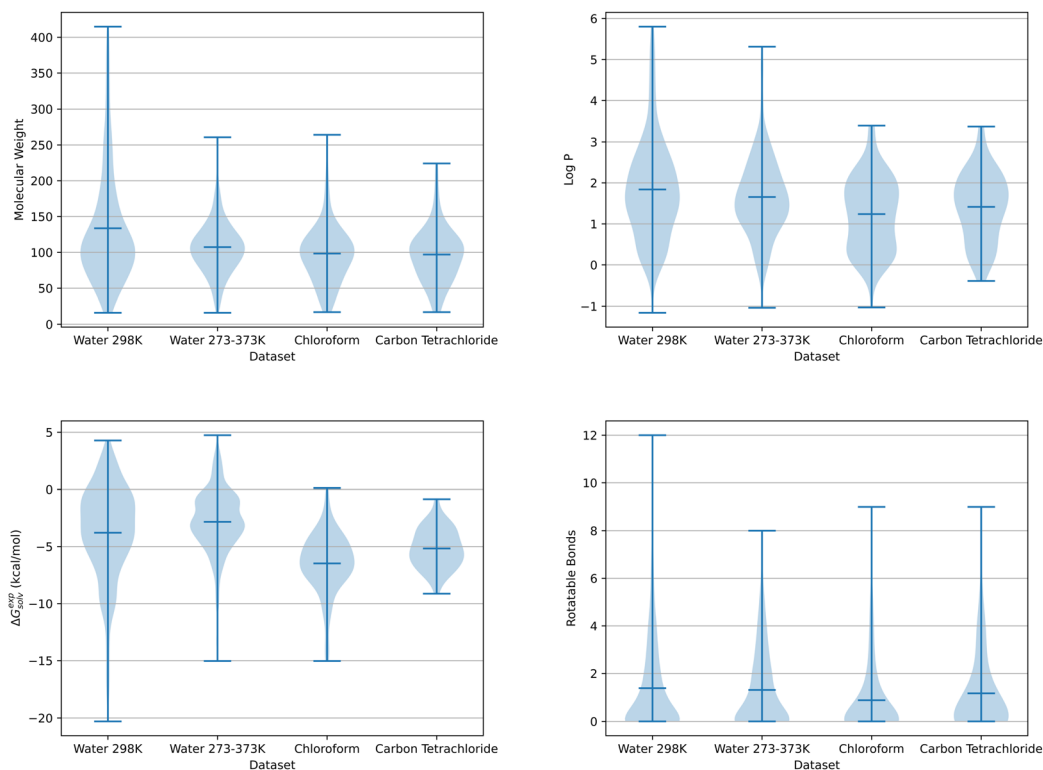


Fig. 1 Violin plots showing solute molecule data for each solvent dataset. The data shown from top left to bottom right is molecular weight, log  $P$ , experimental solvation free energy and the number of rotatable bonds.

combined Chamberlin *et al.* 273–373 K water data and MSD water and organic solvent data, while the other dataset excluded the additional MSD 298 K water data. Fig. 1 shows the distribution of solute molecules across four chemical descriptors for the chloroform, carbon tetrachloride, water and multi-temperature water datasets discussed here. These molecular properties include molecular weight, log  $P$ , experimental solvation free energy and the number of rotatable bonds per solute.

The InChi<sup>34</sup> function within Open Babel<sup>35</sup> was used to generate a unique InChi descriptor for each solute within the MSD and Chamberlin *et al.* datasets. Any duplicate molecules were then removed using the “unique” parameter within Open Babel.

A conformational search was performed on each solute molecule using MacroModel<sup>36</sup> with the OPLS-2005 forcefield<sup>37</sup> to obtain the lowest energy conformer for each solute. Each dataset then underwent 1D-RISM calculations as part of the pyRISM package to obtain free energy descriptors.

### 3.2 1D-RISM calculations

1D-RISM calculations were carried out with pyRISM using the KH closure within a system of 16 384 grid points over 20.48 Å from the solute. Aqueous solvent calculations used the dielectrically consistent reference interaction site model (DRISM),<sup>38,39</sup> while organic solvent calculations applied the extended reference interaction site model (XRISM).<sup>40</sup> Organic solvent calculations were found to converge more consistently with XRISM

than with DRISM. Solvation free energy calculations with the KH, HNC and GF free energy functionals were performed for both aqueous and organic solvent systems. For all calculations it was assumed that solute molecules were embedded within an infinitely dilute aqueous solution. A convergence tolerance of  $10^{-12}$  was set for all calculations, with a minimum tolerance of  $10^{-5}$  if the initial calculation failed to converge. The impact from lowering the minimum convergence tolerance to  $10^{-5}$ , as well as the choice of model for 1D-RISM calculations (DRISM or XRISM) on the quality of SFED generated was found to be negligible. A comparison of tolerance threshold and model choice for 1D-RISM calculations is available from Fig. 1 and 2 in the ESI.†

**3.2.1 Solvent parameters.** The Lue and Blankschtein version of the SPC/E water model (MSPC/E)<sup>41</sup> was used for modelling aqueous solvent. This altered version differs from the original model with the inclusion of modified Lennard–Jones (LJ) potential energy parameters for water based hydrogen, which were adjusted to prevent any possible divergence of the algorithm.<sup>40,42,43</sup> Both organic solvent models were modelled using the general Amber forcefield (GAFF) non-bonded parameters, which were assigned using the Antechamber and tLEaP programs within Amber18.<sup>44</sup> The Lorentz–Berthelot mixing rules<sup>45</sup> were used to generate solute–solvent LJ parameters *i.e.*,  $\sigma_{s\alpha} = (\sigma_s + \sigma_\alpha)/2$  and  $\epsilon_{s\alpha} = \sqrt{\epsilon_s \epsilon_\alpha}$ .

**3.2.2 Solute parameters.** Two sets of solute LJ parameters and atomic charges were tested: GAFF<sup>46</sup> and OPLS-2005.<sup>47</sup> GAFF parameters were assigned using the Antechamber and tLEaP





**Table 1** Breakdown of each descriptor dataset used as input for machine learning models. In total, six datasets were compiled for each solvent<sup>a</sup>

Solvent	Temperature range	Temp. descr.	SFE functional	Solute parameters	Datapoints
Carbon tetrachloride	298 K	No	KH/HNC/GF	GAFF/OPLS	79
Chloroform	298 K	No	KH/HNC/GF	GAFF/OPLS	109
Water	298 K	No	KH/HNC/GF	GAFF/OPLS	521
	273–373 K	No	KH/HNC/GF	GAFF/OPLS	3053
	273–373 K	Yes	KH/HNC/GF	GAFF/OPLS	3053
	298 K	No	KH/HNC/GF	GAFF/OPLS	709
Multi-solvent	273–373 K	No	KH/HNC/GF	GAFF/OPLS	3241
	273–373 K	Yes	KH/HNC/GF	GAFF/OPLS	3241
	273–373 K	No	KH/HNC/GF	GAFF/OPLS	3629
	273–373 K	Yes	KH/HNC/GF	GAFF/OPLS	3629

<sup>a</sup> Each of these six datasets can be separated according to the free energy functional and forcefield used to parameterise solute molecules, while the multi-temperature datasets can also be separated by the inclusion of temperature descriptors. The multi-solvent/multi-temperature datasets can be further separated by the inclusion of additional MSD 298 K water data, which the dataset with 3629 datapoints includes.

programs within Amber18,<sup>44</sup> while Maestro was used to assign OPLS-2005 parameters.<sup>48</sup>

### 3.3 Obtaining RISM solvation free energy descriptors

Solute specific SFE descriptors were obtained as a 1D-RISM calculation output using the pyRISM program. A descriptor set was generated for each free energy functional, totaling three sets per set of solute forcefield parameters. As the grid used to represent the 1D-RISM total and direct correlation functions was very fine, leading to multiple correlated variables, a coarser grid-spacing was used to obtain SFE descriptors. To minimise the inclusion of redundant data and to exclude data at long solute–solvent separations in the region where SFEDs approach zero, only every 40th grid point from  $r = 0 \text{ \AA}$  to  $r = 8 \text{ \AA}$  was considered. This approach produced 160 descriptors per SFE functional for each solute 1D-RISM calculation.

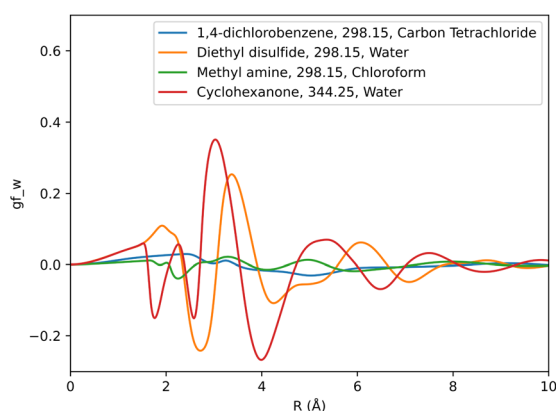
These descriptors were then used as an input to machine learning models to predict experimental solvation free energy. For each solvent a total of six separate SFED datasets were compiled. Each of these SFED datasets contain solute specific descriptors generated from pyRISM calculations involving

combinations of the KH, HNC and GF free energy functionals and OPLS, GAFF based solute parameters. For example, a typical breakdown of the 298 K chloroform dataset would look like: chloroform-298-KH-GAFF, chloroform-298-HNC-GAFF, chloroform-298-GF-GAFF, chloroform-298-KH-OPLS, chloroform-298-HNC-OPLS, chloroform-298-GF-OPLS. Multi-solvent datasets were also compiled for all the relevant free energy functional and forcefield combinations, which included SFED from each solvent. Multi-temperature datasets were tested in two separate formats, with one including an additional descriptor representing the temperature at which each solute's respective experimental SFE was recorded. Table 1 provides a breakdown of each dataset and its contents. Fig. 2 shows descriptors produced from the GF functional for GAFF parameterised solute molecules in carbon tetrachloride, chloroform and water solvent.

### 3.4 Machine learning models

Machine learning models were trained on all six SFED variations for each of the solvent datasets, as well as a combined solvent dataset. Models were validated by nested cross-validation (CV), with hyper parameters tuned by an inner 5-fold CV loop. Final tuned model performance was evaluated by an outer 50-fold Monte Carlo cross-validation loop with a 70% train/30% test split. A stratified sampling approach was taken for multi-temperature and multi-solvent datasets to ensure datapoints were separated by molecule before splitting into training, validation and test sets. More details on hyper parameter selection for our convolutional neural network (CNN), partial least squares (PLS) and random forest (RF) models can be found in Sections 3.4.1, 3.4.2 and 3.4.3 respectively. Each variable was centered by subtracting its mean value in the training data, and scaled by dividing by the standard deviation of its values in the training data. A repository of SFED datasets and scripts to train CNN, PLS and RF models can be found at: <https://doi.org/10.5281/zenodo.7108371>.

**3.4.1 Convolutional neural network.** Convolutional Neural Networks (CNN) were built using the 'sequential' model package in Tensorflow<sup>49</sup> and accessed using Keras<sup>50</sup> with a Python implementation. Several rounds of hyper parameter



**Fig. 2** Graphical representation of SFED functions generated using the GF functional and GAFF forcefield for a range of solute molecules. The temperature and solvent specific to each solute is also noted.



tuning were carried out to determine the best CNN architecture. A simple CNN architecture consisting of single layers of Conv1D-MaxPooling1D-BatchNormalisation with a single Flatten layer and Dense layer output was used as a starting point, with each layer using its default parameters. More complex architectures were subsequently tested and followed a set structure where one or two additional Conv1D-MaxPooling1D-BatchNormalisation blocks could be included. These blocks could be followed by a combination of Flatten-Dense-Dropout before output. From these architectures a range of Conv1D, MaxPooling, Dense and Dropout hyper parameter values were trialled.

The final architecture included three blocks of Conv1D-MaxPooling1D-BatchNormalisation with a subsequent Flatten layer and Dense output layer. Convolutional layers were created using the 'Conv1D' layer package in Keras with 32 output filters, a kernel size of 3 and stride length of 2. No padding was included and the rectified linear activation function (ReLU)<sup>51</sup> was used. Each of the subsequent layers were also taken from Keras, with the max pool size within MaxPooling1D layers set to 2. Default parameters were used for BatchNormalisation and Flatten layers. The loss function and metric was set to 'mse' (mean squared error), with the 'Adam' optimiser.<sup>52</sup> Each model could run for a maximum of 60 epochs with a patience of 20 epochs included through the Keras 'EarlyStopping' callback.

**3.4.2 Partial least squares.** Partial Least Squares (PLS) models were trained using the 'PLSRegression' package within Scikit-learn<sup>53</sup> with a Python implementation. Hyper parameter tuning was carried out to determine the optimal number of components between 1 and 30. A final value of 10 components was chosen.

**3.4.3 Random forest.** Random Forest (RF) models were trained using the 'RandomForestRegressor' package within Scikit-learn with a Python implementation. The Random Forest algorithm has been shown to be insensitive to training parameters, such that increasing the number of trees above 500 has little effect on prediction accuracy.<sup>54,55</sup> Therefore only the node size and minimum sample number per leaf were tested, while the number of trees per model was set to 500. Hyper parameter tuning determined that a node size of 2 and minimum sample number per leaf of 1 to be optimal. The maximum number of randomly selected features to test at each split was set as the square root of the total number of features. A tabulated breakdown of Random Forest model performance can be found in the ESI.†

### 3.5 Statistical analysis

Solvation free energy predictions were evaluated against experimental values of SFE using the coefficient of determination ( $R^2$ ) and root mean squared deviation (RMSD).

$$R^2 = 1 - \frac{\sum_{i=1}^N (y^i - y_{\text{exp}}^i)^2}{\sum_{i=1}^N (y^i - M(y_{\text{exp}}))^2} \quad (14)$$

$$\text{RMSD}(y, y_{\text{exp}}) = \sqrt{\frac{1}{N} \sum_i (y^i - y_{\text{exp}}^i)^2} \quad (15)$$

where index  $i$  goes through a set of  $N$  molecules, and  $y^i$  and  $y_{\text{exp}}^i$  are the predicted and experimental values for molecule  $i$  respectively. The coefficient of determination represents a statistical measure of how well the regression predictions fit the experimental data, and so negative values below 1 are possible for models which fit the data worse than the mean of the experimental data. The total deviation can be separated into two parts: bias (or mean displacement,  $M$ ) and standard deviation (or SDEP,  $\sigma$ ).

$$\text{bias} = M(y - y_{\text{exp}}) = \frac{1}{N} \sum_{i \in S} (y^i - y_{\text{exp}}^i) \quad (16)$$

$$\sigma(y - y_{\text{exp}}) = \sqrt{\frac{1}{N} \sum_{i \in S} (y^i - y_{\text{exp}}^i - M(y - y_{\text{exp}}))^2} \quad (17)$$

The bias provides the systematic error, while the standard deviation gives the random error that is not explained by the model. The bias and standard deviation are connected to the RMSD by:

$$\text{RMSD}(y, y_{\text{exp}})^2 = M(y - y_{\text{exp}})^2 + \sigma(y - y_{\text{exp}})^2 \quad (18)$$

A model which reports an RMSD greater than the standard deviation of the experimental data provides less accurate predictions than the null model provided by the mean of the experimental data.

Statistical analyses were performed in a Python environment using the 'sklearn.metrics' module available in scikit-learn.<sup>53</sup>

## 4 Results and discussion

### 4.1 pyRISM solvation free energy predictions

Solvation free energy calculations were performed using our in-house 1D-RISM solver, with the KH, HNC and GF free energy functionals. Table 2 provides a breakdown of SFEs calculated by the standard 1D-RISM theory (*i.e.* eqn (7)–(9)) separated by solvent, temperature and solute parameters. In-line with previous studies, 1D-RISM free energy functionals are generally unable to accurately predict solvation free energy. The KH and HNC functionals are particularly inaccurate, with a root mean squared deviation (RMSD) exceeding 40 kcal mol<sup>-1</sup> for hydration free energies. As noted in the Theory section, the functional is usually selected to be consistent with the closure used within the 1D-RISM calculation. Therefore, using the HNC functional with the KH closure is not standard practice, but the results are provided here as a direct comparison for the CNN model trained on the related SFEDs; using the HNC functional with the HNC closure also results in large errors.<sup>24</sup> Conversely, the GF functional provides the most accurate predictions, with an average RMSD between OPLS and GAFF of 1.67, 2.31 and 6.86 kcal mol<sup>-1</sup> for carbon tetrachloride, chloroform and water respectively. A similar trend is seen from a comparison of  $R^2$ ,



Table 2 Solvation free energy predictions from pyRISM calculations using the KH, HNC and GF free energy functionals<sup>a</sup>

Solvent	Temperature	R <sup>2</sup>	RMSD	Bias	SDEP	R <sup>2</sup>	RMSD	Bias	SDEP
<b>KH</b>		<b>GAFF</b>				<b>OPLS</b>			
Water	298 K	-109.24	44.13	-39.50	19.69	-135.63	49.21	-44.04	21.95
Water	273–373 K	-201.43	39.37	-36.25	15.35	-219.61	41.10	-38.22	15.11
Chloroform	298 K	-35.72	16.12	-15.04	5.80	-42.16	17.48	-16.45	5.91
Carbon tetrachloride	298 K	-100.96	17.46	-16.55	5.59	-104.14	17.73	-16.82	5.61
<b>HNC</b>		<b>GAFF</b>				<b>OPLS</b>			
Water	298 K	-117.20	45.70	-41.09	20.00	-143.78	50.65	-45.51	22.23
Water	273–373 K	-215.75	40.73	-37.71	15.39	-233.86	42.40	-39.60	15.16
Chloroform	298 K	-40.86	17.21	-16.16	5.93	-47.15	18.46	-17.44	6.04
Carbon tetrachloride	298 K	-121.71	19.16	-18.25	5.83	-125.50	19.45	-18.55	5.86
<b>GF</b>		<b>GAFF</b>				<b>OPLS</b>			
Water	298 K	-1.03	5.98	-3.53	4.83	-3.56	8.99	-7.50	4.96
Water	273–373 K	-2.94	5.49	-4.42	3.26	-5.36	6.98	-6.45	2.66
Chloroform	298 K	0.10	2.52	1.59	1.95	0.39	2.09	0.43	2.04
Carbon tetrachloride	298 K	-0.04	1.76	0.67	1.63	0.16	1.59	0.55	1.49

<sup>a</sup> Predictions are separated by solute parameters, temperature and solvent. Units are in kcal mol<sup>-1</sup>. Organic solvents were modelled with GAFF parameters, and water solvent was modelled with MSPC/E.

which shows that only the GF functional is able to provide reasonable correlation between experimental and predicted SFE. Large negative R<sup>2</sup> values suggest most of these models fit the data worse than the mean of experimental SFE data (more information on the statistical analysis used in this study can be found in Section 3.5). The choice of solute forcefield parameters also impacts prediction accuracy. This is clearest for aqueous solvent calculations performed at 298 K, with OPLS giving a higher RMSD than GAFF by 5.08, 4.95 and 3.01 kcal mol<sup>-1</sup> for KH, HNC and GF respectively. To a lesser extent the same trend is observed for multi-temperature aqueous solvent calculations, with an average decrease in RMSD of 1.63 kcal mol<sup>-1</sup> from OPLS to GAFF over all functionals. Organic solvent models were the least affected by the choice of solute forcefield parameters, leading to an average change in RMSD for chloroform and carbon tetrachloride of 0.47 kcal mol<sup>-1</sup> over all functionals. The noticeable divergence in prediction accuracy between aqueous and organic solvents may be due to differences in dataset size

and composition, rather than a particular failure of RISM to model organic solvents. The 298 K water dataset, with roughly five times more solutes present than in the chloroform or carbon tetrachloride datasets, is considerably larger. Combining this with the greater range of values present within each of the four chemical descriptors shown in Fig. 1, may explain the poor aqueous solvent predictions. Further, although the increase in dataset size for the multi-temperature dataset is not as significant, a drop in 1D-RISM performance may be occurring for calculations outwith 298 K. Fig. 3 shows the correlation plots of calculated solvation free energy against experimental values for the multi-temperature water dataset with the GF functional and GAFF solute parameters.

## 4.2 Convolutional neural network predictions

**4.2.1 Aqueous solvent models.** Convolutional neural network models were trained on each of the three aqueous

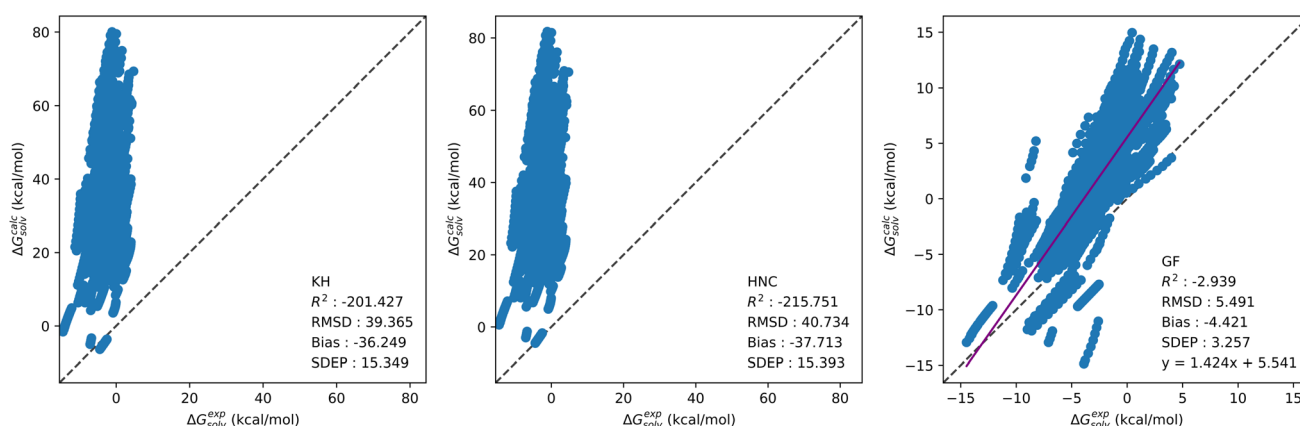


Fig. 3 Correlation plots showing calculated solvation free energy from the standard 1D-RISM theory against experimental values for the best performing multi-temperature water dataset, which used the GF functional and GAFF solute parameters.



**Table 3** Solvation free energy predictions for aqueous solvent using convolutional neural network models trained on KH, HNC and GF calculated descriptors<sup>a</sup>

Solvent	Temp.	Temp. descr.	Datapoints	$R^2$	RMSD	Bias	SDEP	$R^2$	RMSD	Bias	SDEP
<b>KH</b>				<b>GAFF</b>				<b>OPLS</b>			
Water	298 K	No	521	0.95	0.91	0.04	0.89	0.95	0.89	0.01	0.85
	273–373 K	No	3053	0.93	0.66	−0.01	0.65	0.91	0.75	0.02	0.74
	273–373 K	Yes	3053	0.95	0.55	0.01	0.54	0.93	0.65	0.01	0.64
<b>HNC</b>				<b>GAFF</b>				<b>OPLS</b>			
Water	298 K	No	521	0.94	0.94	−0.14	0.91	0.95	0.88	−0.01	0.85
	273–373 K	No	3053	0.94	0.64	−0.01	0.63	0.90	0.75	0.03	0.74
	273–373 K	Yes	3053	0.94	0.57	0.02	0.56	0.91	0.71	0.03	0.70
<b>GF</b>				<b>GAFF</b>				<b>OPLS</b>			
Water	298 K	No	521	0.94	0.96	−0.06	0.91	0.95	0.88	0.02	0.86
	273–373 K	No	3053	0.93	0.65	0.02	0.64	0.92	0.68	0.03	0.67
	273–373 K	Yes	3053	0.94	0.62	0.01	0.61	0.94	0.59	0.01	0.58

<sup>a</sup> Predictions are separated by temperature, inclusion of a temperature descriptor, solute forcefield parameters and number of datapoints. Results for each model are taken from test set predictions. Units are in kcal mol<sup>−1</sup>. Water solvent was modelled with MSPC/E. The standard deviation of each statistic per model is available in the ESI.

solvent SFED datasets, as shown in Table 1, to predict solvation free energy. For each descriptor dataset, six variations were tested: KH, HNC GF generated descriptors and OPLS, GAFF solute parameters, for a total of 18 models. Table 3 provides a breakdown of the test set based performance for each of these models.

The CNN models give significantly more accurate predictions of SFE than the standard 1D-RISM theory. From a direct comparison of SFEs computed by the standard 1D-RISM theory and by the CNN models, taken from Tables 2 and 3 respectively, a significant improvement across all measurements can be seen. In stark contrast to the standard 1D-RISM theory, from which predictions using all SFE functionals give RMSDs greater than 5 kcal mol<sup>−1</sup>, all 18 CNN models achieve an RMSD below 1 kcal mol<sup>−1</sup>. A remarkable consistency can be seen between CNN trained on *kh\_w*, *hnc\_w* and *gf\_w* calculated descriptors. This consistency across functionals suggests important and

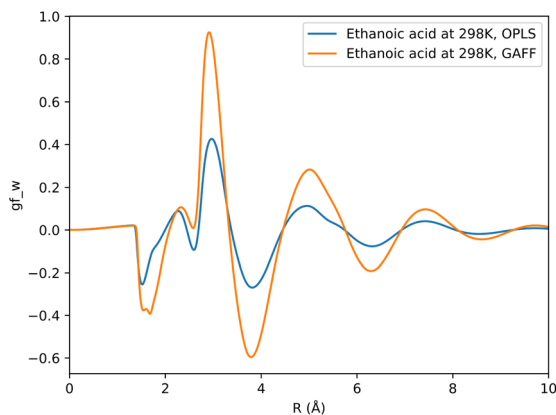
solute specific physical data is present within 1D-RISM free energy functionals that cannot be applied during their routine use, regardless of their accuracy. Indeed the only change in RMSD above 0.1 kcal mol<sup>−1</sup> between functionals occurs for the multi-temperature OPLS water dataset with temperature descriptors, which notes an increase from 0.59 kcal mol<sup>−1</sup> for *gf\_w* to 0.71 kcal mol<sup>−1</sup> for *hnc\_w*.

The choice of forcefield for solute parameters has a marginal impact on prediction accuracy. Between GAFF and OPLS datasets, an average change in RMSD of 0.07, 0.10 and 0.05 kcal mol<sup>−1</sup> is observed for KH, HNC and GF functionals respectively. Fig. 4 shows how the choice of forcefield affects the shape and magnitude of a set of descriptors, where the GF functional has been used to generate these descriptors. From this example it can be noted that both forcefields affect SFED magnitude, while maintaining a similar shape across the 1D-RISM grid, which may suggest that the shape of a given SFED is more important to a CNN for accurate SFE predictions.

Including temperature descriptors in a CNN model appears to have a limited impact on performance across all functionals and forcefields. On average, over all multi-temperature dataset variations, RMSD decreases by 0.07 kcal mol<sup>−1</sup> when temperature descriptors are included.

**4.2.2 Organic solvent models.** Convolutional neural network models were trained on chloroform and carbon tetrachloride solvent SFED datasets, as shown in Table 1, to predict solvation free energy. For each descriptor dataset six variations were tested: KH, HNC GF generated descriptors and OPLS, GAFF solute parameters, for a total of 12 models. Table 4 provides a breakdown of the test set based performance for each of these models.

Similarly to CNN models trained on aqueous solvent datasets, those trained on organic solvents consistently outperformed the standard 1D-RISM theory. The best performing CNN models achieved an RMSD of 0.44 and 0.73 kcal mol<sup>−1</sup> for carbon tetrachloride and chloroform respectively, whereas the



**Fig. 4** Comparison of descriptors generated using the GF functional for ethanoic acid in aqueous solvent, where the solute has been parameterised using either GAFF or OPLS forcefield parameters.





**Table 4** Solvation free energy predictions for chloroform and carbon tetrachloride using convolutional neural network models trained on KH, HNC and GF calculated descriptors<sup>a</sup>

Solvent	Temp.	Temp. descr.	Datapoints	$R^2$	RMSD	Bias	SDEP	$R^2$	RMSD	Bias	SDEP
<b>KH</b>				<b>GAFF</b>				<b>OPLS</b>			
Carbon tet.	298 K	No	79	0.93	0.44	0.06	0.42	0.91	0.45	0.03	0.40
Chloroform	298 K	No	109	0.92	0.74	0.00	0.72	0.90	0.76	-0.01	0.75
<b>HNC</b>				<b>GAFF</b>				<b>OPLS</b>			
Carbon tet.	298 K	No	79	0.93	0.45	0.03	0.43	0.91	0.47	-0.01	0.44
Chloroform	298 K	No	109	0.89	0.78	-0.00	0.76	0.90	0.74	0.03	0.72
<b>GF</b>				<b>GAFF</b>				<b>OPLS</b>			
Carbon tet.	298 K	No	79	0.91	0.47	0.05	0.44	0.90	0.51	-0.01	0.47
Chloroform	298 K	No	109	0.89	0.80	0.01	0.77	0.91	0.73	0.00	0.72

<sup>a</sup> Predictions are separated by temperature, inclusion of a temperature descriptor, solute forcefield parameters and number of datapoints. Results for each model are taken from test set predictions. Units are in kcal mol<sup>-1</sup>. Organic solvents were modelled with GAFF parameters. Carbon tet. refers to Carbon tetrachloride. The standard deviation of each statistic per model is available in the ESI.

RMSD of the most accurate 1D-RISM predictions were three times higher.

From Table 4, the impact of both OPLS and GAFF solute parameters to the overall accuracy of CNN models can be seen, with an average change in RMSD from OPLS to GAFF parameterised solutes of 0.01, 0.03 and 0.05 kcal mol<sup>-1</sup> for *kh\_w*, *hnc\_w* and *gf\_w* respectively. Further, with consistent performance across all functionals the viewpoint that solvation descriptors can be generated from any 1D-RISM free energy functional, regardless of its performance during 1D-RISM calculations, is reinforced. These points again suggest chemically relevant solvation data is present within these 1D-RISM generated descriptors.

The number of solute molecules per solvent dataset appears to have a negligible impact on SFE prediction accuracy. Chloroform and carbon tetrachloride have considerably fewer solutes than water at 298 K: 79 and 109 solutes, compared to 521 for water. However, organic solvent models make SFE predictions to a similar accuracy as their aqueous counterparts.

**4.2.3 Combined solvent models.** Convolutional neural network models were trained on multi-solvent SFE descriptor datasets, as shown in Table 1, to predict solvation free energy. For each descriptor dataset, six variations were tested: KH, HNC, GF generated descriptors and OPLS, GAFF solute parameters, for a total of 30 models. Table 5 provides

**Table 5** Solvation free energy predictions for multi-solvent datasets containing descriptors for all three solvents using convolutional neural network models trained on KH, HNC and GF calculated descriptors<sup>a</sup>

Solvent	Temp.	Temp. descr.	Datapoints	$R^2$	RMSD	Bias	SDEP	$R^2$	RMSD	Bias	SDEP
<b>KH</b>				<b>GAFF</b>				<b>OPLS</b>			
Multi-solvent	298 K	No	709	0.95	0.83	-0.01	0.82	0.95	0.82	-0.02	0.80
	273–373 K	No	3241	0.88	0.94	0.08	0.92	0.89	0.91	0.03	0.88
	273–373 K	Yes	3241	0.88	0.97	0.06	0.95	0.87	1.01	0.03	0.98
	273–373 K	No	3629	0.93	1.01	0.10	1.00	0.91	1.12	0.00	1.10
	273–373 K	Yes	3629	0.93	1.03	0.21	1.00	0.92	1.06	-0.02	1.05
<b>HNC</b>				<b>GAFF</b>				<b>OPLS</b>			
Multi-solvent	298 K	No	709	0.95	0.84	0.06	0.83	0.95	0.80	0.02	0.79
	273–373 K	No	3241	0.89	0.93	0.08	0.90	0.77	1.15	0.05	1.11
	273–373 K	Yes	3241	0.88	0.98	0.11	0.95	0.86	1.04	0.07	1.01
	273–373 K	No	3629	0.93	1.00	0.09	0.98	0.93	0.98	0.00	0.97
	273–373 K	Yes	3629	0.93	0.97	0.12	0.95	0.92	1.08	0.03	1.07
<b>GF</b>				<b>GAFF</b>				<b>OPLS</b>			
Multi-solvent	298K	No	709	0.95	0.87	0.05	0.85	0.95	0.77	0.01	0.76
	273–373K	No	3241	0.90	0.91	0.03	0.89	0.36	1.27	0.07	1.24
	273–373K	Yes	3241	0.88	0.97	0.10	0.93	0.74	1.10	-0.02	1.07
	273–373K	No	3629	0.94	0.97	0.04	0.96	0.94	0.93	0.05	0.91
	273–373K	Yes	3629	0.94	0.95	0.04	0.95	0.93	1.04	0.03	1.03

<sup>a</sup> Predictions are separated by temperature, inclusion of a temperature descriptor, solute forcefield parameters and number of datapoints. Results for each model are taken from test set predictions. Units are in kcal mol<sup>-1</sup>. Organic solvents were modelled with GAFF parameters, and water solvent was modelled with MSPC/E. The standard deviation of each statistic per model is available in the ESI.



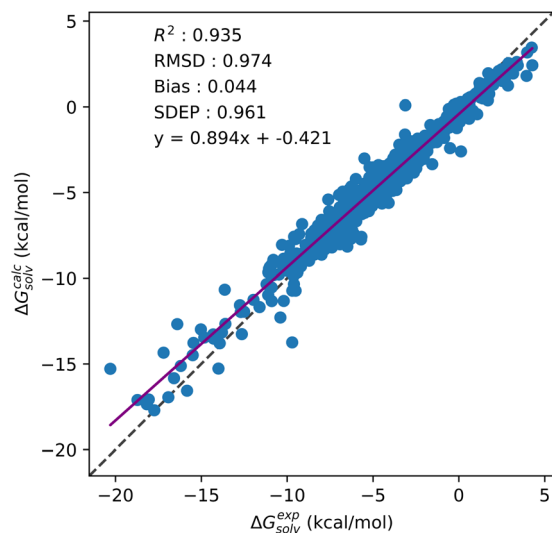


Fig. 5 Correlation plot showing test set solvation free energy predictions against experimental values for the GF and GAFF based multi-solvent and multi-temperature dataset.

a breakdown of the test set based performance for each of these models.

From Table 5, a drop in accuracy is observed for OPLS based multi-temperature datasets when compared to their GAFF counterparts. This drop in model performance is clearest for multi-temperature datasets that do not include the additional 298 K water data from the MSD (3241 datapoint models). For these models, on average between like-for-like multi-temperature datasets, the use of OPLS results in a 0.13 kcal mol<sup>-1</sup> increase in prediction error. Significant variations in correlation between experimental and predicted SFE are also only observed for these OPLS datasets. When CNN are trained on datasets which include the additional 298 K water data, prediction errors only increase by 0.05 kcal mol<sup>-1</sup> on average against their GAFF counterparts. This drop in accuracy is not observed for single solvents, nor is it present in multi-solvent GAFF models, suggesting the use of SFE descriptors generated

from OPLS parameterised solutes in multi-solvent models will interfere in prediction quality. As with the single-solvent aqueous and organic models, negligible changes in model performance are seen across the multi-solvent *kh\_w*, *hnc\_w* and *gf\_w* datasets. As GAFF based models are consistent across KH, HNC and GF, the drop in performance of OPLS based models is likely only caused by the use of OPLS parameters. Including a temperature descriptor for multi-temperature models also has a negligible impact on prediction accuracy. CNN trained on multi-solvent and multi-temperature GAFF datasets predict SFE to below 1 kcal mol<sup>-1</sup> of experiment without any need for re-parameterisation. Fig. 5 provides the correlation plot for the 3629 datapoint multi-solvent and multi-temperature dataset generated using GAFF and GF without any temperature descriptors.

A breakdown of solvation free energy predictions for individual solvents within multi-solvent datasets generated using the GF functional is shown in Table 6. Training CNN on multi-solvent SFED datasets results in an average increase in prediction error for chloroform and carbon tetrachloride solvents of 0.26 and 0.32 kcal mol<sup>-1</sup> respectively against their corresponding single solvent models. A drop in accuracy is not observed for water however, and may be due to there being significantly more experimental data available for water (as compared to the organic solvents). Despite the fact each multi-solvent dataset includes disproportionately more aqueous solvent datapoints than organics solvent, SFE prediction errors for chloroform and carbon tetrachloride only peak at 1.14 and 0.81 kcal mol<sup>-1</sup> respectively. The greatest prediction error for each solvent is below the standard deviation of experimental SFE values at 4.21, 2.85, 2.67 and 1.74 kcal mol<sup>-1</sup> for water 298 K, water 273–373 K, chloroform and carbon tetrachloride respectively.

pyRISM-CNN is capable of predicting the solvation free energies of small organic molecules with comparable accuracy to state-of-the-art methods. MD based free energy perturbation (FEP) calculations have been reported for a dataset of 239 small molecules in water, achieving an average unsigned error (AUE) of 1.10 kcal mol<sup>-1</sup>.<sup>56</sup> The semi empirical universal correction (UC) free energy functional paired with 3D-RISM has been

Table 6 Solvation free energy predictions for individual solvents within the multi-solvent datasets using convolutional neural network models trained on GF calculated descriptors<sup>a</sup>

GF				GAFF				OPLS			
Solvent	Temp.	Temp. descr.	Datapoints	R <sup>2</sup>	RMSD	Bias	SDEP	R <sup>2</sup>	RMSD	Bias	SDEP
Water				0.96	0.83	-0.02	0.81	0.96	0.79	-0.00	0.78
Chloroform	298 K	No	709	0.83	1.03	0.25	0.93	0.89	0.78	0.06	0.73
Carbon tet.				0.74	0.81	0.25	0.68	0.86	0.61	0.03	0.54
Water				0.95	0.95	0.01	0.94	0.95	0.92	0.01	0.91
Chloroform	273–373 K	No	3629	0.80	1.14	0.13	1.08	0.82	1.05	0.11	0.98
Carbon tet.				0.80	0.76	0.12	0.68	0.82	0.71	0.13	0.61
Water				0.95	0.96	-0.02	0.95	0.94	1.07	0.04	1.06
Chloroform	273–373 K	Yes	3629	0.85	1.01	0.10	0.97	0.81	1.13	0.04	1.03
Carbon tet.				0.77	0.80	0.20	0.74	0.81	0.72	0.00	0.62

<sup>a</sup> SFE predictions are separated by temperature, inclusion of a temperature descriptor, solute forcefield parameters and number of datapoints. Results for each model are taken from test set predictions. Units are in kcal mol<sup>-1</sup>. Organic solvents were modelled with GAFF parameters, and water solvent was modelled with MSPC/E. Carbon tet. refers to Carbon tetrachloride.



shown to accurately predict hydration free energies for a set of 504 organic molecules, with an RMSD of  $1.18 \text{ kcal mol}^{-1}$ .<sup>11</sup> The most comprehensive comparisons can be made against the SMD, which has been tested extensively against both aqueous and non-aqueous solvents at 298 K.<sup>7</sup> By solvent, AUE of 0.52, 0.84 and  $0.59 \text{ kcal mol}^{-1}$  were reported for carbon tetrachloride, chloroform and water respectively with the SMD. Although not directly comparable, RMSD values of 0.76, 1.14 and  $0.95 \text{ kcal mol}^{-1}$  were determined with the GF based multi-solvent pyRISM-CNN model. These errors drop to 0.47, 0.8 and  $0.65 \text{ kcal mol}^{-1}$  for carbon tetrachloride, chloroform and water respectively with the GF based single solvent output pyRISM-CNN models.

## 5 Conclusions

Here we have applied pyRISM, a new 1D-RISM solver capable of applying the standard 1D-RISM theory beyond 298 K, to water and a wide range of common organic solvents. Further, by combining pyRISM with a deep learning model, accurate SFE predictions can be made far beyond the capabilities of the standard 1D-RISM theory. Replacing the standard 1D-RISM SFE functionals with a CNN delivers a 40-fold improvement compared to the standard 1D-RISM theory, with consistent prediction errors below  $1 \text{ kcal mol}^{-1}$  as compared to experiment for organic solvents at 298 K and aqueous solvent at 273–373 K. This move from 1D-RISM calculations to pyRISM-CNN predictions requires minimal additional computational expenditure as descriptors can be generated as part of the typical 1D-RISM workflow. Efforts are ongoing to assess the generalisability of pyRISM-CNN and gather experimental SFE data in other organic solvents against which to test and further develop the model.

## Data availability

The pyRISM v0.1.1 code for solving 1D RISM equations and computing solvation free energy density functions can be found as freely available and open-source software at <https://zenodo.org/record/7107645>. Future versions of this software will be released at <https://github.com/2AUK/pyRISM>. Software and datasets used to develop the pyRISM-CNN models are freely available as open-source software at <https://doi.org/10.5281/zenodo.7108371>. The molecular structures and experimental solvation free energy data contained within this repository were initially obtained from: A. V. Marenich, C. P. Kelly, J. D. Thompson, G. D. Hawkins, C. C. Chambers, D. J. Giesen, P. Winget, C. J. Cramer and D. G. Truhlar, Minnesota Solvation Database – version 2012, University of Minnesota, Minneapolis, 2012. A. C. Chamberlin, C. J. Cramer and D. G. Truhlar, *J. Chem. Phys. B*, 2006, **110**, 5665–5675. A. C. Chamberlin, C. J. Cramer and D. G. Truhlar, *J. Chem. Phys. B*, 2008, **112**, 3024–3039.

## Author contributions

Conceptualization: DSP; methodology: DJF, RGMcH, AA; software: DJF, RGMcH, AA; validation/verification: DJF, RGMcH,

AA, DSP; formal analysis: DJF, RGMcH; investigation: DJF, RGMcH, AA; resources: DSP, AA; data curation: DJF, RGMcH; writing – original draft: DJF; writing – review editing: DJF, RGMcH, AA, DSP; visualization: DJF, DSP; supervision: DSP; project administration: DSP, DJF; funding acquisition: DSP.

## Conflicts of interest

There are no conflicts to declare.

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## Notes and references

- 1 L. Xu and M. L. Coote, *J. Phys. Chem. A*, 2019, **123**, 7430–7438.
- 2 M. S. Bodnarchuk, D. M. Heyes, D. Dini, S. Chahine and S. Edwards, *J. Chem. Theory Comput.*, 2014, **10**, 2537–2545.
- 3 F. R. Dutra, C. de Souza Silva and R. Custodio, *J. Phys. Chem. A*, 2021, **125**, 65–73.
- 4 S. Genheden, T. Luchko, S. Gusarov, A. Kovalenko and U. Ryde, *J. Phys. Chem. B*, 2010, **114**, 8505–8516.
- 5 D. S. Palmer, J. L. McDonagh, J. B. O. Mitchell, T. van Mourik and M. V. Fedorov, *J. Chem. Theory Comput.*, 2012, **8**, 3322–3337.
- 6 D. J. Fowles, D. S. Palmer, R. Guo, S. L. Price and J. B. O. Mitchell, *J. Chem. Theory Comput.*, 2021, **17**, 3700–3709.
- 7 A. V. Marenich, C. J. Cramer and D. G. Truhlar, *J. Phys. Chem. B*, 2009, **113**, 6378–6396.
- 8 J. Tomasi, B. Mennucci and E. Cancès, *J. Mol. Struct.*, 1999, **464**, 211–226.
- 9 S.-T. Lin and C.-M. Hsieh, *J. Chem. Phys.*, 2005, **125**, 124103.
- 10 D. L. Mobley, C. I. Bayly, M. D. Cooper, M. R. Shirts and K. A. Dill, *J. Chem. Theory Comput.*, 2009, **5**, 350–358.
- 11 J.-F. Truchon, B. M. Pettitt and P. Labute, *J. Chem. Theory Comput.*, 2014, **10**, 934–941.
- 12 D. S. Palmer, A. I. Frolov, E. L. Ratkova and M. V. Fedorov, *J. Phys.: Condens. Matter*, 2010, **22**, 492101.
- 13 V. Sergiievskiy, G. Jeanmairet, M. Levesque and D. Borgis, *J. Chem. Phys.*, 2015, **143**, 184116.
- 14 M. Misin, M. V. Fedorov and D. S. Palmer, *J. Chem. Phys.*, 2015, **142**, 091105.
- 15 S. Tanimoto, N. Yoshida, T. Yamaguchi, S. L. Ten-no and H. Nakano, *J. Chem. Inf. Model.*, 2019, **59**, 3770–3781.
- 16 D. Roy and A. Kovalenko, *J. Phys. Chem. A*, 2019, **123**, 4087–4093.
- 17 M. Misin, D. S. Palmer and M. V. Fedorov, *J. Phys. Chem. B*, 2016, **120**, 5724–5731.
- 18 M. Misin, P. A. Vainikka, M. V. Fedorov and D. S. Palmer, *J. Chem. Phys.*, 2016, **145**, 194501.



- 19 M. Misin, M. V. Fedorov and D. S. Palmer, *J. Phys. Chem. B*, 2016, **120**, 975–983.
- 20 S. J. Singer and D. Chandler, *Mol. Phys.*, 1985, **55**, 621–625.
- 21 K. Sato, H. Chuman and S. Ten-no, *J. Phys. Chem. B*, 2005, **109**, 17290–17295.
- 22 S. Ten-no, J. Jung, H. Chuman and Y. Kawashima, *Mol. Phys.*, 2010, **108**, 327–336.
- 23 E. L. Ratkova, G. N. Chuev, V. P. Sergiievskiy and M. V. Fedorov, *J. Phys. Chem. B*, 2010, **114**, 12068–12079.
- 24 D. S. Palmer, M. Misin, M. V. Fedorov and A. Llinas, *Mol. Pharmaceutics*, 2015, **12**, 3420–3432.
- 25 A. Ahmad, *2AUK/pyRISM: v0.1.1*, 2022, DOI: [10.5281/zenodo.7107645](https://doi.org/10.5281/zenodo.7107645).
- 26 E. L. Ratkova, D. S. Palmer and M. V. Fedorov, *Chem. Rev.*, 2015, **115**, 6312–6356.
- 27 A. Kovalenko and F. Hirata, *J. Phys. Chem. B*, 1999, **103**, 7942–7957.
- 28 S. Ten-no, *J. Chem. Phys.*, 2001, **115**, 3724–3731.
- 29 F. Hirata, *Molecular Theory of Solvation*, Kluwer Academic Publishers, Dordrecht, The Netherlands, 1st edn, 2003.
- 30 V. P. Sergiievskiy, W. Hackbusch and M. V. Fedorov, *J. Comput. Chem.*, 2011, **32**, 1982–1992.
- 31 A. V. Marenich, C. P. Kelly, J. D. Thompson, G. D. Hawkins, C. C. Chambers, D. J. Giesen, P. Winget, C. J. Cramer and D. G. Truhlar, *Minnesota Solvation Database – version 2012*, University of Minnesota, Minneapolis, 2012.
- 32 A. C. Chamberlin, C. J. Cramer and D. G. Truhlar, *J. Phys. Chem. B*, 2006, **110**, 5665–5675.
- 33 A. C. Chamberlin, C. J. Cramer and D. G. Truhlar, *J. Phys. Chem. B*, 2008, **112**, 3024–3039.
- 34 S. R. Heller, A. McNaught, I. Pletnev, S. Stein and D. Tchekhovskoi, *J. Cheminf.*, 2015, **7**, 23.
- 35 N. M. O’Boyle, M. Banck, C. A. James, C. Morley, T. Vandermeersch and G. R. Hutchison, *J. Cheminf.*, 2011, **3**, 33.
- 36 *Schrödinger Suite 2008, Maestro Version 8.5, MacroModel Version 9.6*, Schrödinger LLC, New York, NY, 2008.
- 37 W. L. Jorgensen, D. S. Maxwell and J. Tirado-Rives, *J. Am. Chem. Soc.*, 1996, **118**, 11225–11236.
- 38 J. Perkyns and B. M. Pettitt, *J. Chem. Phys.*, 1992, **97**, 7656–7666.
- 39 J. Perkyns and B. M. Pettitt, *Chem. Phys. Lett.*, 1992, **190**, 626–630.
- 40 P. H. Lee and G. Maggiora, *J. Phys. Chem.*, 1993, **97**, 10175–10185.
- 41 L. Lue and D. Blankschtein, *J. Phys. Chem.*, 1992, **96**, 8582–8594.
- 42 F. Hirata, P. J. Rossky and B. M. Pettitt, *J. Chem. Phys.*, 1982, **78**, 4133.
- 43 G. N. Chuev, M. V. Fedorov and J. Crain, *Chem. Phys. Lett.*, 2007, **448**, 198–202.
- 44 H. A. D. A. Case, K. Belfon, I. Ben-Shalom, J. Berryman, S. Brozell, D. Cerutti, T. Cheatham, G. C. III, V. Cruzeiro, T. Darden, R. Duke, G. Giambasu, M. Gilson, H. Gohlke, A. Goetz, R. Harris, S. Izadi, S. Izmailov, K. Kasavajhala, M. Kaymak, E. King, A. Kovalenko, T. Kurtzman, T. Lee, S. LeGrand, P. Li, C. Lin, J. Liu, T. Luchko, R. Luo, M. Machado, V. Man, M. Manathunga, K. Merz, Y. Miao, O. Mikhailovskii, G. Monard, H. Nguyen, K. O’Hearn, A. Onufriev, F. Pan, S. Pantano, R. Qi, A. Rahnamoun, D. Roe, A. Roitberg, C. Sagui, S. Schott-Verdugo, A. Shajan, J. Shen, C. Simmerling, N. Skrynnikov, J. Smith, J. Swails, R. Walker, J. Wang, J. Wang, H. Wei, R. Wolf, X. Wu, Y. Xiong, Y. Xue, D. York, S. Zhao, and P. Kollman, *Amber 2018*, 2018.
- 45 M. P. Allen and D. J. Tildesley, *Computer Simulation of Liquids*, Clarendon Press, Oxford, 1987.
- 46 J. Wang, R. M. Wolf, J. W. Caldwell, P. A. Kollman and D. A. Case, *J. Comput. Chem.*, 2004, **25**, 1157–1174.
- 47 J. L. Banks, H. S. Beard, Y. Cao, A. E. Cho, W. Damm, R. Farid, A. K. Felts, T. A. Halgren, D. T. Mainz, J. R. Maple, R. Murphy, D. M. Philipp, M. P. Repasky, L. Y. Zhang, B. J. Berne, R. A. Friesner, E. Gallicchio and R. M. Levy, *J. Comput. Chem.*, 2005, **26**, 1752–1780.
- 48 *Schrödinger Release 2022–3: Maestro*, Schrödinger, LLC, New York, N.
- 49 M. Abadi, A. Agarwal, P. Barham, E. Brevdo, Z. Chen, C. Citro, G. S. Corrado, A. Davis, J. Dean, M. Devin, S. Ghemawat, I. Goodfellow, A. Harp, G. Irving, M. Isard, Y. Jia, R. Jozefowicz, L. Kaiser, M. Kudlur, J. Levenberg, D. Mané, R. Monga, S. Moore, D. Murray, C. Olah, M. Schuster, J. Shlens, B. Steiner, I. Sutskever, K. Talwar, P. Tucker, V. Vanhoucke, V. Vasudevan, F. Viégas, O. Vinyals, P. Warden, M. Wattenberg, M. Wicke, Y. Yu and X. Zheng, *TensorFlow: Large-Scale Machine Learning on Heterogeneous Systems*, 2015, <https://www.tensorflow.org/>.
- 50 F. Chollet, *et al.*, *Keras*, 2015, <https://github.com/fchollet/keras>.
- 51 A. F. Agarap, arXiv:1803.08375, preprint, 2018.
- 52 D. Kingma and J. Ba, *International Conference on Learning Representations*, 2014.
- 53 F. Pedregosa, G. Varoquaux, A. Gramfort, V. Michel, B. Thirion, O. Grisel, M. Blondel, P. Prettenhofer, R. Weiss, V. Dubourg, J. Vanderplas, A. Passos, D. Cournapeau, M. Brucher, M. Perrot and E. Duchesnay, *J. Mach. Learn. Res.*, 2011, **12**, 2825–2830.
- 54 M. Eklund, U. Norinder, S. Boyer and L. Carlsson, *J. Chem. Inf. Model.*, 2014, **54**, 837–843.
- 55 D. S. Palmer and J. B. O. Mitchell, *Mol. Pharmaceutics*, 2014, **11**, 2962–2972.
- 56 D. Shivakumar, J. Williams, Y. Wu, W. Damm, J. Shelley and W. Sherman, *J. Chem. Theory Comput.*, 2010, **6**, 1509–1519.

