

PAPER

 View Article Online
 View Journal | View Issue
Cite this: *Nanoscale*, 2023, **15**, 19757

Impedance spectroscopy of Sb_2Se_3 photovoltaics consisting of $(\text{Sb}_4\text{Se}_6)_n$ nanoribbons under light illumination†

 Jaemin Park,^{‡a} Thomas P. Shalvey,^{‡b} Thomas Moehl,^c Kyoohee Woo,^{id d}
 Jonathan D. Major,^b S. David Tilley^c and Wooseok Yang^{id *a,e}

Sb_2Se_3 , consisting of one-dimensional $(\text{Sb}_4\text{Se}_6)_n$ nanoribbons has drawn attention as an intriguing light absorber from the photovoltaics (PVs) research community. However, further research is required on the performance-limiting factors in Sb_2Se_3 PVs. In this study, we investigated the charge carrier behavior in Sb_2Se_3 PVs by impedance spectroscopy (IS) under light illumination. $(\text{Sb}_4\text{Se}_6)_n$ nanoribbons with two different orientations were used to investigate the effect of crystal orientation on the device performance. Regardless of the $(\text{Sb}_4\text{Se}_6)_n$ orientation, negative capacitance was observed at forward bias, representing a recombination pathway at the $\text{TiO}_2/\text{Sb}_2\text{Se}_3$ interface. A comparison of the recombination resistances and lifetimes of two different Sb_2Se_3 PVs showed that a better interface could be formed by placing the $(\text{Sb}_4\text{Se}_6)_n$ ribbons parallel to the TiO_2 layer. Based on these observations, an ideal structure of the $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface is proposed, which will enhance the performance of Sb_2Se_3 PVs toward its theoretical limit.

 Received 14th August 2023,
 Accepted 27th November 2023

DOI: 10.1039/d3nr04082h

rsc.li/nanoscale

1. Introduction

Antimony selenide (Sb_2Se_3), owing to its unique properties, has recently captured considerable attention as a semiconductor material in thin-film photovoltaics (PVs). In addition to its inexpensive cost and simple composition, the fundamental properties of Sb_2Se_3 , such as its band gap, light absorption coefficient, and mobilities, demonstrate the material's potential as a promising light absorber. The more intriguing nature of Sb_2Se_3 is its one-dimensional (1D) crystal structure comprising $(\text{Sb}_4\text{Se}_6)_n$ nanoribbons packed with weak van der Waals forces. The formation of nanoribbons can be predicted through first-principles simulations. Computational calculations indicate that the surface energies for the $(hk0)$ surfaces are lower than $(hk1)$.¹ This suggests that the most prevalent surfaces in Sb_2Se_3 samples will be the $(hk0)$ surfaces, which are parallel to the $[001]$

direction, leading to the formation of nanoribbons. When the ribbons are placed parallel to the grain boundary (GB), contrary to other semiconductors in which the GBs are detrimental to the performance, the GB of Sb_2Se_3 is known to be electrically benign.¹ Moreover, atomic relaxation—the so-called “self-healing” process—removes the deep trap states within the band gap of Sb_2Se_3 not only at the GBs but also at all crystallographic orientations of Sb_2Se_3 .² Owing to these interesting characteristics, efforts have been devoted to fabricate Sb_2Se_3 -based high-performance devices in the solar cell³ and water splitting⁴ fields. Despite their increasing efficiencies, it is necessary to further enhance their performance, which requires deeper understanding of the performance-limiting factors.

Impedance spectroscopy (IS) is a powerful characterization technique for optoelectronic devices. Its value lies in its ability to separate the contributions of different layers or regions of the device from the performance of the overall device.⁵ For instance, by analyzing frequency-based current-voltage responses with small amplitudes, IS can individually extract information such as resistances, capacitances, and time constants from each of the device regions.⁶ This comprehensive understanding can provide valuable insights for optimizing and improving the performance of various devices by identifying the performance-limiting sources. Consequently, the IS technique has been applied to various materials systems and devices, including (dye-sensitized) solar cells,^{7,8} water-splitting devices,⁹ batteries,¹⁰ fuel cells,¹¹ and sensors.¹²

IS measurement is also widely used by the thin-film solar cells community to demonstrate various photo-physical pro-

^aSchool of Chemical Engineering, Sungkyunkwan University (SKKU), 16419, Republic of Korea. E-mail: wooseok.yang@skku.edu

^bStephenson Institute for Renewable Energy, Physics Department, University of Liverpool, Liverpool L69 7ZF, UK

^cDepartment of Chemistry, University of Zurich, Winterthurerstrasse 190, 8057 Zurich, Switzerland

^dDepartment of Printed Electronics, Korea Institute of Machinery & Materials, Daejeon 305-343, South Korea

^eSKKU Institute of Energy Science and Technology (SIEST), Sungkyunkwan University, Suwon 16419, Republic of Korea

†Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d3nr04082h>

‡These authors contributed equally.

cesses. For instance, in halide perovskite-based solar cells, several models and mechanisms have been established, including charge transfer, recombination-related resistance and capacitance, and the presence of negative capacitance.^{13,14} However, systematic studies using IS of Sb₂Se₃ solar cells have not yet been reported. Most existing research have performed only limited comparisons among samples, while comprehensive exploration of IS data under various light and voltage conditions remains elusive.^{15–19} Notably, the crystallographic orientation of the Sb₂Se₃ thin film has a pivotal role in the solar cell performance because of its anisotropic properties.¹ Electron mobility and GB characteristics can be altered by changing the orientation of the stacked 1D (Sb₄Se₆)_n nanoribbons in Sb₂Se₃ film, which inevitably affects the device efficiencies.²⁰ Therefore, IS characterization of Sb₂Se₃-based solar cells, including the TiO₂-electron transport layers (ETL) with Sb₂Se₃ having different crystallographic orientations, is significantly important as it forms the fundamental basis for improving the overall performance of the device.

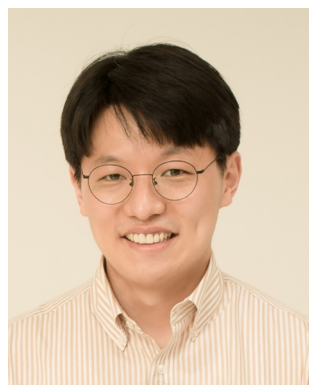
In this study, we first present the result of the in-depth IS characterization of Sb₂Se₃ solar cells. By examining the information obtained from analyzing the IS data of TiO₂-ETL Sb₂Se₃ solar cells, we attained insights about the semiconductor and TiO₂/Sb₂Se₃ interfaces, such as the recombination processes, lifetimes of carriers, and flat band potentials under different light intensities, and applied potentials. In addition, ‘negative capacitance’ was observed in Sb₂Se₃ solar cells; therefore, a recombination-related mechanism was proposed for the first time to explain this phenomenon. Furthermore, IS data was carefully investigated according to the crystallographic orientations, *i.e.*, the orientation of (Sb₄Se₆)_n nanoribbons. More favorable interface characteristics, a larger recombination resistance, and a longer lifetime were observed when the (Sb₄Se₆) ribbons were parallel to the TiO₂ layer. Based on these findings, we have proposed a novel

Sb₂Se₃ solar-cell configuration that shows promising potential for future advancements and development. Our results reveal the effectiveness of the IS technique in identifying the performance-limiting factors under operational conditions, paving the way for advancements of Sb₂Se₃-based solar cells.

2. Results and discussion

The compact film-type Sb₂Se₃-film-based PV with $[hk1]$ oriented crystallographic orientation achieved higher efficiencies than the $[hk0]$ -oriented devices because the carrier transport in Sb₂Se₃ along the $[001]$ direction was more efficient than that along the other directions owing to its anisotropic crystallographic nature.¹ A two-step close space sublimation (CSS) method was employed to avoid the formation of the thermodynamically favorable 1D structure and the rearrangement of (Sb₄Se₆)_n nanoribbons along the $[hk1]$ direction.²¹ A thin seed layer was pre-deposited at a relatively low substrate temperature (350 °C), which then recrystallized during the higher-temperature (450 °C) main growth step.²² Fig. 1a depicts the microstructures of compact Sb₂Se₃ thin films deposited on fluorine-doped tin oxide (FTO) substrate with a TiO₂-ETL by a two-step CSS deposition process. The cross-section scanning electron microscope (SEM) image of the corresponding Sb₂Se₃ film is presented in Fig. S1 (ESI†). The experimental details are described in the Experimental section. The orientation of the crystal can be anticipated by measuring the crystallographic planes, since the nanoribbons act as a fundamental building block of the Sb₂Se₃ structure if this basic repeat unit is not broken. In this regard, most of the literature used X-ray diffraction (XRD) analysis to predict the crystallographic orientation of the Sb₂Se₃ film. These findings align closely with the results obtained from the direct measurement of the orientation of (Sb₄Se₆)_n ribbons through TEM imaging.^{1,23} The crystallographic orientations of the Sb₂Se₃ film were investigated by XRD analysis (Fig. 1b). The sample showed diffraction patterns corresponding to orthorhombic Sb₂Se₃ (Joint Committee on Powder Diffraction Standards, JCPDS 15-0861, space group *Pbnm*). In addition, the dominant $[211]$ and $[2\bar{2}1]$ orientations of Sb₂Se₃ in the XRD patterns (Fig. 1c) confirm the success of the seed-derived growth strategy for preferred $[hk1]$ -oriented film growth by using a kinetically controlled growth.

Sb₂Se₃ PV devices were fabricated using Sb₂Se₃ film, a poly (3-hexylthiophene) (P3HT) hole transfer layer (HTL), and Au electrode (denoted w-seed Sb₂Se₃ PV, inset of Fig. 1d). Fig. 1d displays representative current density–voltage (*J*–*V*) curves of the Sb₂Se₃ PV. Typical *J*–*V* characterizations performed under 1 sun illumination yielded an optimal conversion efficiency of 5.01% for the Sb₂Se₃ PV with an open circuit voltage (*V*_{oc}) of 0.41 V, short circuit current density (*J*_{sc}) of 25.20 mA cm^{–2}, and fill factor (FF) of 47.96%. To provide statistical information on the device performance, 81 individual devices in a 5 × 5 cm² FTO glass were fabricated (Fig. S2, ESI†). To finalize fabrication of the subcells, HTL/Au layers were deposited on



Wooseok Yang

Prof. Wooseok Yang is an Assistant Professor at the School of Chemical Engineering at Sungkyunkwan University (SKKU) in Korea. He received his BS in 2012 from the Department of Ceramic Engineering and his PhD in 2018 from the Department of Materials Science and Engineering at Yonsei University in Korea. From 2019 to 2021, he worked as a postdoctoral researcher at the University of Zurich in Switzerland with Prof. David

Tilley before joining SKKU in 2022. His research focuses on (photo-) electrochemical energy conversion, such as photoelectrochemical water splitting, electrochemical CO₂ capture, and valorization of waste biomass.

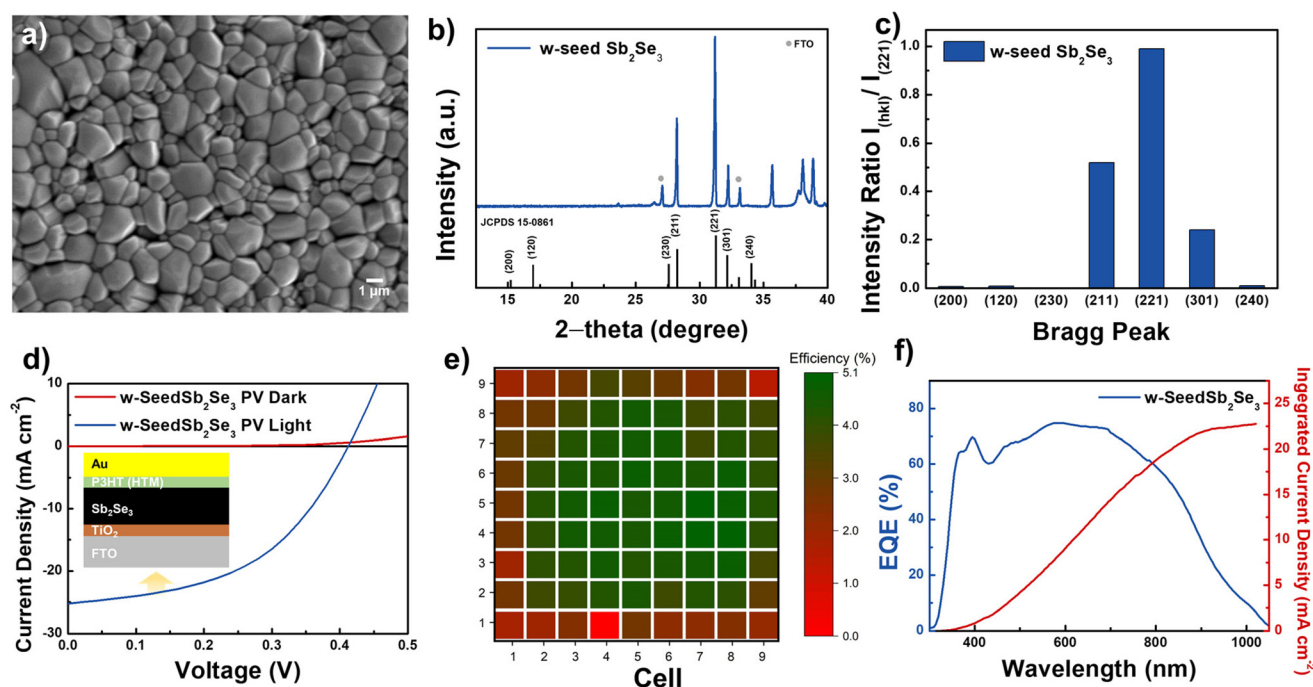


Fig. 1 (a) Top-view SEM images of the w-seed Sb_2Se_3 films. (b) Diffraction patterns and (c) peak intensities ratio ($I_{(hkl)}/I_{(221)}$) of the obtained Sb_2Se_3 films. (d) J - V characteristics of champion Sb_2Se_3 PV. (e) 2D plot of the PCE values of the large area Sb_2Se_3 PV (25 cm^2) for each sub cell. (f) External quantum efficiency (EQE) spectra of a w-seed Sb_2Se_3 -based PV.

them. With optimal depositing condition, we achieved an average V_{OC} of $0.40 \pm 0.05 \text{ V}$, J_{SC} of $20.53 \pm 3.67 \text{ mA cm}^{-2}$, and FF of $43.53 \pm 6.48\%$ from the 81 subcells (Fig. S3, ESI†), indicating the high uniformity of the device parameters in the w-seed Sb_2Se_3 PVs, except for the edges, which are more prone to pinholes (Fig. 1e). Measurement of the external quantum efficiency (EQE) showed that the photocurrent contribution extended to over 1000 nm, which was consistent with the measured optical bandgap.²⁰ In addition, the photocurrent calculated by integrating the EQE over the light spectrum was nearly identical to the values determined *via* J - V characterizations.

IS analysis was performed with the w-seed Sb_2Se_3 PVs under white LED light illumination as a function of the applied potential. Fig. 2a shows the current density with 10% sun white LED, which was calibrated with a Si diode. It should be noted that the current density is not exactly 10% of the value measured under 1 sun illumination with the Xe lamp (Fig. 1d), owing to differences in the wavelength-dependent absorption of the Si diode and Sb_2Se_3 . Fig. 2b–d show the Nyquist plots at each applied potential. When the applied potential was more positive than the V_{OC} ($\sim 0.32 \text{ V}$), an inductive response, which is mathematically equivalent to a “negative capacitance”, was observed at the very low frequency region (Fig. 2b, inset c). Considering the definition of capacitance (C), $C = dQ/dV$, a negative capacitance would mean a decreased charge (Q) when the applied voltage (V) is increased. This negative capacitance was observed in organic PV²⁴ and hybrid perovskite PV¹⁴ devices. Mora-Sero *et al.* reported a

negative capacitance in an inorganic CdSe-based device, which is not observed in a conventional monocrystalline Si solar cell.²⁵ They attributed the negative capacitance to an additional recombination pathway that opens up at forward bias and decreases the charge accumulation ability of the solar cells. It was also reported that a negative capacitance can originate from a time-dependent trap-assisted recombination process.²⁶ Note that a possible trap-mediated recombination process was proposed at the $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface in the previous IS study on Sb_2Se_3 photocathode for photoelectrochemical water splitting.⁹ Thus, the recombination at the $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface, which is activated at forward bias, could be one source of the negative capacitance. The photo-shunt, observed as the crossing point between dark and photocurrents after passing the V_{OC} (Fig. 1d), indicates activation of an additional recombination pathway in the PV device at forward bias.

Another evident feature in the IS spectra of the Sb_2Se_3 PVs is the increasing size of the semicircle as the applied potential approaches the short circuit (Fig. 2b–d). The fact that higher resistance was observed at higher photocurrent regions implies that the resistance originates from a recombination process. In order to elucidate the nature of this resistance, IS spectra with higher light intensities were measured. Fig. 3a shows the current densities with 10, 50, and 100% sun LED light illumination. A small semicircle was observed at higher light intensities, indicating decreasing resistance as the light intensity increases (Fig. S4, ESI†). We fitted the Nyquist plots with the simple Voigt circuit with three or two RC elements (Fig. S5, ESI†). The high frequency resistance, R_{HF} , which is

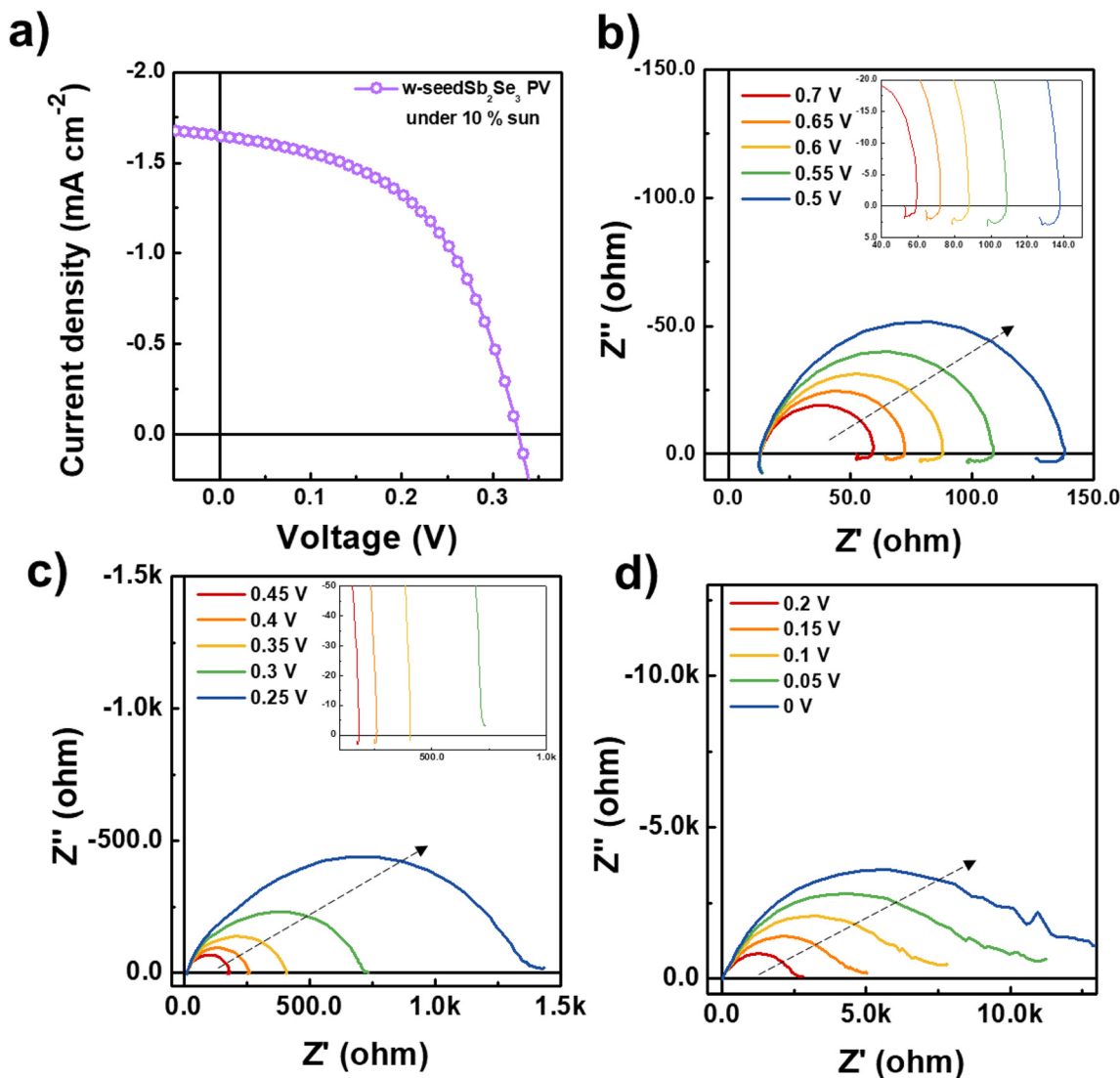


Fig. 2 (a) J - V curve of the w-seed Sb₂Se₃ PV under 10% sun illumination. (b–d) Nyquist plots of the w-seed Sb₂Se₃ PV at the designated applied potential.

almost constant over potentials (Fig. S6†), is considered an artifact and not attributed to real physical phenomena, as we demonstrated in our previous study.⁹ Thus, we focused on the resistance, capacitance, and RC time constant from the low-frequency semicircle, as shown in Fig. 3b–d. As the low-frequency resistance is related to the recombination process, we denoted it as R_{rec} . The light intensity dependency indicates that the resistance to recombination decreases (*i.e.*, the more chance to recombine) at a higher light intensity, providing further support that the resistance originates from the recombination process. The corresponding capacitance exhibited a different trend: a higher capacitance was observed in relation to the light intensity near the V_{OC} region (Fig. 3c). The capacitance originated from the depletion region at the TiO₂/Sb₂Se₃ interface, as evidenced by Mott–Schottky plots representing doping densities close to that of Sb₂Se₃ (Fig. S7 and ESI Note 1, ESI†). The higher doping density under higher light intensity may

reflect the contribution of the photo-generated charges. The RC time constant can be obtained by multiplying the resistance and capacitance, as shown in Fig. 3d. The time constant, which represents the lifetime of the photo-generated charge carriers, was higher under lower light intensity, proving that faster recombination occurs when the population of the photo-generated carrier is high.

To investigate the carrier dynamics depending on the crystallographic orientation of the Sb₂Se₃ film, we prepared a reference Sb₂Se₃ film that was deposited quickly at lower substrate temperature and did not contain a seed layer (denoted w/o-seed Sb₂Se₃ film) for a comparative study. As shown in Fig. 4a and b, the values of the XRD peak intensities of the ($hk0$) planes/peak intensities of (221) planes were higher than those of the w-seed Sb₂Se₃ film, indicating that the more (Sb₄Se₆)_{*n*} nanoribbons that grew along the [001] direction were horizontally laid on the superstrate. As shown in Fig. S8 and

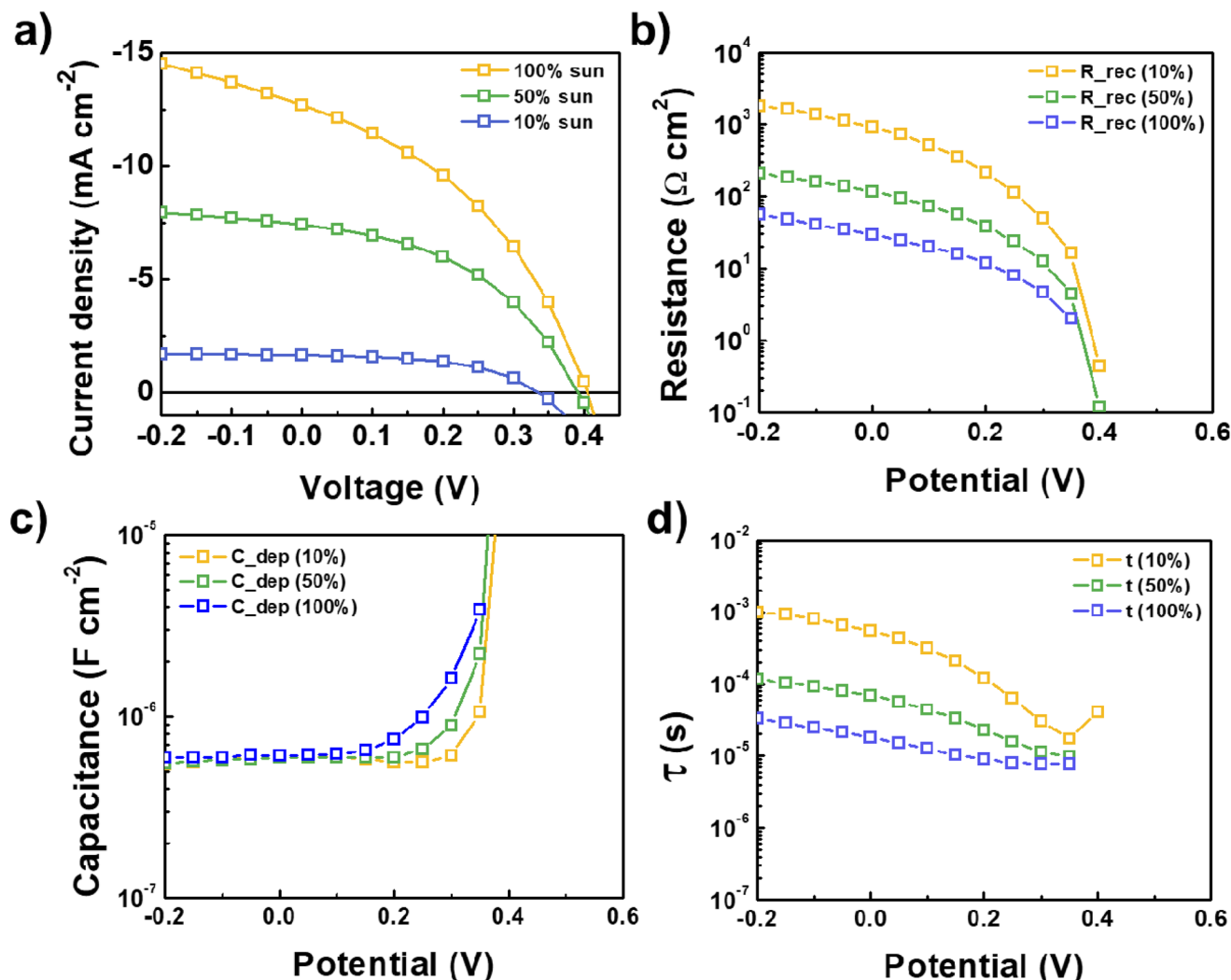


Fig. 3 (a) Steady-state photocurrent densities obtained by the w-seed Sb_2Se_3 PV during IS measurement under 10%, 50%, and 100% sun illumination. (b) Resistances (c) capacitance, and (d) time constants for charge transfer of the w-seed Sb_2Se_3 PV from the IS fitting procedure.

S9 (ESI[†]), attempts to use the w/o-seed Sb_2Se_3 film for PVs resulted in low efficiencies compared to the w-seed Sb_2Se_3 PV, especially for the photocurrent density. It seems that carriers are required to hop across (Sb_4Se_6) ribbons held by van der Waals forces in w/o-seed Sb_2Se_3 containing horizontally laid (Sb_4Se_6) ribbons. Moreover, as shown in the Fig. S10 (ESI[†]), the w/o-seed Sb_2Se_3 film contains more pin-holes than the w-seed Sb_2Se_3 film owing to its narrow thickness. The exposure of the superstrate can result in direct contact between the n-type TiO_2 and p-type P3HT, which acts as a recombination center. Therefore, poor photocurrent arose from the recombination of the photogenerated charge carriers in the w/o-seed Sb_2Se_3 film, owing to its inefficient carrier-transport behavior and decreased shunt resistance. It is crucial to note that the deposition method for the w/o-seed Sb_2Se_3 film was also optimized in our previous study to achieve a planar film-type morphology aiming at attaining a planar film-type morphology.²⁷ This optimization involved meticulous control of the source

and substrate temperatures, preventing the formation of 1D structures. The resultant film demonstrated surface morphology and XRD patterns are determined by the source and substrate temperatures, while thickness is a separated parameter.²⁷ As shown in Fig. 4c, current densities with 10% white light intensities were measured, and IS analysis was conducted with the w/o-seed Sb_2Se_3 PVs under 10% sun light illumination as a function of the applied potential (Fig. 4d–f). When the applied potential was more positive than the V_{OC} ($\sim 0.3 \text{ V}$), a similar propensity was exhibited that inductive negative capacitance was observed at the very low frequency region (Fig. 4d and e inset). Thus, it is believed that trap-mediated recombination at the $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface is also activated in w/o-seed Sb_2Se_3 at forward bias. Moreover, an increasing size of the main semicircle as the applied potential approached the short circuit was also observed in the IS spectra of the w/o-seed Sb_2Se_3 PVs, implying that the resistance originated from the recombination process, as in the w-seed Sb_2Se_3 PVs.

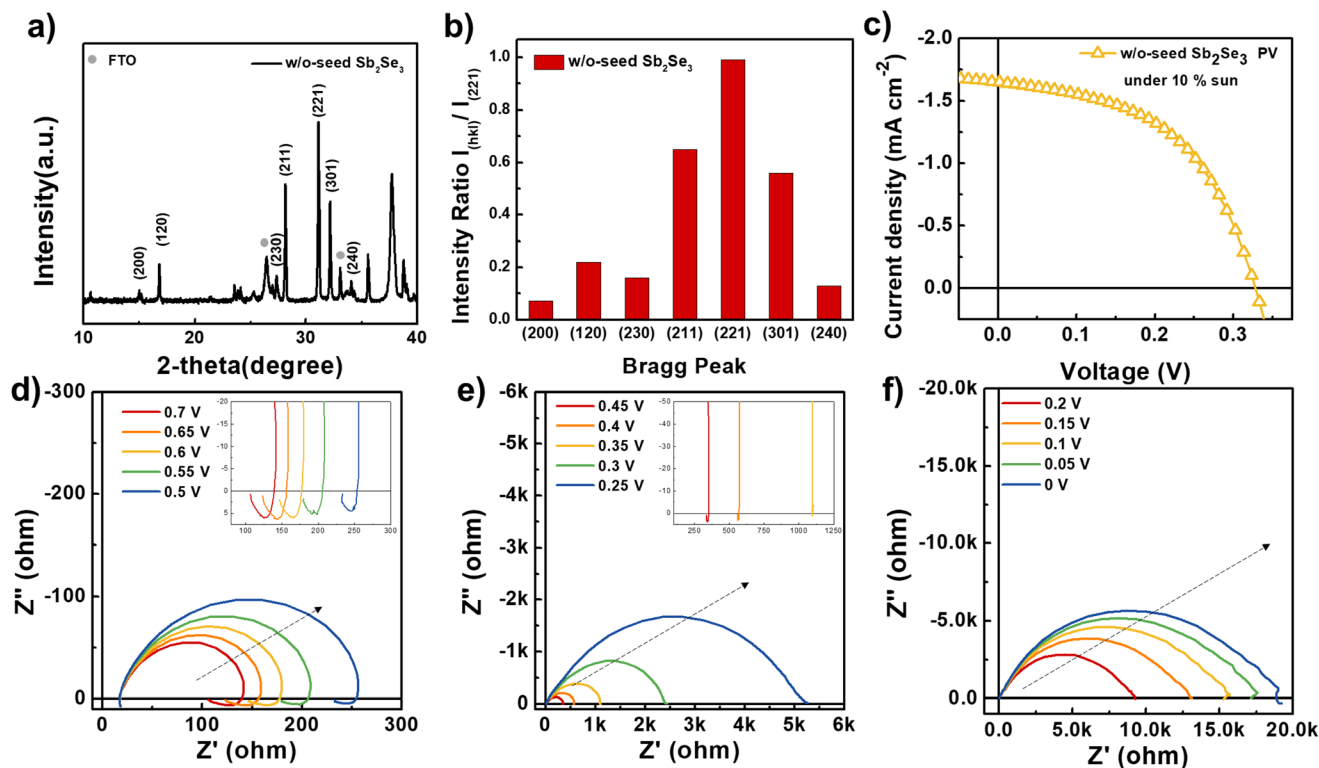


Fig. 4 (a) Diffraction patterns and (b) peak intensities ratio ($I_{(hkl)}/I_{(221)}$) of the obtained w/o-seed Sb_2Se_3 films (c) J - V behavior exhibited by the w/o-seed Sb_2Se_3 PV under 10% sun illumination. (d-f) Nyquist plots of the w/o-seed Sb_2Se_3 PV at the designated applied potential.

As shown in Fig. 5a, current densities with various LED light intensities were measured, and IS spectra with higher light intensities were measured for the w/o-seed Sb_2Se_3 PV. Increased resistance was observed as the applied potential moved toward the short circuit for all light conditions. Identical Voigt circuits (Fig. S5†) were used to fit the Nyquist plots, and the resulting resistance, RC time constants, and capacitance are displayed in Fig. 5b-d, respectively. The resistance and time constant also decreased as the light intensity increased, which confirms that similar recombination-related mechanisms may exist in w/o-seed Sb_2Se_3 PVs. Divergent capacitance profiles are evident around the V_{OC} in the two devices. However, the photo-physical processes significantly change near the V_{OC} . Some R - C elements (*e.g.*, negative capacitance) disappear while a new R - C element (*e.g.*, depletion region and recombination) appears as the applied potential passes the V_{OC} , making it challenging to fit and interpret the Nyquist plots. Thus, there can be some ambiguity and care must be taken to interpret the EIS data near this “transition condition”. In the V_{OC} region, an observation of increasing capacitance with increasing applied voltage was observed for the w-seed Sb_2Se_3 -based device. In contrast, the w/o-seed Sb_2Se_3 -based device exhibited a lower capacitance that did not exhibit an increase. We have validated that this capacitance is associated with the depletion region through Mott-Schottky plots (Fig. S7 and ESI Note 1, ESI†). In an ideal semiconductor junction, an increase in applied voltage leads to a reduction in

the depletion region and a consequent rise in capacitance originated from the depletion region (eqn (1) and (2))

$$W = \sqrt{\frac{2\epsilon_s(\phi_{bi} - V)}{q} \left(\frac{1}{N_a} + \frac{1}{N_d} \right)} \quad (1)$$

$$\frac{1}{C_{\text{depletion}}^2} = \frac{2(\phi_{bi} - V)}{qN\epsilon_s A^2} \quad (2)$$

where, V is the applied voltage, ϵ_s is the permittivity, N is the dopant concentration, and ϕ_{bi} is the built-in potential. Therefore, it is plausible to think that remained capacitance is attributed to the shunt pass effect induced by pinholes in the w/o-seed Sb_2Se_3 film, hindering the reduction of the depletion region under forward bias. Interestingly, a larger recombination resistance was observed in the w/o-seed Sb_2Se_3 PV than in the w-seed Sb_2Se_3 PV, which resulted in it having a 6.7 times longer RC time at 0 V under 100% illumination as shown in Fig. 5c. Given that a larger recombination resistance typically signifies lower recombination, this phenomenon may seem contradictory to the fact that the sample with w/o-seed Sb_2Se_3 exhibited lower performance. It should be noted that, under current measurement conditions, there is an absence of additional impedance components stemming from bulk recombination, and the $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface remains unaffected by film thickness. Consequently, it is plausible to conclude that our impedance spectroscopy observations

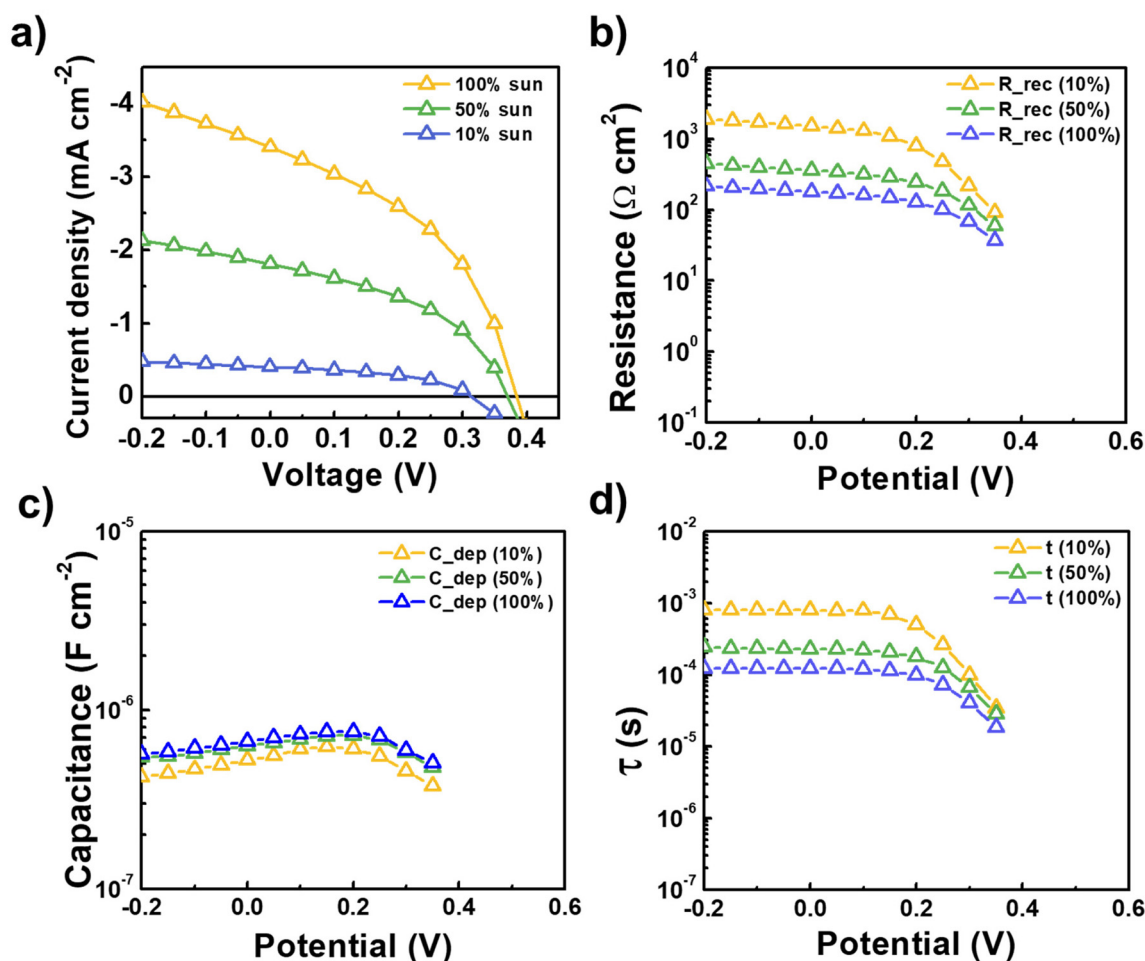


Fig. 5 (a) Steady-state photocurrent densities of the w/o-seed Sb_2Se_3 -based PV during IS measurement under 10%, 50%, and 100% sun illumination. (b) Resistances, (c) capacitance, and (d) time constants for the charge transfer of w/o-seed Sb_2Se_3 PV from the IS fitting procedure.

between the two Sb_2Se_3 solar cells are indeed independent of the thickness of the Sb_2Se_3 films.

A possible explanation for the longer lifetime in the w/o-seed Sb_2Se_3 sample is in its $\text{TiO}_2/\text{Sb}_2\text{Se}_3$ interface. As the time constant can be obtained by multiplying the resistance and capacitance originating from the depletion region at the $\text{TiO}_2/\text{Sb}_2\text{Se}_3$, it is reasonable to assume that the calculated time constant represents the lifetime of the photo-generated charge carriers at the $\text{TiO}_2/\text{Sb}_2\text{Se}_3$ interfaces at the given applied potential. In the w/o-seed Sb_2Se_3 film, containing an ($hk0$) plane, the (Sb_4Se_6) ribbons lie parallel to the TiO_2 layer. Hence, a high-quality $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface with Wan der Waals planes that has fewer trap sites can be formed. On the contrary, more recombination centers can be created at the interface of the w-seed Sb_2Se_3 film and TiO_2 owing to the dangling bond at the ($hk1$) planes. As a result, the photo-generated carriers in the interfaces have a shorter lifetime, even though they are readily collected in a vertical direction within the covalently bonded (Sb_4Se_6) ribbons along the [001] direction within the Sb_2Se_3 film that only contains ($hk1$) planes. In other words, the interface of the w/o-seed Sb_2Se_3 film and TiO_2 itself

is superior, as the nanoribbons are combined with TiO_2 horizontally. However, a lower performance is observed because of the decreased shunt resistance and severe bulk recombination derived from the pin-holes and inefficient charge transport due to the carrier hopping among the ribbons. The decreased shunt resistance of the w-seed Sb_2Se_3 PV is in accordance with the decreased dark current (Fig. S11†). Therefore, a bilayer structure is the most suitable.

A bilayer structure comprises a bottom thin layer of (Sb_4Se_6) ribbons arranged horizontally to the TiO_2 layer and a top film with vertically aligned (Sb_4Se_6) ribbons. This design ensures efficient charge transport through covalent bonds and provides sufficient thickness to prevent pinholes. Such a configuration can effectively transport the charge carriers without recombination and also passivate the interface defect (Fig. 6). However, there will be an interface between the two layers, which can possibly affect the charge transport. The practical feasibility of having vertically aligned ribbons coexist with 100% parallel-oriented nanoribbons presents a significant challenge. Consequently, there is a notable likelihood of additional film growth within certain existing vertically aligned

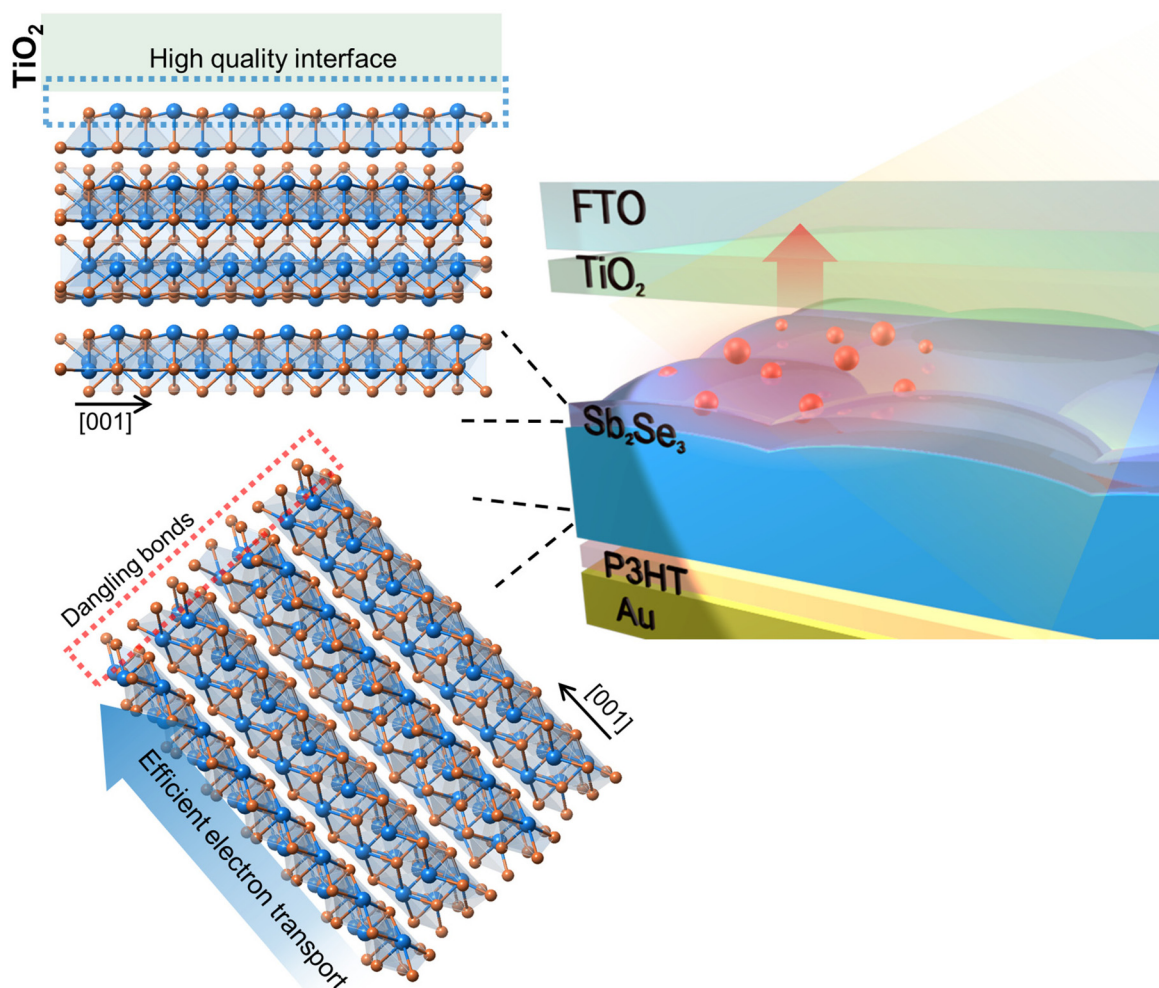


Fig. 6 Schematic illustration of the proposed novel bilayer Sb_2Se_3 configuration that can effectively transport the charge carriers without recombination and also form a high quality interface with the TiO_2 layer.

nanoribbons through the formation of new covalent bonds at dangling sites. Furthermore, the $\text{Sb}_2\text{Se}_3/\text{Sb}_2\text{Se}_3$ interface is anticipated to exhibit relative freedom from defects attributed to differences in lattice parameters. Therefore, we believe that this interface would have a comparatively lesser adverse impact on charge transport than the interface between Sb_2Se_3 and TiO_2 . Nevertheless, a more in-depth analysis is required to comprehend the implications of the additional interface and the improved $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface on performance once the device is implemented.

3. Conclusions

In summary, we performed IS characterization of Sb_2Se_3 PVs with two different (Sb_4Se_6) ribbon orientations under illumination. The w-seed Sb_2Se_3 , with the $[hk1]$ orientations, exhibited a better uniformity and performance, particularly with regard to J_{SC} and efficiency. Both samples exhibited negative capacitance behaviors at large forward bias, indicating that a trap mediated recom-

bination pathway can be activated at the $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface regardless of the crystallographic orientation. Additionally, both samples exhibited a recombination-related resistance, manifested by the potential and light intensity dependencies. The w-seed Sb_2Se_3 revealed a smaller recombination resistance and shorter lifetime, presumably owing to the terminated covalent bonds at the $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface. Based on these observations, we have proposed an ideal $\text{Sb}_2\text{Se}_3/\text{TiO}_2$ interface comprising two successive layers with different ribbon directions—parallel to the TiO_2 layer for a longer lifetime and vertical to the TiO_2 layer for better charge separation.

4. Experimental section

Fabrication of the Sb_2Se_3 PV

The TiO_2 layer was deposited on the FTO substrate using CSS. The preparation method for the Sb_2Se_3 thin films is identical to that provided in a previous report.²¹ w-seed Sb_2Se_3 was deposited in a two-step process, whereby an initial seed layer

(~100 nm) was deposited with $T_{\text{source}} = 400\text{ }^{\circ}\text{C}$ and $T_{\text{sub}} = 350\text{ }^{\circ}\text{C}$ followed by a higher temperature growth step with $T_{\text{source}} = 480\text{ }^{\circ}\text{C}$ and $T_{\text{sub}} = 400\text{ }^{\circ}\text{C}$. w/o-seed Sb_2Se_3 was deposited without a seed layer. In this case, the source tray was quickly heated to $T_{\text{source}} = 480\text{ }^{\circ}\text{C}$ without intentional heating of the substrate; however, due to the close proximity of the source tray to the substrate (~10 mm), the substrate typically reaches ~350 $^{\circ}\text{C}$ during deposition. Poly(3-hexylthiophene-2,5-diyl) (P3HT) hole transport layer (HTL) was deposited on Sb_2Se_3 film by dynamically spin coating 100 μl of a 10 mg ml^{-1} solution in chlorobenzene at 3000 rpm for 30 s. Finally, an Au electrode with a thickness of 50 nm was deposited onto the HTL by thermal evaporation through a mask to define individual 0.1 cm^2 contacts. The resulting phase evolution of the Sb_2Se_3 films was examined using a Rigaku SmartLab diffractometer with a rotating Cu anode and Ge (220) two-bounce monochromator. The surface morphology and cross section image of the samples was obtained via field emission SEM (JEOL 7001F).

Device characterization

The photovoltaic performance of the Sb_2Se_3 PV was measured under simulated AM 1.5G illumination using a TS space systems class AAA solar simulator. A standard Si reference cell was used to calibrate the 1-sun intensity. Current density–voltage (J – V) scanning was performed with a dwell time of 200 ms at each point over a range from –1 to 1 V at a rate of 0.28 V s^{-1} . The EQE was measured without an additional light bias using a Bentham PVE300 system in the wavelength range 300–1100 nm.

Impedance measurements for Sb_2Se_3 PV

IS measurements were performed with a BioLogic SP-300. The peak-to-peak AC voltage amplitude was $V_{\text{pp}} = 15\text{ mV}$ (which gave a root mean square voltage $V_{\text{rms}} = 5.3\text{ mV}$), the bias potential steps were approximately 50 mV, and the equilibrium time at each bias potential step was 17 s. A frequency range of 7 MHz to 0.1 Hz was used. In addition, a white light bias by LEDs (SP-12-W5, cool white Luxeon Rebel) was used for the IS measurements under illumination. The IS spectra were fitted with Z_{view} software from Scribner.

Author contributions

J. Park analyzed the data, conducted experiment, and wrote the manuscript. T. P. Shalvey optimized and prepared Sb_2Se_3 PVs. T. Moehl helped analyze impedance data and manuscript preparation. K. Woo assisted data analysis. J. D. Major contributed to prepared Sb_2Se_3 PV. S. D. Tilley helped contributed to the data analysis and manuscript preparation. W. Yang designed and supervised the project, directed the research, and contributed to the writing of the manuscript.

Conflicts of interest

The authors declare no competing financial interests.

Acknowledgements

This work was supported by the National R&D Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Science and ICT (grant no. RS-2023-0025177) and Korea Institute of Machinery and Materials (grant no. NK242B, NK243H). JDM and TS acknowledge funding provided by EPSRC (grant no. EP/N014057/1, EP/T006188/1, and EP/W03445X/1).

References

- 1 Y. Zhou, L. Wang, S. Y. Chen, S. K. Qin, X. S. Liu, J. Chen, D. J. Xue, M. Luo, Y. Z. Cao, Y. B. Cheng, E. H. Sargent and J. Tang, *Nat. Photonics*, 2015, **9**, 409.
- 2 K. P. McKenna, *Adv. Electron. Mater.*, 2021, **7**, 2000908.
- 3 R. F. Tang, X. M. Wang, W. T. Lian, J. L. Huang, Q. Wei, M. L. Huang, Y. W. Yin, C. H. Jiang, S. F. Yang, G. C. Xing, S. Y. Chen, C. F. Zhu, X. J. Hao, M. A. Green and T. Chen, *Nat. Energy*, 2020, **5**, 587.
- 4 W. Yang, J. Park, H. C. Kwon, O. S. Hutter, L. J. Phillips, J. Tan, H. Lee, J. Lee, S. D. Tilley, J. D. Major and J. Moon, *Energy Environ. Sci.*, 2020, **13**, 4362.
- 5 Y. Y. Proskuryakov, K. Durose, M. K. Al Turkestani, I. Mora-Sero, G. Garcia-Belmonte, F. Fabregat-Santiago, J. Bisquert, V. Barrioz, D. Lamb, S. J. C. Irvine and E. W. Jones, *J. Appl. Phys.*, 2009, **106**, 044507.
- 6 I. Mora-Sero, G. A. Garcia-Belmonte, P. P. Boix, M. A. Vazquez and J. Bisquert, *Energy Environ. Sci.*, 2009, **2**, 678.
- 7 F. Fabregat-Santiago, J. Bisquert, G. Garcia-Belmonte, G. Boschloo and A. Hagfeldt, *Sol. Energy Mater. Sol. Cells*, 2005, **87**, 117.
- 8 E. von Hauff, *J. Phys. Chem. C*, 2019, **123**, 11329.
- 9 W. Yang, T. Moehl, E. Service and S. D. Tilley, *Adv. Energy Mater.*, 2021, **11**, 2003569.
- 10 D. Andre, M. Meiler, K. Steiner, C. Wimmer, T. Soczka-Guth and D. U. Sauer, *J. Power Sources*, 2011, **196**, 5334.
- 11 S. M. R. Niya and M. Hoorfar, *J. Power Sources*, 2013, **240**, 281.
- 12 B. Pejic and R. De Marco, *Electrochim. Acta*, 2006, **51**, 6217.
- 13 H. C. Kwon, W. Yang, D. Lee, J. Ahn, E. Lee, S. Ma, K. Kim, S. C. Yun and J. Moon, *ACS Nano*, 2018, **12**, 4233.
- 14 F. Ebadi, N. Taghavinia, R. Mohammadpour, A. Hagfeldt and W. Tress, *Nat. Commun.*, 2019, **10**, 1574.
- 15 J. Zhou, D. Meng, T. H. Yang, X. T. Zhang, Z. Q. Tang, Y. Cao, J. Ni, J. J. Zhang, Z. Y. Hu and J. B. Pang, *Appl. Surf. Sci.*, 2022, **591**, 153169.
- 16 L. P. Guo, S. N. Vijayaraghavan, X. M. Duan, H. G. Menon, J. Wall, L. Y. Kong, S. Gupta, L. Li and F. Yan, *Sol. Energy*, 2021, **218**, 525.
- 17 C. Liu, K. Shen, D. X. Lin, Y. Cao, S. D. Qiu, J. Z. Zheng, F. X. Bao, Y. Y. Gao, H. B. Zhu, Z. Q. Li and Y. H. Mai, *ACS Appl. Mater. Interfaces*, 2020, **12**, 38397.

- 18 Y. F. Liu, B. Li, X. Y. Liang, T. Liu, S. F. Wang and Z. Q. Li, *J. Alloys Compd.*, 2023, **932**, 167313.
- 19 G. Li, Z. Q. Li, X. Y. Liang, C. S. Guo, K. Shen and Y. H. Mai, *ACS Appl. Mater. Interfaces*, 2019, **11**, 828.
- 20 C. Chen, D. C. Bobela, Y. Yang, S. Lu, K. Zeng, C. Ge, B. Yang, L. Gao, Y. Zhao, M. C. Beard and J. Tang, *Front. Optoelectron.*, 2017, **10**, 18.
- 21 L. J. Phillips, C. N. Savory, O. S. Hutter, P. J. Yates, H. Shiel, S. Mariotti, L. Bowen, M. Birkett, K. Durose, D. O. Scanlon and J. D. Major, *IEEE J. Photovolt.*, 2019, **9**, 544.
- 22 W. Yang, J. H. Kim, O. S. Hutter, L. J. Phillips, J. Tan, J. Park, H. Lee, J. D. Major, J. S. Lee and J. Moon, *Nat. Commun.*, 2020, **11**, 861.
- 23 R. E. Williams, Q. M. Ramasse, K. P. McKenna, L. J. Phillips, P. J. Yates, O. S. Hutter, K. Durose, J. D. Major and B. G. Mendis, *ACS Appl. Mater. Interfaces*, 2020, **12**(19), 21730.
- 24 K. Shimizu, Y. Tanaka, Y. Noguchi and H. Ishii, *J. Phys.: Conf. Ser.*, 2017, **924**, 012012.
- 25 I. Mora-Sero, J. Bisquert, F. Fabregat-Santiago, G. Garcia-Belmonte, G. Zoppi, K. Durose, Y. Proskuryakov, I. Oja, A. Belaidi, T. Dittrich, R. Tena-Zaera, A. Katty, C. Levy-Clement, V. Barrioz and S. J. C. Irvine, *Nano Lett.*, 2006, **6**, 640.
- 26 E. Ehrenfreund, C. Lungenschmied, G. Dennler, H. Neugebauer and N. S. Sariciftci, *Appl. Phys. Lett.*, 2007, **91**, 012112.
- 27 O. S. Hutter, L. J. Phillips, P. J. Yates, J. D. Major and K. Durose, *2018 IEEE 7th World Conf. Photovolt. Energy Conversion*, 2018, 0027.