


 Cite this: *RSC Adv.*, 2023, 13, 32039

# Formation of charge-transfer complexes in ionic crystals composed of 1,3-bis(dicyanomethylidene)indan anion and viologens bearing alkyl chains†

Erika Saito, Ryohei Yamakado, \* Taichi Yasuhara, Hiroto Yamaguchi, Shuji Okada and Tsukasa Yoshida \*

The relationship between charge-transfer (CT) properties and the molecular arrangement formed from  $\pi$ -electronic ion pairs remains unclear because of the limited variety of  $\pi$ -electron anions. This study addressed this issue by synthesising a series of ion pair assemblies composed of viologen dications with diverse alkyl chains as  $\pi$ -electron cations and 1,3-bis(dicyanomethylidene)indan anion ( $\text{CMI}^-$ ) as a stable  $\pi$ -electron anion. We obtained seven ionic crystals and identified their assembled structures using single-crystal X-ray analysis. These structures are categorized into three types: “columnar”, “slipped columnar” and “independent”. The CT properties were characterised using UV-Vis absorption spectroscopy, which revealed that the CT absorption bands were dependent on the alkyl chain length. This intriguing variation in the CT transitions can be explained by the differences in the type of assembled structure.

 Received 6th October 2023  
 Accepted 25th October 2023

DOI: 10.1039/d3ra06782c

[rsc.li/rsc-advances](https://rsc.li/rsc-advances)

## Introduction

Organic functional cocrystals can modulate optical properties, electrical conductivity and other properties based on the combination of molecules.<sup>1</sup> Electron-rich donors and electron-poor acceptors can be merged to obtain charge-transfer (CT) complex cocrystal materials, which are a common type of cocrystals. These materials have been reported to have potential applications in semiconductors<sup>2</sup> and ferroelectric materials.<sup>3</sup> It is widely acknowledged that the arrangement of  $\pi$ -electron-based assemblies has a significant impact on the functions of CT complexes. Specifically, CT complexes can exhibit conductivity or insulating behaviour depending on their arrangement.<sup>4</sup> Thus, it is crucial to elucidate the molecular framework and arrangement for the development of organic functional cocrystals based on  $\pi$ -conjugated systems. The formation of CT complexes can be facilitated by utilising covalent and non-covalent bonds, such as hydrogen bonds.<sup>5</sup> Furthermore, the formation of CT complexes using ionic interactions *via* the addition of a charge has also attracted attention. However, most previous studies have examined  $\pi$ -electron systems with ionic substituents, such as sulfonates and carboxylates, and only a few studies have investigated CT complexes using  $\pi$ -electron systems with delocalised charges. This is primarily due

to the limited number of reports on  $\pi$ -electron anions, such as 7,7,8,8-tetracyanoquinodimethane (TCNQ) radical anions,<sup>6</sup> pentacyanocyclopentadienide,<sup>7</sup> and tris-(7*H*-dibenzo[*c,g*]fluorenylidene-methyl)methide ion (Kuhn's anion),<sup>8</sup> compared to  $\pi$ -electron cations.<sup>9</sup> Additionally, controlling the molecular arrangement in the crystal structure *via*  $\pi$ -electronic interactions is a challenging task. Meanwhile, peripheral modifications of ionic  $\pi$ -electron systems have led to the development of “charge-by-charge assemblies”, where anions and cations are stacked alternately, and “charge segregated assemblies”, where each charge species is stacked separately.<sup>10</sup> Therefore, appropriate molecular modifications can be used to comprehend different interactions, including CT interactions, and construct various assembled forms consisting of ion pairs having the same main skeleton.

This study investigated using 1,3-bis(dicyanomethylidene)indan anion ( $\text{CMI}^-$ ), prepared *via* the self-deprotonation of 1,3-bis(dicyanomethylidene)indan ( $\text{CMI}$ ), as a  $\pi$ -electron anion for the formation of ionic CT complexes.<sup>11</sup> The fabrication of various ion-pair crystals using  $\text{CMI}^-$  has been previously reported.<sup>12</sup> Viologen, a well-known electron acceptor,<sup>13</sup> was selected as the counter cation due to its molecular structure, which allows for facile modification and possesses high symmetry. Good packing and substantial CT properties can be achieved by selecting a combination of molecules with symmetric shapes and similar molecular sizes.<sup>14</sup> In this study, new organic CT ion-pair crystals consisting of viologen dications with alkyl chains ranging from methyl to octyl (C1–C8) and  $\text{CMI}^-$  were synthesised<sup>15</sup> and their optical properties were characterised. The correlation between the crystal structure and the optical properties of the organic ion pairs was elucidated.

Department of Organic Materials Science, Graduate School of Organic Materials Science, Yamagata University, Yonezawa 992-8510, Japan. E-mail: yamakado@yz.yamagata-u.ac.jp

† Electronic supplementary information (ESI) available. CCDC 2268733–2268739. For ESI and crystallographic data in CIF or other electronic format see DOI: <https://doi.org/10.1039/d3ra06782c>



## Results and discussion

### Synthetic procedures

Salts of 1,3-bis(dicyanomethylidene)indan anions ( $\text{CMI}^-$ ) and *N*-alkyl viologen cations ( $\text{C}1^{2+}$ – $\text{C}8^{2+}$ ) were prepared *via* ion exchange using the sodium salt of  $\text{CMI}^-$  ( $\text{Na}^+\text{-CMI}^-$ ) and the chloride, bromide and iodide salts of the cations (Scheme 1). Their identifications were performed using  $^1\text{H}$  and  $^{13}\text{C}$  NMR spectroscopy (see ESI, Fig. S1–S12†).

### Single-crystal X-ray analysis

All synthesised compounds were screened using several crystallisation methods (*e.g.*, vapour diffusion and vapour

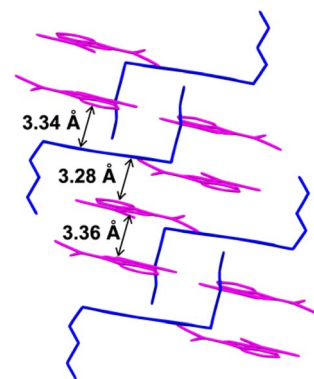
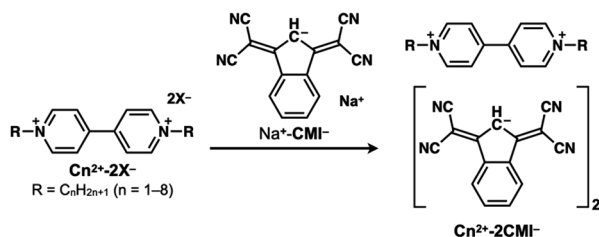


Fig. 3 Single-crystal X-ray structure of  $\text{C}5^{2+}\text{-}2\text{CMI}^-$  wherein the magenta- and blue-coloured parts represent the anion and cation species, respectively.



Scheme 1 Preparation of ion pairs comprising 1,1'-dialkyl-4,4'-bipyridinium cations ( $\text{C}1^{2+}$ – $\text{C}8^{2+}$ ) and 1,3-bis(dicyanomethylidene)indan anion ( $\text{CMI}^-$ ).

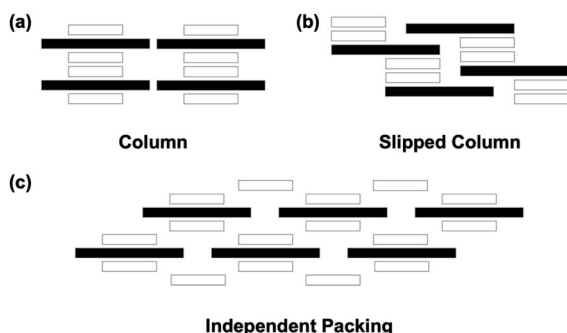


Fig. 1 Ion pair stacking motifs: (a) columnar structure; (b) slipped columnar structure; (c) independent packing structure. The black and white rectangles indicate cation and anion, respectively.

evaporation), and seven single crystals ( $\text{C}1^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}2^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}3^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}4^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}5^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}6^{2+}\text{-}2\text{CMI}^-$  and  $\text{C}8^{2+}\text{-}2\text{CMI}^-$ ) appropriate for X-ray crystallographic analysis were obtained (Table S1†). All ionic crystals were obtained as 1 : 2 complexes of viologen cations and  $\text{CMI}^-$  without other ions, and the packing structures are categorized into three types: “columnar structure”, “slipped columnar structure” and “independent structure” (Fig. 1–4).

The ionic crystals of  $\text{C}1^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}2^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}3^{2+}\text{-}2\text{CMI}^-$  and  $\text{C}4^{2+}\text{-}2\text{CMI}^-$  were crystallised in the  $C2/c$ ,  $P\bar{1}$ ,  $P\bar{1}$  and  $P\bar{1}$  space groups, respectively. Columnar structures comprising alternating anion–cation–anion stacking were observed in the stacking structures of these four ion pairs (Fig. 1a, 2 and S13–16†). These columnar structures exhibited symmetric anion–cation–anion-type stacking for  $\text{C}1^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}2^{2+}\text{-}2\text{CMI}^-$  and  $\text{C}4^{2+}\text{-}2\text{CMI}^-$ , whereas asymmetric anion–cation–anion-type stacking was observed for  $\text{C}3^{2+}\text{-}2\text{CMI}^-$ . The stacking distances between the average ring planes, calculated from the related  $\text{sp}^2$  atoms of  $\text{CMI}^-$  (9 atoms) and bipyridinium moiety (12 atoms), were estimated as 3.75 Å, 3.31 Å, 3.37/3.35 Å and 3.26 Å in  $\text{C}1^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}2^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}3^{2+}\text{-}2\text{CMI}^-$  and  $\text{C}4^{2+}\text{-}2\text{CMI}^-$ , respectively. The rotations of the long axis of  $\text{CMI}^-$  in relation to the long axis of the cation were  $30.5^\circ$ ,  $46.0^\circ$  and  $87.8^\circ$  for  $\text{C}1^{2+}\text{-}2\text{CMI}^-$ ,  $\text{C}2^{2+}\text{-}2\text{CMI}^-$  and  $\text{C}4^{2+}\text{-}2\text{CMI}^-$ , respectively. In  $\text{C}3^{2+}\text{-}2\text{CMI}^-$ , one  $\text{CMI}^-$  was observed to be rotated  $74.4^\circ$  relative to the long axis of  $\text{C}3^{2+}$ ,

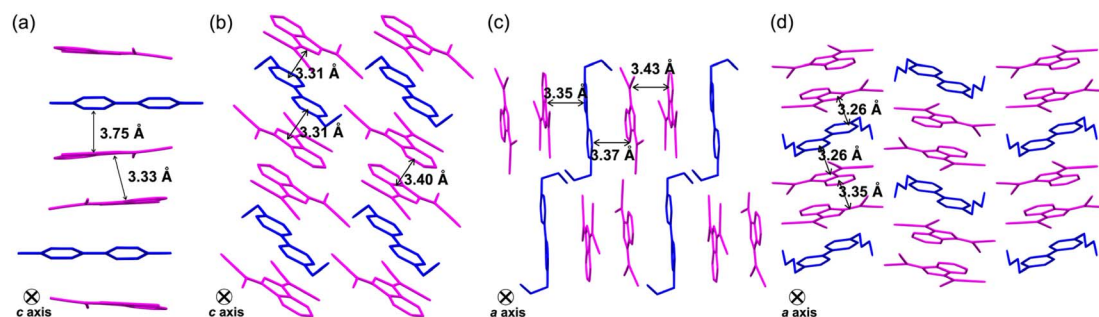


Fig. 2 Single-crystal X-ray structures of (a)  $\text{C}1^{2+}\text{-}2\text{CMI}^-$ , (b)  $\text{C}2^{2+}\text{-}2\text{CMI}^-$ , (c)  $\text{C}3^{2+}\text{-}2\text{CMI}^-$  and (d)  $\text{C}4^{2+}\text{-}2\text{CMI}^-$  wherein the magenta- and blue-coloured parts represent the anion and cation species, respectively.



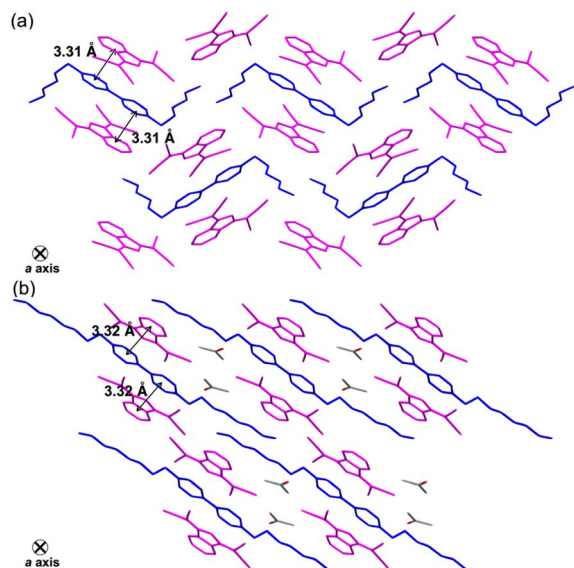


Fig. 4 Single-crystal X-ray structures of (a)  $C6^{2+}-2CMI^{-}$  and (b)  $C8^{2+}-2CMI^{-}$  wherein the magenta- and blue-coloured parts represent the anion and cation species, respectively.

whereas the other  $CMI^{-}$  was rotated  $37.6^{\circ}$  relative to the long axis of  $C3^{2+}$  (Fig. 5 and S20<sup>†</sup>). Conversely, two  $CMI^{-}$  were observed to exhibit face-to-face stacking within the column (Fig. 6). The anionic dimer was observed to exhibit anti-directional stacking, in which the dipole moments cancelled each other out, with intermolecular distances of 3.33 Å, 3.40 Å and 3.35 Å for  $C1^{2+}-2CMI^{-}$ ,  $C2^{2+}-2CMI^{-}$  and  $C4^{2+}-2CMI^{-}$ , respectively. Notably, the anion–anion stacking of  $C3^{2+}-2CMI^{-}$  exhibited a rotational deviation of approximately  $90^{\circ}$ , resulting in an intermolecular distance of 3.43 Å.

The ionic crystal of  $C5^{2+}-2CMI^{-}$  was crystallised in the triclinic  $P\bar{1}$  space group. The stacking structure of  $C5^{2+}-2CMI^{-}$

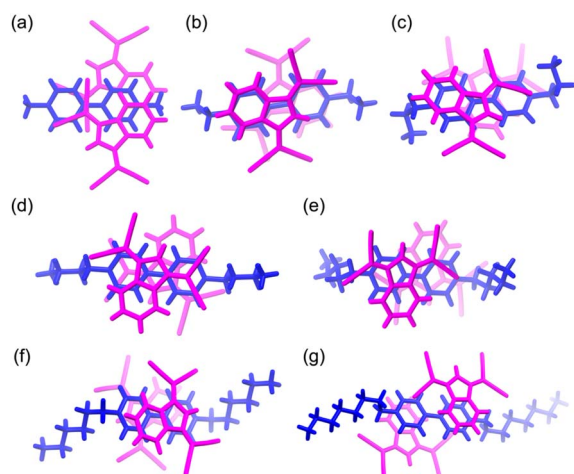


Fig. 5 The stacking structures of anion–cation–anion (top view) of (a)  $C1^{2+}-2CMI^{-}$ , (b)  $C2^{2+}-2CMI^{-}$ , (c)  $C3^{2+}-2CMI^{-}$ , (d)  $C4^{2+}-2CMI^{-}$ , (e)  $C5^{2+}-2CMI^{-}$ , (f)  $C6^{2+}-2CMI^{-}$  and (g)  $C8^{2+}-2CMI^{-}$  wherein the magenta- and blue-coloured parts represent the anion and cation species, respectively.

revealed a slipped columnar structure consisting of alternating anion–cation–anion-type stacking along the  $a$  axis (Fig. 1b, 3 and S17<sup>†</sup>). The columnar structure exhibited asymmetrical anion–cation–anion stacking. The distances between the facing-conjugated planes of the adjacent cations and anions were measured to be 3.28/3.34 Å. Additionally, the orientation of the long axis of  $CMI^{-}$  in relation to the long axis of the cation was calculated to be  $74.2^{\circ}$  (Fig. 5 and S20<sup>†</sup>). In contrast, the distance between the facing  $\pi$ -conjugated planes of the anions was recorded as 3.36 Å. Anion dimers exhibit anti-directional stacking, in which the dipole moments cancel each other (Fig. 6).

The ionic crystals of  $C6^{2+}-2CMI^{-}$  and  $C8^{2+}-2CMI^{-}$  were crystallised in the space groups  $P2_1/c$  and  $P\bar{1}$ , respectively, and  $C8^{2+}-2CMI^{-}$  contained isopropanol as the solvent. Despite the presence of anion–cation–anion-type stacking, the observed packing structures of both ion pairs lacked a columnar structure (independent packing; Fig. 1c, 4, S18 and S19<sup>†</sup>). The distances between the facing  $\pi$ -conjugated planes of the adjacent cation–anions were 3.31 Å and 3.32 Å for  $C6^{2+}-2CMI^{-}$  and  $C8^{2+}-2CMI^{-}$ , respectively. Furthermore, the orientations of the long axis of  $CMI^{-}$  in relation to the long axis of the cation were calculated to be  $25.7^{\circ}$  and  $76.4^{\circ}$  for  $C6^{2+}-2CMI^{-}$  and  $C8^{2+}-2CMI^{-}$ , respectively (Fig. 5 and S20<sup>†</sup>).

Thus, ion-pair crystals comprising  $Cn^{2+}$  and  $CMI^{-}$  formed aggregates based on the anion–cation–anion structure. Moreover, it was established that the ion-pair crystals exhibited three distinct modes of assembly, determined by the length of the alkyl chain in  $Cn^{2+}$ . In  $C1^{2+}-2CMI^{-}$ ,  $C2^{2+}-2CMI^{-}$ ,  $C3^{2+}-2CMI^{-}$ , and  $C4^{2+}-2CMI^{-}$ , characterised by short alkyl chains ( $n = 1-4$ ), a highly ordered columnar structure was formed. Upon increasing the alkyl chain length ( $n = 5$ ), the regularity of the structure was slightly impaired, resulting in the formation of a slipped columnar structure. This observation implies that the packing of the alkyl chains inhibits the electrostatic interactions between the ion pairs. In  $C6^{2+}-2CMI^{-}$  and  $C8^{2+}-2CMI^{-}$ , which supported longer alkyl chains ( $n = 6$  and  $8$ ), the packing of the alkyl chains fully suppressed the formation of a columnar

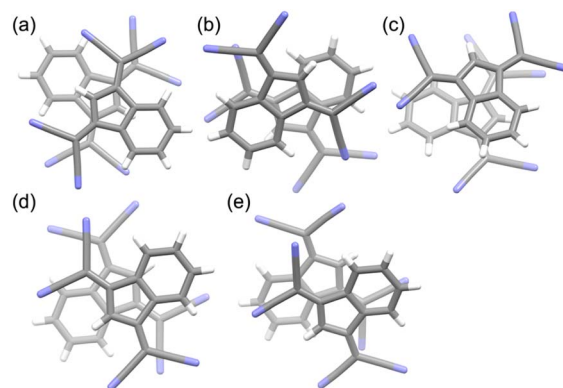


Fig. 6 The stacking structures of anion–anion (top view) of (a)  $C1^{2+}-2CMI^{-}$ , (b)  $C2^{2+}-2CMI^{-}$ , (c)  $C3^{2+}-2CMI^{-}$ , (d)  $C4^{2+}-2CMI^{-}$  and (e)  $C5^{2+}-2CMI^{-}$ . Atom colour code: grey and blue refer to carbon and oxygen, respectively.



structure, and the anion–cation–anion structure existed independently. Interestingly, attention to the dihedral angles between the pyridinium moieties in  $Cn^{2+}$  tended towards  $0^\circ$  when the alkyl chain lengths were even, whereas they exhibited a range of  $2.45\text{--}4.45^\circ$  when the alkyl chain lengths were odd. Further investigation of this observable even–odd effect is imperative.

### Optical properties

UV-Vis absorption spectrophotometry was performed by dissolving the sample in methanol to examine the interaction between  $Cn^{2+}$  and  $CMI^-$  in a polar solvent (Fig S21†). The analysis of  $Na^+ \cdot CMI^-$  revealed two distinctive absorption bands. In the long-wavelength region, the initial absorption was observed at 440 nm, and the cut-off wavelength was determined to be approximately 670 nm in the visible range. This complicated absorption band with a maximum at 582 nm indicates that  $CMI^-$  exists as the anion form without protonation in methanol. In contrast, the absorption spectra of  $Cn^{2+} \cdot 2CMI^-$  in methanol exhibited a similar absorption band with a maximum at 582 nm, indicating a substantial two-fold enhancement in its molar absorption coefficient. This result indicates the absence of an interaction between  $CMI^-$  and  $Cn^{2+}$  in methanol.

Diffuse reflectance spectra were measured using powder samples to reveal the interactions between ion pairs (Fig. S22†). In  $Cn^{2+} \cdot 2CMI^-$  ( $n = 1\text{--}4$ ), a significant redshift of the absorption edge compared to that of  $Na^+ \cdot CMI^-$  was observed (Fig. 7a). The presence of an absorption edge in the near-infrared (NIR) region suggests the formation of CT complexes with  $Cn^{2+}$  as the

acceptor and  $CMI^-$  as the donor. In contrast, in  $Cn^{2+} \cdot 2CMI^-$  ( $n = 5\text{--}8$ ), only a small redshift of the absorption edge compared to that of  $Na^+ \cdot CMI^-$  was observed (Fig. 7b). These results correlated with the packing arrangement of the crystal structures. In particular,  $Cn^{2+} \cdot 2CMI^-$  ( $n = 1\text{--}4$ ), which forms a columnar structure, demonstrated a remarkable CT absorption band because of the formation of a proficient stacking structure of anions and cations. In contrast, CT absorption was hardly observed in the slipped columnar structure with limited overlap between the anions and cations, as well as in the independent packing structures of the anion–cation–anion arrangements. Thus, the packing structure and the associated optical characteristics can be regulated by manipulating the length of the alkyl chain.

The optical energy bandgap ( $E_g$ ) of optical functional materials is a critical parameter for producing optoelectronic application of these materials. Therefore, the optical bandgaps of the ionic crystals were calculated using the Tauc plots (Table 1). Based on the data of diffuse reflection spectra of powder sample, the Tauc plots were prepared by the following eqn (1).

$$(\alpha h\nu)^m = A(h\nu - E_g) \quad (1)$$

$\alpha$  is UV-Vis absorption coefficient ( $=2.303 \log A/t$ ,  $t$ : thickness: 1),  $h$  is Planck constant,  $\nu$  is frequency,  $m$  is direct transitions (using for direct allowed transition,  $m = 1/2$ ),  $A$  is absorbance and  $E_g$  is optical energy bandgap.  $(\alpha h\nu)^{1/2}$  was plotted as a function of the  $h\nu$ , and  $h\nu$  is also photon energy, can calculate as  $1240/\lambda$ . Calculated  $E_g$  by the Tauc plots of  $Cn^{2+} \cdot 2CMI^-$  were consistent with the photon energy at the absorption edge. Among  $Cn^{2+} \cdot 2CMI^-$ , the order of the bandgap from the regression curve of the Tauc plot is as follow;  $C1^{2+} \cdot 2CMI^- < C2^{2+} \cdot 2CMI^- < C6^{2+} \cdot 2CMI^- < C4^{2+} \cdot 2CMI^- < C3^{2+} \cdot 2CMI^- < C7^{2+} \cdot 2CMI^- < C8^{2+} \cdot 2CMI^- < C5^{2+} \cdot 2CMI^-$ . The bandgap of  $C6^{2+} \cdot 2CMI^-$  was estimated to be low due to the presence of an extended absorption tail in the NIR range. However, the bandgap obtained from the Tauc plots clearly indicates a distinct boundary between  $C4^{2+} \cdot 2CMI^-$  and  $C5^{2+} \cdot 2CMI^-$ .  $Cn^{2+} \cdot 2CMI^-$  ( $n = 1\text{--}4$ ) exhibit low bandgap values, approximately ranging from 1.30 to 1.45, suggesting the formation of a CT band. Thus, the presence of an elaborate columnar structure, consisting of the donor and acceptor, promotes the formation

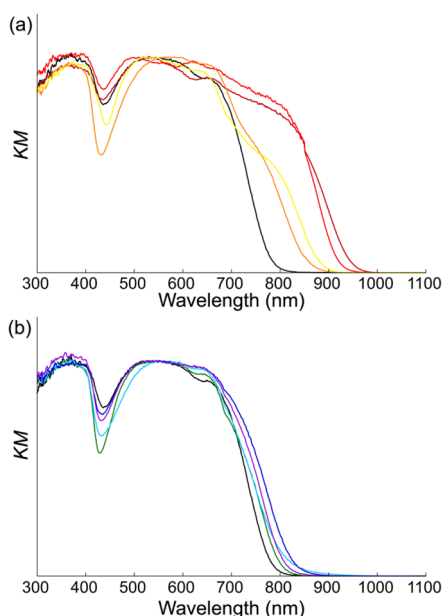


Fig. 7 Comparison of diffuse reflectance spectra of  $Cn^{2+} \cdot 2CMI^-$  (a)  $n = 1\text{--}4$  and (b)  $n = 5\text{--}8$  with  $Na^+ \cdot CMI^-$ . Spectra line colour code: black, dark red, red, orange, yellow, green, light blue, blue and purple refer to  $Na^+ \cdot CMI^-$ ,  $C1^{2+} \cdot 2CMI^-$ ,  $C2^{2+} \cdot 2CMI^-$ ,  $C3^{2+} \cdot 2CMI^-$ ,  $C4^{2+} \cdot 2CMI^-$ ,  $C5^{2+} \cdot 2CMI^-$ ,  $C6^{2+} \cdot 2CMI^-$ ,  $C7^{2+} \cdot 2CMI^-$  and  $C8^{2+} \cdot 2CMI^-$ , respectively.

Table 1 The optical bandgap ( $E_g$ ) of  $Cn^{2+} \cdot 2CMI^-$  calculation by the Tauc plots

Compound	Bandgap $E_g$ [eV]	Wavelength of absorption edge [nm]
$C1^{2+} \cdot 2CMI^-$	1.27	974
$C2^{2+} \cdot 2CMI^-$	1.30	957
$C3^{2+} \cdot 2CMI^-$	1.37	904
$C4^{2+} \cdot 2CMI^-$	1.35	916
$C5^{2+} \cdot 2CMI^-$	1.47	842
$C6^{2+} \cdot 2CMI^-$	1.34	927
$C7^{2+} \cdot 2CMI^-$	1.41	878
$C8^{2+} \cdot 2CMI^-$	1.44	864





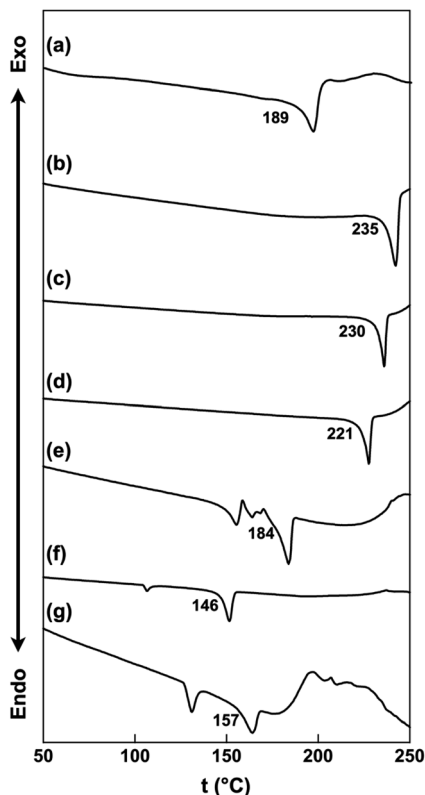


Fig. 8 DSC thermograms of (a)  $C2^{2+}-2CMI^-$ , (b)  $C3^{2+}-2CMI^-$ , (c)  $C4^{2+}-2CMI^-$ , (d)  $C5^{2+}-2CMI^-$ , (e)  $C6^{2+}-2CMI^-$ , (f)  $C7^{2+}-2CMI^-$  and (g)  $C8^{2+}-2CMI^-$ . Onset temperatures ( $^{\circ}C$ ) of melting points are labeled.

of the CT band. On the other hand, for  $Cn^{2+}-2CMI^-$  ( $n = 5-8$ , excepted for  $C6^{2+}-2CMI^-$ ), larger bandgap values were calculated compared to  $Cn^{2+}-2CMI^-$  ( $n = 1-4$ ), approximately ranging from 1.47 to 1.54. In  $C6^{2+}-2CMI^-$ , the transition derived from CT was observed even without the columnar structure, which was probably due to the small angle between the long axes of the cation and anion resulting in the substantial cation–anion overlap. However, it is noteworthy that the absorption derived from the CT transition was noticeably smaller compared to those of  $Cn^{2+}-2CMI^-$  ( $n = 1-4$ ), which form columnar structures.

### Thermal properties

The thermal properties of the ion pairs in their bulk states were examined using differential scanning calorimetry (DSC). The melting points were observed at 189, 235, 230, 221, 184, 146 and  $157^{\circ}C$  for  $C2^{2+}-2CMI^-$ ,  $C3^{2+}-2CMI^-$ ,  $C4^{2+}-2CMI^-$ ,  $C5^{2+}-2CMI^-$ ,  $C6^{2+}-2CMI^-$ ,  $C7^{2+}-2CMI^-$  and  $C8^{2+}-2CMI^-$ , respectively, and  $C1^{2+}-2CMI^-$  decomposed under its melting point (Fig. 8 and S23<sup>†</sup>). The melting point order was approximately dependent on the length of the alkyl chain.

## Conclusions

Ion pairs consisting of  $CMI^-$  and viologen dication with alkyl chains ranging from methyl to octyl groups were prepared *via*

ion exchange using the sodium salt of  $CMI^-$  as well as the chloride, bromide, and iodide salts of the cations. Seven ionic crystals were obtained, and their assembled structures were revealed *via* single-crystal X-ray analysis. In all the ionic crystals, anion–cation–anion-type stacking formations were observed. The assembled structures were classified into three types: “columnar”, “slipped columnar”, and “independent”, as shown in Fig. 1. The  $Cn^{2+}-2CMI^-$  crystals ( $n = 1-4$ ) were sorted into columnar structures, whereas the  $C5^{2+}-2CMI^-$  crystal was sorted into slipped columnar structures. The crystals of  $Cn^{2+}-2CMI^-$  ( $n = 6$  and  $8$ ), including the long alkyl chains, were sorted into independent structures. These results indicate a correlation between the crystal structure and the optical properties. Ion pairs forming columnar and slipped columnar structures exhibited absorption bands derived from the CT transition, whereas the CT absorption band was not observed for the independent structure. Therefore, we successfully generated numerous assemblies by altering only the alkyl chain length while keeping the primary framework of the  $\pi$ -electron system unchanged. Moreover, we elucidated the correlation between the arrangement of anions and cations and the development of charge-transfer complexes. Systematic studies on the crystal structures of organic ionic species are just beginning to emerge, and further detailed analyses are ongoing and worth investigating.

## Author contributions

E. Saito, T. Yoshida and R. Yamakado conceived and designed the experiments. T. Yasuhara and H. Yamaguchi conducted preliminary examinations. S. Okada assisted the measurement for the optical properties. E. Saito and R. Yamakado wrote the manuscript. All the authors have read and commented on the manuscript.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

A part of this work was supported by Program for Advancing Strategic International Networks to Accelerate the Circulation of Talented Researchers, “Advanced Next Generation Energy Leadership (R2601, FY2014-2016)” JSPS KAKENHI Grant-in-Aid for Scientific Research (B) (JP 18H02068), JSPS Fellows (JP 21J22405 and 22KJ0332), Japan-Austria Bilateral Joint Research Project and Japan-Ukraine Bilateral Joint Research Project from Japan Society for the Promotion of Science (JSPS). The synchrotron radiation experiments were performed at the BL02B1 (2022B1944) of SPring-8 with the approval of the Japan Synchrotron Radiation Research Institute (JASRI) with the supports by Dr Y. Nakamura, JASRI/SPring-8 and Dr A. Matsunaga, Yamagata University. The NMR analysis was performed using a JEOL JNM-EC500, JEOL Ltd introduced by the subsidy program for development of advanced research infrastructure.



We thank Mr Junya Kimura Yamagata University, for preliminary examinations.

## Notes and references

- (a) Y. Huang, Z. Wang, Z. Chen and Q. Zhang, *Angew. Chem., Int. Ed.*, 2019, **58**, 9696–9711; (b) M. Jiang, C. Zhen, S. Li, X. Zhang and W. Hu, *Front. Chem.*, 2021, **9**, 764628.
- D. Jerome, *Chem. Rev.*, 2004, **104**, 5565–5592.
- (a) F. Kagawa, S. Horiuchi, H. Matsui, R. Kumai, Y. Onose, T. Hasegawa and Y. Tokura, *Phys. Rev. Lett.*, 2010, **104**, 227602; (b) A. S. Tayi, A. K. Shveyd, A. C. Sue, J. M. Szarko, B. S. Rolczynski, D. Cao, T. J. Kennedy, A. A. Sarjeant, C. L. Stern, W. F. Paxton, W. Wu, S. K. Dey, A. C. Fahrenbach, J. R. Guest, H. Mohseni, L. X. Chen, K. L. Wang, J. F. Stoddart and S. I. Stupp, *Nature*, 2012, **488**, 485–489.
- H. Alves, A. S. Molinari, H. Xie and A. F. Morpurgo, *Nat. Mater.*, 2008, **7**, 574–580.
- A. Das and S. Ghosh, *Angew Chem. Int. Ed. Engl.*, 2014, **53**, 2038–2054.
- (a) A. L. Sutton, B. F. Abrahams, D. M. D'Alessandro, R. W. Elliott, T. A. Hudson, R. Robson and P. M. Usov, *CrystEngComm*, 2014, **16**, 5234–5243; (b) A. L. Sutton, B. F. Abrahams, D. M. D'Alessandro, T. A. Hudson, R. Robson and P. M. Usov, *CrystEngComm*, 2016, **18**, 8906–8914; (c) K. Sambé, N. Hoshino, T. Takeda, T. Nakamura and T. Akutagawa, *Cryst. Growth Des.*, 2020, **20**, 3625–3634.
- (a) E. L. Goff and R. B. LaCount, *J. Am. Chem. Soc.*, 1963, **85**, 1354–1355; (b) O. W. Webster, *J. Am. Chem. Soc.*, 1965, **87**, 1820–1821; (c) C. Richardson and C. A. Reed, *Chem. Commun.*, 2004, 706–707; (d) R. J. Less, M. McPartlin, J. M. Rawson, P. T. Wood and D. S. Wright, *Chem. –Eur. J.*, 2010, **16**, 13723–13728; (e) Y. Bando, Y. Haketa, T. Sakurai, W. Matsuda, S. Seki, H. Takaya and H. Maeda, *Chem. –Eur. J.*, 2016, **22**, 7843–7850.
- (a) R. Kuhn and D. Rewicki, *Angew. Chem., Int. Ed.*, 1967, **6**, 635–636; (b) K. Okamoto, T. Kitagawa, K. Takeuchi, K. Komatsu, T. Kinoshita, S. Aonuma, M. Nagai and A. Miyabo, *J. Org. Chem.*, 1990, **55**, 996–1002.
- J.-B. Ming, X. Li, J. Li, W. Yu and W. Wang, *Cryst. Growth Des.*, 2022, **22**, 1212–1220.
- (a) Y. Haketa and H. Maeda, *Chem. Commun.*, 2017, **53**, 2894–2909; (b) Y. Haketa and H. Maeda, *Bull. Chem. Soc. Jpn.*, 2018, **91**, 420–436; (c) Y. Haketa, K. Urakawa and H. Maeda, *Mol. Syst. Des. Eng.*, 2020, **5**, 757–771.
- (a) Y. Son and S. Kim, *Dyes Pigm.*, 2005, **64**, 153–155; (b) J. Fujisawa, *Chem. Phys. Lett.*, 2014, **608**, 355–359.
- Y. Tanaka, K. Ichijo, S. Kodama, S. Aoyama, T. Yoshida, R. Yamakado and S. Okada, *Cryst. Growth Des.*, 2019, **19**, 5811–5818.
- (a) R. J. Mortimer, *Chem. Soc. Rev.*, 1997, **26**, 147–156; (b) H. Chen, V. Brasiliense, J. Mo, L. Zhang, Y. Jiao, Z. Chen, L. O. Jones, G. He, Q. H. Guo, X. Y. Chen, B. Song, G. C. Schatz and J. F. Stoddart, *J. Am. Chem. Soc.*, 2021, **143**, 2886–2895.
- P. Huo, Y.-H. Li, L.-J. Xue, T. Chen, L. Yu, Q.-Y. Zhu and J. Dai, *CrystEngComm*, 2016, **18**, 1904–1910.
- E. Saito, T. Yasuhara, Y. Tanaka, R. Yamakado, S. Okada, T. Nohara, A. Masuhara and T. Yoshida, *ECS Trans.*, 2019, **88**, 301–311.

