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## 1. Introduction

Nowadays, humidity sensors are indispensable informationgathering equipment in the age of artificial intelligence, which plays an important role in the field of detecting ambient humidity, controlling various industrial production and monitoring the human respiratory rate.<sup>1-3</sup> To meet the needs of high performance such as linearity, fast response/recovery ability, low hysteresis and pronounced sensitivity, a wide range of humidity sensors, including capacitance, impedance, frequency, and other types, have been developed using sensing materials such as ceramics, polymers, and carbon-based materials.4,5 The utilization of nanoparticles from metal oxide semiconductors (SnO<sub>2</sub>, Al<sub>2</sub>O<sub>3</sub>, TiO<sub>2</sub>, etc.) as building humidity sensing materials is advantageous due to mechanical flexibility and excellent electrical/sensing properties.<sup>6-8</sup>

Tin oxide  $(SnO_2)$ , a kind of wide band-gap (3.6 eV) n-type semiconductor which has the advantages of high chemical stability, simple manufacturing and a favorable adsorption

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# Investigation on acceptor-donor co-doped SnO<sub>2</sub> nanoparticles enriched with oxygen vacancies: a capacitive humidity sensor for respiration detection

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In this work, we develop a novel capacitive humidity sensor based on Al-Si acceptor-donor co-doped SnO<sub>2</sub> for real-time monitoring of ambient humidity and human respiration. XRD measurements reveal that all samples exhibit a tetragonal rutile phase and the crystallite size of SnO<sub>2</sub> decreases with increasing Al-Si content. The high intensity of the Raman peak at 762 cm<sup>-1</sup> confirms the presence of bridging mode oxygen vacancies in (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub>. The EPR results show that the amount of singly ionized oxygen vacancies increases after the introduction of Al-Si. Both types and amounts of oxygen vacancy defects are particularly sensitive to the adsorption of water molecules. Moreover, according to DFT calculations, the contribution of the Si 3s orbital and Al 3s orbital to the band edge verifies the formation of acceptor-donor complexes in Al-Si co-doped SnO<sub>2</sub>. The humidity sensing results reveal that the  $(Al + Si)_{0.02}Sn_{0.98}O_2$  humidity sensor shows high sensitivity (S = 839), low hysteresis (1.94%) and fast response/recovery times (25 s/5 s). The respiratory intervals during shallow, medium and deep breathing states of (Al + Si)0.02Sn0.98O2 were measured at 2.8 s, 3.8 s and 4.5 s, respectively. The chemical mechanism for the enhancement of humidity sensing performance corresponding to the oxygen vacancy defects induced by Al-Si interplay is proposed.

> capacity for water molecules, remains the focus of attention in the outline and advancement of high sensitivity humidity sensors.<sup>9</sup> However, the intrinsic SnO<sub>2</sub>-based humidity sensor suffers from long response and recovery times, as well as high hysteresis. In order to further enhance the sensing performance, such as the linearity and sensitivity, a number of elements including Zn,<sup>10</sup> Cd,<sup>11</sup> Al,<sup>12</sup> Sb,<sup>13</sup> N<sup>14</sup> and P<sup>14</sup> have been selected as dopants in SnO2-based sensors. Among these doping elements, Al and Si have an edge on being resourceful, inexpensive, non-toxic and environmentally friendly. In addition, the radii of  $Al^{3+}$  (0.51 Å) and  $Si^{4+}$  (0.4 Å) are considerably smaller than Sn<sup>4+</sup> (0.69 Å), which can lead to lattice distortion and trigger the formation of substitutional or interstitial defects. Generally, due to the fact that the Al acceptor has one less valence electron compared to the Si donor, the crystalline structure inherently introduces an electronic hole in the SiO<sub>2</sub> energy band.<sup>15,16</sup> The interplay of Al and Si ions can generate a significant number of positive charges and create a local electric field within the crystal, thereby facilitating the efficient separation of electron-hole pairs. Previous studies have reported that the incorporation of an Al acceptor and a Si donor into metal oxides is responsible for increasing the oxygen vacancy and electron content.<sup>17,18</sup> Si behaves as a wellbehaved shallow donor within the range of conducting oxides,





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while Al acts as an acceptor to facilitate p-type conductivity. Both the high activity of the oxygen vacancy and the significant contribution of electrons are considered to be responsible for enhancing sensing performance. The oxygen vacancy defects provide additional active surface adsorption sites for water molecules on the SnO<sub>2</sub> surface, which plays a crucial role in promoting the humidity sensitivity of the material.<sup>19</sup> Therefore, the incorporation of Al–Si into SnO<sub>2</sub> can enhance the humidity sensitivity by increasing the number of oxygen vacancy defects.

In this study, we present a successful fabrication of a novel capacitive humidity sensor based on Al–Si acceptor–donor codoped  $SnO_2$  nanoparticles, enriched with oxygen vacancy defects. To explore the interactions between the Al–Si acceptor–donor complex and water molecules, as well as understand their impact on the sensing properties, the oxygen vacancy defects and the electronic structure were investigated by EPR and DFT. The capacitive humidity sensor based on the acceptor–donor complex demonstrates the potential for simultaneous detection of ambient humidity and human breath rate.

## 2. Experimental

#### 2.1 Synthesis of Al-Si co-doped SnO<sub>2</sub>

All the starting materials used in the experiment were of analytical grade and used directly without further purification. Fig. 1 depicts the synthesis of pure and Al–Si co-doped SnO<sub>2</sub> nanoparticles *via* a simple sol–gel method. In a typical procedure, 0.05 mol SnCl<sub>2</sub>·2H<sub>2</sub>O was dissolved in the solution mixed with 100 mL ethanol and 100 mL deionized water. Thereafter,  $(C_2H_5O)_4Si$  and Al $(NO_3)_3$ ·9H<sub>2</sub>O were incorporated into Sn precursor solutions to prepare (Al + Si)<sub>x</sub>Sn<sub>1-x</sub>O<sub>2</sub> (x = 0, 0.5, 1, 1.5, 2 mol%) and then stirred at 300 K for 6 h on a magnetic stirrer.  $(C_2H_5O)_4Si$  and Al $(NO_3)_3$ ·9H<sub>2</sub>O were added into the solutions in

a 1:1 ratio to ensure the doping concentration of Al and Si. Then, a certain amount of ammonia water was dropped into the resultant solutions to regulate the pH to 8, until white turbid precipitates appeared at the bottom of solution. The reaction products were aged, filtered, and washed 3 times before drying for 1 hour in an oven at 100 °C. Moreover, the prepared samples were kept in a furnace and annealed at 500 °C for 60 min. Finally, Al–Si co-doped SnO<sub>2</sub> nanoparticles with different dopant contents were synthesized, respectively.

#### 2.2 Characterization and calculation methods

The lattice structure of the SnO<sub>2</sub> sample was determined using an X-ray diffractometer (Ultima IV, Rigaku, Japan). The test range of  $2\theta$  was  $20^{\circ}$ – $80^{\circ}$ . The  $\lambda$  was 0.15406 nm and the scan rate was  $5^{\circ}$  min<sup>-1</sup>. The morphology and structure of the prepared sample were observed by employing the SUPRA55 field emission scanning electron microscope, which was operated at 15 kV. Raman spectroscopy measurements were performed using a LabRAM UV-VIS-NIR system with an excitation wavelength of  $\lambda$  = 590 nm. The investigation of oxygen vacancies in the samples was carried out using electro paramagnetic resonance (EPR) spectroscopy with a Bruker A300-10/12 spectrometer operating at a field modulation frequency of 100 kHz under ambient conditions. By employing the density functional theory (DFT) module offered by Q-flow (Lonxun Quantum), we performed electronic structure calculations of SnO2 within the DFT framework. The exchange-correlation functional was modeled using the Perdew-Burke-Ernzerhof potential within the GGA+U scheme. To conduct the calculations, we employed energy cutoffs of 120 Ry for the plane-wave expansion of the wave functions and the charge density. For Brillouin zone sampling, we utilized a  $2 \times 2 \times 6$  Monkhorst–Pack k-point grid. In order to account for the modeling of different impurities in oxides, we considered



Fig. 1 Preparation of the pure and Al-Si co-doped SnO<sub>2</sub> and procedure for testing humidity sensing performance.

not only substitution (S) impurities but also interstitial (I) impurities, as well as their combinations with oxygen vacancies.

# 2.3 Humidity sensing performance testing and respiratory detection

The testing procedure for the pristine and Al-Si co-doped SnO<sub>2</sub> humidity sensors is shown in Fig. 2. In the testing procedures, a polyimide substrate (10 mm  $\times$  10 mm) with 15 pairs of interdigital electrodes was used to measure the humidity sensing performance. Then pure and Al-Si co-doped SnO<sub>2</sub> nanoparticles (2 mg) were mixed with ethanol (20 µL) to form a paste and coated on the surface of the electrode. Then, the electrodes were placed in the oven with a drying temperature of 70  $^{\circ}$ C for 10 min. After that, all prepared electrodes were sequentially placed in different air tight glass humidity chambers, which were added with saturated aqueous solution of LiCl, MgCl<sub>2</sub>, K<sub>2</sub>CO<sub>3</sub>, NaCl, KCl and KNO<sub>3</sub>. The above saturated salt solutions have 11%, 34%, 43%, 75%, 84% and 95% RH, respectively. The corresponding variations in the capacitance of sensing materials were recorded using an LCR digital instrument at room temperature with a working voltage of 1 V and a frequency range of 120 Hz-10 kHz. Fig. 2 shows the circuit and the physical diagram of the humidity sensor designed to enable respiratory monitoring across various states of human respiration. When the humidity sensor is exposed to different relative humidity (RH) levels, significant variations in the brightness of the LED can be observed. When a person breathes onto the humidity sensor, the presence of moisture influences the capacitance of the sensor, leading to an increase in the current within the circuit. Consequently, the LED reaches the breakover voltage and turns on.

## 3. Results and discussion

#### 3.1 Structural and morphological properties

As shown in Fig. 3(a), the crystalline structure and phase purity of pure and Al–Si co-doped  $\text{SnO}_2$  were determined by X-ray diffraction patterns. All the diffraction peaks observed are in agreement with the standard tetragonal rutile  $\text{SnO}_2$  phase (JCPDS#41-1445). The strong diffraction peaks located at  $2\theta$  of 26.6°, 33.9° and 51.8° are respectively assigned to (110), (101) and (211) planes. There is no phase related to Al–Si oxides in the diffraction pattern, which confirms that the Al–Si ions are



Fig. 2 Circuit diagram for the humidity sensor system for monitoring breathing (LED) and ambient humidity (LCR).

homogeneously incorporated into  $\text{SnO}_2$  lattice. The intensity of diffraction peak (101) exhibits a sharp increase with increasing Al–Si content, which shows the reduction of crystallinity for Al–Si co-doped samples.<sup>20</sup> This evidence suggests the substitution of dopants ions for Sn ions in the  $\text{SnO}_2$  lattice.

We calculated the crystallite size *D* and micro strain  $\varepsilon$  using the following relations:<sup>18</sup>

$$D = \frac{K\lambda}{\beta\cos\theta} \tag{1}$$

$$\varepsilon = \frac{\beta \cos \theta}{4} \tag{2}$$

where  $\lambda$  is the wavelength of X-ray radiation and  $\beta$  is the full width at half maximum (FWHM) of the peaks at the diffracting angle  $\theta$ . From Fig. 3(b), it is evident that the crystallite size decreases from 44 to 25 nm with an increase in the doping rate. Simultaneously, the microstrain increases from  $4.5 \times 10^{-2}$  to  $8.0 \times 10^{-2}$ . This can be interpreted by the fact that the smaller Al<sup>3+</sup>/Si<sup>4+</sup> ions can induce the interstitial or substitutional defects in the SnO<sub>2</sub> lattice.<sup>21</sup> During the annealing process of Al–Si co-doped SnO<sub>2</sub> nanoparticles, the smaller ionic radii of Al<sup>3+</sup> and Si<sup>4+</sup> ions allow them to occupy the octahedral sites leads to an increase in compression stresses and lattice distortion. Therefore, the crystalline structural properties of SnO<sub>2</sub> have been tuned by incorporating the Al–Si with different concentrations.

Fig. 4(a)–(f) presents the FESEM images and elemental mapping of the  $(Al + Si)_{0.02}Sn_{0.98}O_2$  nanoparticles. The surface morphology of the nanoparticles shows the spherical grain with a size ranging from 21.5 to 39.2 nm. The elemental compositions of  $(Al + Si)_{0.02}Sn_{0.98}O_2$  nanoparticles are studied by elemental mapping. The results show that the materials are composed of Al, Si, Sn and O, and the four elements are distributed homogeneously in  $(Al + Si)_{0.02}Sn_{0.98}O_2$ , which further confirms the introduction of dopants.

#### 3.2 Oxygen vacancy defect analysis

Raman spectroscopy was used to investigate the structural defects and crystalline disorder induced by dopant ions in the SnO<sub>2</sub> lattice. Fig. 5 shows the Raman spectra of (Al +  $Si_rSn_{1-r}O_2$  (x = 0, 0.01 and 0.02) samples measured in the range of 200-900 nm and the corresponding plots of these Raman peak fitting by the Gaussian function. The typical lattice vibration mode at the  $\Gamma$  point of the Brillouin zone of SnO<sub>2</sub> are described as  $A_{1g} + A_{2g} + A_{2u} + B_{1g} + B_{2g} + 2B_{1u} + E_g + 3E_u$ .<sup>22</sup> Among these modes  $A_{\rm 1g},~B_{\rm 1g},~B_{\rm 2g}$  and  $E_{\rm g}$  are Raman active modes,  $A_{2u}$  and  $E_u$  are IR active modes, while  $A_{2g}$  and  $B_{1u}$  are inactive in both IR and Raman regions. Fig. 5(b) shows the B<sub>1u</sub>  $(238 \text{ cm}^{-1})$ ,  $E_g (503 \text{ cm}^{-1})$ ,  $A_{1g} (627 \text{ cm}^{-1})$ ,  $A_{2u} (688 \text{ cm}^{-1})$  and  $B_{2g}$  (762 cm<sup>-1</sup>) bands in the pure SnO<sub>2</sub>. But in (Al +  $Si)_{0.01}Sn_{0.99}O_2$  and  $(Al + Si)_{0.02}Sn_{0.98}O_2$ , the  $A_{1g}$  band exhibits a shift towards the higher wavenumber side, which is due to the incorporation of Al-Si ions into the SnO<sub>2</sub> lattice, as shown in Fig. 5(c) and (d). This observed blue shift of the  $A_{1g}$  modes from



Fig. 3 (a) X-ray diffraction patterns of pure and Al–Si co-doped SnO<sub>2</sub>. (b) Plots of crystallite size and microstrain versus doping concentrations of (Al + Si)<sub>x</sub>Sn<sub>(1-xi</sub>O<sub>2</sub>.



Fig. 4 (a) and (b) FESEM images of  $(Al + Si)_{0.02}Sn_{0.98}O_2$  nanoparticles and (c)–(f) (Al + Si) elemental mapping images.

627 cm<sup>-1</sup> to 634 cm<sup>-1</sup> can be attributed to the substitution of Sn<sup>4+</sup> with smaller Al<sup>3+</sup> and Si<sup>4+</sup> ions, as well as the decreased crystallite size resulting from phonon confinement.<sup>23</sup> In Fig. 5(d), the Raman peak at 507 cm<sup>-1</sup> in (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub> is observed to shift to a lower wavenumber compared to pure SnO<sub>2</sub> due to the rearrangement of the internal defect equilibrium. Moreover, in the Raman modes of (Al + Si)<sub>0.01</sub>Sn<sub>0.99</sub>O<sub>2</sub> and (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub>, here a small peak appearing at 324 cm<sup>-1</sup> and 363 cm<sup>-1</sup> represents E<sub>u</sub> transverse optical (TO) vibration and that in pure SnO<sub>2</sub> was absent.<sup>24</sup> This is due to

interstitial Si (Si<sub>i</sub>) in the tin oxide lattice, and the high disordering and low crystallinity of the Al–Si co-doped  $SnO_2$  samples. In the Al–Si addition the following equations may apply.<sup>25</sup>

$$Al_2O_3 \xrightarrow{SnO_2} 2Al'_{Sn} + 3O_O + V_O^{\bullet\bullet}$$
(3)

$$\operatorname{SiO}_2 \xrightarrow{\operatorname{SnO}_2} \operatorname{Si}_{\operatorname{Sn}} + 2O_O$$
 (4)

$$Al_2O_3 \xrightarrow{SiO_2} 2Al'_{Si} + 3O_O + V_O^{\bullet\bullet}$$
(5)

$$O_{O} \rightarrow V_{O}^{\bullet \bullet} + 2e^{-} + \frac{1}{2}O_{2} \tag{6}$$

In the case of substituting Al<sup>3+</sup> ions for Sn<sup>4+</sup> and Si<sup>4+</sup> ions, more oxygen vacancy defects  $V_{\Omega}^{\bullet\bullet}$  are formed in the  $(Al + Si)_{0.02}$ - $Sn_{0.98}O_2$  lattice. Fig. 6 shows the distribution of different kinds of oxygen vacancies like bridging, sub-bridging or in-plane oxygen vacancy of SnO<sub>2</sub> on the surface layer. For different types of oxygen vacancy defects, the Raman mode shows different positions and intensities due to the external strain on the rutile structure. Generally, the peak of  $SnO_2$  around 762 cm<sup>-1</sup> is attributed to the presence of bridging or sub-bridging mode of vacancy, while the peak at around 688  $\text{cm}^{-1}$  of the A<sub>2u</sub> mode is corresponding to the presence of in-plane oxygen vacancy.<sup>24</sup> The bridging and sub-bridging modes of oxygen vacancies are sensitive to the adsorption of the water molecules and the conduction of charge carriers.<sup>26</sup> As shown in Fig. 5(b) and (c), the Raman peak at 762 cm<sup>-1</sup> of pure SnO<sub>2</sub> and (Al + Si)<sub>0.02</sub>-Sn<sub>0.98</sub>O<sub>2</sub> indicates the presence of a similar concentration of bridging or sub-bridging oxygen vacancies in both materials. However, in Fig. 5(d), there is a notable increase in the intensity of the Raman peak around 762  $\text{cm}^{-1}$  of  $(\text{Al} + \text{Si})_{0.02}\text{Sn}_{0.98}\text{O}_2$ , which suggests that the bridging or sub-bridging oxygen vacancy content is enhanced significantly. Hence, the incorporation of Al<sup>3+</sup> and Si<sup>4+</sup> into the SnO<sub>2</sub> lattice can potentially increase the presence of bridging oxygen vacancies and enhance the adsorption ability towards water molecules. Besides, other fitting small peaks including the  $E_g$  mode at 503 cm<sup>-1</sup> in the Raman spectra may imply the presence of



Fig. 5 Raman spectra of pure and Al-Si co-doped SnO<sub>2</sub> (a) general survey spectrum, (b)-(d) deconvolution of the Raman spectrum.



Fig. 6 The illustration of different modes of oxygen atoms on the  $SnO_2$  surface.

defects like tin interstitials  $(Sn_i)$  or substitutional defects  $(Al_{Si})$ . It can be concluded that the structural defects have been introduced by doping Al–Si ions in the SnO<sub>2</sub> lattice structure.

Electron paramagnetic resonance (EPR) is a highly sensitive technique used to investigate the spin dynamics of paramagnetic centers, including those originating from Al and Si ion incorporated into SnO<sub>2</sub> and oxygen vacancies defects. The EPR patterns of pure SnO<sub>2</sub> and  $(AI + Si)_{0.02}Sn_{0.98}O_2$  are shown in Fig. 7(a). The *g* factor (EPR parameter) can be calculated to determine the type of defects in the SnO<sub>2</sub> particles. The following equation is commonly used to calculate the *g* factor:<sup>27</sup>

$$g = \frac{h\nu}{\beta H} \tag{7}$$

where  $h = 4.135 \times 10^{-15}$  eVs is the Planck constant, v = 9.48 GHz is the microwave frequency (Hz),  $\beta = 5.788 \times 10^{-5}$  eV T<sup>-1</sup> is the electron Bohr magnetron and H is the applied magnetic field. The strong symmetrical EPR signal observed at 3514 G (g = 2.003) in pure SnO<sub>2</sub> and (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub> can be ascribed to unpaired electrons originating from the presence of oxygen vacancies. As demonstrated in Fig. 7(b), the concentration of paramagnetic centers is the primary factor influencing the EPR signal intensity and the intensity of the EPR signal increases with the doping of Al and Si. This can be attributed to the increase in the number of oxygen vacancies in (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub>. The number of the spins corresponding to oxygen vacancies was estimated by using the following formula<sup>28</sup>

$$N_{\rm spins} = 0.285 I (\Delta H)^2 \tag{8}$$

where *I* represents the peak-to-peak height of the EPR signal (in arbitrary units), and  $\Delta H$  represents the line width of the signal (in mT). The calculated  $N_{\text{spins}}$  of pure SnO<sub>2</sub> and



Fig. 7 (a) and (b) EPR spectrum and intensity comparison of pure SnO<sub>2</sub> and (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub>.

 $(Al + Si)_{0.02}Sn_{0.98}O_2$  are 86.5 and 132.9, respectively. As Al and Si are incorporated into SnO<sub>2</sub>, the calculated  $N_{spins}$  of  $(Al + Si)_{0.02}Sn_{0.98}O_2$  increase significantly, which is closely related to a 53.6% increase in oxygen vacancies.

#### 3.3 DFT calculations

In order to determine the correlation between the theoretical results and the experimental findings, we performed DFT modeling to investigate the electronic structure of Al–Si co-doped SnO<sub>2</sub>. Fig. 8(a)–(d) shows the schematic representation of the optimized atomic structure of the (100) surface of (a) pure SnO<sub>2</sub> (b) with a single substitutional Al and one oxygen vacancy (Al<sub>Sn</sub> + V<sub>O</sub>) (c) with one interstitial Si (Si<sub>i</sub>) and (d) substitutional Al, substitutional Si and one oxygen vacancy (Al<sub>Sn</sub> + Si<sub>Sn</sub> + V<sub>O</sub>). In Table 1, the presence of various combinations of substitutional and interstitial defects results in the



Fig. 8 The schematic representation of the optimized atomic structure of (a) pure SnO<sub>2</sub>, (b)  $AI_{Sn} + V_O$ , (c) Si<sub>i</sub> and (d)  $AI_{Sn} + Si_{Sn} + V_O$ .

variations in Fermi energy, band gap and lattice parameters. In the case of the pure  $\text{SnO}_2$ , the band gap (0.7046 eV) estimated from PBE was an underestimate compared to the experimental values due to its lack of consideration of the discontinuity in the exchange–correlation potential.<sup>29</sup> However, this shortcoming has no bearing on our study as we focus on the relative changes in the electronic structure when Al and Si are introduced into  $\text{SnO}_2$ .

To better understand the formation of oxygen vacancy defects upon the incorporation of Al-Si ions in SnO2, we investigate the DOS and PDOS corresponding to the pure and defective SnO<sub>2</sub>. As shown in Fig. 9(a), the density of partial wave states analysis indicates that the conduction band is primarily contributed by O 2p orbitals and the valence band is mainly composed of Sn 5s and Sn 5p orbitals. It is worth noting that the energy band gap of  $SnO_2$  (0.2344 eV) narrowed significantly for the contributions of Al 3s orbitals to the valence band (Fig. 9b). Therefore the introduction of Al into the system plays the key role in the formation of oxygen vacancies and enhances the carrier transfer between the energy band. Fig. 9(c) shows the contributions of Si 3s orbitals to the electronic structure where the band gap of Si doped  $SnO_2$  (1.3455 eV) has been broadened. This is due to the presence of excess electrons in the system. As seen in Fig. 9(d), the Sn atoms in the SnO<sub>2</sub> crystalline lattice were simultaneously substituted by an Al and a Si atom and oxygen vacancy was introduced, leading to a decrease of band gap to 1.0208 eV.

In general, Al defects and Si defects can act as acceptors and donors, respectively, and they interact with each other through complex mechanisms.<sup>30</sup> Li *et al.* reported that co-doping Mo

Table 1 Fermi energy, band gap and lattice parameters for pure and defective  $\mbox{SnO}_2$ 

	Fermi energy (eV)	Band gap (eV)	a (Å)	b (Å)	c (Å)
Pure SnO <sub>2</sub>	4.9084	0.7046	9.637	9.637	3.233
$Al_{sn} + V_{O}$	5.0725	0.2344	9.407	9.407	3.203
Sii	7.6627	1.3455	8.629	9.243	3.266
$Al_{Sn}+Si_{sub}+V_{O} \\$	5.9109	1.0208	9.241	9.209	3.140



Fig. 9 TDOS and PDOS for (a) pure SnO<sub>2</sub>, (b)  $AI_{Sn} + V_O$ , (c) Si<sub>i</sub> and (d)  $AI_{Sn} + Si_{Sn} + V_O$ .

and C can effectively shift the valence band edge while leaving the conduction band edge nearly unaffected.<sup>31</sup> As shown in Fig. 10, the substitution of Al<sup>3+</sup> for Sn<sup>4+</sup> ions was consistently accompanied by the introduction of oxygen vacancies, where Al<sup>3+</sup> ions played an acceptor role and the Si dopants state functions as temporary trap centers.<sup>32</sup> The Si 3s orbitals and Al 3s orbitals have important contributions to the valence band which indicates that the crystalline and electronic structure has been changed with Al-Si co-doping in the SnO<sub>2</sub>. In the case of isovalent doping, the introduction of Si into SnO<sub>2</sub> is favorable for generating additional electrons and forming stable donorlike defects (Si<sub>Sn</sub>) in the SnO<sub>2</sub> structure. The substitution of Sn<sup>4+</sup> ions with dopant ions induces the disruption of the original electronic structure, which subsequently leads to the rearrangement of the internal defect equilibrium. Therefore, the contribution of Si 3s orbital and Al 3s orbital to the band edge verifies the presence of acceptor-donor complexes in Al-Si codoped SnO<sub>2</sub>.

#### 3.4 Humidity sensing properties and respiration detection

The sensing properties can be associated with the change in the capacitance of humidity sensors exposed to humidity at different levels. Fig. 11(a) shows the capacitance variation of the pure and Al–Si co-doped  $SnO_2$  humidity sensors in different RH at a testing frequency of 1 kHz. It can be observed that the capacitance of all the humidity sensors increases with increasing



Fig. 10 A schematic band diagram with defect states.

humidity level, which means water molecules are gradually adsorbed on the  $SnO_2$  surface. The sensitivity of the humidity sensor can be calculated using the following formula:<sup>33</sup>

$$S = \frac{(C_1 - C_2)}{C_2}$$
(9)

where C1 and C2 represent the capacitance values at 95% RH



Fig. 11 (a) Capacitance relative humidity curves of pure and Al-Si co-doped  $SnO_2$  humidity sensors. (b) Capacitance frequency curves of (Al + Si)\_{0.02}Sn\_{0.98}O\_2 in different RH.

and 11% RH, respectively. The calculated sensitivity of pure and Al-Si co-doped SnO<sub>2</sub> humidity sensors corresponds to 27, 40, 38, 77 and 839, respectively. It can be found that the sensitivity of the humidity sensor is enhanced as the Al-Si codoping concentration increases. In these samples (Al +  $Si_{0.02}Sn_{0.98}O_2$  (S = 839) exhibits better linearity and sensing performance due to the 31 times sensitivity of the pure SnO<sub>2</sub> sensor (S = 27). This phenomenon can be attributed to more oxygen vacancies formed on the surface of the doped sample where the interplay of Al and Si ions play a significant role. Fig. 11(b) reveals the effect of different frequencies on the capacitance of  $(Al + Si)_{0.02}Sn_{0.98}O_2$  in the range from 11% RH to 95% RH. It can be observed that the capacitance of all the samples decreases as the frequency increases. This phenomenon occurs because the polarization speed of water molecules struggles to keep up with rapid changes in the electric field direction at high working frequencies. As a result, the capacitance decreases and its dependence on relative humidity (RH) diminishes. It can be clearly seen that the variations in the capacitance were clearly distinguishable at low frequencies while it became ambiguous at higher frequencies. The capacitance of the sensor increases with higher relative humidity (RH). In the case of  $(Al + Si)_{0.02}Sn_{0.98}O_2$ , an increased number of oxygen vacancies on the surface allows for more adsorption of water molecules. This leads to the formation of effective ionic conduction, ultimately resulting in an increase in capacitance. As verified above in EPR and Raman analysis, the presence of oxygen vacancies was associated with the Raman peaks at 688 cm<sup>-1</sup> and 762 cm<sup>-1</sup> and higher intensity of the EPR signal. Furthermore, the electronic structure has been tuned by the contribution of Si 3s orbitals and Al 3s orbitals to the valence band after Al–Si co-doping, which has noticeable implications on the humidity sensing properties. As a result, the capacitance of the humidity sensing response.

Humidity hysteresis is a key aspect to evaluate the dependability and stability of a humidity sensor. Fig. 12(a) and (b) reveals the humidity hysteresis for pure  $\text{SnO}_2$  and  $(\text{Al} + \text{Si})_{0.02}\text{Sn}_{0.98}\text{O}_2$  sensors in the adsorption and desorption cycle. During the desorption cycle, the capacitance of the humidity sensor is higher compared to the adsorption process. This is due to the delayed separation of water molecules from the  $\text{SnO}_2$ surface, resulting in a prolonged desorption process.<sup>34</sup> The equation for the calculated hysteresis of the humidity sensor



Fig. 12 Humidity hysteresis characteristics of (a) pure  $SnO_2$  and (b) (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub> humidity sensors.

is as follows:

$$H(\% RH) = \frac{C_{\rm A} - C_{\rm D}}{S} \tag{10}$$

where  $C_A$  and  $C_D$  are the capacitance values obtained at the same RH in the desorption and adsorption process, respectively, and *S* is the sensitivity of the sensor. The maximum humidity hysteresis for pure SnO<sub>2</sub> and (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub> are 4.47% and 1.94% RH, respectively. The (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub> sensor exhibits a significantly reduced hysteresis compared to pure SnO<sub>2</sub>, which indicates that the addition of Al<sup>3+</sup> and Si<sup>4+</sup> ions accelerates the desorption process on the SnO<sub>2</sub> surface. However, it is important to note that the fabricated SnO<sub>2</sub> humidity sensor may encounter challenges related to variations in environmental humidity and temperature, which can impact the response and decay time of the device, leading to less than optimal results.

Fig. 13 illustrates the response/recovery time of the humidity sensors based on pure and Al-Si co-doped SnO<sub>2</sub>, covering a humidity range from 11% RH to 95% RH. The response/ recovery time of the humidity sensor are defined as the duration required for sensor's capacitance variation to reach 90% of its total capacitance change. The  $(Al + Si)_{0.02}Sn_{0.98}O_2$  sensor exhibits a rapid response time of 25 s and a recovery time 5 s while the corresponding response/recovery times of (Al + Si)<sub>0.01</sub>Sn<sub>0.99</sub>O<sub>2</sub> and pure SnO<sub>2</sub> are 37 s/6 s and 54 s/14 s. The fast response/recovery times of (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub> are mainly due to the fact that the substitution of Al<sup>3+</sup> and Si<sup>4+</sup> into SnO<sub>2</sub> can enhance the adsorption activity of SnO<sub>2</sub>. It can be concluded that the recovery capability of the SnO<sub>2</sub> sensor depends on Al and Si doping concentrations. Table 2 presents a comparison of the sensing performance for the humidity sensors based on different sensing materials. The results highlight that the  $(AI + Si)_{0.02}Sn_{0.98}O_2$  humidity sensor exhibits several advantages, including faster response and recovery times, as well as lower hysteresis. These findings suggest that the inclusion of



Fig. 13 Response/recovery time of pure and Al–Si co-doped  $SnO_2$  humidity sensors at the humidity ranging from 11% to 95% RH.

 $\label{eq:table_$ 

Material	Response time (s)	Recovery time (s)	Hysteresis (%)	Ref.
SnO <sub>2</sub> /SiO <sub>2</sub>	14	16	2	35
ZrP	57	9	5.2	36
Al/SnO <sub>2</sub>	100	88	_	37
$Na_{1/3}Sr_{1/3}Tb_{1/3}Cu_{3}Ti_{4}O_{12}$	179	220	10.3	38
$(Al + Si)_{0.02}Sn_{0.98}O_2$	25	5	1.94	This work

Al–Si dopants plays a crucial role in enhancing the adsorption of water molecules on the SnO<sub>2</sub> surface, thus improving the overall performance of the sensor.

Respiratory detection plays a crucial role in the prevention of respiratory diseases and fitness tracking. Fig. 14(a)-(d) shows the response testing of pure SnO<sub>2</sub> and Al-Si co-doped SnO<sub>2</sub> sensors towards different breathing rates. The instantaneous capacitance of the humidity sensor was measured using the LCR testing instrument. The changes in capacitance were highly variable when the person breathed in different rates or depths and then according to it, the breathing was distinguished into shallow, medium and deep respiration. As seen from Fig. 14(b)-(d), the capacitance values of the Al-Si co-doped samples change more quickly than that of SnO<sub>2</sub> under different respiratory rates. For pure SnO<sub>2</sub>, the respiratory intervals for response/recovery time at shallow (20-250 pF), medium (250-400 pF) and deep (400-650 pF) breathing states were 4.7 s, 5.5 s and 6.3 s respectively. After increasing the content of Al-Si (as shown in Fig. 14(b)-(d)), the average response intervals decreased by approximately 1.8 seconds in different breathing states compared to pure  $SnO_2$ . In the case of the (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub> sensor, it was observed that the average response time in the shallow breathing state was shorter, indicating a quicker response. This suggests that the introduction of Al<sup>3+</sup> and Si<sup>4+</sup> into SnO<sub>2</sub> can enhance the response and recovery ability of the material. This finding demonstrates the promising potential of Al-Si co-doped SnO<sub>2</sub> for promptly monitoring breathing rates and depths. Consequently, the  $(Al + Si)_{0.02}Sn_{0.98}O_2$  sensor exhibited outstanding humidity sensing performance, further establishing its potential as an advanced respiratory sensor.

#### 3.5 Humidity sensing mechanism

The chemical mechanisms for the enhancement in the humidity sensitivity induced by structural defects such as oxygen vacancies were proposed. As shown in Fig. 15, the introduction of dopant ions induces the formation of substitutional defects  $Al_{Sn}$ ,  $Si_{Sn}$  and the oxygen vacancies in the crystalline structure of  $SnO_2$ . The introduction of the Al–Si dopant ions leads to the generation of oxygen vacancies. This increased the presence of oxygen vacancy defects provides additional active surface adsorption sites for water molecules on the surface of Al–Si co-doped  $SnO_2$ . The oxygen vacancies accelerate water adsorption and desorption reactions, which is beneficial for hopping transportation of protons on the  $SnO_2$  surface.<sup>39</sup> The water molecules adsorption/ desorption reaction formula on the surface of  $SnO_2$  are as



Fig. 14 The respiratory intervals of (a) pure and (b–d) Al–Si co-doped SnO<sub>2</sub> for response/recovery time at shallow, medium and deep breathing states.

follows:40

$$\frac{1}{2}O_2 + \alpha \cdot e^- + S \Leftrightarrow O_s^{-\alpha}$$
(11)

$$H_2O + 2 \cdot Sn_{Sn} + O_S^{-\alpha} \Leftrightarrow 2 \cdot (SnOH) + \alpha \cdot e^- + S$$
(12)

where  $O_S^{-\alpha}$  represents a surface adsorbed oxygen species with = 1 indicating singly ionized oxygen species or 2 for doubly ionized oxygen species; e<sup>-</sup> represents an electron in the conduction band surface; and *S* represents a surface adsorption site. In the case of Al–Si co-doping, the oxygen vacancies originating from Al and Si ions incorporated into SnO<sub>2</sub> can consume excess electrons, and



Fig. 15 Structure diagram of the defects  $Al_{\text{Sn}}$  +  $Si_{\text{Sn}}$  +  $V_{\text{O}}$  reacting with water molecules.

since then, the oxygen vacancies with one trapped electron increased significantly. Both the oxygen vacancies and the electrons can serve as charge carriers to facilitate electronic conduction and enhance the capacitance. As a result, the oxygen vacancies can accelerate the ionisation of water molecules into conductive H<sup>+</sup> and bind the water molecules near oxygen vacancies into  $H_3O^+$  ( $H_2O \rightarrow H^+ + OH^-$ ,  $H^+ + H_2O \rightarrow H_3O^+$ ), thereby enhancing the sensors sensitivity at different humidity levels and rapidly increasing its capacitance.<sup>41</sup> Hence, the introduction of Al and Si ions is responsible for the enhancement of the sensitivity and linearity of the sensor.

## 4. Conclusions

In this study, we successfully synthesized a novel capacitive humidity sensor utilizing acceptor–donor co-doped SnO<sub>2</sub> nanoparticles for the detection of ambient humidity and human respiration. The Raman spectra reveal an increased presence of bridging mode oxygen vacancies in  $(AI + Si)_{0.02}Sn_{0.98}O_2$ . Additionally, the EPR results demonstrate a higher concentration of singly ionized oxygen vacancies following the introduction of Al–Si. Furthermore, DFT calculations provide evidence that Si 3s orbital and Al 3s orbital contribute to the band edge, confirming the formation of acceptor–donor complexes in Al– Si co-doped  $SnO_2$ . The humidity sensing experiment showed that the  $(AI + Si)_{0.02}Sn_{0.98}O_2$  humidity sensor exhibited high sensitivity (S = 839), low hysteresis (1.94%), fast response/ recovery times (25 s/5 s) and quick respiratory intervals at three breathing states. We also discussed the chemical mechanisms for the enhancement in the humidity sensitivity induced by oxygen vacancies in (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub>. Hence, the (Al + Si)<sub>0.02</sub>Sn<sub>0.98</sub>O<sub>2</sub> sensor exhibits excellent humidity sensing characteristics for detecting ambient humidity and human respiratory rate.

# Research data policy and data availability statements

All data generated or analysed during this study are included in this published article.

## Author contributions

All authors contributed to the study conception and design. Yuchuan Ding performed material preparation, data collection, analysis, and the creation of new software used in the work. Yong Chen made substantial contributions to the conception and design of the work, and Mao-Hua Wang revised the work critically for important intellectual content. The first draft of the manuscript was written by Yuchuan Ding and all authors commented on the previous versions of the manuscript.

# Conflicts of interest

The authors declare that they have no conflict of interest.

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