

# PCCCP

Physical Chemistry Chemical Physics

Accepted Manuscript

This article can be cited before page numbers have been issued, to do this please use: C. J. Gibb, M. M. Majewska, G. J. Strachan, D. Pocięcha, J. M. D. Storey, E. Gorecka and C. T. Imrie, *Phys. Chem. Chem. Phys.*, 2024, DOI: 10.1039/D4CP04076G.



This is an Accepted Manuscript, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this Accepted Manuscript with the edited and formatted Advance Article as soon as it is available.

You can find more information about Accepted Manuscripts in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this Accepted Manuscript or any consequences arising from the use of any information it contains.

# Twist-bend liquid crystal phases and molecular structure: the role of methoxybiphenyl

View Article Online  
DOI: 10.1039/D4CP04076G

Calum J Gibb <sup>a†</sup>, Magdalena M Majewska <sup>b</sup>, Grant J Strachan <sup>b</sup>, Damian Pocięcha <sup>b</sup>, John MD Storey <sup>a</sup>, Ewa Gorecka <sup>b</sup>, and Corrie T Imrie <sup>a\*</sup>

<sup>a</sup> Department of Chemistry, School of Natural and Computing Sciences, University of Aberdeen, Meston Building, Aberdeen, AB24 3UE (UK)

<sup>†</sup> Current address: School of Chemistry, University of Leeds, Leeds, LS2 9JT (UK)

<sup>b</sup> University of Warsaw, Faculty of Chemistry, ul. Zwirki i Wigury 101,02-089, Warsaw, Poland

\*Author for correspondence: e-mail [c.t.imrie@abdn.ac.uk](mailto:c.t.imrie@abdn.ac.uk)

## Abstract

The synthesis and characterisation of the 4-[[4-((6-[4-(4-methoxyphenyl)phenyl]hexyl)oxy)phenyl)methylidene]amino] phenyl 4-alkyloxybenzoates is reported. These are referred to using the acronym MeOB6OIBeOm in which *m* denotes the number of carbon atoms in the terminal alkyloxy chain and is varied from one to ten. All ten members exhibit an enantiotropic conventional nematic (N) phase. In addition, for *m*=1-9, the twist-bend nematic (N<sub>TB</sub>) phase was observed on cooling the N phase. The N-isotropic (I) and N<sub>TB</sub>-N transition temperatures decrease on increasing the length of the terminal chain and this is more pronounced for the former. This supports the view that the N<sub>TB</sub>-N transition is predominantly shape driven and this depends largely on the length and parity of the flexible spacer. The transitional behaviour of the MeOB6OIBeOm series is compared with that of the corresponding dimers based instead on a cyanobiphenyl fragment, the CB6OIBeOm series. The rich smectic polymorphism exhibited by the CB6OIBeOm series is extinguished for the MeOB6OIBeOm series. The CB6OIBeOm series shows higher values of the N-I transition temperature than the corresponding members of the MeOB6OIBeOm series whereas the values of the N<sub>TB</sub>-N transition temperatures are rather similar for corresponding members of the two series. This again suggests that the N<sub>TB</sub>-N transition is largely shape driven whereas the mixed core interaction plays a more distinct role in driving N phase formation. The promotion of



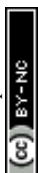
smectic behaviour in the CB6OIBeOm series is attributed to the strong tendency of the cyanobiphenyl fragment to adopt anti-parallel associations.

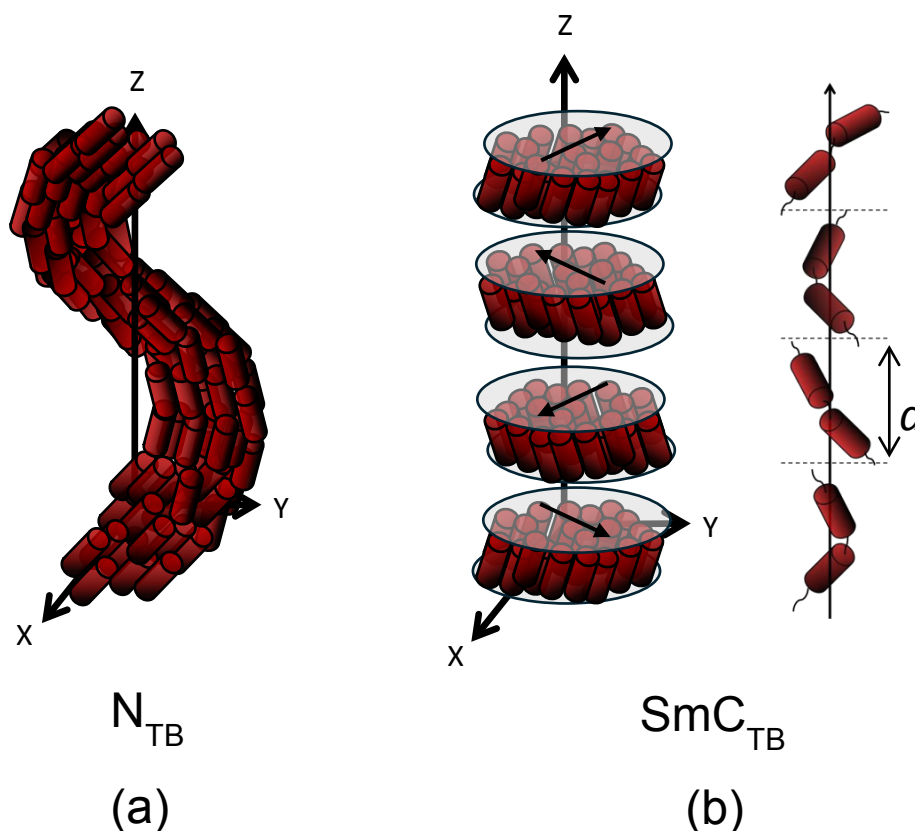
### Keywords:

Liquid crystal dimers; nematic phase; twist-bend nematic phase; twist-bend smectic phase; intercalated; methoxybiphenyl; cyanobiphenyl.

### Introduction

Reports of spontaneous chiral symmetry breaking in soft matter systems composed of achiral molecules always garner significant attention among the scientific community. The twist-bend nematic ( $N_{TB}$ ) phase provided the first example of such behaviour in a fluid with no positional order<sup>1,2,3</sup>. First predicted by Meyer<sup>4</sup> and later independently by Dozov<sup>5</sup> a decade before its experimental discovery, in the  $N_{TB}$  phase the director adopts a heliconical structure in which it is tilted at arbitrary angle with respect to the helical axis (**Fig. 1(a)**). Chirality emerges spontaneously giving doubly degenerate domains of opposite handedness. Molecular chirality removes this degeneracy, and the chiral twist-bend nematic phase is obtained<sup>6,7,8,9</sup>. A striking feature of the  $N_{TB}$  phase is that the pitch length of the helix is very short and just a few molecular lengths. At the core of Dozov's prediction was the assumption that bent molecules have a strong natural tendency to pack into bent structures. Such arrangements, however, cannot fill space and so bend must be accompanied by a second deformation, namely twist in the case of the  $N_{TB}$  phase. The requirement for a bent molecular structure was realised by using odd-membered liquid crystal dimers and these now constitute the overwhelming majority of the twist-bend nematogens reported to date<sup>10-16</sup>. A liquid crystal dimer consists of molecules containing two semi-rigid mesogenic units, the interactions between which drive the formation of liquid crystal phases, separated by a flexible spacer, most commonly an alkyl chain. The parity of this spacer determines whether the molecules is on average linear or bent<sup>17</sup>.





**Fig. 1.** Sketches of the molecular arrangement in (a) the twist-bend nematic phase ( $N_{TB}$ ) and (b) in the twist-bend smectic phase ( $SmC_{TB}$ ).

In his seminal paper Dozov also predicted the existence of twist-bend smectic phases<sup>5</sup> and these have now also been found, again by studying odd-membered liquid crystal dimers<sup>18</sup>. These have been termed the twist-bend smectic C ( $SmC_{TB}$ ) phases in which the director again forms a helical structure having a very short pitch length (**Fig. 1(b)**)<sup>18</sup>. Several structural variants of the  $SmC_{TB}$  phase have been discovered<sup>19</sup> although this new class of liquid crystal phase has been observed for rather few molecules, all having similar molecular architectures<sup>20-23</sup>.

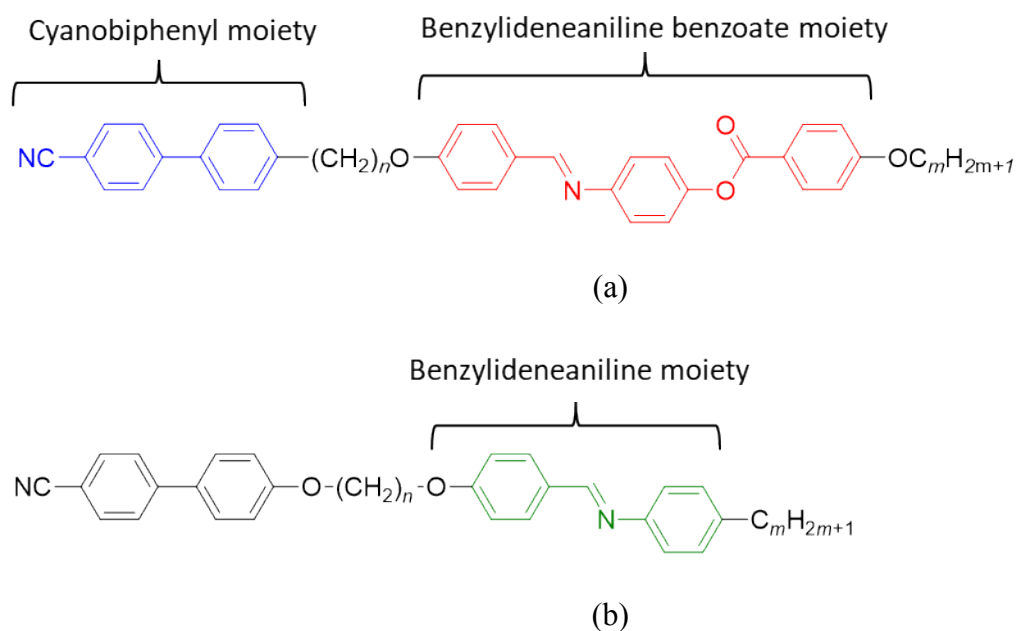
The  $SmC_{TB}$  phase was first observed in the CB6OIBeOm series<sup>18</sup> in which two differing mesogenic units are attached via the spacer and this class of materials are referred to as non-symmetric liquid crystal dimers. The CB6OIBeOm series contains cyanobiphenyl (CB) and



benzylideneaniline benzoate (IBeO) mesogenic units linked via a hexyloxy spacer (6O) and  $m$  represents the number of carbon atoms in the terminal alkyloxy chain (**Fig. 2(a)**). The CB6OIBeO $m$  series displays fascinating liquid crystalline behaviour that depends on the length of the terminal chain ( $m$ ). For terminal chains shorter than the length of the spacer, the N<sub>TB</sub> phase was observed below the conventional nematic (N) phase. If the terminal chain exceeded the length of the spacer, the N<sub>TB</sub> phase was extinguished and replaced by a rich smectic polymorphism including the smectic A (SmA), biaxial smectic A (SmA<sub>b</sub>) and SmC<sub>TB- $\alpha$</sub>  phases. Recent studies have further highlighted the importance of the relative lengths of the spacer and terminal chain in determining phase behaviour in this class of materials<sup>24</sup>. We have also shown that the molecular bend angle plays a critical role in determining phase behaviour<sup>25</sup>. Empirically, albeit based on a rather small collection of molecules having similar structures, it appears that the N<sub>TB</sub> phase is extinguished if the terminal chain length exceeds that of the spacer and twist-bend smectic phases are instead observed.

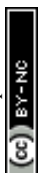
View Article Online  
DOI: 10.1039/D4CP04076G





**Fig. 2.** The general structures of (a) the CB6OIBeOm series<sup>18</sup> in which the cyanobiphenyl mesogenic unit is highlighted in blue and the benzylideneaniline benzoate moiety in red, and (b) the CBO $n$ O□ $m$  series<sup>26</sup> in which the benzylideneaniline mesogenic unit is highlighted in green.

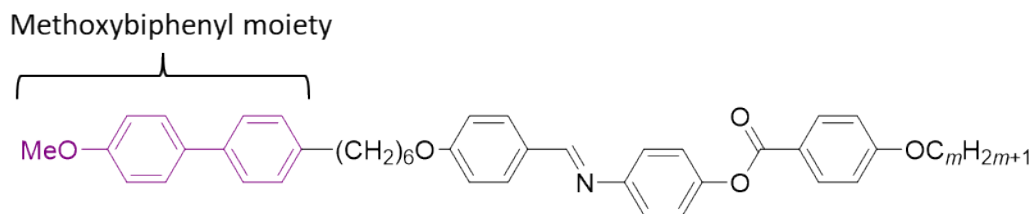
The driving force for the formation of the smectic phases observed in these materials has recently been suggested to be anti-parallel associations between the benzylideneaniline benzoate mesogenic units and the structure further stabilised by anti-parallel associations between the cyanobiphenyl moieties<sup>24</sup>. In consequence, the smectic mesophases exhibited by these dimers show layer periodicities comparable to full molecular length,  $d \sim l$ . This contrasts with the behaviour reported for cyanobiphenyl-based non-symmetric dimers containing benzylideneaniline-based units for which the relative lengths of the spacer and terminal chain controls the ability of the mesogenic units to interact leading to smectic phases with differing local packing arrangements. For example, for the CBO $n$ O. $m$  series (**Fig. 2(b)**) when the terminal chain ( $m$ ) is shorter than the spacer ( $n$ ) favourable interactions between the unlike mesogenic groups promote the formation of intercalated structures in which the layer



periodicity is approximately half the molecular length,  $d \sim l/2$ <sup>26,27,28</sup>. This structure is only stable if the terminal chain can be accommodated in the structural volume determined by the length of the spacer. Increasing the terminal chain destabilises the intercalated arrangements and instead the molecules pack into interdigitated structures. These have bilayer periodicities with  $d \sim 2l$  and are stabilised by the antiparallel association of the cyanobiphenyl units that serve to minimise dipolar energy while the smectic behaviour is driven by molecular inhomogeneity<sup>26,27,29</sup>.

We have seen that the cyanobiphenyl unit is thought to play an important role in determining the quite different liquid crystal behaviour of the CB6OIBeOm (**Fig. 2(a)**)<sup>18</sup> and CBOm series (**Fig. 2(b)**)<sup>26</sup>. Of particular significance is the strong tendency of the cyanobiphenyl group to form anti-parallel associations<sup>30</sup>. To test these suggestions, here we report a series of non-symmetric dimers containing the benzylideneaniline benzoate mesogenic unit but replace the cyanobiphenyl moiety with the methoxybiphenyl mesogenic unit (**Fig. 3**). We refer to this series using the acronym MeOB6OIBeOm series in which MeOB refers to methoxybiphenyl. We selected the methoxybiphenyl unit because it is less polar and polarizable than the cyanobiphenyl moiety and also reduces the structural anisotropy of the molecule. These differences manifest in lower transition temperatures compared to the corresponding cyanobiphenyl-based dimers<sup>31,32,33</sup>. Critically for this study, the methoxybiphenyl moiety does not show a tendency to exhibit anti-parallel dimers and this should significantly effect the ability of these dimers to exhibit the fascinating smectic phase behaviour seen for the CB6OIBeOm series. In terms of molecular design, a hexyloxy spacer was selected to provide the requisite molecular bend for twist-bend phases to be observed<sup>26,27,29</sup>. We compare the transitional properties of the MeOB6OIBeOm and CB6OIBeOm<sup>18</sup> in order to further understand the role played by molecular structure in determining the fascinating twist-bend phase behaviour.





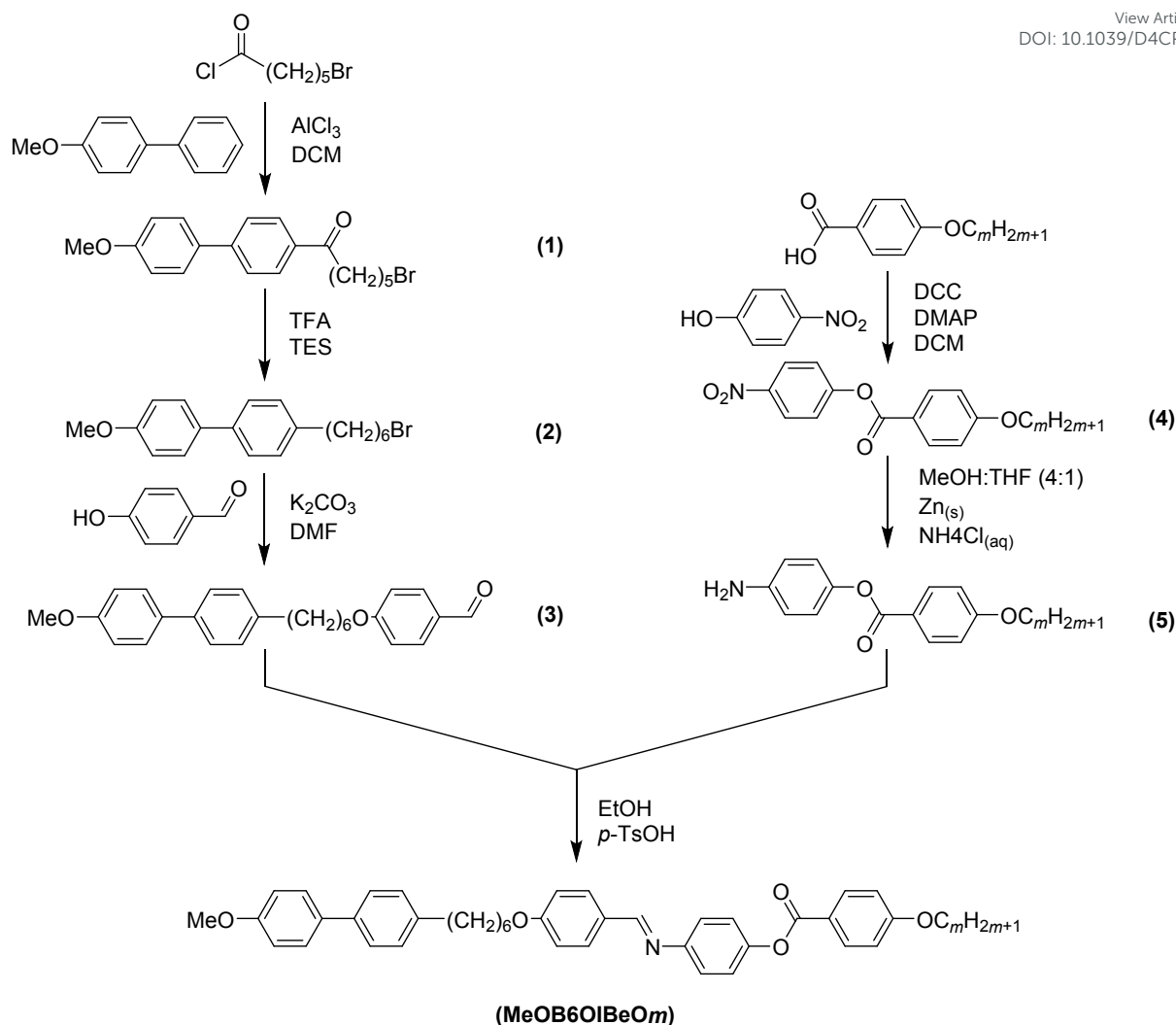
**Fig. 3.** The general molecular structure of the MeOB6OIBeOm, where  $m = 1-10$ . The methoxybiphenyl mesogenic unit is highlighted in purple.

## Experimental

The synthesis of the MeOB6OIBeOm series is shown in **Scheme 1**. The synthesis begins with a Friedel-Crafts acylation reaction between 6-bromohexanoyl chloride and 4-methoxybiphenyl forming 6-bromo-1-(4'-methoxy-[1,1'-biphenyl]-4-yl)hexan-1-one (**1**).<sup>34</sup> The ketone functional group is reduced through treatment with triethyl silane (TEA) and trifluoroacetic acid (TFA) forming 4-(6-bromohexyl)-4'-methoxy-1,1'-biphenyl (**2**).<sup>35</sup> A standard Williamson ether synthesis is performed using 4-hydroxybenzaldehyde, generating 4-((6-(4'-methoxy-[1,1'-biphenyl]-4-yl)hexyl)oxy)benzaldehyde (**3**).<sup>33</sup> (**3**) is coupled with the appropriate 4-aminophenyl-4'-alkoxy benzoate ester (**5**)<sup>18,24</sup> using a standard Schiff-base synthesis forming the 4-[[4-({6-[4-(4-methoxyphenyl)phenyl]hexyl}oxy)phenyl]methylidene}amino]phenyl 4-alkylbenzoates. Complete synthetic details, along with structural, and purity analysis for final products and their intermediates are provided in the accompanying ESI.



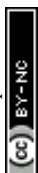




**Scheme 1.** The synthetic scheme for the synthesis of the MeOB6OIBeOm series ( $m = 1-10$ ).

## Results and Discussion

The transitional properties of the MeOB6OIBeOm series are listed in **Table 1**. All ten members exhibit an enantiotropic nematic phase identified by the observation of schlieren optical textures containing both 2- and 4-point brush singularities and which flashed when subjected to mechanical stress (**Fig. 4(a)**). On cooling the N phase for  $m = 1-9$ , a blocky, rope-like texture was obtained, characteristic of the N<sub>TB</sub> phase. The N<sub>TB</sub> phase was enantiotropic in nature for  $m = 1-6$  and monotropic for  $7-9$ . Representative DSC traces obtained on heating and cooling for MeOB6OIBeO7 are shown in **Fig. 4(b)**. On heating, a strong endotherm associated with the Cr-N transition and a weaker endotherm associated with the N-I transition

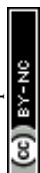


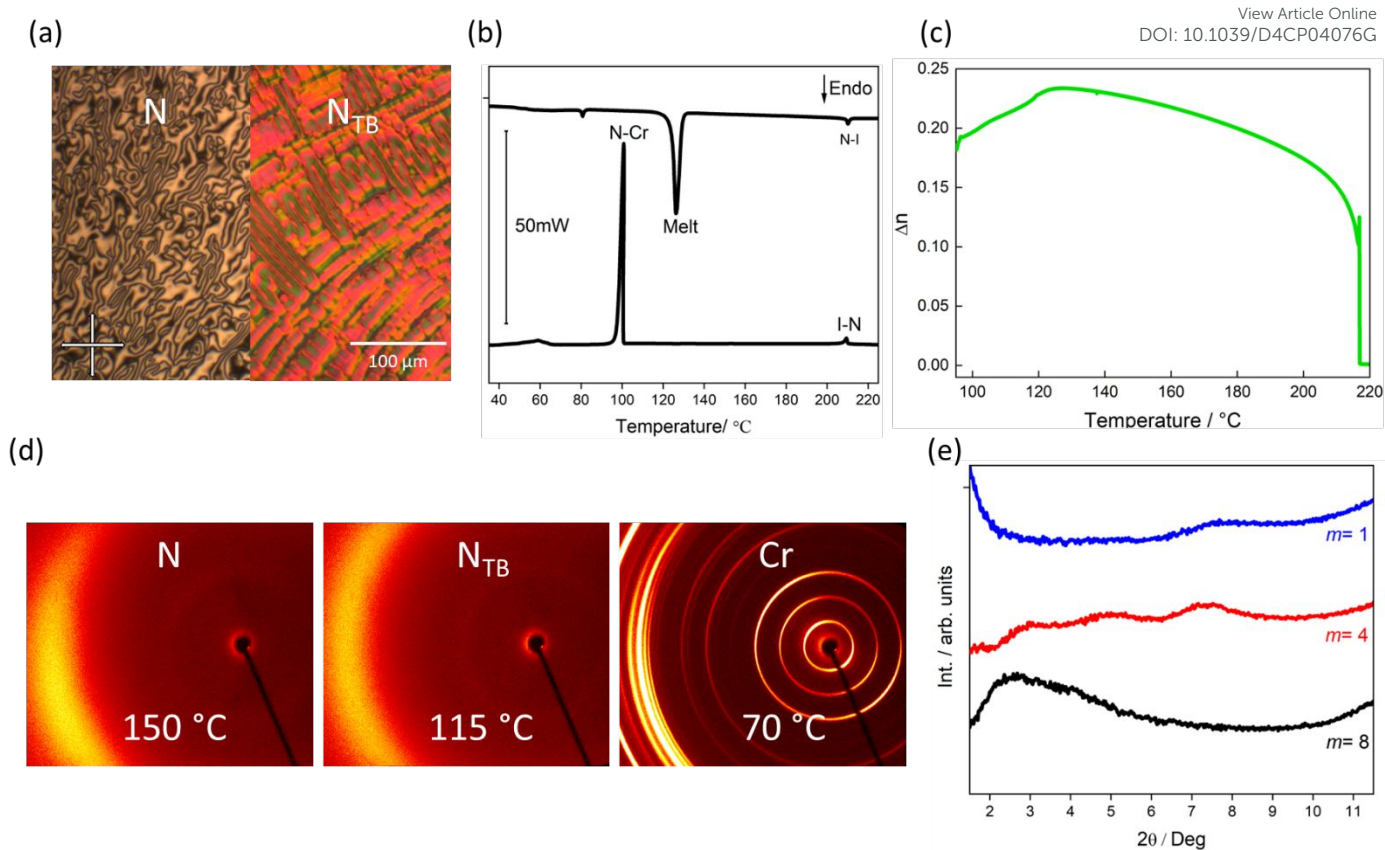
are observed. The corresponding exotherms are observed in the cooling trace. The  $N_{TB-N}$  transition is not detected using DSC given that the  $N_{TB-N}$  phase transition becomes weakly first order as the temperature range of the preceding nematic increases<sup>36</sup> and the MeOB6OIBeO $m$  series shows wide temperature range nematic phases. **Fig. 4(c)** shows the temperature dependence of the optical birefringence ( $\Delta n$ ) measured for MeOB6OIBeO8 and a continuous increase in the order parameter is seen on cooling the N phase. At the  $N_{TB-N}$  transition a characteristic decrease is observed in  $\Delta n$  due to the averaging of optical anisotropy resulting from the formation of the short-pitched helix<sup>37,38</sup>.

**Table 1.** Transition temperatures ( $^{\circ}\text{C}$ ) and associated scaled entropy changes, in brackets, for the MeOB6OIBeO $m$  series.

$m$	m.p.		$N_{TB-N}$		N-I	
	T/ $^{\circ}\text{C}$	$\Delta S/R$	T/ $^{\circ}\text{C}$	$\Delta S/R$	T/ $^{\circ}\text{C}$	$\Delta S/R$
<b>1</b>	136.3	17.7	140.0 <sup>†</sup>	$\approx 0$	260.2	0.44
<b>2</b>	123.3	15.7	138.5 <sup>†</sup>		262.2	0.54
<b>3</b>	133.8	15.8	135.5 <sup>†</sup>		248.6	0.38
<b>4</b>	131.4	16.1	135.0 <sup>†</sup>		243.0	0.38
<b>5</b>	131.0	14.1	133.0 <sup>†</sup>		236.5	0.35
<b>6</b>	128.4	16.0	130.0 <sup>†</sup>		229.7	0.36
<b>7</b>	127.9	16.9	[126.5] <sup>†</sup>		222.0	0.34
<b>08</b>	128.1	16.9	[125.5] <sup>†</sup>		219.3	0.34
<b>9</b>	127.2	17.0	[123.0] <sup>†</sup>		212.8	0.28
<b>10</b>	128.2	17.6	-		-	210.5

<sup>†</sup> microscope observation, [ ] monotropic transition.





**Fig. 4:** (a) Characteristic POM textures for a sample sandwiched between untreated glass slides of the N and N<sub>TB</sub> phases exhibited by MeOB6OIBeO7, (b) DSC thermogram for MeOB6OIBeO9, (c) the temperature dependence of the optical birefringence,  $T(\Delta n)$ , for MeOB6OIBeO8 measured using green light,  $\lambda=532$  nm in a 1.7  $\mu\text{m}$  cell treated for planar alignment, (d) 2D X-ray diffraction patterns for MeOB6OIBeO1 in the N [ $T=150$  °C] (left), N<sub>TB</sub> [ $T=115$  °C] (centre) and Cr [ $T=70$  °C] (right) phases, and (e) the low-angle XRD intensity profiles of the N<sub>TB</sub> phases shown by homologues with  $m=1, 4$  and  $8$  recorded at 115 °C, 128 °C, and 120 °C, respectively, of the MeOB6OIBeO $m$  series. The curves have been shifted along the vertical axis for clarity.

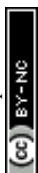
The phase assignments made using POM for the MeOB6OIBeO $m$  series were confirmed using X-ray diffraction (XRD). The diffraction patterns obtained in both the N and N<sub>TB</sub> phases showed only diffuse scattering in both the small and wide-angle regions indicative of short-range translational ordering (**Fig. 4(d)**). The local periodicity the nematic phases may be calculated from the positions of the peak in the small angle region (**Fig. 4(e)**) and these are given in **Table 2**. The  $d/l$  ratio for  $m=3$  is approximately 0.33 suggesting an intercalated structure. For  $m=4$  three broad signals are observed corresponding to  $d/l$  ratios of 0.33, 0.50



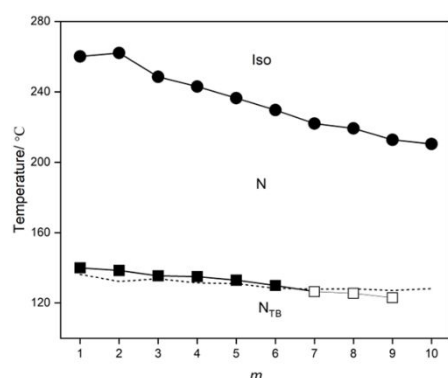
and 0.84. This suggests a frustrated local structure involving three different packing arrangements. For  $m=8$  a single diffuse peak is again observed with an associated  $d/l$  ratio of 0.84. These data suggest that the local structure in both nematic phases evolves from being intercalated for short terminal chains, experiences frustration for intermediate chain lengths, and adopts an approximately monolayer arrangement for long terminal chain lengths. We note, however, that this change in the local structure cannot be discerned from the dependence of the transition temperatures on  $m$  (**Fig. 5**) that show a regular dependence. Similar behaviour has been reported for the CB6O. $m$  series<sup>39</sup>. Furthermore, there appears to be no simple relationship between local structure and the helical pitch length in the  $N_{TB}$  phase. Both locally bilayer and intercalated types of  $N_{TB}$  phases form helical structures with similarly short pitch lengths. The frustration in the local packing mode (as observed, for example, in MeOB6OIBeO9) may be understood as a coexistence of both intercalated and bilayer local structures, but it appears that the breaking of the head and tail equivalence, which is responsible for the bilayer structure, does not affect the formation of the heliconical structure.

**Table 2:** Molecular end-to-end separation in the N and  $N_{TB}$  phases estimated from the position of the small angle peaks in the X-ray patterns,  $d$ , the estimated molecular lengths,  $l$ , and the  $d/l$  ratios for selected homologues of the MeOB6OIBeO $m$  series.

$m$	$d / \text{\AA}$	$l / \text{\AA}$	$d/l$
<b>1</b>	11.7	39.0	0.33
<b>4</b>	12.6	35.5	0.33
	19.3		0.50
	33.0		0.84
<b>8</b>	37.0	44.1	0.83



Increasing the length of the terminal chain sees essentially a linear decrease in the values of  $T_{N-I}$ , (**Fig. 5**), and this may be attributed to the dilution of mesogenic group interactions due to the increasing volume fraction of alkyl chains. The absence of an odd-even effect superimposed on the decreasing trend in  $T_{N-I}$  suggests that the effect of the change in molecular shape arising from the disposition of the methyl group with respect to the major axis of the benzylideneaniline benzoate mesogenic unit on  $T_{N-I}$  is much smaller than the dilution effect. The dependence of  $T_{N_{TB-N}}$  on  $m$  is similar but the decrease in  $T_{N_{TB-N}}$  on passing from  $m=1$  to 9 is only 17 K whereas that associated with  $T_{N-I}$  is 50 K. The smaller change in  $T_{N_{TB-N}}$  on increasing terminal chain length suggests that the dilution of interactions between the mesogenic groups has less of an effect on the  $N_{TB-N}$  transition than on the  $N-I$  transition. This is consistent with the view that the formation of the  $N_{TB}$  phase is predominantly shape-driven<sup>1,36,40</sup> and that this depends largely on the flexible spacer (**Fig. 5**).



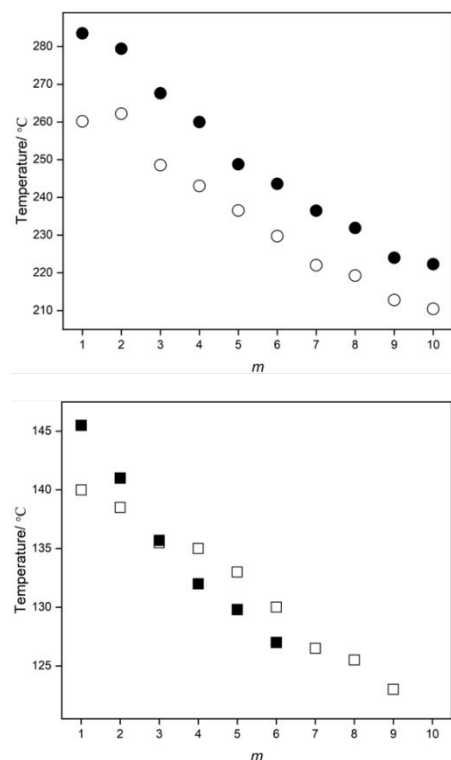
**Fig. 5.** Temperature dependence of the transition temperatures on increasing terminal chain length ( $m$ ) for the MeOB6OIBeOm series. The melting points are indicated by the dotted line. Open symbols represent monotropic phase transition.

The phase behaviour of the MeOB6OIBeOm series differs dramatically from that of the CB6OIBeOm series described earlier<sup>18</sup>. Indeed, the rich smectic polymorphism seen for the latter series is completely extinguished in the former. The  $T_{N-I}$  and  $T_{N_{TB-N}}$  transition temperatures for these two series are compared in **Fig. 6**. As we discussed earlier, replacing a nitrile by a methoxy group reduces  $T_{N-I}$  and this is clearly evident in **Fig. 6(a)**. Over all ten members,  $T_{N-I}$  is higher on average by 13.3 K for the CB6OIBeOm series. This difference is largest for the shorter chain lengths and is higher than the difference seen between the CB6O.Om and MeOB6O.Om series<sup>33</sup>. By comparison the values of  $T_{N_{TB-N}}$  for the two series are very similar and cross-over on increasing  $m$  (**Fig. 6(b)**). The larger difference seen for the



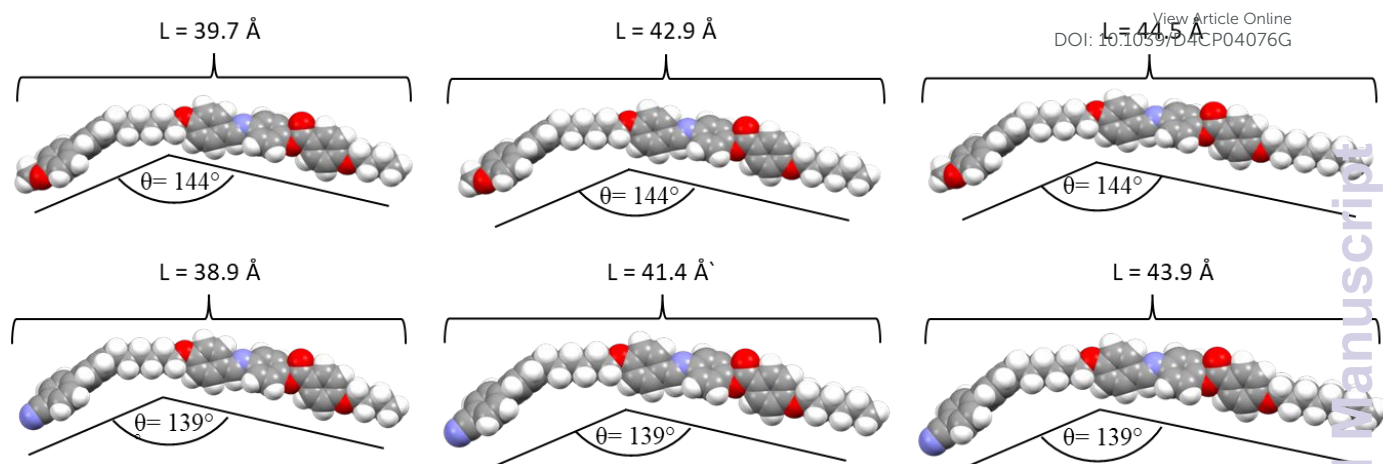
values of  $T_{N-I}$  for these two series may suggest that the interaction between the unlike mesogenic units does play a role in driving nematic phase formation, whereas as we noted earlier the  $N_{TB}-N$  transition is widely believed to be predominantly shape driven and the shapes of these dimers are rather similar (**Fig. 7**).

View Article Online  
DOI: 10.1039/D4CP04076G



**Fig 6.** Comparison of the  $T_{N-I}$  [top] and  $T_{N_{TB}-N}$  [bottom] for the MeOB6OIBeOm [open] and CB6OIBeOm<sup>18</sup> [solid] series.





**Fig. 7:** Geometry optimised structures of the MeOB6OIBeOm [top] and CB6OIBeOm [bottom] series showing the similarities in molecular structure on increasing the length of the terminal chain ( $m$ ). The structures were obtained by quantum mechanical density functional theory (DFT) calculations at the B3LYP 6-31G(d) level of theory<sup>41,42,43</sup>.

Finally, we consider the absence of smectic behaviour for the MeOB6OIBeOm series. The smectic phases exhibited by the CB6OIBeOm series are thought to be stabilised by favourable, presumably quadrupolar, interactions between neighbouring pairs of anti-parallel benzylideneaniline benzoate mesogenic units with further stabilisation of the structure coming from anti-parallel associations between the cyanobiphenyl moieties concentrated at the layer interfaces<sup>24</sup>. Methoxybiphenyl fragments do not have the same strong tendency to align in antiparallel associations and this will destabilise the smectic structure. This difference in how the molecular fragments interact results in the observation of the  $N_{TB}$  phase for longer terminal chain lengths for the MeOB6OIBeOm series compared to the CB6OIBeOm series for which the emergent SmA phase precludes the observation of the  $N_{TB}$  phase. This contrasting behaviour highlights the importance of the cyanobiphenyl moiety in stabilising the fascinating twist-bend smectic phases.

## Conclusions

We have reported a series of non-symmetric dimers containing methoxybiphenyl and benzylideneaniline benzoate mesogenic units linked by a hexyloxy spacer, the MeOB6OIBeOm series. Unlike the corresponding materials containing a cyanobiphenyl moiety<sup>18</sup>, all smectic polymorphism is lost and N and  $N_{TB}$  phases are observed across all members of the MeOB6OIBeOm series. This dramatic change in phase behaviour is attributed



to the importance of anti-parallel associations between the cyanobiphenyl mesogenic units in stabilising lamellar structures in this sub-class of non-symmetric liquid crystal dimers. This observation now has an important role to play in the design of new materials expected to exhibit these fascinating and rare twist-bend liquid crystal phases.

View Article Online

DOI: 10.1039/D4CP04076G

### Acknowledgements

D.P. gratefully acknowledges financial support from the National Science Centre (Poland) under the grant no. [2021/43/B/ST5/00240].

### References

1. M. Cestari, E. Frezza, A. Ferrarini and G.R. Luckhurst, *J. Mater. Chem.*, 2011, **21**, 12303-12308.
2. V. Borshch, Y.K. Kim, J. Xiang, M. Gao, A. Jakli, V.P. Panov, et al. *Nat. Commun.*, 2013, **4**, 2635.
3. D. Chen, J.H. Porada, J.B. Hooper, A. Klitnick, Y. Shen, M.R. Tuchband, et al. *Proc Natl Acad Sci U S A.*, 2013, **110**, 15931-15936.
4. R.B. Meyer. *Structural problems in liquid crystal physics, Les houches summer school in theoretical physics*. In: Molecular Fluids, Balian R.G. and Wil, G., editors. New York: Gordon and Breach; 1976. p. 273-373.
5. I. Dozov, *Europhys. Lett.*, 2001, **56**, 247-253.
6. R. Walker, *Liq. Cryst. Today*, 2020, **29**, 2-14.
7. R. Walker, D. Pocięcha, J.M.D. Storey, E. Gorecka and C.T. Imrie, *Chem. Eur. J.*, 2019, **25**, 13329-13335.
8. A. Ozegovic, A. Knezevic, J. Novak, S. Segota, P. Davidson and A. Lesac, *Chemphyschem.*, 2024, **25**, DOI 10.1002/cphc.202400065.
9. A. Ozegovic, J. Hobbs, R. Mandle, A. Lesac and A. Knezevic A, *J. Mater. Chem. C*, 2024, **12**, 13985-13993.
10. R.J. Mandle RJ, *Molecules*, 2022, **27**, 2689.
11. R.J. Mandle, *Soft Matter*, 2016, **12**, 7883-7901.
12. E.E. Pocock, R.J. Mandle and J.W. Goodby, *Soft Matter*, 2018, **14**, 2508-2514.
13. P.A. Henderson and C.T. Imrie, *Liq. Cryst.*, 2011, **38**, 1407-1414.
14. N. Tufaha, C.J. Gibb, J.M.D. Storey and C.T. Imrie, *Liq. Cryst.*, 2023, **50**, 1362-1374.
15. G.J. Strachan, M.M. Majewska, D. Pocięcha, J.M.D. Storey and C.T. Imrie, *Chemphyschem.*, 2023, **24**, .





16. E. Cruickshank, R. Walker, G.J. Strachan, C.H.F. Goode, M.M. Majewska, D. Pocięcha et al. *J. Molec. Liq.*, 2023, **391**, 123226. View Article Online  
DOI: 10.1039/D4CP04076G
17. C.T. Imrie and P.A. Henderson, *Chem. Soc. Rev.*, 2007, **36**, 2096-2124.
18. J.P. Abberley, R. Killah, R. Walker, J.M.D. Storey, C.T. Imrie, M. Salamonczyk, et al., *Nat. Commun.*, 2018, **9**, 228.
19. C.T. Imrie, R. Walker, J.M.D. Storey, E. Gorecka and D. Pocięcha, *Crystals*, 2022, **12**, 1245.
20. A.F. Alshammari, D. Pocięcha, R. Walker, J.M.D. Storey, E. Gorecka and C.T. Imrie, *Soft Matter*, 2022, **18**, 4679-4688.
21. D. Pocięcha, N. Vaupotic, M. Majewska, E. Cruickshank, R. Walker, J.M.D. Storey et al., *Adv. Mater.*, 2021, **33**, 2103288.
22. A.F. Alshammari, A. Zattarin, A. Pearson, E. Cruickshank, M. Majewska, D. Pocięcha et al., *Liq. Cryst.*, 2024, DOI 10.1080/02678292.2024.2334342.
23. M. Salamonczyk, N. Vaupotic, D. Pocięcha, R. Walker, J.M.D. Storey, C.T. Imrie et al., *Nat. Commun.*, 2019, **10**, 1922.
24. C.J. Gibb, M. Majewska, D. Pocięcha, J.M.D. Storey, E. Gorecka and C.T., *Chemphyschem.*, 2024, DOI 10.1002/cphc.202300848.
25. C.J. Gibb, J.M.D. Storey and C.T. Imrie, *Liq. Cryst.*, 2024, **51**, 1956-1963.
26. J.L. Hogan, C.T. Imrie and G.R. Luckhurst, *Liq. Cryst.*, 1988, **3**, 645-650.
27. C.T. Imrie, *Liq. Cryst.*, 2006, **33**, 1449-1454.
28. G.S. Attard, R.W. Date, C.T. Imrie, G.R. Luckhurst, S.J. Roskilly, J.M. Seddon et al., *Liq. Cryst.*, 1994, **16**, 529-581.
29. A.E. Blatch, I.D. Fletcher and G.R. Luckhurst, *Liq. Cryst.*, 1995, **18**, 801-809.
30. D.A. Dunmur, *Liq. Cryst.*, 2015, **42**, 678-687.
31. J.P. Abberley, J.M.D. Storey and C.T. Imrie, *Liq. Cryst.*, 2019, **46**, 2102-2114.
32. D.A. Paterson, R. Walker, J.P. Abberley, J. Forestier, W.T.A. Harrison, J.M.D. Storey et al., *Liq. Cryst.*, 2017, **44**, 2060-2078.
33. J.P. Abberley, S.M. Jansze, R. Walker, D.A. Paterson, P.A. Henderson, A.T.M. Marcelis et al., *Liq. Cryst.*, 2017, **44**, 68-83.
34. D.A. Paterson, M. Gao, Y.K. Kim, A. Jamali, K.L. Finley, B. Robles-Hernández et al., *Soft Matter*, 2016, **12**, 6827-6840.
35. S.M. Jansze, A. Martínez-Felipe, J.M.D. Storey, A.T.M. Marcelis and C.T. Imrie, *Angew. Chem. Int. Ed.*, 2015, **54**, 643-646.



36. D.A. Paterson, J.P. Abberley, W.T. Harrison., J.M. Storey and C.T. Imrie, *Liq. Cryst.*, [View Article Online DOI: 10.1059/D4CP04076G](#) 2017, **44**, 127-146.
37. D. Pocięcha, C.A. Crawford, D.A. Paterson, J.M.D. Storey, C.T. Imrie, N. Vaupotic et al., *Phys. Rev. E*, 2018, **98**, 052706.
38. C. Meyer, G.R. Luckhurst and I. Dozov, *J. Mater. Chem. C*, 2015, **3**, 318-328.
39. R. Walker, D. Pocięcha, G.J. Strachan, J.M.D. Storey, E. Gorecka and C.T. Imrie, *Soft Matter*, 2019, **15**, 3188-3197.
40. D.A. Paterson, C.A. Crawford, D. Pocięcha, R. Walker, J.M.D. Storey, E. Gorecka et al., *Liq. Cryst.*, 2018, **45**, 2341-2351.
41. M.J. Frisch, G.W. Trucks, H.B. Schlegel, G.E. Scuseria, M.A. Robb, J.R. Cheeseman JR et al. *Gaussian 09 (Revision D.01)*, Wallingford CT: Gaussian Inc., 2016.
42. R. Dennington, T. Keith and J. Millam, *Gauss View, Version 5.5*, Shawnee Mission, KS: Semichem Inc; 2009.
43. C.F. Macrae, I. Sovago, S.J. Cottrell, P.T.A. Galek, P. McCabe, E. Pidcock et al. *J. Appl. Cryst.*, 2020, **53**, 226-235.



# Twist-bend liquid crystal phases and molecular structure: the role of methoxybiphenyl

[View Article Online](#)

DOI: 10.1039/C4CP04076G

Calum J Gibb <sup>a†</sup>, Magdalena M Majewska <sup>b</sup>, Grant J Strachan <sup>b</sup>, Damian Pocięcha <sup>b</sup>, John MD Storey <sup>a</sup>, Ewa Gorecka <sup>b</sup>, and Corrie T Imrie <sup>a\*</sup>

<sup>a</sup> Department of Chemistry, School of Natural and Computing Sciences, University of Aberdeen, Meston Building, Aberdeen, AB24 3UE (UK)

<sup>†</sup> Current address: School of Chemistry, University of Leeds, Leeds, LS2 9JT (UK)

<sup>b</sup> University of Warsaw, Faculty of Chemistry, ul. Zwirki i Wigury 101,02-089, Warsaw, Poland

\*Author for correspondence: e-mail [c.t.imrie@abdn.ac.uk](mailto:c.t.imrie@abdn.ac.uk)

- The data supporting this article have been included as part of the Supplementary Information.

