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***The Effect of H₂ Occupancy Modes in Small and Large Cages of
H₂-Tetrahydrofuran Hydrates on the Hydrates Stability and
H₂ Storage Capacity***

Ruyi Zheng ^{1,2}, Sohaib Mohammed ¹, Yang Jia ¹, Rituparna Hazra ¹ and Greeshma Gadikota^{1,*}

¹ School of Civil and Environmental Engineering, Cornell University, Ithaca, NY 14853

² Frontier Institute of Science and Technology, Xi'an Jiaotong University, Xi'an, Shaanxi
710054, China

*Corresponding author: Greeshma Gadikota, email: gg464@cornell.edu



Abstract

Hydrogen storage as hydrates is considered as one of the most environmentally benign approaches to store hydrogen as it requires only water and traces of promoters. However, advancing scalable controls on hydrogen hydrates formation is hindered by the limited understanding of the structure, dynamics and energetics of hydrogen and promoters in the hydrate cages. In this study, molecular dynamic simulation configurations with different occupancy modes of H₂ and Tetrahydrofuran (THF) in the hydrate cages are investigated under the following scenarios: *i*) two H₂ molecules occupying the small cages, *ii*) occupancy of H₂ molecules in the THF-free large cages, and *iii*) co-occupancy of H₂ and THF in one large cage. Exploring these scenarios reveals the impact of occupancy modes on the dynamic motion of guest and water molecules and on the hydrate structure stability. The results show that the occupancy of two H₂ molecules in the small cages reduced the stability of hydrate structure, triggered the inter-cage hopping of H₂ molecules through pentagonal faces, and increased the probability of hydrogen bond formation between THF and cage H₂O molecules. The thermodynamic stability of hydrate cages is increased when the THF-free large cages are occupied by H₂ molecules and the tetrahedral feature of H₂ distribution in the large cages is enhanced when the number of H₂ in one cage increased from two to three. When the large cages are co-occupied by H₂ and THF, the inter-cage migration of H₂ originated from large cages demonstrated two different features, i.e., ballistic motion ($\text{MSD} \propto t^2$) due to the tunneling migration behavior in the initial stage and diffusive motion in the late stage. The ballistic migration of H₂ molecules is more favorable for achieving a higher storage capacity. The decay rate of THF orientation is reduced when the interaction between THF and H₂O molecules is stronger. The insights provided by this study are crucial to advance the understanding of hydrogen storages as hydrates for a sustainable energy future.



Keywords: H₂-THF hydrates; Storage capacity; Cage occupancy; Inter-cage migration; Ballistic motion

1. Introduction

Hydrogen (H₂) is a promising energy carrier with a high energy density (~142 MJ/kg), making it an essential vector in the integrated energy system, such as balancing the seasonal fluctuation of energy demand and supply. However, safe and efficient H₂ storage is challenged by its high flammability, low density, and small molecular size. Numerous approaches have been proposed for hydrogen storage, e.g., geological storage, high-pressure-tank storage, and hydrogen storage materials including hydrate-based storage.(1–7) Hydrate-based H₂ storage is a less explored but transformative approach owing to: (i) the ease of recovering the stored H₂ with high purities, (ii) the formation of hydrates at mild pressure and temperature, (iii) environmentally benign features of hydrates compared to other H₂ storage materials, and (iv) stable trapping of H₂ molecules in the hydrate cages, which reduces the risk of H₂ leakage.

Clathrate hydrates are ice-like crystalline solids formed by guest molecules and cages of H₂O molecules that are connected through hydrogen bonds and in which the guest molecules are trapped. Depending on the guest species and thermodynamic conditions, various structures of clathrate hydrates might be formed.(8) In this context, sII structure has been extensively detected in H₂-containing hydrates.(9,10) A unit cell of sII hydrate is composed of 16 small cage and 8 large cages, as shown in **Figure 1**. Each small cage is made of 12 pentagonal faces and each large cage consists of 12 pentagonal and 4 hexagonal faces. The radius of the small and large cages is 3.91 and 4.73 Å respectively, which results in an average lattice parameter of a sII unit cell around 17.31 Å. Both the large and small hydrate cages can accommodate multiple H₂ molecules, which



makes it conducive for storing H₂ in the hydrate form. It has been reported that H₂ storage capacity of 5.3 wt% is achieved for the H₂ hydrates formed in pure water.(11)

The formation of pure H₂ hydrates requires high pressure and low temperature conditions. For instance, the temperature of 140 K is required to form H₂ hydrates at the ambient pressure, while 300 MPa is needed to form pure H₂ hydrates at 280 K(7,11) (see **Figure 1**). Therefore, it is essential to enable H₂ hydrate formation at moderate conditions to make this approach economically viable. Thermodynamic promoters have been considered to enhance the formation of H₂ hydrates under mild conditions as they increase the stability of hydrate structure by occupying the large cages.(12–14) As indicated in **Figure 1**, the addition of THF to H₂ hydrates could reduce the hydrate formation pressure from 300 MPa to 5 MPa at a temperature around 280 K.(9) The presence of thermodynamic promoters greatly improves the economic viability of the hydrate-based H₂ storage approach, however, at the cost of reduced storage capacity. Some representative findings about the storage capacity of H₂ in pure H₂ hydrates and binary H₂-THF hydrates can be found in **Table S1** in *Supplementary Information*. Generally, the storage capacity is higher when more H₂ molecules can be stably accommodated in the small and large hydrate cages. However, the conclusions on the occupancy modes of H₂ in the small and large cages (e.g., single or multiple H₂ occupying one hydrate cage) demonstrate great discrepancy, especially for the binary H₂-THF hydrates. Some researchers suggested that high storage capacity of H₂ is achievable in the H₂-THF hydrates because small cages could accommodate two H₂ molecules or that multiple H₂ molecules might occupy the large cages in the presence or absence of promoters. For instance, Lee and co-workers inferred from NMR results that the small cages in H₂-THF hydrates are occupied by two H₂ molecules and the H₂ occupancy in the large cages could be increased by reducing the concentration of THF, which resulted in a tunable H₂ storage capacity up to 4.0 wt%.(15) By using



ice powders and solid THF in the experiments, Sugahara and co-workers(16) observed a maximum H₂ storage capacity of 3.4 wt%, with the small cages occupied by one H₂ and multiple occupancy of H₂ molecules in the large ones. Nishikawa and co-workers detected H₂ occupancy in large cages of H₂-THF hydrates using *in-situ* Raman spectra without the quenching procedure.(17) On the contrary, it was concluded by other researchers that the small cages in H₂-THF hydrates could only accommodate one H₂ molecule.(18–20) For instance, by using similar experimental procedures as Lee and co-workers,(15) Strobel and co-workers concluded that the storage capacity of H₂ in the H₂-THF hydrates increased with pressure but asymptotically approached 1.0 wt%, with one H₂ in the small cage and the large cages being empty.(20)

To reveal the mechanisms of different occupancy modes of H₂ in the small and large cages, theoretical modeling and molecular dynamic (MD) simulations are conducted to understand the influence of promoters on the storage capacity of H₂. Some research findings support the multiple occupancy of H₂ in the small and large cages by analyzing the variation in energy or hydrate cage volume. For example, Patchkovskii and Tse concluded that the occupancy of two H₂ molecules in the small cages and four H₂ in the large ones favors the stability of the H₂ hydrate structure due to the increased surface contact between H₂ and cage H₂O molecules.(21) Sebastianelli and co-workers also suggested that the strong attractive interaction between H₂ and cage H₂O contributed to the negative ground-state energy when two H₂ molecules occupied the small cage.(22) Alavi and co-workers suggested that the configuration energy did not change significantly with the occupancy modes of H₂ and THF based on classical MD simulation.(23) By using *ab initio* molecular scale simulations, Tachikawa and co-workers concluded that the small and large cages could hold up to 2 and 5 H₂ molecules, respectively.(24) Koh and co-workers revealed that the small cages only needed to expand by around 3% in volume to provide thermodynamically stable



room for two H₂ molecules.(25) On the other hand, some researchers argued against the multiple occupancy of H₂ in the small and large cages.(26) It was suggested that the H₂ molecules are tightly confined in the repulsive potential field with two H₂ molecules occupying the small cages.(27) Liu and co – workers also indicated that the occupancy of two H₂ molecules in the small cage is less favorable and the most stable occupancy mode is the single occupancy of H₂ in the small cage and single H₂ + THF in the large cage.(28) Kang and co – workers concluded that the small cage is occupied by one H₂ molecule, whereas the large cages free of THF could accommodate two to three H₂ molecules.(29)

Despite considerable experimental observations and simulation results, several knowledge gaps associated with the storage capacity of H₂ in H₂-THF hydrates remain. The three key points with respect to the storage capacity of H₂ in the hydrate forms that need to be resolved are: (i) feasibility of occupancy of two H₂ molecules in one small hydrate cage; (ii) whether the large cages free of THF can be occupied by H₂ when the concentration of THF is lower than the stoichiometric value, i.e., 5.56 mol%; (iii) the viability of co-occupancy of THF and H₂ molecule in the same large cages. Also, the influence of hydrogen bonds between THF and cage H₂O molecules on hydrate stability and dynamic motion of H₂ and THF molecules in the small and large cages needs refinement. In this work, the classical MD simulation approach is applied to systematically analyze the abovementioned three factors by combining the dynamic motion of guest and water molecules with the interaction energy, as illustrated in **Figure 1**, which is crucial for the understanding of mechanisms underlying the storage capacity of H₂ in the hydrate form. The cage occupancy modes investigated in this work are based on those either suggested based on previous experimental observations or the ones that are reported viable in simulation work and under realistic conditions. Resolving the knowledge gaps articulated above will unlock the



mechanisms underlying the role of promoters in enabling H₂ storage as hydrates at moderate temperatures and pressures.

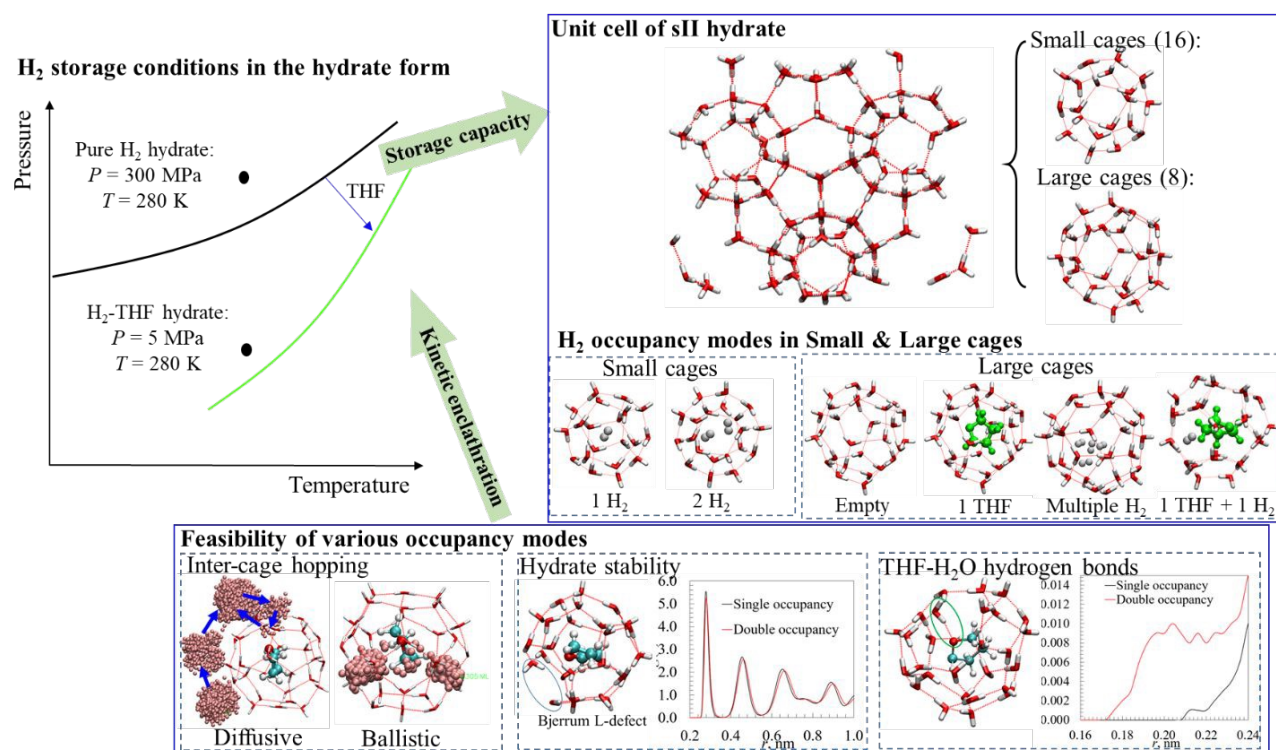


Figure 1. Schematic representation of the approach and methodology used to probe the mechanisms underlying the influence of THF on H₂ hydrate formation.

2. Methodology

Eight different configurations of binary H₂-THF hydrates with various occupancy modes of small and large cages are investigated in this work using MD simulations, as listed in **Table 1**. The occupancy modes of H₂ and THF in each configuration have been reported by either experiments or simulations. *HS1-THF8* configuration denotes that each of the 16 small cages in a unit cell of sII hydrates is occupied by one H₂ molecule, whereas all the 8 large cages are occupied by THF molecules. (9,18–20) *HS2-THF8* configuration is composed with 2 H₂ molecules occupying each small cage and all large cages occupied by THF. (15) By comparing these two configurations,



insights into the impact of various occupancy modes in the small cages on the dynamic motion and structural change are obtained. Configuration *HS2-THF7-HL4* represents the occupancy modes reported by Lee and co-workers,⁽¹⁵⁾ with 7 large cages in a unit cell occupied by THF and the other large one occupied by four H₂ molecules. Each of the small cages is occupied by two H₂ molecules. In *HS1-THF7*, each of the small cages accommodates one H₂ molecule and seven of the large cages in a unit cell are occupied by THF, which denotes the observations by Strobel and co-workers that H₂ molecules could not enter the large cages even when the concentration of THF was lower than the stoichiometric value of 5.56 mol%.⁽²⁰⁾ As reported by Kang and co-workers, the large cages free of THF could be occupied by 2-3 H₂ molecules, which are investigated in Configurations *HS1-THF7-HL2* and *HS1-THF7-HL3*, respectively.⁽²⁹⁾ *HS1-(THF+HL)8* denotes the occupancy mode suggested by Kaur and Ramachandran, i.e., each small cage accommodates one H₂ molecule and each large cage is co-occupied by one THF and one H₂ molecule.⁽³⁰⁾ Liu and co-workers reported a same occupancy mode as *HS1-(THF+HL)8*, but at a lower temperature condition of 140 K.⁽²⁸⁾ To investigate the impact of pressure and temperature conditions, Configuration *HS1-(THF+HL)8-140K* is designed, with the pressure and temperature of 0.1 MPa and 140 K, respectively. The pressure and temperature conditions in other configurations are 12 MPa and 270 K, respectively.

Table 1. Cage occupancy details of H₂ and THF in the small and large cages of a unit sII cell

Configuration *	Cage occupancy		Reference	Summary / Novel findings
	Small cages (16)	Large cages (8)		
HS1-THF8	Each by one H ₂	Each by one THF	9, 18-20	No inter-cage hopping of H ₂ molecules; Hydrate structure is stable but with low storage capacity
HS2-THF8	Each by two H ₂	Each by one THF	15	Active inter-cage hopping of H ₂ ; High probability of hydrogen bond formation between THF and H ₂ O; Reduced stability of hydrate structure; Slower decay of THF
HS2-THF7-HL4	Each by two H ₂	7 by THF, 1 by four H ₂	15	



				orientation due to stronger THF-H ₂ O interactions
HS1-THF7	Each by one H ₂	7 by THF, 1 empty	20	Thermodynamic stability of hydrate is enhanced by occupying three H ₂ molecules in the THF-free large cages
HS1-THF7-HL2	Each by one H ₂	7 by THF, 1 by two H ₂	29	
HS1-THF7-HL3	Each by one H ₂	7 by THF, 1 by three H ₂	29	
HS1-(THF+HL)8			30	The inter-cage motion of H ₂ molecules evolved from ballistic to diffusive regime with co-occupancy of THF and H ₂ in large cages
HS1-(THF+HL)8-140K ^a	Each by one H ₂	All by one THF + one H ₂	28	

Note: * HS, HL, and THF stand for H₂ molecules in the small cages, large cages, and occupancy of THF molecules in the large cages, respectively; ^a $P = 0.1$ MPa, $T = 140$ K.

The molecular dynamic simulation package of GROMACS 2020.6(31) is used. According to the sensitivity analysis on the size of simulation box, we did not detect a significant change in the pairwise energy between water and H₂ molecules in the large and small cages with the number of unit cells in the simulation box increasing from 2×2×2 to 5×5×5 (see **Figure S1**). Therefore, a simulation box containing 2×2×2 unit cells of sII hydrates is adopted in this work, with dimensions of 3.462×3.462×3.462 nm. The periodic boundary condition is employed in all three directions of the simulation box. The coordinates of O and H atoms of cage H₂O molecules determined by Takeuchi and co-workers(32) are applied in this work as the configuration of hydrate cages. The center of mass of H₂ (or H₂ clusters in the cases of multiple occupancy) and promoter molecules are placed in the center of the hydrate cages. The interaction between guest and water molecules are regulated by the force fields. In this work, the TIP4P/Ice potential function(33) is applied for H₂O molecules, which has been reported reliable in reproducing the phase boundary conditions of gas hydrates.(34) The three-site model developed by Alavi et al. is employed for H₂ molecules.(35) As for THF molecules, the general AMBER force field (GAFF) is adopted, with the partial charges and atom positions optimized.(12) The H₂O, H₂, and THF molecules are regarded as rigid and constrained using the LINCS algorithm in GROMACS. The parameters of the potential functions are summarized in **Table S2** in the *Supplementary Information*.



A typical simulation algorithm includes an energy minimization step performed to the initial simulation configuration, followed by 100 ps NPT ensemble on the optimized configuration to reach equilibrium state, which is further continued for 1.0 ns in the production stage. The Nose-Hoover thermostat and Berendsen barostat are used for the temperature and pressure coupling, respectively. A time step of 0.2 fs is applied in the equilibrium and production stages. The non-bonded interactions, e.g., van der Waals force and electrostatic forces are modeled using 12-6 Lennard-Jones (LJ) and Columbic models, respectively.(36)

3. Results and Discussions

3.1 Simulation validation

To validate the developed MD simulation models in this work, the distribution and average distance of H₂ molecules in the small and large cages of pure H₂ hydrates are compared with those in the literature. In the pure H₂ hydrates, the small cages are occupied by two H₂ and the large ones by four H₂ molecules, with a simulation pressure and temperature of 0.1 MPa and 100 K respectively, which is within the phase boundary conditions.(35) The initial configuration of the pure H₂ hydrates in the simulation box and that at the end of simulation time are illustrated in **Figure S2** in the *Supplementary Information*.

The two-dimensional (2D) density maps of H₂ molecule distribution in the small and large cages of pure H₂ hydrates are shown in **Figure 2**. The number density is truncated at 20 nm⁻² to highlight the distribution feature of H₂. The distribution of two H₂ molecules in each small cages separated into two regions away from the cage centers, while that for the four H₂ molecules accumulated in four regions, demonstrating a tetrahedral structure in three dimensions (see **Figure S3** in the *Supplementary Information*). With a closer observation of the 3D distribution profile of the four H₂ molecules in **Figure S3**, it is found that the vertices of the tetrahedral structure pointed



to the four hexagonal faces of the large cages, which is consistent with Burnham and co-workers.(37) Also, the average distance between the two H₂ molecules in the small cages was 2.5-2.6 Å and that between the four H₂ molecules in the large cage is 3.0-3.1 Å, consistent with previous studies.(21,27,28,37) Further, the lattice parameter obtained in this work (1.726 nm) is slightly higher than 1.699 nm obtained by Alavi and co-workers(35) at the same cage occupancy mode and simulation pressure-temperature conditions. This higher lattice parameter is a result of a different force field of H₂O, i.e., SPC/E applied by Alavi and co-workers. The H₂ distribution and hydrate cage structure results indicate reliability of the developed simulation model.

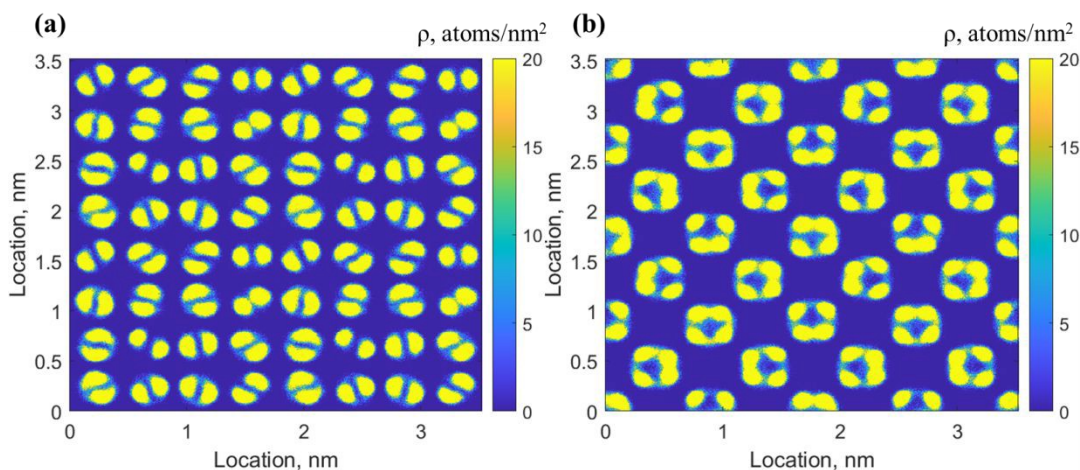


Figure 2. 2D density maps of H₂ molecules distribution in the small (a) and large (b) cages at $T = 100$ K, $P = 0.1$ MPa. The density maps are averaged over 1.0 ns of the simulation time.

3.2 Occupancy of two H₂ in small cages

The impact of different occupancy modes of H₂ in the small cages on the dynamic motion and distribution of H₂, THF and cage water molecules, e.g., probability of hydrogen bond formation between THF-oxygen (Ot) and water-hydrogen (Hw), pair-wise interaction energy between different molecules, and stability of hydrate structure is analyzed by comparing the simulation results of various configurations in **Table 1**. It is noted that the impact of H₂ occupancy in small



cages is not significantly influenced by the occupancy modes of H₂ and THF in the large cages. For instance, the findings by comparing results of Configurations *HS1-THF8* and *HS2-THF8* are similar to those by comparing results of Configurations *HS1-THF7-HL3* and *HS2-THF7-HL4*. Therefore, the simulation results of Configurations *HS1-THF8* and *HS2-THF8* are presented as an example to analyze the impact of occupancy mode of H₂ in the small cages.

Impact on dynamic motion of H₂ molecules

Occupancy of two H₂ in the small cages makes the translational motion of H₂ more active, as indicated by the increasing MSD in **Figure 3**. The MSD curve of H₂ molecules showed a steady profile when the small cages were occupied by one H₂ molecule, whereas it increased linearly at an occupancy of two H₂ molecules, with a diffusion constant of 1.7×10^{-8} cm²/s. This active motion might be attributed to the strong repulsive force between H₂ molecules. The pair-wise potential energy between H₂ was 1.415 kJ/mol with two H₂ molecules occupying the small cages. As a comparison, the interaction between H₂ molecules is attractive when the small cages are occupied by one H₂ molecules, with a pair-wise potential energy between H₂ molecules of -0.075 kJ/mol. The strong repulsive force between H₂ when the small cages are occupied by two H₂ molecules results in the migration of H₂ molecules from cage to cage. This inter-cage hopping is reflected on the 2D density maps of H₂ distribution, as indicated by the non-zero density in the region between cages denoted by the red circles in **Figure 4 (b)**. No inter-cage hopping of H₂ is detected when the small cages are occupied by one H₂ molecule, which is consistent with the observation by Frankcombe et al and Cao et al.(38,39)



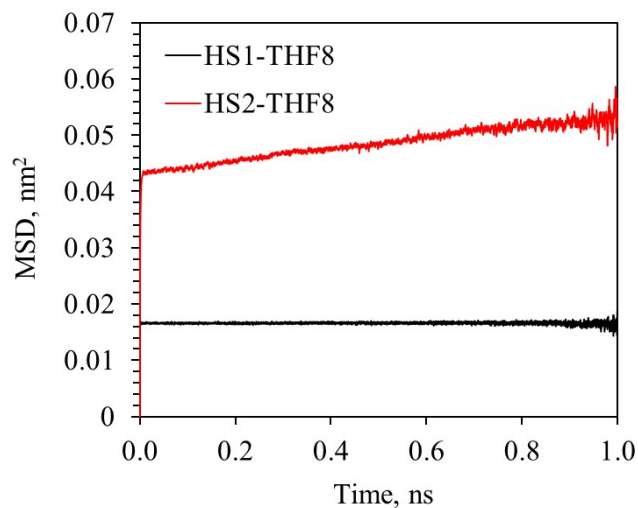


Figure 3. MSD curves of H_2 with small cages occupied by one H_2 (*HS1-THF8*) and two H_2 molecules (*HS2-THF8*).

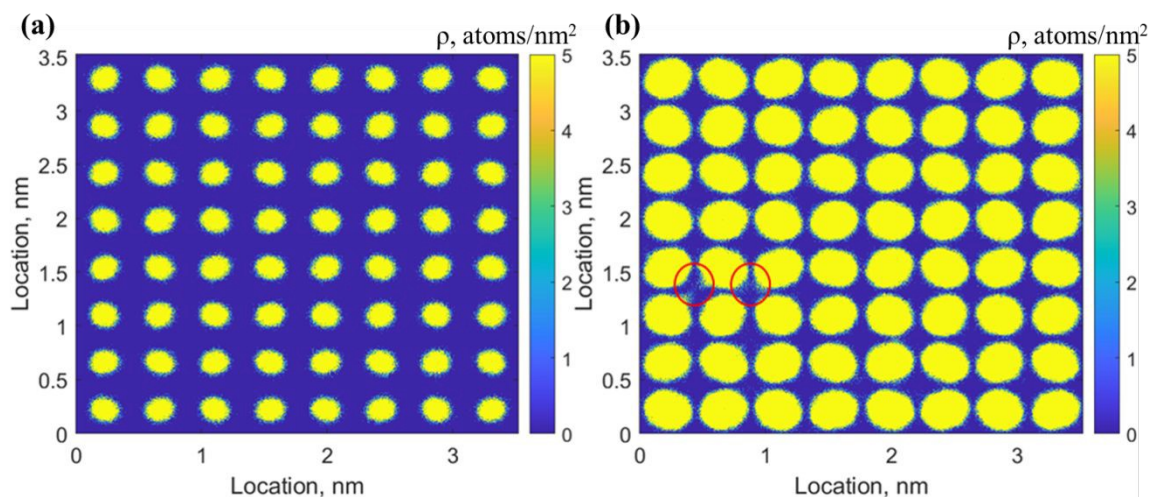


Figure 4. 2D density maps of H_2 distribution in Configurations *HS1-THF8* (a) and *HS2-THF8* (b).

With a closer look into the evolution of the H_2 migration, it is found that the migration may occur to only one of the H_2 molecules occupying a single cage, as shown in **Figure 5**. The **Figures 5 (a1) – 5 (a2)** and **5 (b1) – 5 (b2)** demonstrate the cumulative positions of the first and second H_2 molecules at $t = 0$ and $t = 1.0$ ns, respectively. It can be observed that the first H_2 molecule is



retained in the original cage (marked by the green label) throughout the simulation time of 1.0 ns. As for the second H₂ molecule, its inter-cage hopping extended to 4 cages in addition to the original one, with the trajectory described by the blue arrows. It is clearly shown that this H₂ molecule first migrated to a neighboring small cage and the inter-cage hopping from the original small cage to the first cage was through a pentagonal face. Then it continued migrating to two other neighboring small cages through the pentagonal faces. The fourth cage where this H₂ molecule appears is the center large cage which is stably occupied by a THF molecule throughout the simulation time. Two observations are noted in this process: *i*) H₂ migrates from cage to cage and the migration occurs mostly in the small cages; *ii*) H₂ hops into the large cage occupied by the THF molecule through a pentagonal face. For pure H₂ sII hydrates, it is suggested that the energy barrier for the H₂ hopping through a pentagonal face was much higher than a hexagonal one.(24,38,40) However, it is also observed that the H₂ migrated to a neighboring cage through a pentagonal face.(41) The temporary breaking of the hydrogen bonds in the pentagonal faces may facilitate inter-cage hopping of H₂ molecules and some of these hydrogen bonds restored after the migration.(42) Therefore, the observation of H₂ migration through pentagonal faces in Configuration *HS2-THF8* may be attributed to the distorted cage network by the newly formed hydrogen bond between the Ot-Hw in THF and cage water molecules.



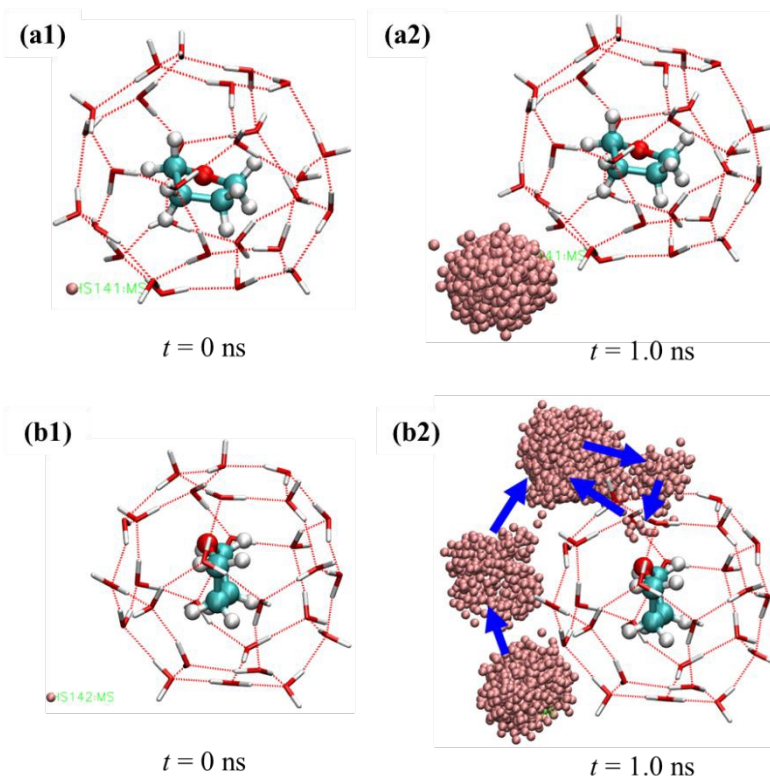


Figure 5. Cumulative positions of the first (a1 and a2) and the second (b1 and b2) H_2 molecule in one small cage of Configuration *HS2-THF8*.

Impact on stability of hydrate structure

The lattice parameter of hydrate cages increases by 1.03% when the number of H_2 molecules occupying the small cages increase from one to two, corresponding to a volume expansion of 3%, which is consistent with the conclusion by Koh and co-workers.⁽²⁵⁾ Even though an expansion of 3% in the volume of hydrate cages is not pronounced, it has a significant impact on the configuration of hydrate cages, as shown in the RDF curve of oxygen (Ow) pairs of water molecules in **Figure 6 (a)**. It is seen that the peaks become less sharp and shift to a higher radius when the small cages are occupied by two H_2 molecules, indicating a reduced stability of hydrate structure.⁽³⁶⁾ This reduced stability of hydrate structure may be attributed to the much higher



probability of hydrogen bonds formation between O in THF (Ot) and H in cage H₂O (Hw) molecules when the small cages are occupied by two H₂ molecules, as demonstrated in **Figure 6 (b)**. As suggested by Alavi and co-workers,^(43–45) hydrogen bonding between THF and cage water molecules is assumed when the distance of Ot-Hw is shorter than 2.1 Å. The hydrogen bond formed between THF and cage water molecules is marked by the black circle in **Figure 6 (b)**. The formation of hydrogen bond between THF and H₂O is accompanied by the distortion of hydrogen bond network in the hydrate cages, as shown in the blue circle in **Figure 6 (b)**, due to the formation of the Bjerrum L-defect, i.e., the absence of the hydrogen atom between two Ow atoms. It is believed that the guest-cage hydrogen bonds enhance the rotational dynamics of water molecules due to the reduced energy barrier for water rotation in the presence of the Bjerrum L-defect.⁽⁴⁶⁾

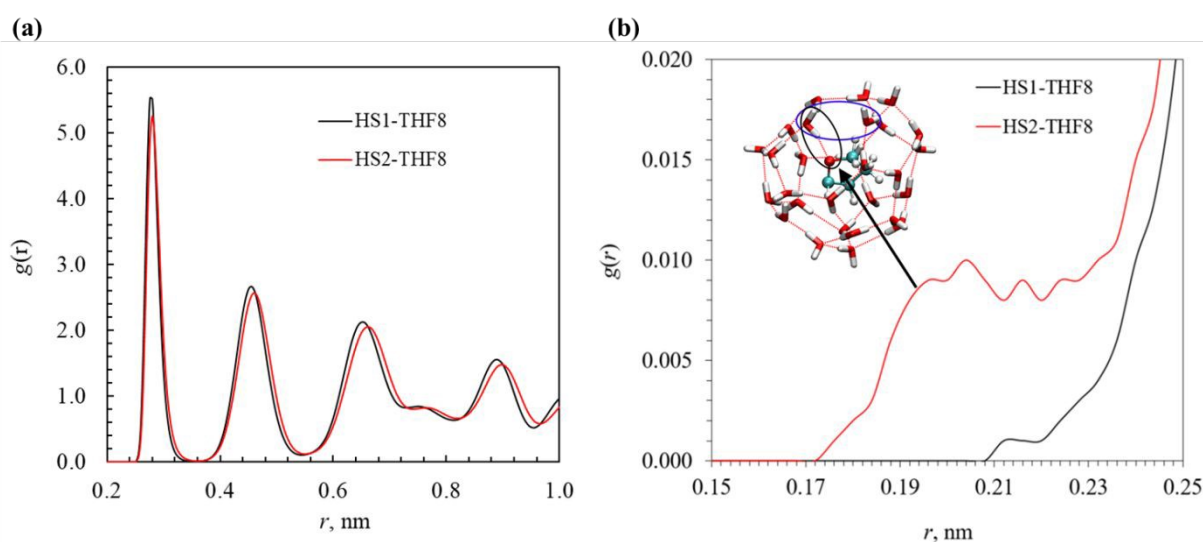


Figure 6. RDF curves of (a) Ow-Ow pairs in H₂O molecules and (b) Ot-Hw pairs in THF and H₂O molecules respectively with different occupancy mode of H₂ in the small cages.

The enlarged volume of hydrate cages by the increased number of H₂ in the hydrate cages also results in slightly more active translational motion of THF. Interestingly, it is found that the decay rate of THF orientation decreases with the increased occupancy of H₂ in the small cages (see



Figure S4 in *Supplementary Information*). Alavi and Ripmeester(47) suggested that the hydrogen bond between THF and cage H₂O molecules slows down the decay of THF orientation, which is consistent with the observation in this work by comparing the probability of hydrogen bond formation as demonstrated in **Figure 6(b)**. The probability of hydrogen bond formation between THF and cage H₂O molecules is greatly elevated when the small cages are occupied by two H₂ molecules, as indicated by the larger area below the RDF curves of O-H pairs around a distance of 0.21 nm. However, this relationship between THF-H₂O hydrogen bonding and the decay of THF orientation is not true when comparing the Configurations *HS1-(THF+HL)8* and *HS1-(THF+HL)8-140K*, which would be further discussed in Section 3.4.

The long-range electrostatic and short-range LJ interactions between water (W), H₂ in the small cages (HS), and THF molecules for in Configurations *HS1-THF8* and *HS2-THF8* is also compared. It is noted that the pair-wise interaction energy (i.e., sum of long-range and short-range interaction energies) is averaged by dividing with the number of molecules for convenient comparison. From the energetics perspective, it is seen that the strongest interaction is between water molecules due to the hydrogen bonding, followed by the interaction between THF and cage water molecules owing to the large size of THF and temporary hydrogen bonds between THF and water molecules. The pair-wise interaction energy between water, H₂, and THF with different occupancy modes in the small cages are compared in **Figure 7**. It is seen that the pair-wise interaction energy of W-HS, HS-THF, and THF-THF are negative due to the overall attractive forces. The absolute value of W-HS and THF-THF interaction energy decreased with the occupancy ratio of H₂ in the small cages due to the expansion of hydrate cages. The absolute value of HS-THF interaction energy is higher because the two H₂ molecules in one small cage are distributed closer to cage edge as shown in **Figure 4 (b)**, which shortened the distance between H₂



and THF. It is noticeable that the interaction energy between the H₂ molecules turned from negative to positive when the occupancy of H₂ in the small cage increased from one to two. As discussed in Section 3.2, this repulsive force between H₂ molecules facilitated the inter-cage hopping of H₂ molecules.

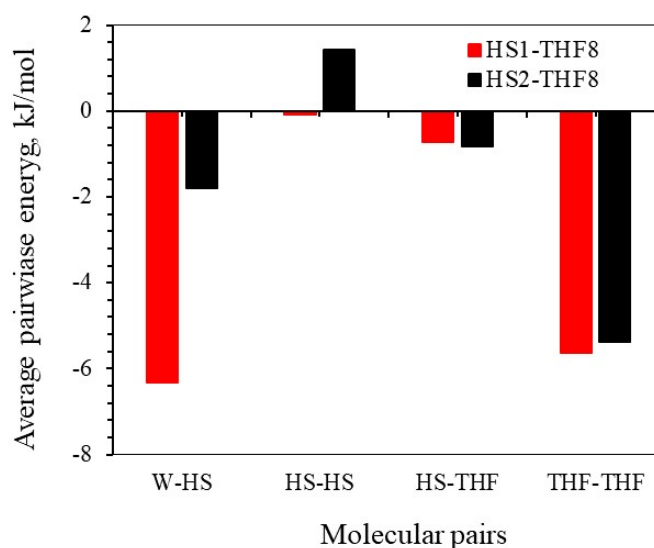


Figure 7. Pair-wise interaction energy between water (W), H₂ in the small cages (HS), and THF in Configurations *HS1-THF8* and *HS2-THF8*.

3.3 Occupancy of H₂ in THF-free large cages

Configurations *HS1-THF7*, *HS1-THF7-HL2*, and *HS1-THF7-HL3* are compared to investigate the impact of H₂ occupancy in the large cage free of THF. It is found that the number of H₂ molecules in the large cage has no noticeable influence on the probability of hydrogen bond formation between THF and cage water molecules. The RDF curves of Ow-Ow pairs in cage H₂O molecules are demonstrated in **Figure 8 (a)**. It is seen that the tetrahedral structure of the water cage is slightly enhanced with more H₂ occupying the THF-free large cages, which is consistent with the slightly less active translational motion of H₂O molecules (see **Figure S5** in



Supplementary Information). Therefore, the occupancy of three H₂ molecules in the THF-free large cages results in slightly more stable hydrate structure. The occupancy modes of H₂ in the THF-free large cages influence the distribution of H₂ molecules, as shown in **Figure 8 (b)**. It is seen that the H₂ molecules became more orderly distributed in the large cages with the number of H₂ molecules increased from two to three. When observing the 2D density maps of H₂ distribution in these two configurations, it is found that the tetrahedral feature of H₂ in the large cages becomes more pronounced when the number of H₂ molecules is increased from two to three, which may suggest that the tetrahedral distribution of H₂ in the large cages is favorable for the stability of hydrate structure.

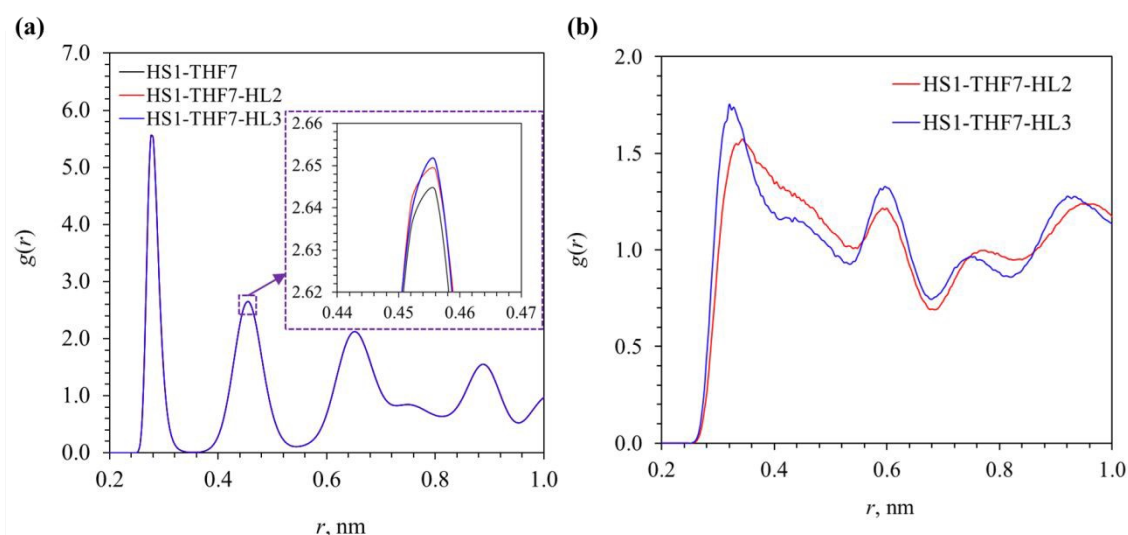


Figure 8. RDF of Ow-Ow pairs in cage H₂O molecules with various occupancy modes of H₂ in the THF-free large cages.

It is interesting to note that even though the occupancy of three H₂ molecules in the THF-free large cages contributes to the thermodynamic stability of hydrate structure, the tunability of H₂ storage in the THF-free large cages is not always captured in the experiments.⁽²⁰⁾ This observation may be because of the high energy barrier to kinetically enclathrate H₂ molecules into the large



cages. It has been found that the occupancy of H₂ in the THF-free large cages is feasible when ice powders(16) are used to form the H₂-THF hydrates or gas exchange approach is applied to form H₂ hydrates, e.g., enclathrating H₂ into hydrate cages by exchanging with N₂.(48,49) The potential mechanisms for these methods of enhancing the storage capacity of H₂ in the large cages may be due to the reduced energy barrier induced by the pre-existence of cavities in the ice powders and hydrate cages. Therefore, it is important to develop new promoters and techniques to overcome the kinetic barriers for H₂ molecules enclathrating in the promoter-free large cages, which is essential to enhance the storage capacity of hydrate-based hydrogen storage technology.

3.4 Co-occupancy of H₂ and THF in one large cage

It is seen in **Figure 5** that the probability of co-occupancy of H₂ and THF in one large cage is low based on the short time the H₂ molecule is retained in the center large cage, although co-occupancy of THF and small gas molecules, e.g., H₂ and He in the same large cages has been suggested viable according to the ab initio or grand canonical Monte Carlo simulations.(30,50) In this section, Configurations *HSI-(THF+HL)8* and *HSI-(THF+HL)8-140K* are applied to explore the feasibility of co-occupancy of H₂ and THF in one large cage by observing its impact on the dynamic motion of guest and cage water molecules and the hydrate cage stability.

The 2D density maps of H₂ molecules originally occupying the small and large cages are demonstrated in **Figure 9**. By comparing with the typical 2D density maps of one or two H₂ molecules in small cages in **Figure 4**, the 2D density maps of H₂ distribution in **Figure 9 (a)** suggested that some of the small cages are occupied by two H₂ molecules during the simulation. Given that all the small cages are originally occupied by one H₂ molecule, the extra one H₂ molecule migrates from the large cages, which can be confirmed in the 2D density maps of H₂ molecules originated from the large cages in **Figure 9 (b)**. In **Figure 9 (b)**, it is seen that the H₂



molecules originated from the large cages not only showed up in the small cages throughout the simulation time, but also demonstrated bar-like feature in the density map. With a closer observation of one H_2 molecule initially co-occupying a large cage with THF, it turns out this bar-like distribution of the H_2 molecule is a result of the tunneling behavior of H_2 through the hexagonal faces of the large cages, as shown in **Figure 10**. The example in **Figure 10** demonstrates that this H_2 molecule goes back and forth through one hexagonal face in the first 0.7 ns and then moves to another hexagonal face and continues the tunneling behavior to the end of simulation.

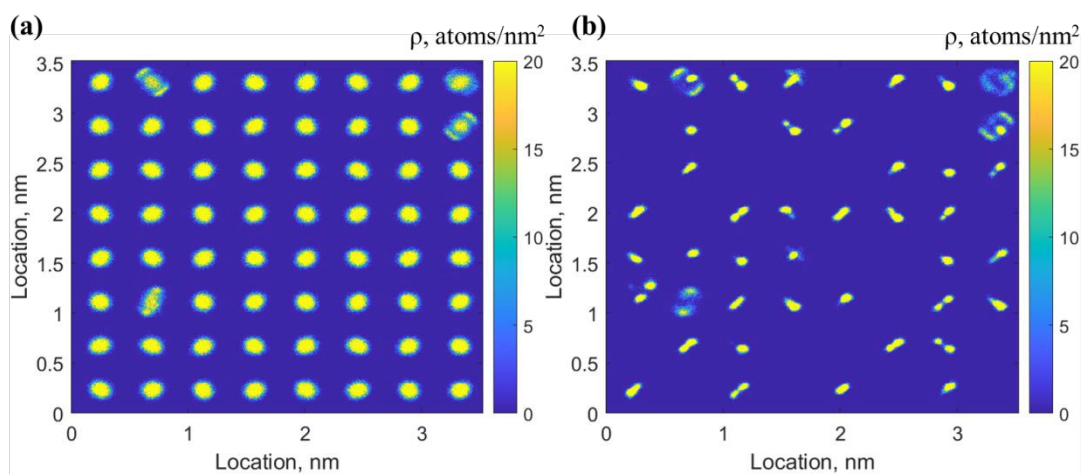


Figure 9. 2D density maps of H_2 distribution originally seated in the small (a) and large (b) cages in Configuration *HS1-(THF+HL)8*.

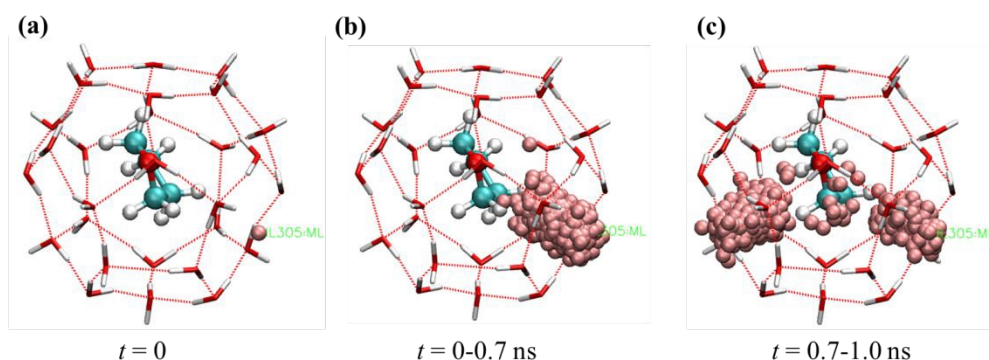


Figure 10. Evolution of the distribution of one H_2 molecule originated from a large cage in Configuration *HS1-(THF+HL)8*.



The impact of pressure and temperature conditions is explored using Configuration *HS1-(THF+HL)8-140K*, where the temperature and pressure are 140 K and 0.1 MPa, respectively. From the RDF curves of Ow-Ow pairs in cage H₂O molecules in **Figure 11**, it is seen that the hydrate cage structure is more stable at $P = 0.1$ MPa and $T = 140$ K compared with that at $P = 12$ MPa and $T = 270$ K. The peaks in the RDF curve not only became sharper, which indicates more orderly distribution of water molecules in the hydrate cages, they also shifted slightly to left due to the shrinkage of hydrate cages at lower temperature, which is consistent with the smaller lattice parameter (by 1.04%) at the lower temperature conditions.

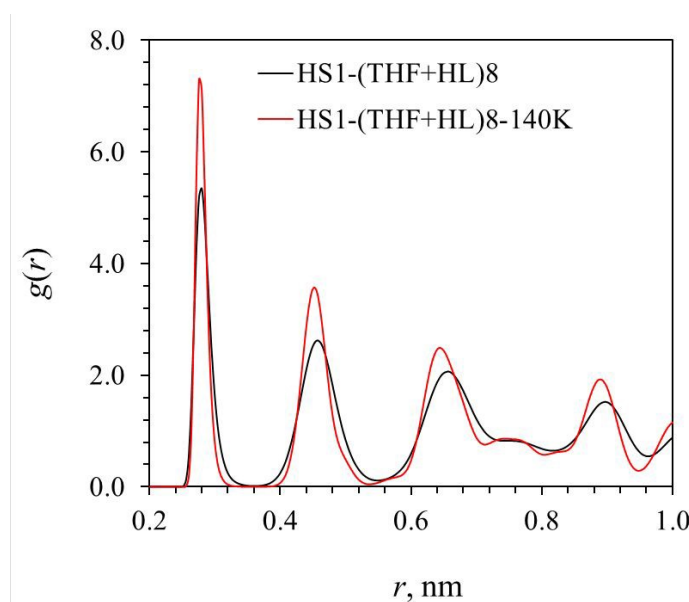


Figure 11. RDF curves of Ow-Ow pairs in cage H₂O molecules with co-occupancy of THF and H₂ in large cages at different pressure-temperature conditions (12 MPa and 270 K for the black curve; 0.1 MPa and 140 K for the red curve).

The impact of the lower temperature and lower pressure condition on the dynamic motion of H₂ molecules initially co-occupying the large cages with THF can be found in **Figure 12**. The MSD curves in **Figure 12 (a)** shows that the motion of H₂ molecules is greatly reduced when the



temperature is decreased to 140 K. However, different from the MSD curves of H₂ in Configuration *HS2-THF8* (see **Figure 3**, where the MSD increased linearly, indicating a diffusive motion), the MSD curves in **Figure 12** suggest a ballistic motion(51,52) of H₂ molecules ($\text{MSD} \propto t^2$), which is more pronounced for Configuration *HS1-(THF+HL)8-140K*. By observing the 2D density map of these H₂ molecules in **Figure 12(b)**, it is seen that the migration of H₂ molecules is featured with mostly tunneling behavior, which was also through the hexagonal faces as shown in **Figure S6** in the *Supplementary Information*. Given the fact that the ballistic motion feature is more noticeable in Configuration *HS1-(THF+HL)8-140K* (see **Figure 12 (a)**) and that the tunneling behavior is more remarkable in the migration of these H₂ molecules, it is indicated that the tunneling motion of H₂ molecules originally co-occupying the large cages with THF leads to the ballistic behavior in the MSD curves.

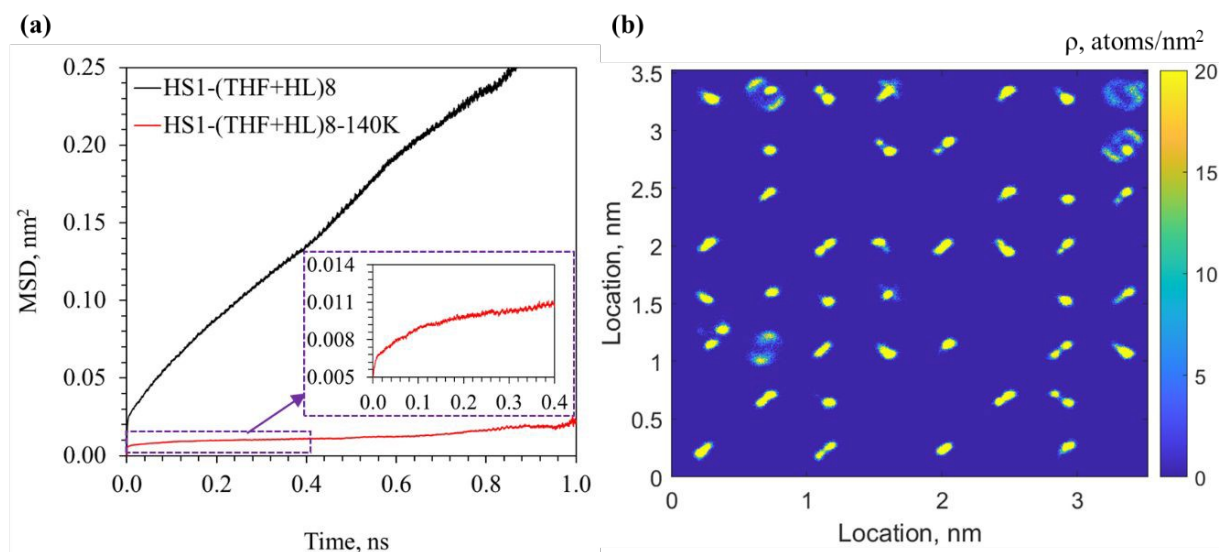


Figure 12. MSD curves of H₂ molecules originally co-occupying the large cages with THF (a) and 2D density map of H₂ distribution in Configuration *HS1-(THF+HL)8-140K* (b).



To further demonstrate the relation between different H₂ inter-cage migration features (i.e., diffusion and tunneling) and features in the MSD curves (i.e., $\text{MSD} \propto t$ and $\text{MSD} \propto t^2$), the simulation time of Configuration *HS1-(THF+HL)8* is prolonged to 50 ns, with the MSD curves and the 2D density maps of H₂ molecule distribution illustrated in **Figure 13**. It is shown that the $\text{MSD} \propto t^2$ relation is captured in the initial stage, ended up with a linear relation. By observing the 2D density maps of H₂ molecules at two different stages, i.e., in the first and last 2 ns, it is found that the probability of tunneling motion of H₂ molecules is greatly reduced in the last 2 ns (See **Figure 13 (b)**), which corresponds to the $\text{MSD} \propto t$ feature in the MSD curve. Therefore, it is concluded that the $\text{MSD} \propto t^2$ feature in the MSD curve of H₂ molecules results from the tunneling motion of H₂ molecules through the hexagonal faces, whereas the inter-cage hopping of H₂ molecules through diffusion leads to a linear feature in the MSD curves, which also applies to inter-cage hopping of H₂ molecules initially occupying the small cages (e.g., Configuration *HS2-THF8*). Also, the probability of hydrogen bond formation between THF and H₂O molecules is low when the ballistic motion dominates as shown in **Figure 14 (b)**, which means that the hydrate structure is less likely distorted in the ballistic regime. In contrast, when the hydrate structure is less stable, the ballistic motion of H₂ molecules in the initial stage will be gradually replaced by the diffusive motion. These results may indicate that the stable co-occupancy of H₂ and THF in one large cage is possible, accompanied by the ballistic motion of H₂ molecules. In other words, the ballistic motion is favorable for a higher storage capacity of hydrogen in the hydrate form, due to the possible co-occupancy of H₂ and THF in one hydrate cage. However, it requires a more stable hydrate structure, i.e., at higher pressure or lower temperature conditions.



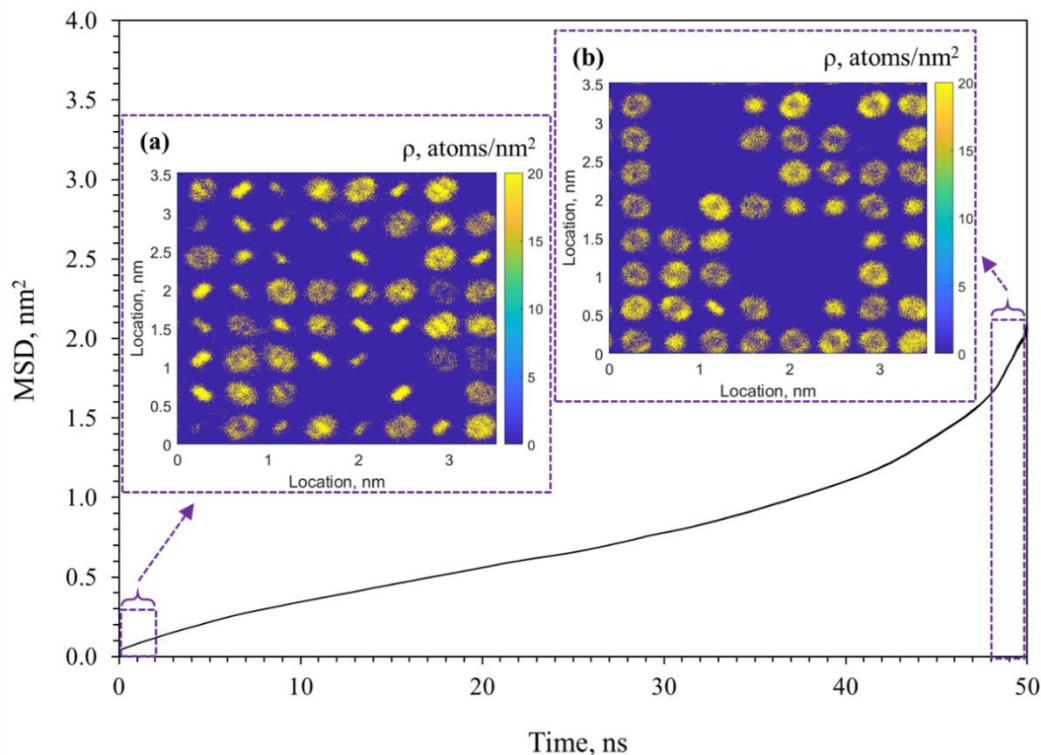


Figure 13. MSD curve of H₂ molecules originally co-occupying the large cages with THF and 2D density maps of H₂ distribution in the first 2 ns (a) and last 2 ns (b) with the simulation time of Configuration *HSI-(THF+HL)8* prolonged to 50 ns.

The impact of temperature and pressure on the decay rate of THF orientation is demonstrated in **Figure 14 (a)**. It is seen that the decay rate is significantly reduced in Configuration *HSI-(THF+HL)8-140K*. As discussed in Section 3.2, it is suggested that the slower decay of THF orientation in H₂-THF hydrates is due to the higher probability of hydrogen bond formation between THF and cage water molecules.⁽⁴²⁾ However, it is observed in **Figure 14 (b)** that the probability of hydrogen bond formation in Configuration *HSI-(THF+HL)8-140K* is greatly reduced. Therefore, the relationship between stronger hydrogen bonds of THF-H₂O and smaller decay rate of THF orientation is not valid. To explore the cause of slower decay of THF orientation,



the pair-wise interaction energy between THF and cage water molecules in different configurations is compared. It is found that the decay rate of THF orientation is regulated by the interaction energy between THF and H₂O molecules in general. For instance, the pair-wise interaction energy of THF-H₂O is -54.95 kJ/mol and -57.52 kJ/mol respectively for Configurations *HS1-(THF+HL)8* and *HS1-(THF+HL)8-140K*. The higher interaction force between THF and H₂O arises from the the short-range LJ interaction, which may be attributed to the shorter distance between THF and H₂O with the shrunk cages. Another example is that the pair-wise interaction energy of THF-H₂O for Configurations *HS1-THF8* and *HS2-THF8* is -55.45 kJ/mol and -56.37 kJ/mol, respectively. This stronger interaction between THF and H₂O results in slower decay of THF orientation (see **Figure S4**).

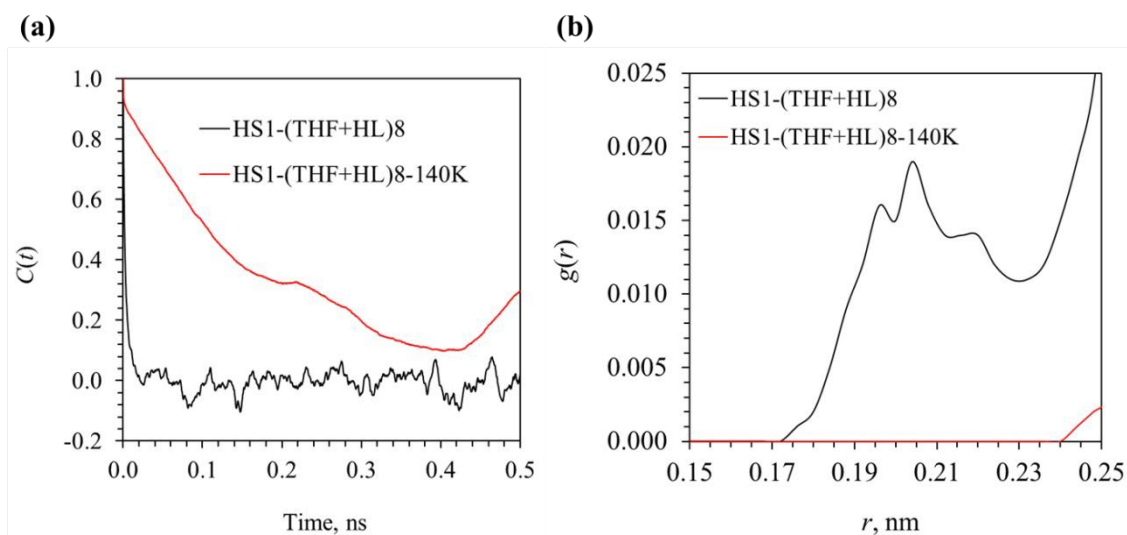


Figure 14. RACF curve of THF (a) and RDF of O-H pairs in THF and H₂O molecules (b).

4. Conclusions

The influence of different occupancy modes of H₂ in the small and large cages on the dynamic motion of gas and water molecules and hydrate structure stability to inform the storage capacity of H₂ in H₂-THF hydrate, is elucidated using classical molecular dynamics (MD) simulations.



Eight configurations with occupancy modes determined or suggested by previous experiments and simulation research are used to investigate the impact of (i) occupancy of two H₂ molecules in the small cages, (ii) occupancy of H₂ molecules in the THF-free large cages, and (iii) co-occupancy of H₂ and THF in one large cage. The following conclusions are drawn based on the configurations designed to address key knowledge gaps in hydrogen storage as hydrates:

(1) Occupancy of two H₂ molecules in small cages: The presence of two H₂ molecules in the small cages induced strong repulsive forces between the two H₂ molecules within a confined space, leading to active inter-cage hopping. The inter-cage migration occurred predominantly among small cages via pentagonal faces, with occasional migration into THF-occupied large cages. This phenomenon indicates that the distortion of the hydrate framework due to THF-water hydrogen bonding reduces the energy barrier for H₂ migration through pentagonal faces. Structurally, the occupancy of two H₂ molecules caused a 1.03% increase in lattice parameter, corresponding to a 3% expansion in cage volume. This expansion weakened the stability of the hydrate framework, as reflected by broader RDF peaks and disrupted hydrogen bonding among water molecules.

(2) Occupancy of H₂ molecules in THF-Free large cages: The occupancy of two or three H₂ molecules in large cages free of THF resulted in enhanced tetrahedral arrangements of H₂ molecules, particularly when three H₂ molecules were present. The tetrahedral distribution is energetically favorable and contributes to slightly improved hydrate stability. The findings suggest that higher occupancy in THF-free large cages can stabilize the hydrate framework, albeit the experimental realization of such configurations may face kinetic challenges. Approaches such as using ice powders or gas exchange methods could overcome these barriers by reducing the energy required for enclathration.



(3) Co-occupancy of H₂ and THF in large cages: The inter-cage migration of H₂ molecules when co-occupying large cages with THF demonstrated two motion features, i.e., ballistic motion ($\text{MSD} \propto t^2$) due to the tunneling behavior in the initial stage and diffusive motion ($\text{MSD} \propto t$) in the late stage. Prolonged simulations confirmed that the ballistic regime is associated with greater hydrate stability due to reduced hydrogen bond distortion. The probability of hydrogen bond formation between THF and H₂O is low when the tunneling migration of H₂ molecules is dominant in the inter-cage hopping, which may indicate that the tunneling migration behavior is enabled by a more stable hydrate structure, and is more favorable for achieving a higher storage capacity of H₂ hydrates due to the co-occupancy mode. The decay rate of THF orientation is regulated by the interaction energy between THF and cage water molecules. A stronger interaction between THF and H₂O molecules leads to slower decay of THF orientation.

By systematically analyzing the dynamic motion of H₂ and THF molecules and their interactions with the hydrate framework, this study bridges longstanding knowledge gaps in understanding hydrogen storage capacity and hydrate stability. The insights into the effects of occupancy modes on molecular dynamics and structural integrity provide valuable guidance for optimizing hydrate-based hydrogen storage systems under moderate temperature and pressure conditions, paving the way for advancing sustainable hydrogen storage solutions that align with the goals of a sustainable energy future.

Author contributions

Ruyi Zheng: writing – original draft, concept, data curation, formal analysis, investigation, methodology, validation; Sohaib Mohammed: writing – review & editing, formal analysis, methodology, validation; Yang Jia: writing – editing, validation; Rituparna Hazra: writing –



review & editing, methodology; Greeshma Gadikota: writing – concept, review & editing, supervision, conceptualization, resources, formal analysis.

Conflicts of interest

G. G. is the co-founder of Carbon To Stone, a company focused on commercializing technologies for industrial carbon management.

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Data Availability Statement

The data supporting this article have been included in the main manuscript and as part of the Supplementary Information.

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