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# Theoretical study on lithium storage performance of V-doped $\text{Ti}_2\text{CO}_2$ MXene†

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With the increasing application of lithium-ion batteries, the demand for high energy density, high-rate performance and high stability lithium-ion batteries is becoming more and more urgent.  $\text{Ti}_2\text{CO}_2$  MXene, as a two-dimensional material with multilayer atomic structure and multiple active sites, has great advantages in lithium-ion battery electrode materials. However, the original  $\text{Ti}_2\text{CO}_2$  MXene has been unable to meet the requirements of lithium-ion batteries due to its semiconductor properties. Doping is an effective means to regulate the conductivity and electrochemical properties of  $\text{Ti}_2\text{CO}_2$  and improve the capacity of lithium-ion batteries and other energy storage devices. Hence, we use first-principles calculations to study the effect of V atom doping on the adsorption and diffusion of Li on the MXene surface. The density of states (DOS) and partial density of states (PDOS) of  $\text{TiVCO}_2$  and  $\text{Ti}_2\text{CO}_2$  MXene indicated the transition of their conductive types from semiconductors to conductors. In addition, we observed that  $\text{TiVCO}_2$  has higher electrical conductivity and ion transport speed than the original  $\text{Ti}_2\text{CO}_2$  MXene, and at the same time, Li atoms can be adsorbed well on the surface of MXene and show a lower diffusion energy barrier. Therefore,  $\text{TiVCO}_2$  is expected to become the anode material for the next generation of lithium-ion batteries and has good lithium storage performance.

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## 1. Introduction

Due to the irreversible consumption of non-renewable energy sources and the advancements in the portable electronics and electric vehicle industries, the increasing demand for energy storage devices has spurred extensive research interest in the development of new storage systems.<sup>1</sup> Lithium-ion batteries (LIBs) have garnered significant attention as the primary power source for portable electric vehicles and implantable medical systems, owing to their small size and high energy density.<sup>2–4</sup> However, the safety and cost issues associated with LIBs have hindered their development. Consequently, the search for or development of new materials with high-rate performance, enhanced safety, improved cycling performance and stability, and higher capacity represents a major challenge for LIBs. In this context, the discovery of graphene<sup>5</sup> and its exceptional properties has led to the exploration of numerous two-dimensional (2D) materials, such as germanene,<sup>6</sup> silicene,<sup>7</sup> hexagonal boron nitride (h-BN),<sup>8</sup> black phosphorus,<sup>9</sup> and

transition metal dichalcogenides (TMDC).<sup>10</sup> While most 2D materials have predominantly been applied in electronic devices due to their semiconductor characteristics, such as p–n junctions and field-effect transistors, there has also been research into their application in energy storage devices like supercapacitors, thermoelectric devices, and batteries. Notably, lithium-ion batteries based on 2D materials have garnered increasing attention due to their rapid ion diffusion properties. On the other hand, MXenes, as new members of the 2D materials family, have been demonstrated to be promising candidates for electrode materials in lithium-ion batteries and lithium-ion capacitors.<sup>11</sup> Furthermore, to date, numerous different types of MXene materials have been synthesized through chemical exfoliation methods, and it is anticipated that many MXene materials are stable and experimentally obtainable.<sup>12</sup>

MXene was a novel two-dimensional material that had recently been recognized as one of the best candidate anode materials for lithium-ion batteries. MXenes were derived from MAX phases with a closed-packed hexagonal structure and  $P6_3/mmc$  space group symmetry.<sup>13</sup> MXene nanosheets were typically described by the general formula “ $\text{M}_{n+1}\text{X}_n\text{T}_x$ ”, where “M” represents a transition metal, “X” denotes carbon and/or nitrogen, “ $\text{T}_x$ ” stands for surface terminal groups such as –F, –O, and –OH, and  $n$  can take values from 1 to 3.<sup>14</sup> Currently, extensive theoretical and experimental research has been conducted on MXene materials. Their electronic structure,

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electrical conductivity, surface chemistry, and electrochemical properties have been widely studied. The exceptional conductivity, unique layered structure, ultra-high specific surface area, and excellent ion transport capabilities of MXenes have been well-documented. However, most research to date has focused predominantly on single-component MXene materials. There was relatively little research on MXene electrode materials that incorporate transition metal impurity atoms.

Shunsuke *et al.*<sup>15</sup> studied the formation energy of Na/Li adsorption on  $\text{Ti}_2\text{CT}_x$  ( $T = \text{O}, \text{OH}, \text{F}$ ) sheets, showing that the Li- $\text{Ti}_2\text{CO}_2$  system was more stable than others. Michael Ashton *et al.*<sup>16</sup> reported a reversible capacity of  $192 \text{ mA h g}^{-1}$  and a diffusion barrier of  $0.63 \text{ eV}$  for  $\text{Ti}_2\text{CO}_2$  MXene through systematic density functional theory calculations. Yierpan Aierken *et al.*<sup>17</sup> found, based on first-principles calculations, that the charge transfer of Li on mono- $\text{Ti}_2\text{CO}_2$  is  $0.912e$ , which is higher than that of most MXene materials, indicating strong binding between Li atoms and the  $\text{Ti}_2\text{CO}_2$  monolayer, resulting in stronger ion interactions in the  $\text{Ti}_2\text{CO}_2$ -based system. Additionally, theoretical studies show that  $\text{Ti}_2\text{CO}_2$  has a low migration barrier ( $0.27 \text{ eV}$ ) and a high theoretical weight capacity ( $348.9 \text{ mA h g}^{-1}$ ). However, according to previous studies and our own calculations,  $\text{Ti}_2\text{CO}_2$  MXene behaves as a semiconductor, which makes it unfavorable for use as a negative electrode material for lithium-ion batteries.

In this context, this paper constructed a doped MXene material  $\text{TiVCO}_2$  based on the geometric lattice structure of  $\text{Ti}_2\text{CO}_2$  using the Special Quasi-random Structure (SQS) method in the ATAT software. First-principles calculations were used to systematically study the electronic properties and electrochemical performance of V-doped  $\text{Ti}_2\text{CO}_2$  MXene materials. To understand the electronic conductivity of  $\text{TiVCO}_2$  MXene materials, the total density of states (DOS) and partial density of states (PDOS) of  $\text{Ti}_2\text{CO}_2$  and  $\text{TiVCO}_2$  MXene materials were calculated. At the same time, the adsorption energy and diffusion barriers of Li atoms on the surface of MXene materials were also calculated to study their adsorption and migration properties. Furthermore, this paper systematically studies the open-circuit voltage (OCV), binding energy, and maximum theoretical capacity of Li on the surface of MXene using density functional theory (DFT). By comparing the electrochemical performance data with the original  $\text{Ti}_2\text{CO}_2$  MXene materials, this paper explores the feasibility of  $\text{TiVCO}_2$  MXene as an anode material for lithium-ion batteries. The theoretical calculation data indicate that  $\text{TiVCO}_2$  MXene materials are promising candidates for new lithium-ion battery anode materials.

## 2. Calculation method

Using density functional theory (DFT), all calculations were performed with the Vienna *ab initio* simulation package (VASP).<sup>18–20</sup> The Perdew–Burke–Ernzerhof (PBE) functional within the Generalized Gradient Approximation (GGA) was used to describe the interactions between nuclei and electrons.<sup>21,22</sup> The valence electronic states of  $3s^23p^63d^34s^1$  for Ti,  $3s^23p^64s^23d^3$  for V,  $2s^22p^2$  for C,  $2s^22p^4$  for O, and  $1s^22s^1$  for Li were taken into account. The mcsqs script in the Alloy Theoretic

Automated Toolkit (ATAT) software was used to generate the Special Quasi-random Structure (SQS) of V-doped MXene materials. An energy cutoff of  $520 \text{ eV}$  was used for the plane wave expansion of the electronic wave functions. All calculations used a  $3 \times 3 \times 1$  Monkhorst–Pack  $k$ -point grid. The convergence criteria for atomic forces and energy were set to  $0.02 \text{ e \AA}^{-1}$  and  $10^{-5} \text{ eV}$ , respectively. To avoid interactions due to the use of periodic boundary conditions, a real space gap of approximately  $15 \text{ \AA}$  was set between adjacent slabs. Considering the weak interaction systems we are dealing with, van der Waals interactions were considered using the semi-empirical DFT-D2 method. Additionally, the Climbing-Image Nudged Elastic Band (CI-NEB) method implemented in VASP was used to determine the diffusion barriers of Li atoms on MXene surfaces.

The adsorption energy ( $E_{\text{ad}}$ ) of Li atoms on the MXene surface is defined as follows:<sup>23</sup>

$$E_{\text{ad}} = E(\text{MXene-Li}) - E(\text{MXene}) - E(\text{Li})$$

where  $E_{\text{MXene-Li}}$  represents the total energy of the system with a single adsorbed Li atom;  $E_{\text{MXene}}$  is the total energy of the pristine MXene monolayer;  $E_{\text{Li}}$  is the total energy of a single Li atom in the bulk bcc phase.

The open-circuit voltage (OCV) can be calculated as:<sup>24</sup>

$$\text{OCV} = \frac{E_{\text{MXene}} + xE_{\text{Li}} - E_{\text{MXene-xLi}}}{x}$$

where  $E_{\text{MXene-xLi}}$  is the total energy of the system with multiple adsorbed Li atoms, and  $x$  is the number of Li atoms adsorbed on the surface.

The binding energy per Li atom on the MXene surface is defined as:

$$E_{\text{b}} = \frac{E_{\text{MXene-xLi}} - E_{\text{MXene}} - xE_{\text{Li}}}{x}$$

where  $E_{\text{MXene-xLi}}$  is the total energy of the system with multiple adsorbed Li atoms, and  $x$  is the number of Li atoms adsorbed on the surface.

The theoretical capacity is defined as:<sup>25</sup>

$$C = \frac{nzF}{(M_{\text{MXene}} + nM_{\text{Li}})}$$

where  $n$  is the number of adsorbed Li atoms,  $z$  is the valence state of Li,  $F$  is the Faraday constant ( $26\,801 \text{ mA h mol}^{-1}$ ), and  $M_{\text{MXene}}$  and  $M_{\text{Li}}$  are the molar masses of MXene and Li, respectively.

## 3. Results and discussion

### 3.1 Structure models

First, utilized the “mcsqs” script in the Alloy Theoretic Automated Toolkit (ATAT) software to generate special quasi-random structures (SQS) for V-doped MXene materials. In the ATAT software, stack five layers in the sequence of O–M (Ti, V)–C–M–O to establish a monolayer geometric model with  $\text{TiVCO}_2$  structure and electronic properties. In this work, a total of 26 different structures were predicted (Fig. S1†), the formation energy of different structures was calculated (Fig. S2†), and the



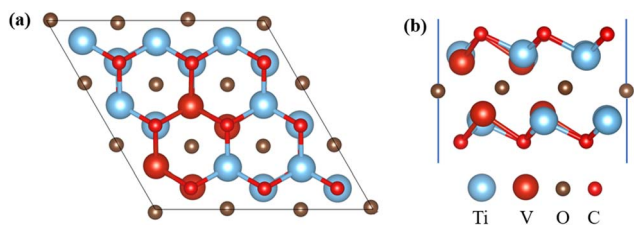


Fig. 1 Geometric model of TiVCO<sub>2</sub>. Top (a) and side (b) view.

most stable structure was selected. This model, based on the Ti<sub>2</sub>CO<sub>2</sub> structure, replaced four Ti atoms in two layers of Ti atoms with V atoms, resulting in optimized lattice constants of  $a = b = 9.0127 \text{ \AA}$  for TiVCO<sub>2</sub>. To investigate Li adsorption on its surface, we constructed a  $3 \times 3 \times 1$  supercell consisting of 45 atoms, as depicted in Fig. 1.

### 3.2 Electronic structural properties

The electronic conductivity of a material is related to its electronic properties. In order to understand the effect of vanadium doping on electronic properties, the DOS and PDOS of Ti<sub>2</sub>CO<sub>2</sub> and TiVCO<sub>2</sub> were calculated, as shown in Fig. 2. Among them, the Fermi level was set to 0 eV, which is represented by a black dotted line in the figure. The total DOS is represented by gray lines, and the PDOS of Ti, V, C, and O are represented by green, light blue, dark blue, and brown respectively. As can be seen from the figure, the DOS at the Fermi level of Ti<sub>2</sub>CO<sub>2</sub> is zero, while the DOS at the Fermi level of TiVCO<sub>2</sub> after incorporating vanadium is not zero, which indicated that the material has metallic conductivity after incorporating vanadium. This is mainly due to the contribution of the d orbital of V to the density of states at the Fermi level. Transformed from a semiconductor into a conductor, the conductivity was significantly enhanced. In addition, the unpaired electrons of doping atoms will cause spin polarization of the material, which is manifested as spin-up and spin-down DOS diagram asymmetry. From

Fig. 2b, it can be seen that the spin-up and spin-down DOS diagrams are asymmetric at energies of  $-4.6 \text{ eV}$ ,  $-3.1 \text{ eV}$ ,  $0.6 \text{ eV}$ , and  $2.2 \text{ eV}$ . What is more noteworthy was that in the energy range of  $-6 \text{ eV}$  to  $-4 \text{ eV}$ , the PDOS of the C 2p orbital and the V 3d orbital overlap significantly, which indicated that the C 2p orbital and the V 3d orbital are highly hybridized, forming a V–C bond. In the conduction band range, the peak at an energy of  $0.7 \text{ eV}$  was mainly contributed by the V 3d orbital, which showed that the incorporation of vanadium atoms makes a huge contribution to the improvement of the conductivity of Ti<sub>2</sub>CO<sub>2</sub>. In addition, it can be seen from the charge density difference diagram (Fig. S3†) of TiVCO<sub>2</sub> after V doping that the introduction of V has an effect on the charge distribution of MXene, which may be the reason for the change of its electronic conductivity. According to the calculation results of DOS, TiVCO<sub>2</sub> has metallic conductivity, which may make it suitable as an anode material for lithium-ion batteries.

### 3.3 Adsorption sites and adsorption energy

As an electrode material, the electrochemical performance of TiVCO<sub>2</sub> is closely related to the adsorption and diffusion behavior of surface Li atoms. First, it is necessary to confirm the optimal adsorption position of Li atoms on the surface of MXene material. For general Ti<sub>2</sub>CO<sub>2</sub> materials with only one transition metal and one functional group, there are usually three non-equivalent adsorption positions, as shown in Fig. 3a, which are the top of the C atom (fcc) and the top of the Ti atom (hcp). and the top of the oxygen atoms.

Similarly, MXene materials with multiple transition metals can also be classified according to the above positions. However, since the incorporation of vanadium atoms destroys the symmetry of Ti<sub>2</sub>CO<sub>2</sub>, the surface environment of V-doped Ti<sub>2</sub>CO<sub>2</sub> was more complex than that of the original Ti<sub>2</sub>CO<sub>2</sub> MXene, and the adsorption positions directly above Ti, C, and O atoms were no longer equivalent. Therefore, for TiVCO<sub>2</sub>, it was necessary to consider the possibility of Li atom adsorption

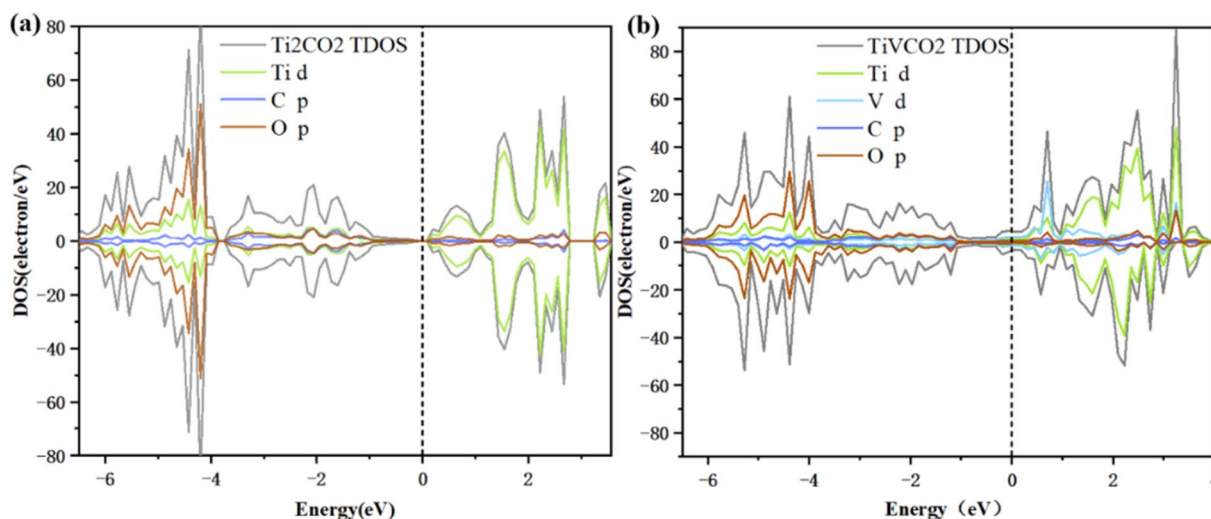


Fig. 2 Total density of states (DOS) and partial density of states (PDOS) of (a) Ti<sub>2</sub>CO<sub>2</sub> and (b) TiVCO<sub>2</sub>. The Fermi level was set to 0 eV.





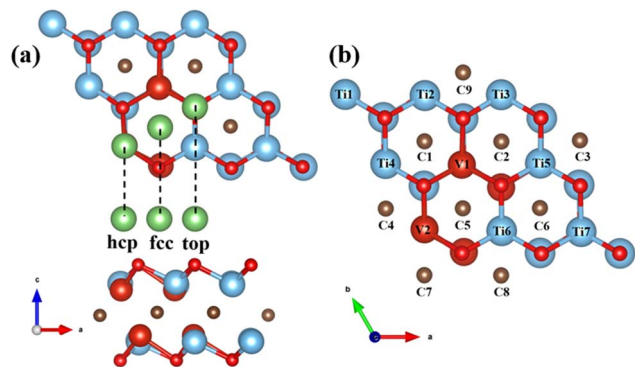


Fig. 3 (a) Top view and side view of three possible lithium-ion adsorption sites on MXene, (b) possible lithium-ion adsorption sites on TiVCO<sub>2</sub>.

Table 1 Adsorption sites and adsorption energies of Li atoms on the surface of TiVCO<sub>2</sub> MXene

Adsorption site	Adsorption energy (eV)	Adsorption site	Adsorption energy (eV)
C1	−2.44	Ti1	−1.93
C2	−2.45	Ti2	−2.03
C3	−2.16	Ti3	−2.04
C4	−2.44	Ti4	−2.00
C5	−2.57	V1	−2.45
C6	−2.24	Ti5	−1.93
C7	−2.45	V2	−2.45
C8	−2.25	Ti6	−1.95
C9	−2.30	Ti7	−2.03

directly above each atom. However, due to the particularity of the position directly above the oxygen atom, when the potential around the oxygen atom is not symmetrical, it is usually unstable as an adsorption position. Therefore, the position directly above the oxygen atom was not taken into account in this article. Other adsorption positions are shown in the Fig. 3b.

The calculation results are shown in Table 1. The adsorption energy of Li atoms on the surface ranges from −1.92 to −2.57 eV, all of which are negative values. The more negative the adsorption energy, the more stable the system between the adsorbed atoms and the surface. Therefore, Li atoms can be stable adsorbed on the MXene surface. It should be emphasized that the adsorption energy of Li atoms above C atoms is greater than that of transition metal atoms, indicating that Li atoms will be preferentially adsorbed at the fcc position. This result is similar to the adsorption performance of pristine MXene–Li ions. The reason for this phenomenon may be that the pore size of the “fcc” site is larger, which is more suitable for adsorbing Li ions.

In order to describe the adsorption mechanism in more detail, the charge transfer between them was analyzed using charge difference and Bader charge analysis. The charge difference calculation formula is:

$$\Delta\rho = \rho(\text{MXene-Li}) - \rho_{\text{MXene}} - \rho_{\text{Li}}$$

in the formula,  $\rho_{\text{MXene-Li}}$ ,  $\rho_{\text{MXene}}$  and  $\rho_{\text{Li}}$  are the charge densities of Li-adsorbed TiVCO<sub>2</sub> MXene, original TiVCO<sub>2</sub> MXene and single Li atom respectively. The charge difference diagram of a single Li atom adsorption system is shown in Fig. 4a–d and blue and yellow areas represent charge loss and charge accumulation, respectively. The charge transfer between Li atoms and the adsorption surface was calculated using the formula, and the specific charge transfer amount was obtained through Bader charge analysis. The charge transfer amounts of Li atoms adsorbed on the four adsorption sites of C5, Ti3, V1, and V2 are 0.911, 0.923, 0.917, and 0.916 $|e|$  respectively. In Fig. 4a–d, it can be clearly seen that the charges above the Li atoms are transferred to the material surface and localized between the Li atoms and the material surface, with a tendency to form bonds. Therefore, the differential charge and Bader charge analysis results showed that Li is adsorbed on the surface of TiVCO<sub>2</sub> MXene through chemical bonds.

### 3.4 Diffusion of lithium ions on MXene surface

The migration rate of Li atoms is an important criterion for evaluating the rate performance of electrode materials, and the migration rate of Li atoms on electrode materials depends on the diffusion energy barrier between the two most stable adsorption

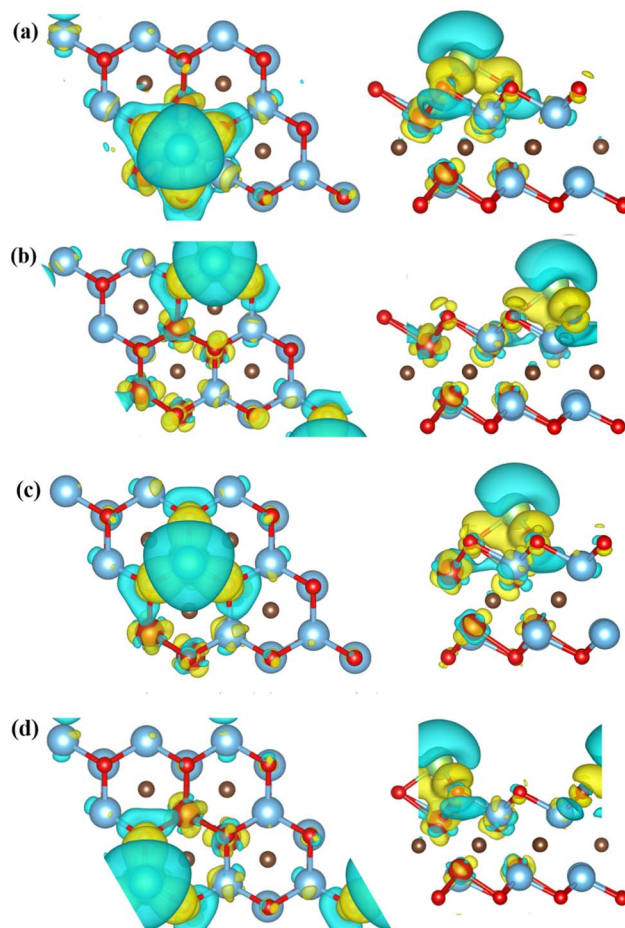


Fig. 4 Differential charge diagrams of Li atoms adsorbed on (a) C5, (b) Ti3, (c) V1, (d) V2 respectively.



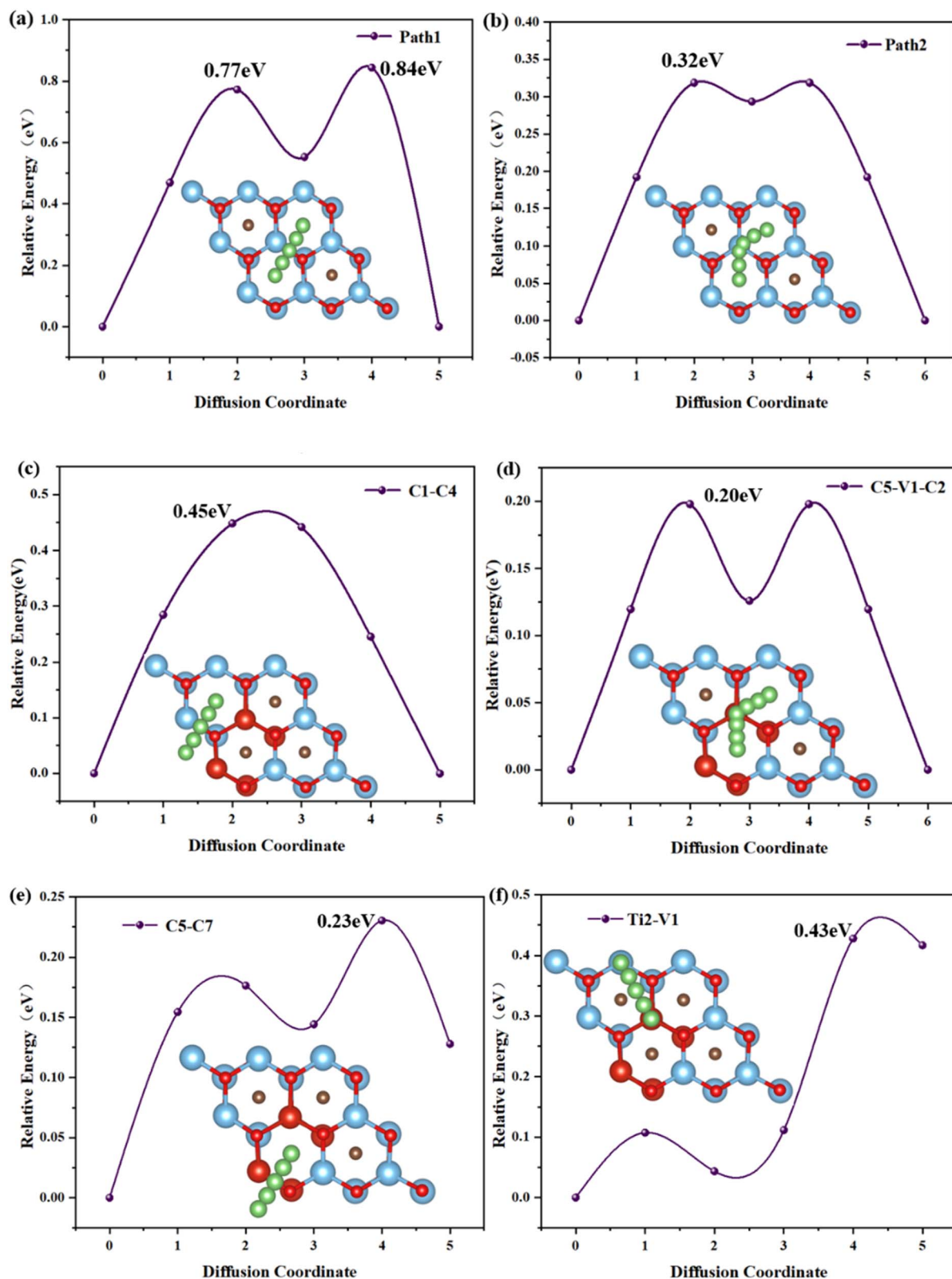


Fig. 5 Diffusion energy barrier and path of lithium ions on the surface of  $\text{Ti}_2\text{CO}_2$  (a and b) and  $\text{TiVCO}_2$  (c-f) monolayers.

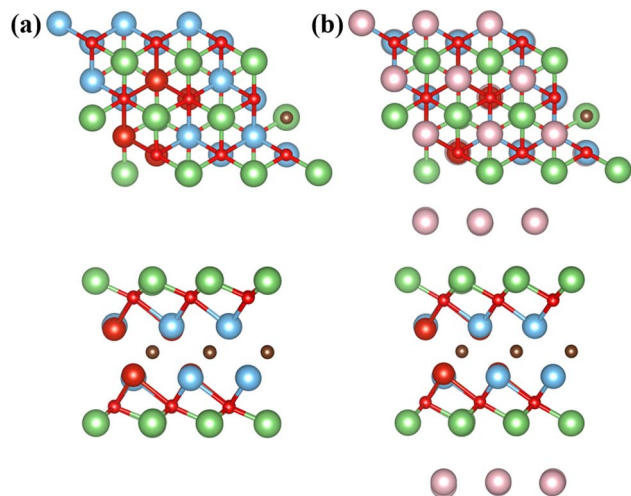


Fig. 6 Top view and side view of optimized MXene, (a) single layer Li and (b) double layer Li on both sides. The green and pink balls represent the first and second layers of lithium atoms respectively.

sites of Li atoms. Considering that the incorporation of vanadium atoms results in a complex surface potential environment of  $\text{TiVCO}_2$  MXene, the article considered the diffusion paths of multiple Li atoms on the  $\text{TiVCO}_2$  surface and obtained four representative diffusion paths through optimization calculations. As shown in Fig. 5c–f, according to the positions marked previously, the diffusion paths are  $\text{C1} \rightarrow \text{C4}$ ,  $\text{C5} \rightarrow \text{V1} \rightarrow \text{C2}$ ,  $\text{C5} \rightarrow \text{C7}$  and  $\text{Ti2} \rightarrow \text{V1}$  respectively. For comparison, the article also considers the diffusion path of Li atoms on the original  $\text{Ti}_2\text{CO}_2$  MXene material. The diffusion path was selected along the high symmetry lines between energetically favorable adsorption sites on the surface, respectively path1 and path2, as shown in Fig. 5a and b. It can be found through a calculation that the diffusion energy barrier of Li in the migration path designed in this article was in the range of 0.20–0.45 eV. Among them, the  $\text{C5} \rightarrow \text{V1} \rightarrow \text{C2}$

diffusion path has the lowest energy barrier (0.20 eV), and the  $\text{C1} \rightarrow \text{C4}$  diffusion path has the highest energy barrier (0.45 eV). For original  $\text{Ti}_2\text{CO}_2$  MXene, the diffusion barriers of Li atoms on its surface are 0.84 eV and 0.32 eV respectively, which are higher than  $\text{TiVCO}_2$  MXene. This result showed that vanadium doping reduced the diffusion barrier of lithium atoms on the surface of MXene, making it suitable as an anode material for lithium-ion batteries.

### 3.5 Electrochemical performance

Based on the above analysis, the article believes that Li atoms can be stably adsorbed on the MXene surface, and the energy barrier for Li-ion diffusion is low. However, the above results only consider the adsorption and migration of one Li atom on the surface of  $\text{TiVCO}_2$  MXene, and the adsorption process of multiple Li atoms can better describe the situation of the material during use. Therefore, the article calculated the open circuit voltage (OCV), binding energy ( $E_b$ ) and theoretical lithium storage capacity to evaluate the cumulative effect of multiple lithium atoms on the surface of  $\text{TiVCO}_2$  MXene. According to the adsorption energy results discussed above, lithium atoms were preferentially adsorbed at the C position with more negative adsorption energy (Fig. 6a), and the second layer was adsorbed above the transition metal atoms, as shown in Fig. 6b. The article provided images of the changes in OCV and  $E_b$  of Li atoms on the surface of  $\text{Ti}_2\text{CO}_2$  and  $\text{TiVCO}_2$  MXene as a function of specific capacity, respectively, as shown in Fig. 7a and b. It can be easily seen from the figure that the initial value of the specific capacity of  $\text{TiVCO}_2$  was larger than that of  $\text{Ti}_2\text{CO}_2$ . However, as the specific capacity increases, the specific capacity of  $\text{TiVCO}_2$  gradually approaches that of  $\text{Ti}_2\text{CO}_2$ . From Fig. 7b, contrary to the open circuit voltage, the greater the number of adsorbed Li atoms, the smaller the negative value of the binding energy. The negative value of the binding energy illustrated the stability of Li atom adsorption.

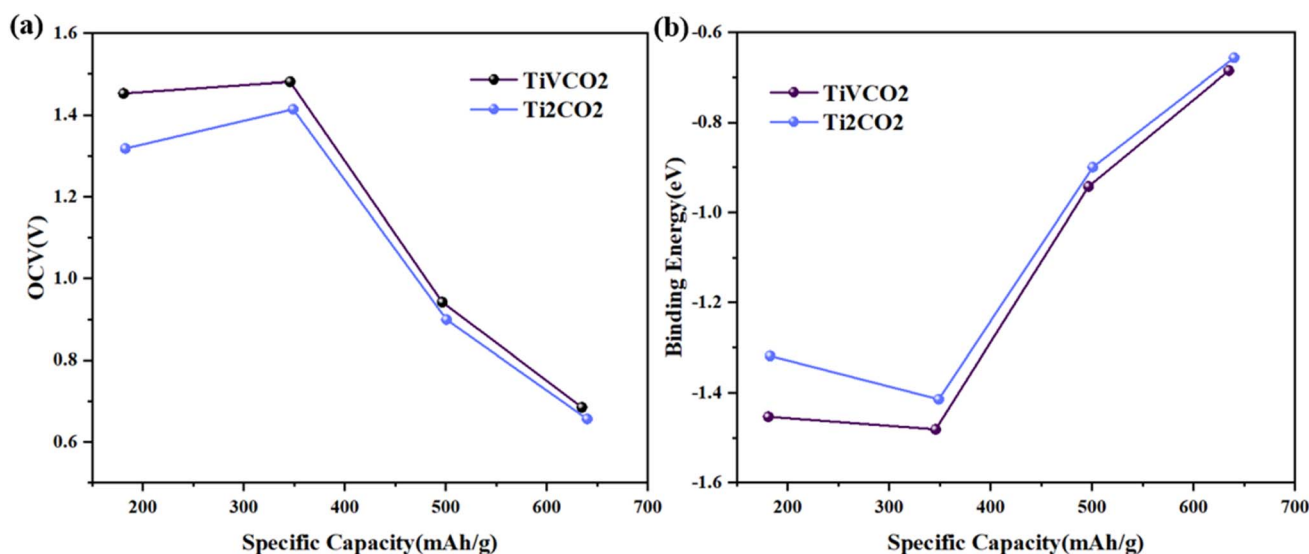


Fig. 7 (a) Change curve of open circuit voltage with specific capacity of  $\text{Ti}_2\text{CO}_2$  and  $\text{TiVCO}_2$ . (b) The relationship between the binding energy of each Li atom and the specific capacity of MXene.





## 4. Conclusion

This paper explores the possibility of MXene material ( $\text{TiVCO}_2$ ) as an anode material for lithium-ion batteries through computational simulation methods using VASP, ATAT and other software as carriers, based on generalized first principles, and using methods such as charge difference and Bader charge analysis. Through horizontal comparison with original MXene ( $\text{Ti}_2\text{CO}_2$ ), the lithium storage performance of high-entropy MXene ( $\text{TiVCO}_2$ ) as an anode material for lithium-ion batteries was studied. By calculating the DOS and PDOS of  $\text{Ti}_2\text{CO}_2$  and  $\text{TiVCO}_2$ , this paper systematically studied their electronic structure and bonding mechanism. The non-zero DOS at the Fermi level and the increase in the DOS above the Fermi level indicated that after the incorporation of V atoms. The electronic structure has changed, from the original semiconductor to a conductor, and the conductivity has been significantly improved. In addition, this article studies the adsorption performance by calculating and analyzing the adsorption sites and adsorption energies of Li atoms on the surface of MXene materials. The negative adsorption energies indicated the stability of adsorption. At the same time, this paper uses charge difference and Bader charge analysis methods to further study the adsorption mechanism of Li on MXene. On average, each Li atom transfers  $0.917|e|$  electrons to the surface of the MXene material, further indicating the adsorption of Li on its surface. By designing paths and using the climbing transition state calculation method (CI-NEB) to find the minimum diffusion energy barrier of Li on the MXene surface, and studying its surface diffusion characteristics,  $\text{TiVCO}_2$  has a lower diffusion barrier (0.20 eV) than  $\text{Ti}_2\text{CO}_2$ , indicating that The  $\text{TiVCO}_2$  MXene material has better ion mobility and is the guarantee of high-speed electrode materials.

In summary, this article further proves that  $\text{TiVCO}_2$  MXene has better conductivity, adsorption and diffusion properties than the original MXene ( $\text{Ti}_2\text{CO}_2$ ). High conductivity, high ion transport rate and low diffusion barrier make it a good alternative anode material for lithium-ion batteries with very impressive lithium storage performance.

## Data availability

Data can be obtained from the authors. Data supporting the findings of this study are available from the corresponding authors upon reasonable request.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

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## References

- 1 B. Dunn, H. Kamath and J.-M. Tarascon, *Science*, 2011, **334**, 928–935.
- 2 J.-M. Tarascon and M. Armand, *Nature*, 2001, **414**, 359–367.
- 3 F. Du, H. Tang, L. Pan, T. Zhang, H. Lu, J. Xiong, J. Yang and C. J. Zhang, *Electrochim. Acta*, 2017, **235**, 690–699.
- 4 J. O. Besenhard and G. Eichinger, *J. Electroanal. Chem. Interfacial Electrochem.*, 1976, **68**, 1–18.
- 5 K. S. Novoselov, A. K. Geim, S. V. Morozov, D. Jiang, Y. Zhang, S. V. Dubonos, I. V. Grigorieva and A. A. Firsov, *Science*, 2004, **306**, 666–669.
- 6 P. Vogt, P. De Padova, C. Quaresima, J. Avila, E. Frantzeskakis, M. C. Asensio, A. Resta, B. Ealet and G. Le Lay, *Phys. Rev. Lett.*, 2012, **108**, 155501.
- 7 E. Bianco, S. Butler, S. Jiang, O. D. Restrepo, W. Windl and J. E. Goldberger, *ACS Nano*, 2013, **7**, 4414–4421.
- 8 L. Song, L. Ci, H. Lu, P. B. Sorokin, C. Jin, J. Ni, A. G. Kvashnin, D. G. Kvashnin, J. Lou, B. I. Yakobson and P. M. Ajayan, *Nano Lett.*, 2010, **10**, 3209–3215.
- 9 H. Liu, A. T. Neal, Z. Zhu, Z. Luo, X. Xu, D. Tománek and P. D. Ye, *ACS Nano*, 2014, **8**, 4033–4041.
- 10 Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, *Nat. Nanotechnol.*, 2012, **7**, 699–712.
- 11 M. Naguib, J. Come, B. Dyatkin, V. Presser, P.-L. Taberna, P. Simon, M. W. Barsoum and Y. Gogotsi, *Electrochem. Commun.*, 2012, **16**, 61–64.
- 12 B. Anasori, M. R. Lukatskaya and Y. Gogotsi, *Nat. Rev. Mater.*, 2017, **2**, 16098.
- 13 Q. Zhong, Y. Li and G. Zhang, *Chem. Eng. J.*, 2021, **409**, 128099.
- 14 M. Naguib, M. Kurtoglu, V. Presser, J. Lu, J. Niu, M. Heon, L. Hultman, Y. Gogotsi and M. W. Barsoum, *Adv. Mater.*, 2011, **23**, 4248–4253.
- 15 S. Kurahashi, S. Arabnejad, H. Ushiyama and K. Yamashita, *J. Comput. Chem., Jpn.*, 2019, **18**, 84–94.
- 16 M. Ashton, R. G. Hennig and S. B. Sinnott, *Appl. Phys. Lett.*, 2016, **108**, 023901.
- 17 Y. Aierken, C. Sevik, O. Gülseren, F. M. Peeters and D. Çakır, *J. Mater. Chem. A*, 2018, **6**, 2337–2345.
- 18 W. Yi, G. Tang, X. Chen, B. Yang and X. Liu, *Comput. Phys. Commun.*, 2020, **257**, 107535.
- 19 V. Wang, N. Xu, J.-C. Liu, G. Tang and W.-T. Geng, *Comput. Phys. Commun.*, 2021, **267**, 108033.
- 20 K. Yuan, P. Hao, X. Li, Y. Zhou, J. Zhang and S. Zhong, *New J. Chem.*, 2021, **45**, 15234–15239.
- 21 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865–3868.



- 22 G. Kresse and D. Joubert, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 1758–1775.
- 23 S. Das, S. U. D. Shamim, M. K. Hossain, F. Ahmed, M. A. Hossain and M. O. Rahman, *Appl. Surf. Sci.*, 2022, **600**, 154173.
- 24 Q. Tang, Z. Zhou and P. Shen, *J. Am. Chem. Soc.*, 2012, **134**, 16909–16916.
- 25 Y. Xie, Y. Dall'Agnese, M. Naguib, Y. Gogotsi, M. W. Barsoum, H. L. Zhuang and P. R. C. Kent, *ACS Nano*, 2014, **8**, 9606–9615.

