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## Single-molecule detection of a terrylenediimide-based near-infrared emitter†

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Environment-sensitive fluorescent agents with near-infrared (NIR) emission are in great demand owing to their applications in biomedical and quantum technologies. We report a novel NIR absorbing ( $\lambda_{\text{max}}^{\text{Abs}} = 734$  nm) and emitting ( $\lambda_{\text{max}}^{\text{Fl}} = 814$  nm) terrylenediimide (TDI) based donor–acceptor chromophore (TDI-TPA<sub>4</sub>), exhibiting polarity-sensitive single-photon emission. By virtue of the charge transfer (CT) character, ensemble level measurements revealed solvatochromism and NIR emission ( $\phi_{\text{Fl}} = 26.2\%$ ), overcoming the energy gap law. The CT nature of the excited states is further validated by state-of-the-art fragment-based excited state theoretical analysis. To mimic the polarity conditions at the single-molecule level, TDI-TPA<sub>4</sub> was immobilized in polystyrene (PS; low polar) and poly(vinyl alcohol) (PVA; high polar) matrices, which enables tuning of the energy levels of the locally excited state and charge-separated (CS) state. Minimal blinking and prolonged survival time of the TDI-TPA<sub>4</sub> molecule in the PS matrix, in contrast to the PVA matrix, possibly confirms the implication of the energy gap law and polarity sensitivity of TDI-TPA<sub>4</sub>. The existence of the CT state in nonpolar and CS state in polar solvents was confirmed by transient absorption measurements in the femtosecond regime. The current work sheds light on the design principle for NIR single-photon emitting organic chromophores for deep tissue imaging and probing the nanoscale heterogeneity.

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## Introduction

Near-infrared single-photon emitters are an emerging class of compounds for quantum information science, including computing, sensing, and communication.<sup>1–5</sup> Single-molecule level scrutiny of organic near-infrared (NIR) emitting molecules has recently gained much attention.<sup>6–9</sup> The efficiency of the NIR organic emitters is limited by the energy gap law; the intrinsic forbidden nature of emission from states with low energy gap results in negligible fluorescence quantum yield.<sup>10–13</sup> The donor–acceptor systems with charge transfer (CT) character<sup>14</sup> are better alternatives to circumvent the limitations possessed by conventional  $\pi$ -conjugated NIR molecules by narrowing the energy gap for NIR emission.<sup>15–18</sup> The emission from the CT state is limited due to the forbidden nature of electronic transition as a result of spatial separation between the highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO).<sup>19,20</sup> However, hybridizing the local excited (LE) state with the CT state results in a favored transition dipole moment, subsequently leading to NIR emission.<sup>21–23</sup> Exciton delocalization in the aggregates is

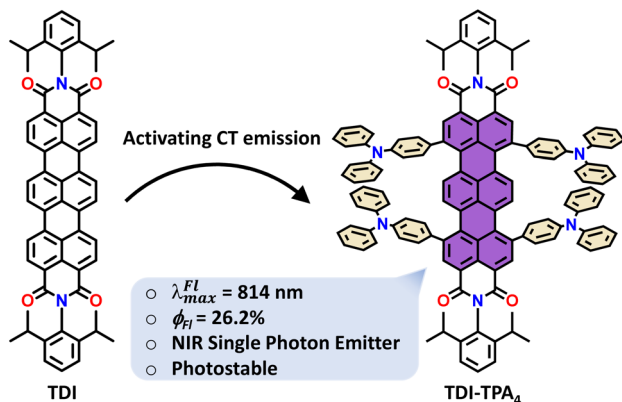
one of the strategies implemented in Pt(II) complexes and bay-alkylated quaterylene to successfully overcome the energy gap law for enhanced NIR fluorescence quantum yield at the ensemble level.<sup>24,25</sup> Recently, Erker and Basché validated the energy gap law for NIR single-photon emitters by investigating the dielectric dependence of dibenzoterrylene ( $\lambda_{\text{max}}^{\text{Fl}} = 741$  nm) at the ensemble and single-molecule levels.<sup>26</sup> In comparison to the previously investigated NIR single-photon emitters (Table S1†), probing the dielectric dependence of CT-mediated NIR emitting chromophores at the single-molecule level remains elusive. Polarity dependence of fluorescence at the nanoscale level can be probed using chromophores with emissive CT in the visible range, which was first demonstrated by Higgins and co-workers<sup>27</sup> based on the model developed by Marcus.<sup>28,29</sup> Monitoring the intrinsic NIR fluorescence fluctuations at the single-molecule level in different dielectric environments is beneficial for probing temporal inhomogeneity and acquiring autofluorescence-free contrasting images of (bio)specimens.<sup>30–38</sup>

By virtue of the limited blinking and extended survival time before photobleaching,<sup>39</sup> terrylenediimide (TDI) is known to be a suitable candidate for membrane labeling<sup>40</sup> and single-molecule investigations.<sup>41–44</sup> Our longstanding interest in monitoring CT excitons in rylencarboximide derivatives at the ensemble<sup>45–47</sup> and single-molecule levels<sup>48,49</sup> motivated us to explore higher analogs, such as TDI-functionalized chromophores.<sup>50–56</sup> Herein, we designed and synthesized

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Scheme 1 Design strategy and chemical structures of TDI and TDI-TPA<sub>4</sub>.

a NIR absorbing and emitting donor–acceptor chromophore based on TDI (acceptor) with triphenylamine moiety (TPA) as a donor (TDI-TPA<sub>4</sub>) and explored its potential as a polarity-sensitive NIR single-photon source (Scheme 1). The striking contrast of fluorescence intensity time traces in different dielectric environments was demonstrated by single-molecule thin film measurements, and femtosecond excited state dynamics investigations corroborated the photoinduced processes at the ensemble level.

## Results and discussion

### Synthesis

Novel TDI-TPA<sub>4</sub> was synthesized *via* a Suzuki cross-coupling reaction between tetrabromo-terrylenediimide (TDI-Br<sub>4</sub>) and 4-(diphenylamino)phenylboronic acid (Schemes S1 and S2†). TDI and TDI-Br<sub>4</sub> were synthesized and characterized as per the previously reported procedures (Scheme S2 and Fig. S1–S8†).<sup>50</sup> The Suzuki coupling reaction between TDI-Br<sub>4</sub> and 4-(diphenylamino)phenylboronic acid was carried out under inert conditions in the presence of potassium carbonate (K<sub>2</sub>CO<sub>3</sub>), dioxane, and a catalytic amount of tetrakis(triphenylphosphine)palladium(0) followed by heating of the reaction mixture at 85 °C. TDI-TPA<sub>4</sub> was isolated and purified using column chromatography (ethyl acetate/hexane, 1 : 10) to yield purple-colored TDI-TPA<sub>4</sub> with a 40% yield (Scheme S2†). The desired compound was successfully characterized by <sup>1</sup>H and <sup>13</sup>C NMR, IR spectroscopy, and HRMS mass spectrometry (see Fig. S9–S13†).

### Optical properties

The electronic absorption spectra of the TDI-TPA<sub>4</sub> and TDI were recorded in hexane (HEX,  $\epsilon = 1.88$ ,  $C = 5\text{--}10 \mu\text{M}$ ) (Fig. 1). The absorption spectrum of TDI-TPA<sub>4</sub> spans through the visible range and extends to the NIR window (400–850 nm) with absorption bands at 734, 597, and 522 nm with a molar extinction coefficient  $\epsilon_{\text{max}} = 3.54 \times 10^4 \text{ M}^{-1} \text{ cm}^{-1}$  (at 734 nm) corresponding to the S<sub>0</sub> to S<sub>1</sub> transition (Table 1 and Fig. S14, S15†). TDI exhibited three distinct vibronic progressions with an absorption maximum at 634 nm corresponding to the S<sub>0</sub> to

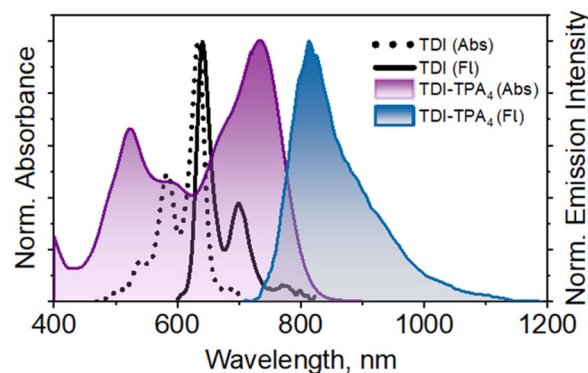


Fig. 1 Normalized UV-vis absorption and fluorescence emission spectra of TDI and TDI-TPA<sub>4</sub> in hexane.

S<sub>1</sub> electronic transition. The fluorescence emission spectrum of TDI-TPA<sub>4</sub> consists of a NIR band with a maximum at 814 nm when excited at 591 nm in HEX (Fig. 1). The remarkable red-shift in the fluorescence emission spectrum of TDI-TPA<sub>4</sub> as compared to TDI ( $\Delta\lambda_{\text{max}}^{\text{FI}} = 173 \text{ nm}$  or  $3315.16 \text{ cm}^{-1}$ ) corroborate the existence of charge transfer between the four TPA moieties and the TDI core. The solvent-dependent absorption measurements of TDI-TPA<sub>4</sub> depicted a red-shift ( $\Delta\lambda_{\text{max}}^{\text{Abs}} = 20 \text{ nm}$ ) with increasing solvent polarity from HEX to ACN (Table 1 and Fig. S14–S17†). The solvent-dependent fluorescence spectra of TDI-TPA<sub>4</sub> ( $C = 5\text{--}10 \mu\text{M}$ ) depict a bathochromic shift along with significant fluorescence quenching upon increasing polarity from HEX ( $\epsilon = 1.88$ ,  $\phi_{\text{FI}} = 26.2\%$ ), THF ( $\epsilon = 7.58$ ,  $\phi_{\text{FI}} = 17.5\%$ ) to ACN ( $\epsilon = 37.50$ ,  $\phi_{\text{FI}} = 10.2\%$ ) (Fig. S18† and Table 1). The decreased fluorescence quantum yield of TDI-TPA<sub>4</sub> in hexane compared to TDI ( $\phi_{\text{FI}} = 90\%$ )<sup>50</sup> is indicative of other non-radiative decay channels upon photoexcitation. The substantial Stokes shift of TDI-TPA<sub>4</sub> in the polar solvents ( $\Delta\tilde{\nu}_{\text{Stokes}} = 227 \text{ nm}$  or  $3097.90 \text{ cm}^{-1}$  in THF and  $\Delta\tilde{\nu}_{\text{Stokes}} = 231 \text{ nm}$  or  $3110.16 \text{ cm}^{-1}$  in ACN) when compared to TDI-TPA<sub>4</sub> in HEX ( $\Delta\tilde{\nu}_{\text{Stokes}} = 80 \text{ nm}$  or  $1338.77 \text{ cm}^{-1}$ ) results from the lowering of the energy gap, leading to the reduced fluorescence quantum yield, which is in accordance with the energy gap law. The emission spectra recorded by exciting at 734, 597, and 522 nm depicted intact emission maximum at  $\lambda_{\text{max}}^{\text{FI}} = 814 \text{ nm}$  (Fig. S19†). The excitation spectra collected at the respective emission maximum of TDI-TPA<sub>4</sub> in HEX, THF, and ACN resemble the absorption spectra, validating the emissive nature of the CT state (Fig. S20†).<sup>47</sup> Time-correlated single photon counting (TCSPC) measurement of TDI-TPA<sub>4</sub> was performed in HEX ( $\lambda_{\text{ex}} = 640 \text{ nm}$  and monitored at 814 nm), suggesting the existence of a single emissive component with a fluorescence lifetime  $\tau_{\text{FI}} = 1.37 \text{ ns}$  (Fig. S21†). The concentration-dependent and temperature-dependent absorption measurements of TDI-TPA<sub>4</sub> were performed in THF to understand the possibility of aggregate formation. The absorption spectra of TDI-TPA<sub>4</sub> were recorded for different concentrations ( $C = 0.01 \text{ mM}$  to  $0.20 \text{ mM}$ ) and exhibited constant  $I_0/I_1$  ratios, ruling out the possibility of aggregation (Fig. S22 and Table S2†). Monitoring the temperature-dependent absorption spectra of TDI-TPA<sub>4</sub> ( $C = 5\text{--}10 \mu\text{M}$ )



**Table 1** Fluorescence quantum yields ( $\phi_{\text{Fl}}$ ) and fluorescence time constants ( $\tau$ ), and driving force for charge separation ( $\Delta G_{\text{CS}}$ ), driving force for charge recombination ( $\Delta G_{\text{CR}}$ ), and rate constants for radiative ( $k_{\text{r}}$ ), nonradiative ( $k_{\text{nr}}$ ), processes of TDI-TPA<sub>4</sub>

Solvent	$\lambda_{\text{max}}^{\text{Abs}}$ (nm)	$\lambda_{\text{max}}^{\text{Fl}}$ (nm)	$\epsilon_{\text{max}} \times 10^4$ ( $\text{M}^{-1} \text{cm}^{-1}$ )	$\phi_{\text{Fl}}$ (%) $\pm 5$	$\tau^a$ (ps)	$\Delta G_{\text{CS}}$ (eV)	$\Delta G_{\text{CR}}$ (eV)	$k_{\text{r}} \times 10^9$ ( $\text{s}^{-1}$ )	$k_{\text{nr}} \times 10^9$ ( $\text{s}^{-1}$ )
HEX ( $\epsilon = 1.88$ )	734	814	3.54	26.2	1210.0	0.20	-1.77	0.216	0.6
THF ( $\epsilon = 7.58$ )	750	977	3.33	17.5	1.5	-0.21	-1.36	116.6	550.0
ACN ( $\epsilon = 37.50$ )	754	985	3.13	10.2	0.5	-0.33	-1.24	204.0	1796.0

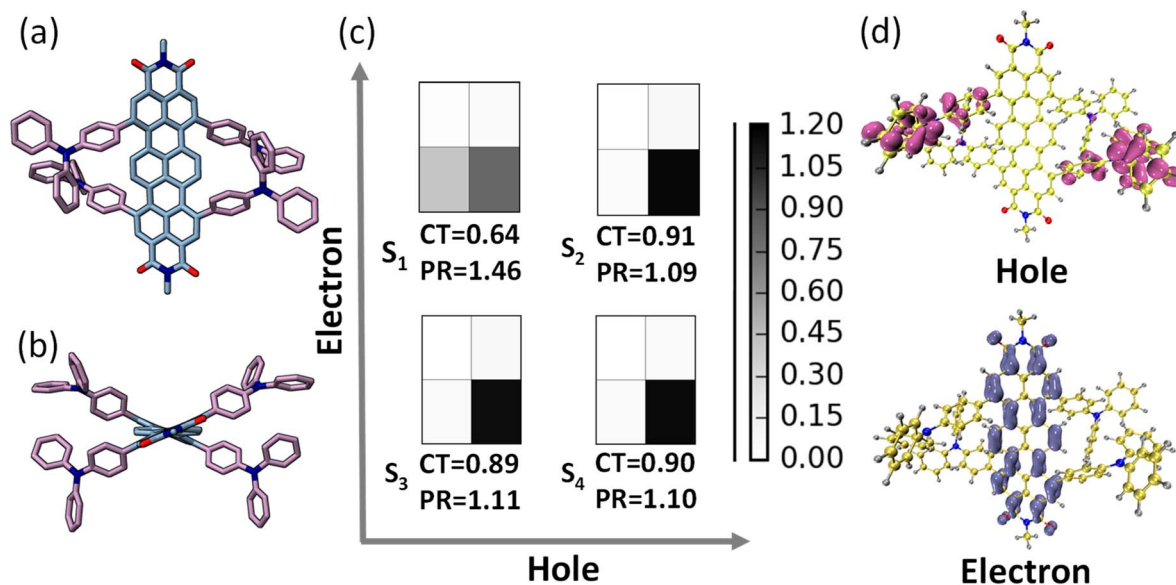
<sup>a</sup> From global fits of fsTA spectra.  $\Delta G_{\text{CS}}$  and  $\Delta G_{\text{CR}}$  from Rehm–Weller analysis.  $k_{\text{r}}$  and  $k_{\text{nr}}$  are calculated as per the reported procedure from the literature.<sup>26,57</sup>

from 10 °C to 60 °C depicted negligible changes in the spectral feature, further strengthening the monomeric nature of TDI-TPA<sub>4</sub> in the solution state (Fig. S23†).

### Theoretical investigations

The ground state optimization of the TDI-TPA<sub>4</sub> was performed using density functional theory (DFT) in vacuum using the Gaussian 16 program.<sup>58</sup> 2,6-Diisopropylphenyl group at the imide position has been replaced with a methyl group to reduce computational cost. For ground state and first singlet excited state optimization, a range-separated hybrid exchange–correlation functional CAM-B3LYP (using the Coulomb-attenuating method) in combination with a def2-SVP basis set was used to account for the charge transfer excitations in TDI-TPA<sub>4</sub>. The optimized structure of TDI and TDI-TPA<sub>4</sub> revealed a considerable contortion in the TDI core with a dihedral angle of 18.8° and 18.6° as a result of bulky TPA substitution in the bay region (Fig. 2a, b and S24†). The propeller-like TPA moieties are oriented at angles 115.3° and 115.1° with respect to the ortho

position of the TDI core to decrease steric repulsion (Fig. S24 and S25†). Time-dependent DFT was employed to compute the vertical excitation energies and oscillator strengths at the CAM-B3LYP/def2-SVP level of theory. The singlet excited states  $S_1$  (2.105 eV),  $S_2$  (2.885 eV), and  $S_3$  (2.886 eV) are identified as bright states with considerable oscillator strength in TDI-TPA<sub>4</sub> (Table S3†). Further investigation into the character of the excited states was carried out using fragment-based excited-state analysis developed by Plasser, implemented in TheoDOR.<sup>59,60</sup> Becke, 3-parameter, Lee–Yang–Parr functional with D3, an empirical correction for dispersion forces and def2-SVP basis set was employed to perform the calculations. The degree of excitation energy delocalization in various fragments is determined by the participation ratio (PR). TDI core and four TPA groups were categorized as two fragments for the fragment analysis of the TDI-TPA<sub>4</sub> (Fig. S26†). The nature of the excited state is represented by the charge transfer (CT) value where  $\text{CT} > 0.8$  denotes the CT state, and  $\text{CT} = 0$  depicts the Frenkel exciton state (FE). A hole–electron correlation plot is used to pictorially represent the nature of the excited states. The  $S_1$  state of TDI-TPA<sub>4</sub> is identified



**Fig. 2** Ground state optimized geometry of TDI-TPA<sub>4</sub> shown in (a) top view and (b) side view at CAM-B3LYP/def2-SVP level of theory (2,6-diisopropylphenyl group has been replaced with methyl group to reduce computational cost. Hydrogen atoms are omitted for clarity). (c) Hole–electron correlation plots showing the nature for the  $S_1$ ,  $S_2$ ,  $S_3$ , and  $S_4$  excited states. (d) Corresponding hole–electron isosurface distribution of  $S_2$  state (isovalued = 0.001). The CT and PR values are shown to define the nature of excitations.



as a partial CT state (hybridized feature of CT and LE) with a CT value of 0.64 (PR = 1.46). Conversely, transitions  $S_0 \rightarrow S_2$ ,  $S_0 \rightarrow S_3$ , and  $S_0 \rightarrow S_4$  possessed CT values of 0.91, 0.89, and 0.90, respectively, corresponding to a strong CT character (Fig. 2c). In the hole–electron isosurface of the  $S_2$  state, the electron is confined to the TDI core while the hole is confined to the TPA fragments, establishing the CT nature of TDI-TPA<sub>4</sub> (Fig. 2d).<sup>64</sup> To obtain further insights into the electron density distribution within the TDI-TPA<sub>4</sub> molecule, the natural transition orbital (NTO) isosurface analysis for  $S_1$ ,  $S_2$ ,  $S_3$ , and  $S_4$  excited states was carried out. For the  $S_1$  state, the electron density is distributed on both the TDI core and TPA fragment, and the lowest unoccupied NTO is localized on the TDI core, corroborating the hybrid nature of the CT and LE. The predominant CT character of higher  $S_2$ ,  $S_3$ , and  $S_4$  excited states is further substantiated by the spatially disjoint electron density distribution of the highest occupied NTO on electron donating TPA fragment and the lowest unoccupied NTO on the TDI core (Fig. S27†).

### Single-molecule fluorescence measurements

Exploitation of CT-based organic chromophores as a single-molecule source remained challenging due to quenching of local excited state fluorescence *via* electron transfer.<sup>62</sup> Considerable efforts have been reported in the literature to monitor single-molecule electron transfer of rylene diimide dendrimers in the visible range.<sup>63–67</sup> The fluorescence lifetime imaging (FLIM, Fig. 3a and e) and fluorescence intensity trajectories (FITs) were measured using the MicroTime 200 (MT200) time-resolved fluorescence microscope from PicoQuant. The inverted microscope (Olympus IX83) is equipped with a piezo-scanning stage (P-733.2CD, PI). The sample was excited using a 512 nm pulsed laser (LDH-D-C-512, PicoQuant) with

a repetition rate of 20 MHz and scanned with a dwell time of 0.4 ms. To acquire the single-molecule emission in the NIR region, the intensity traces were recorded using an SPCM-AQRH single-photon avalanche detector (SPAD) having spectral range 400–1100 nm with 70% maximum detection efficiency at 700 nm and 180 nm diameter of active area.

Single-molecule fluorescence (SMF) measurements of TDI-TPA<sub>4</sub> were carried out in two different dielectric polymer media such as polystyrene (PS,  $\epsilon = 2.6$ ) and poly(vinyl alcohol) (PVA,  $\epsilon = 7.8$ ) nonpolar and polar analogs, respectively, akin to the ensemble measurements.<sup>64</sup> FLIM images represent discrete TDI-TPA<sub>4</sub> molecules in the PS and PVA matrices (Fig. 3a and e). FITs depict the real-time fluorescence fluctuations of single molecules in their local environment.<sup>68</sup> The FITs under investigation include molecules showing blinking (reversible switching between on/bright and off/dark states) or single-step photobleaching (irreversible switch to off state). In the PS matrix, TDI-TPA<sub>4</sub> depicted (81 single molecules) prolonged photon count and seldom blinking in agreement with the photophysical process occurring in the low dielectric ensemble measurements (Fig. 3b). The single-molecule fluorescence lifetime distribution of TDI-TPA<sub>4</sub> in PS showed monoexponential lifetime distribution (Fig. 3c) peaking at  $\sim 2$ –3 ns (Fig. 3d). Frequent blinking and low fluorescence emission were observed for TDI-TPA<sub>4</sub> single molecules in the polar PVA matrix (119 single molecules). The contrasting blinking dynamics in matrices of distinct dielectric constants corroborated photoinduced charge separation is the major contributor to the dark states in TDI-TPA<sub>4</sub>. The feasibility of charge separation being less favored in low polar PS matrix accounts for the high photostability and relatively high rate of radiative process in TDI-TPA<sub>4</sub> compared to PVA.<sup>67</sup> Whereas charge separation is facilitated as an additional nonradiative pathway in the PVA matrix,

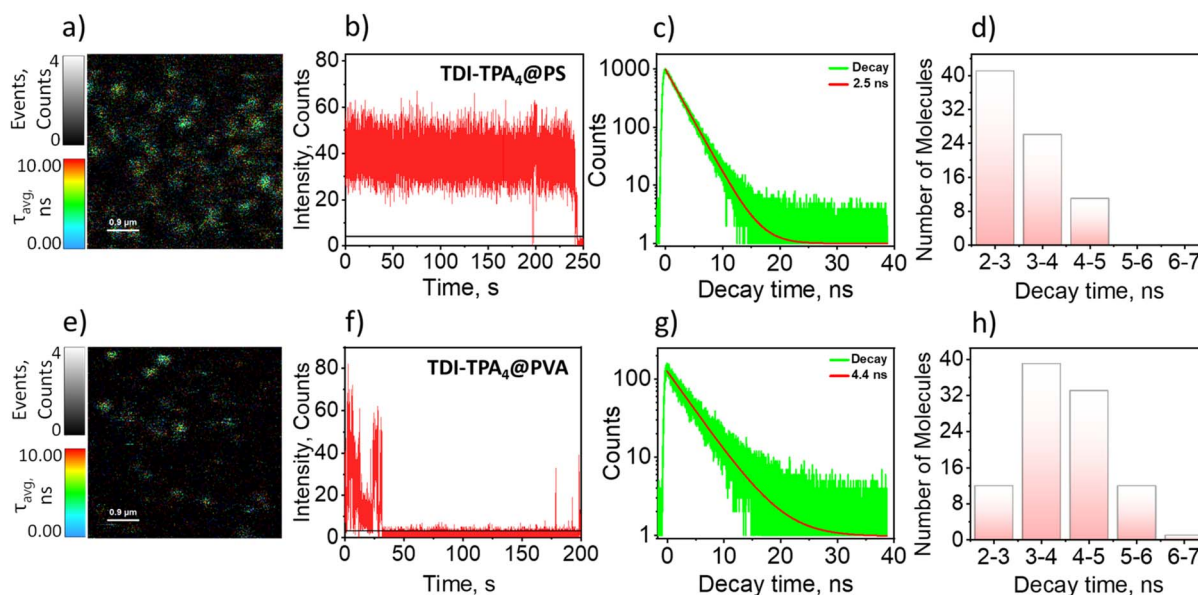


Fig. 3 Representative fluorescence lifetime images with average lifetime ( $\tau_{\text{avg}}$ ) and events scale, of TDI-TPA<sub>4</sub> in (a) PS and (e) PVA matrices. Fluorescence intensity trajectories (FITs) of TDI-TPA<sub>4</sub> in (b) PS and (f) PVA matrices. Representative single-molecule fluorescence decay measurements of TDI-TPA<sub>4</sub> in (c) PS and (g) PVA matrices. Fluorescence lifetime distribution plot of TDI-TPA<sub>4</sub> in (d) PS and (h) PVA matrices.





corroborating the frequent blinking and relatively low fluorescence emission in TDI-TPA<sub>4</sub> (Fig. 3f). The single-molecule fluorescence lifetime of TDI-TPA<sub>4</sub> in the PVA matrix is spread over the 2–6 ns range (Fig. 3g and h). The broad distribution of lifetime implies the existence of a wide range of rate constants for electron transfer in the PVA matrix possibly due to the rigidity and residual low polar local environment offered by toluene molecules in PVA layer (see Section 1.10 in ESI†).<sup>27,69,70</sup> An erratic blinking and irreversible photobleaching is reported as a consequence of the stabilization of the CS state in the polar nano-environment.<sup>71</sup> The excited state stabilization results in the lowering of transition energy at the single-molecule domain, leading to the opening up of new nonradiative channels implied by the energy gap law.

Antibunching experiments of TDI-TPA<sub>4</sub> in PS and PVA matrices were carried out using an 800 nm short-pass filter. A doublet peak was observed at zero delay point, which is characteristic of the afterglow (fluorescence from one detector initiating the other) from the silicon detector (Fig. S28†).<sup>72</sup> Since there is a considerable contribution from afterglow above 700 nm, reliable antibunching data were not acquired. However, the contribution from the afterglow photons in the FIT measurements is buried in the background signal (Fig. S28†). The single-molecule investigation of the present work exclusively contains spatially separated single molecules that undergo blinking or single-step photobleaching. (Representative FITs; Fig. S29 and S30†).

### Cyclic voltammetry and Rehm–Weller analysis

Having established the existence of bright and dark states of TDI-TPA<sub>4</sub> in nonpolar and polar matrices, respectively at the single-molecule domain, we evaluated the thermodynamic feasibility of electron transfer at the ensemble level using cyclic voltammetry (CV) and differential pulse voltammetry (DPV). CV and DPV measurements were performed against Ag/Ag<sup>+</sup> in dry dichloromethane (DCM) under a nitrogen atmosphere to examine the redox properties of TDI-TPA<sub>4</sub> using ferrocene/ferrocenium (Fc/Fc<sup>+</sup>) couple as a reference (Fig. S31†). Two distinct reversible reduction waves ( $E_{1(\text{red})} = -1.01$  V and  $E_{2(\text{red})} = -1.14$  V) and three reversible oxidation waves ( $E_{1(\text{ox})} = 0.44$  V,  $E_{2(\text{ox})} = 0.59$  V and  $E_{3(\text{ox})} = 0.82$  V) of TDI-TPA<sub>4</sub> confirms profound electronic communication between the TDI core and TPA moiety. Rehm–Weller formalism provides the feasibility associated with the electron transfer process in the donor–acceptor system in terms of Gibbs free energy changes for charge separation ( $\Delta G_{\text{CS}}$ ).<sup>73</sup> The positive  $\Delta G_{\text{CS}} \approx 0.20$  eV indicates the infeasibility of charge separation of TDI-TPA<sub>4</sub> in HEX. The  $\Delta G_{\text{CS}} \approx -0.21$  eV in THF and  $\Delta G_{\text{CS}} \approx -0.33$  eV in ACN suggests the feasibility of photoinduced electron transfer in medium and high polar solvents, respectively. The electrochemical energy gap ( $E_{\text{g}}$ ) is estimated to be 1.45 eV from experimentally observed LUMO and HOMO energy levels (Table S4†).

### Femtosecond transient absorption measurements

Solvent-dependent transient absorption measurements provide insights into the excited state dynamics of TDI-TPA<sub>4</sub> in different

dielectric environments. Femtosecond transient absorption spectra (fsTA) of TDI-TPA<sub>4</sub> were recorded in HEX ( $\epsilon = 1.88$ ), THF ( $\epsilon = 7.58$ ), and ACN ( $\epsilon = 37.50$ ) having optical density 0.2–0.3 using ~100 fs laser pulse exciting at 520 nm (Fig. 4). A Spectra-Physics Mai Tai SP mode-locked laser (800 nm, 86 MHz) was used as a seed for Spectra-Physics Spitfire Ace regenerative amplifier (1 kHz, 5.5 mJ). Using TOPAS, a portion of the amplified 800 nm output beam was converted into a 520 nm pump pulse. The magic angle polarization (~54.7°) between pump and probe pulses was used to ensure an isotropic signal of the sample. Singular value decomposition (SVD) and global analysis of the  $\Delta A$  versus time-based wavelength-based three-dimensional map of the fsTA spectra were carried out to extract the kinetic components responsible for the total spectra. The fsTA spectra of TDI-TPA<sub>4</sub> in HEX displayed negative ground state bleach (GSB) and stimulated emission (SE) band from ~552–850 nm accompanied by a positive excited state absorption (ESA) band ranging from ~876–1400 nm (Fig. 4a). The one-component sequential kinetic model A → GS (GS = ground state) was employed to acquire the evolution-associated spectra (EAS) of TDI-TPA<sub>4</sub> in HEX. The deconvolution of the fsTA spectra revealed a single principal component with a decay constant of 1.21 ns in HEX, attributed to a CT state that matches the fluorescence lifetime obtained from TCSPC measurement. The inaccessibility of the CS state from the CT state in HEX is probably due to the large energy barrier for charge separation from the CT state, which is evident from the positive  $\Delta G_{\text{CS}} \approx 0.20$  eV calculated from the Rehm–Weller analysis.

Transient absorption spectra recorded for TDI-TPA<sub>4</sub> in a moderately polar solvent THF exhibited a negative band ~550–827 nm corresponding to GSB and SE. The positive ESA consists of two bands: a low intense band (~829–906 nm) and a broad intense absorption band ranging from ~934–1396 nm (Fig. 4b). A two-component sequential kinetic model A → B → GS was employed to elucidate the excited state process in THF. Upon deconvolution, fsTA spectra result in two principal components centered at 1215 nm (A) and 1196 nm (B), which can be attributed to a CT state and a CS state, respectively. The negative  $\Delta G_{\text{CS}} \approx -0.21$  eV calculated from the Rehm–Weller analysis validates the feasibility of populating the CS state in THF.<sup>74</sup> The CT state decays with a time constant of 1.5 ps, leading to the CS state, which subsequently decays to the ground state within 149.2 ps. The 3D plot of fsTA confirms the evolution of a new species in the later time delay in the ~829–906 nm region, which resembles the spectroscopic signature of chemically oxidized absorption of TPA reported in the literature (Fig. S32 and S33†).<sup>74</sup> The chemical oxidation experiments were carried out for the TPA in toluene and ACN using antimony(v) chloride, further confirming the characteristic absorption feature of the TPA radical cation in this wavelength regime (Fig. S34†).

The transient absorption spectra of TDI-TPA<sub>4</sub> were recorded in ACN to probe excited state dynamics in a high dielectric environment. The negative GSB and SE featured in the ~550–835 nm region, and the positive ESA band spanned from the ~836–1445 nm region. The two-component sequential kinetic model A → B → GS is used to comprehend the photoinduced



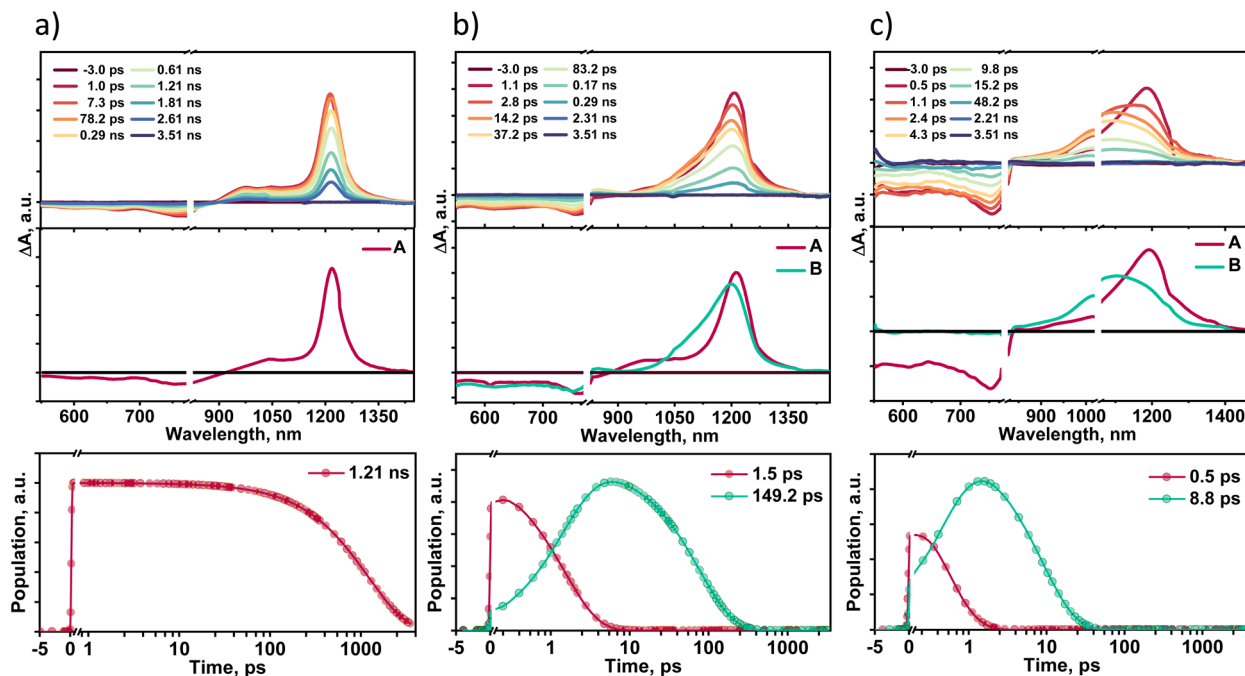


Fig. 4 (Top row) Femtosecond transient absorption spectra of TDI-TPA<sub>4</sub> in (a) HEX, (b) THF, and (c) ACN showing the excited state dynamics upon photoexcitation ( $\lambda_{\text{ex}} = 520$  nm). (Middle row) Species-associated spectra reconstructed from global analysis of TDI-TPA<sub>4</sub> with A  $\rightarrow$  GS model for HEX (left), the A  $\rightarrow$  B  $\rightarrow$  GS model for THF (center), and A  $\rightarrow$  B  $\rightarrow$  GS model for ACN (right). (Bottom row) Relative population profiles of the excited states in TDI-TPA<sub>4</sub> in HEX, THF, and ACN.

processes of TDI-TPA<sub>4</sub> in ACN. The first component (A) at 1195 nm undergoes ultrafast decay to the second component (B) with a rate constant  $k_{A \rightarrow B} = (0.5 \text{ ps})^{-1}$ . The formation of the second component at 1130 nm occurs in the later time delay, which further relaxes to the ground state with the rate constant of  $k_{B \rightarrow \text{GS}} = (8.8 \text{ ps})^{-1}$ . The second band at 1130 nm matches with the chemically reduced anion formed from TDI-TPA<sub>4</sub> in the ACN upon treatment with bis(cyclopentadienyl)cobalt(II) as a reducing agent (Fig. S35<sup>†</sup>). The A and B components can be attributed to CT and CS states, respectively, which are in agreement with the feasibility of charge separation  $\Delta G_{\text{CS}} \approx -0.33$  eV in ACN. The spectroelectrochemistry measurements further reinforced the evolution and spectral signature of TDI-TPA<sub>4</sub> radical anion by applying electric potential. Electrochemical reduction of light-irradiated TDI-TPA<sub>4</sub> solution in DCM exhibited an emergence of a new band at  $\sim 983$  nm upon applying the 3.5 V potential, which confirms the existence of the TDI-TPA<sub>4</sub> radical anion (Fig. S36<sup>†</sup>). The fluorescence time constant ( $\tau$ ),  $\Delta G_{\text{CS}}$ , driving force for charge recombination ( $\Delta G_{\text{CR}}$ ) and rate constants ( $k$ ) of TDI-TPA<sub>4</sub> in different solvents are summarized in Table 1. The remarkable increase in the nonradiative rate with the decrease in the transition energy validates the energy gap law in TDI-TPA<sub>4</sub> at the ensemble level. The selected kinetic traces superimposed with the global analysis fitted curves at different wavelengths of fsTA in HEX, THF, and ACN are shown to display the fitting quality in Fig. S37<sup>†</sup>. To evaluate the possibility of triplet state formation upon charge recombination proposed by El-Sayed,<sup>75</sup> nanosecond transient absorption spectroscopy (nsTA) measurements were performed

in HEX, THF, and ACN by exciting the sample at 532 nm. The absence of a long-lived species with a positive absorption feature ( $\tau_{\text{T}} > 8$  ns) in the nsTA spectra further confirms the rapid charge recombination in TDI-TPA<sub>4</sub> to a singlet state in THF and ACN (Fig. S38<sup>†</sup>).

The fsTA thin film measurements of TDI-TPA<sub>4</sub> in PS and PVA polymer matrices were carried out to comprehend the excited state dynamics in low polar and high polar polymer environments upon photoexciting at 520 nm. The samples for thin film experiments were prepared by drop-casting the TDI-TPA<sub>4</sub> and polymer solution on a cleaned round quartz cuvette. The thin film thickness of TDI-TPA<sub>4</sub> in PS and PVA matrices were measured to be  $\sim 430 \pm 5$  nm and  $470 \pm 5$  nm respectively, using profilometry (KLA Tencor D600 stylus surface profiler; Fig. S39<sup>†</sup>). In the low polar PS matrix, the TDI-TPA<sub>4</sub> exhibited negative GSB at  $\sim 561$ – $771$  nm accompanied by ESA at  $\sim 1046$ – $1349$  nm region (Fig. S40<sup>†</sup>). The wavelength range of the GSB and ESA of TDI-TPA<sub>4</sub> in the PS matrix resembles the spectral signatures of TDI-TPA<sub>4</sub> in HEX (Fig. 4a). The deconvolution of the fsTA spectra of TDI-TPA<sub>4</sub> in PS thin film was carried out using a sequential kinetic model A  $\rightarrow$  GS having a time constant of  $\tau_{A \rightarrow \text{GS}} = 0.87$  ns, which is attributed to the CT state. A faster kinetic decay of TDI-TPA<sub>4</sub> was observed in the PS thin films compared to the fsTA spectra of TDI-TPA<sub>4</sub> in the HEX solution. The accelerated decay kinetics of TDI-TPA<sub>4</sub> in PS compared to HEX is possibly due to the higher dielectric environment offered by PS. The fsTA spectra of TDI-TPA<sub>4</sub> in the PVA matrix depict an ESA band centered at  $\sim 1240$  nm, which is characteristic of the ESA band observed in the polar solvents (THF and ACN).



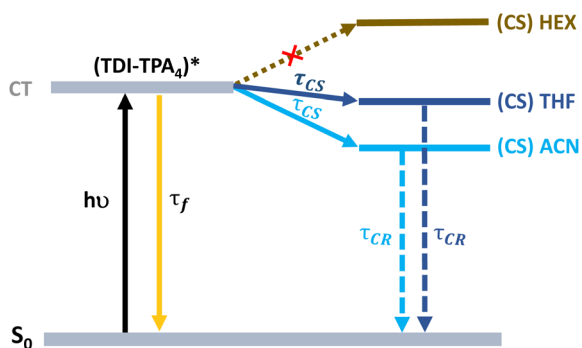


Fig. 5 Jablonski diagram showing the excited state deactivation mechanism of TDI-TPA<sub>4</sub> in HEX, THF, and ACN.

Deconvolution of the total spectra using  $A \rightarrow B \rightarrow GS$  sequential kinetic model resulted in two principal components with decay constants of 0.3 ps followed by 4.7 ps, which can be attributed to the CT and CS state of TDI-TPA<sub>4</sub> in the PVA matrix, respectively (Fig. S40<sup>†</sup>). In thin film fsTA measurements, higher concentration of TDI-TPA<sub>4</sub> (~5 mM) reduces the interaction with residual toluene molecules per unit area. As a result, the influence of low polar environment offered by residual solvent molecules may not be detectable in the ensemble level transient absorption measurements. The thin film fsTA measurements of TDI-TPA<sub>4</sub> in PVA matrix correspond to the excited state dynamics in the high dielectric environment, consistent with observations from high-polarity solution state fsTA measurements. This indeed affirms the capability of single-molecule techniques to probe unique events that are averaged out in ensemble level measurements.

A proposed Jablonski diagram illustrating the excited state dynamics of TDI-TPA<sub>4</sub> in various dielectric environments is shown in Fig. 5. In the HEX or PS matrix, the deactivation from the CT state is identified as a major pathway as a result of a large energy barrier in achieving the CS state. In a moderately polar solvent THF and polar ACN or in a PVA matrix, the access to the CS state from the CT state is facilitated by the stabilization of the CT state in a higher dielectric environment.

## Conclusions

In summary, a novel NIR absorbing and emitting TDI-TPA<sub>4</sub> was explored at bulk and the single-molecule level. The solvatochromism and a significant red-shift in the emission spectrum exhibited by TDI-TPA<sub>4</sub> with respect to TDI ( $\Delta\lambda_{\text{max}}^{\text{Fl}} = 173 \text{ nm}$ ) revealed CT character at the ensemble level and is supported by theoretical analysis. Quantum chemical calculations, including fragment-based excited state analysis of TDI-TPA<sub>4</sub>, suggested the nature of the first singlet excited state ( $S_1$ ) as a hybrid state of both CT and LE states and  $S_2$ ,  $S_3$ , and  $S_4$  excited states as CT states. Single-molecule fluorescence measurements in PS and PVA unambiguously confirmed the dielectric dependence and implication of the energy gap law of the TDI-TPA<sub>4</sub> to serve as a NIR emitting polarity probe in the nanoscale environment. In the low polar PS matrix, relatively

high photostability and photon count were observed in comparison to the frequent blinking and photobleaching in the polar PVA matrix. The thermodynamic feasibility of charge separation of TDI-TPA<sub>4</sub> in THF and ACN ( $\Delta G_{\text{CS}} \approx -0.21$  and  $\Delta G_{\text{CS}} \approx -0.33 \text{ eV}$ ) is validated by Rehm-Weller formalism. The evolution of radical anion and cation of TDI-TPA<sub>4</sub> in polar solvent is characterized by redox titration, spectroelectrochemistry, and femtosecond transient absorption measurements. Our findings manifest the potential of TDI-TPA<sub>4</sub> donor-acceptor chromophore as a polarity-sensitive NIR single-photon emitter and pave the way for a wide range of futuristic applications.

## Data availability

All the experimental data were provided in the ESI.<sup>†</sup>

## Author contributions

The manuscript was written through the contributions of all authors.

## Conflicts of interest

There are no conflicts to declare.

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