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Near-infrared emitting single- and mixed-ligand MOFs that can be excited in the visible region: synthesis, crystal structures and sensitization of Nd³⁺, Er³⁺ and Yb³⁺†

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Near-infrared (NIR) luminescent lanthanide(III) (Ln³⁺) metal-organic frameworks (MOFs) constitute a class of materials possessing unique advantages for a wide range of applications due to their photophysical and structural properties. In this work, we present the synthesis and characterization of two series of Ln³⁺-based MOFs emitting in the NIR region (Ln-MOFs) constructed using 2,6-naphthalene dicarboxylic acid (H₂NDC) and its two amino derivatives, 4-aminonaphthalene-2,6-dicarboxylic acid (H₂ANDC) and 4,8-diaminonaphthalene-2,6-dicarboxylic acid (H₂DANDC). The first series comprises mixed-ligand and mixed-metal MOFs with the general formula [La_{1-x}Ln_x(NDC)_{1-y}(ANDC)_yCl(DMF)] (Ln = Nd³⁺, Yb³⁺, x = 0.005–1, y = 0–0.2), while the second series features MOFs corresponding to the formulae [La_{1-x}Ln_x((NDC)_{1-y}(ANDC)_y)_{1.5}(DMF)₂].DMF (Ln = Nd³⁺, Er³⁺, Yb³⁺, Lu³⁺, x = 0.1 or 1, y = 0.45 or 1) and [Ln(DANDC)_{1.5}(DMF)₂].DMF (Ln = Nd³⁺, Er³⁺, Yb³⁺). Detailed studies of the photophysical properties of Ln-MOFs have demonstrated that MOF scaffolds provide an efficient sensitization of NIR-emitting Ln³⁺ ions (η_{sens} up to 38% for Nd³⁺ analogues). In addition, Ln-MOFs possess long Ln³⁺ luminescence lifetimes (up to 12 μ s for Yb³⁺ analogues) and exhibit Ln³⁺ luminescence quantum yields comparable with the highest values reported today for NIR-emitting Ln-MOFs. We discuss and compare the photophysical processes that lead to Ln³⁺ sensitization in two series of Ln-MOFs. Most importantly, the studied Ln-MOFs are among the few examples of Ln³⁺-based NIR-emitting MOFs where Ln³⁺ emission can be sensitized upon excitation in the visible range (up to 450 nm).

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Introduction

Near-infrared (NIR) luminescent materials have been the subject of intense research over the last few decades primarily in the fields of telecommunications, light emitting diodes, lasers, sensing and biological imaging.^{1–4} Thanks to recent major advances in light detection technology, analytical techniques based on luminescence are reaching sensitivity levels comparable to those attained by methods based on radioactive isotopes.^{5–8} As a result, optical imaging methods, such as fluorescence microscopy, have gained an increasing importance.^{9–11} For instance, today, NIR fluorescence imaging offers molecular specificity, real-time detection, and tissue penetration in the centimeter range.^{12–14} Currently, the

field of NIR bioimaging agents primarily involves highly conjugated organic fluorophores, which, despite significant success, suffer from important limitations associated with the nature of their fluorescence due to spin-allowed transitions, including small Stokes shifts and short emission lifetimes (typically in the ns range) comparable to those of the biological background (auto-fluorescence), thus precluding the use of temporal discrimination for the removal of this unwanted signal.¹⁵ Furthermore, organic NIR fluorescent probes very often exhibit strong and rapid photobleaching. They exhibit broad emission bands within the NIR-I imaging window (650–900 nm), while those covering the NIR-II window (1000–1700 nm) are relatively rare.^{16,17} In contrast, NIR-emitting trivalent lanthanide ions (Ln³⁺) such as Nd³⁺, Er³⁺ and Yb³⁺ display long-lived emission signals (up to several μ s), and sharp, atomic-like emission profiles, the wavelengths of which are largely unaffected by the environment around the Ln³⁺ centre.¹⁸ Additionally, Ln³⁺ emission bands lie within the NIR-I and NIR-II imaging windows and thus offer an attractive complementary alternative to organic fluorophores.^{19–22} The forbidden nature of most f–f transitions (typical

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absorption coefficients $< 1\text{--}10\text{ M}^{-1}\text{ cm}^{-1}$) implies the generation of a low luminescence intensity through direct Ln^{3+} excitation, and thus can only be achieved by the use of high-power sources such as lasers. This limitation can be tackled by the indirect population of f-f excited states *via* energy transfer from highly absorbing organic chromophores which are either in sufficiently close proximity or directly coordinated to Ln^{3+} , an approach known as the “antenna effect” or “ Ln^{3+} sensitisation”.^{18,23} Despite the existence of a large number of antenna ligands for the efficient sensitisation of visible-emitting Ln^{3+} , the emission quantum yields of NIR-emitting Ln^{3+} complexes are generally modest (typically $\ll 1\%$) because of strong non-radiative deactivation of low-lying f-f excited states by the overtones of high-energy oscillators such as O–H, N–H and C–H bonds that are often present in the vicinity of the Ln^{3+} centre. Approaches to alleviate this limitation include the use of deuterated or fluorinated ligands,^{24–27} the synthesis of metallacrowns^{28–30} where the Ln^{3+} is encapsulated in a d-metal containing scaffold devoid of high energy oscillators, and the study of dendritic systems.^{31–34}

Ln^{3+} metal–organic frameworks (Ln-MOFs) constitute a highly promising class of NIR-emitting materials as they allow incorporation of a large number of antenna ligands and emitting ions per unit volume within a rigid framework. This densification strategy has been proven to be successful for generating sufficiently strong and easily recordable signals despite the limited emission of each individual Ln^{3+} centre.^{35–38} An additional attractive feature of Ln-MOFs is the combination of Ln^{3+} luminescence properties and the porosity of their structures, enabling them to potentially host a variety of chemical species. Thus, Ln-MOFs are highly promising platforms for numerous applications such as luminescence sensing,^{39–41} ratiometric thermometry,^{42–46} anticounterfeiting,⁴⁷ encoding applications,^{48–54} bioimaging and biosensing.^{36,55–59} For the latter applications, it is a major advantage to ensure that the Ln^{3+} emission is sensitised in the visible/NIR region in order to avoid unwanted effects caused by UV irradiation to biological systems and to maximise the penetration of the excitation light

within biological tissues and fluids. It is thus important to gain a deeper understanding of the various processes involved in the sensitization of NIR-emitting Ln^{3+} for the rational design of luminescent probes.

In this work, we describe the design, synthesis and photo-physical properties of two series (series 1 and 2) of NIR-emitting MOFs incorporating three Ln^{3+} of different natures, Nd^{3+} , Er^{3+} and Yb^{3+} as emitting centres, and the ligands 2,6-naphthalene dicarboxylate (NDC^{2-}), 4-aminonaphthalene-2,6-dicarboxylate (ANDC^{2-}) and 4,8-diaminonaphthalene-2,6-dicarboxylate (DANDC^{2-}) as Ln^{3+} sensitizers. Series 1 and 2 differ by their structures. In addition to MOFs containing only one type of ligand and metal ion, we have also synthesized mixed-ligand and mixed-metal Ln-MOFs for both series to perform a more detailed systematic investigation of the photophysical processes leading to Ln^{3+} sensitization and of the quenching pathways due to intermetallic Ln–Ln energy transfer. The Ln-MOFs of series 1 can be described by the general formula $[\text{La}_{1-x}\text{Ln}_x(\text{NDC})_{1-y}(\text{ANDC})_y\text{Cl}(\text{DMF})]$ (shorthand 1- Ln_xANDC_y , $\text{Ln} = \text{Nd}^{3+}$, Yb^{3+} , $x = 0.005\text{--}1$, $y = 0\text{--}0.2$; DMF = dimethylformamide) and those of series 2 can be defined by the general formulae $[\text{La}_{1-x}\text{Ln}_x((\text{NDC})_{1-y}(\text{ANDC})_y)_{1.5}(\text{DMF})_2]\cdot\text{DMF}$ (shorthand 2- Ln_xANDC_y , $\text{Ln} = \text{Nd}^{3+}$, Er^{3+} , Yb^{3+} , Lu^{3+} , $x = 0.1$ or 1 , $y = 0.45$ or 1) and $[\text{Ln}(\text{DANDC})_{1.5}(\text{DMF})_2]\cdot\text{DMF}$ (shorthand 2- LnDANDC , $\text{Ln} = \text{Nd}^{3+}$, Er^{3+} , Yb^{3+}) (Table 1). We demonstrate that the NIR Ln^{3+} emission can be sensitized upon excitation of the ligands at wavelengths up to 450 nm. We discuss how the structural characteristics of these Ln-MOFs impact the energy transfer and Ln^{3+} sensitization processes. Overall, the Ln-MOFs synthesized in this work show promising potential for the design of NIR optical imaging agents for biological applications.

Results and discussion

Synthesis

We synthesized two series of Ln-MOFs ($\text{Ln} = \text{La}^{3+}$, Nd^{3+} , Er^{3+} , Yb^{3+}) assembled using the carboxylic acids H_2NDC , H_2ANDC and

Table 1 List of synthesized Ln-MOFs discussed in the main text^a

Series	Metal (mol%)	Ligand (mol%)	Code	Series	Metal (mol%)	Ligand (mol%)	Code
1	Yb	H_2ANDC	20	1	0.5		1-Yb _{0.005} ANDC _{0.2}
					1	100	1-Yb _{0.01} ANDC _{0.2}
					2	100	1-Yb _{0.02} ANDC _{0.2}
					5	100	1-Yb _{0.05} ANDC _{0.2}
					10	100	1-Yb _{0.1} ANDC _{0.2}
					15	100	1-Yb _{0.15} ANDC _{0.2}
					20	100	1-Yb _{0.2} ANDC _{0.2}
					30	Nd	1-Yb _{0.3} ANDC _{0.2}
					50	Er	1-Yb _{0.5} ANDC _{0.2}
					100	Yb	2-Yb ₁ ANDC ₁
					10		2-Yb _{0.1} ANDC ₁
2	Lu	100	100	2	Nd	100	2-LuANDC
					Er	100	
					Nd	100	
					Yb	100	
					Er	100	2-ErDANDC
					Nd	100	2-NdDANDC
		100	2-YbDANDC				

^a General formulae of the Ln-MOFs: $[\text{La}_{1-x}\text{Ln}_x(\text{NDC})_{1-y}(\text{ANDC})_y\text{Cl}(\text{DMF})]$ (shorthand 1- Ln_xANDC_y , $\text{Ln} = \text{Nd}^{3+}$, Yb^{3+} , $x = 0.005\text{--}1$, $y = 0\text{--}0.2$; DMF = dimethylformamide), $[\text{La}_{1-x}\text{Ln}_x((\text{NDC})_{1-y}(\text{ANDC})_y)_{1.5}(\text{DMF})_2]\cdot\text{DMF}$ (shorthand 2- Ln_xANDC_y , $\text{Ln} = \text{Nd}^{3+}$, Er^{3+} , Yb^{3+} , Lu^{3+} , $x = 0.1$ or 1 , $y = 0.45$ or 1) and $[\text{Ln}(\text{DANDC})_{1.5}(\text{DMF})_2]\cdot\text{DMF}$ (shorthand 2- LnDANDC , $\text{Ln} = \text{Nd}^{3+}$, Er^{3+} , Yb^{3+}).



H₂DANDC (the doubly deprotonated forms of the organic ligands are referred in the main text as NDC, ANDC and DANDC). Ln-MOFs of both series were prepared using modified methods previously described by our group.⁶⁰ The synthesis of series 1 Ln-MOFs was performed using Ln³⁺ chlorides as metal sources in combination with the mixtures of H₂NDC/H₂ANDC; the molar ratio of H₂ANDC : H₂NDC was varied from 0 to 0.25 while the total ligand to metal molar ratio was kept at 1:1. The reaction was conducted at 110 °C during 24 h in DMF. In addition, we synthesized Yb³⁺-doped MOFs by adding specified amounts of a standard solution of the dopant in DMF to the reaction mixture. The Ln-MOFs of series 1 can be described by the following formula [La_{1-x}Ln_x(NDC)_{1-y}(ANDC)_yCl(DMF)] (Ln = Nd³⁺, Yb³⁺, x = 0.005–1, y = 0–0.2). Series 2 features Nd³⁺, Er³⁺ and Yb³⁺ Ln-MOFs of the general formulae ([La_{1-x}Ln_x((NDC)_{1-y}(ANDC)_y)_{1.5}(DMF)₂]-DMF (x = 0.1 or 1, y = 0.45 or 1) and [Ln(DANDC)_{1.5}(DMF)₂]-DMF). They were synthesized by the reaction between the Ln³⁺ nitrates, as the metal sources, with the ligands H₂ANDC and H₂DANDC in a 1:1.5 molar ratio in a DMF/H₂O/EtOH 10:1:1 solvent mixture at 70 °C for 72 h. A detailed list of the synthesized Ln-MOFs is provided in Table 1.

Crystal structures from X-ray diffraction on single crystals

Ln-MOFs of series 1 crystallize in the orthorhombic space group *Pnma*. We have previously reported the crystal structure of the visible light-emitting 1-LaNDC MOF,⁶⁰ and here we will only briefly describe the one of 1-NdNDC. The crystal structures of series 1 and 2 Ln-MOFs are summarized in Fig. 1. The Nd³⁺ ion is coordinated to nine atoms, six carboxylate oxygen atoms from the NDC ligand, one oxygen atom from a terminal DMF molecule and two chloride anions. The coordination

polyhedron is best described as a tricapped trigonal prism. The metal centres are bridged by the chloride ions and the carboxylate oxygens to form infinite zig-zag chains along the *a* axis. Neighbouring chains are connected by bridging NDC ligands to form a 3D framework with rhombic channels along the *a* axis, which are filled with terminally coordinated DMF molecules.

Single-crystal X-ray analysis of Ln-MOFs from series 2 revealed two closely related structures, associated with the early (*e.g.* Nd³⁺) and the late (*e.g.* Yb³⁺) lanthanide(III) ions. The 2-Nd₁ANDC₁ MOF crystallizes in the triclinic space group *P* $\bar{1}$. The Nd³⁺ centre is coordinated with nine oxygen atoms, namely seven from the carboxylate groups of the ANDC ligand and two from the terminally coordinated DMF molecules. The coordination [NdO₉] polyhedron is best described as a distorted capped square antiprism. Two adjacent Nd³⁺ polyhedra form an edge-sharing [Nd₂O₁₆] dimer (Fig. 1d). Neighbouring dimers are bridged by ANDC ligands with hapticities μ_2 : η^1 – η^2 and μ_2 : η^1 – η^1 forming a three-dimensional framework with parallelepiped channels along the *a* axis, which are occupied by terminally coordinated and guest DMF molecules.

The Ln-MOFs 2-Yb₁ANDC₁ and 2-LuANDC share the general structural characteristics with 2-Nd₁ANDC₁, with a few subtle differences attributed to the lanthanide(III) contraction.⁶¹ First, the Yb³⁺ and Lu³⁺ centres are coordinated to a total of eight oxygen atoms, six from carboxylate groups of ANDC and two from the terminally coordinated DMF molecules, with the [LnO₈] coordination polyhedra (Ln = Yb³⁺, Lu³⁺) being best described as distorted square antiprisms. Two adjacent polyhedra are bridged by four μ_2 : η^1 – η^1 carboxylate groups from ANDC ligands to form the dimeric SBU which, in this case, comprises two distinct [LnO₈] coordination polyhedra (Fig. 1e).

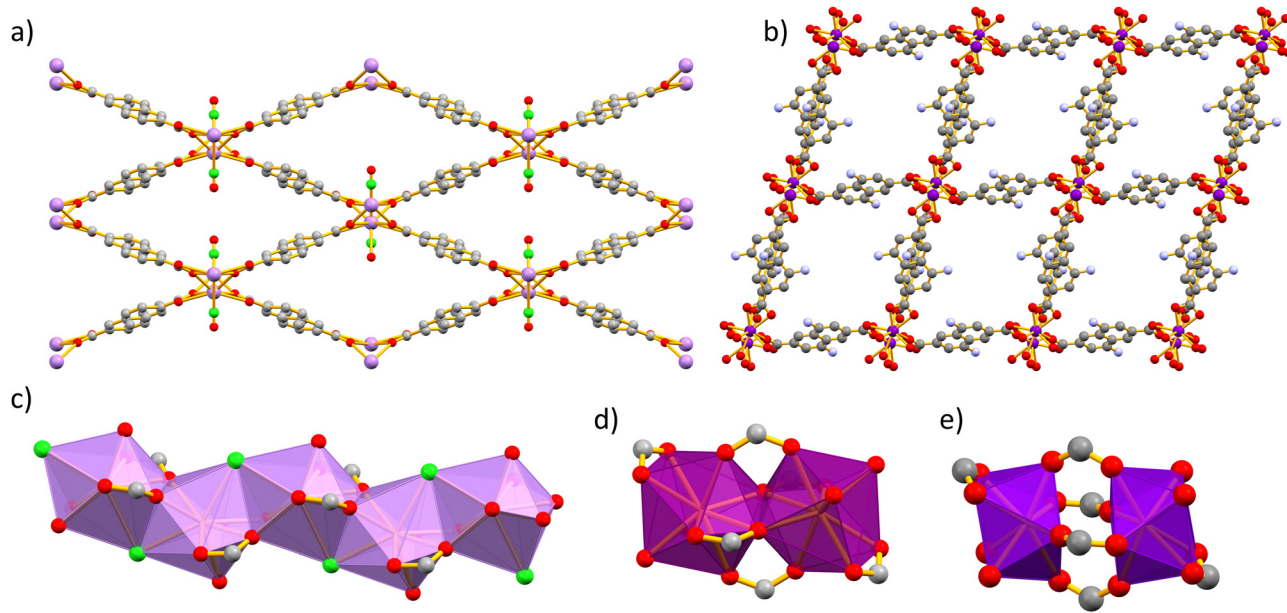


Fig. 1 Crystal structures obtained from X-ray diffraction on single crystals of (a) series 1 Ln-MOFs viewed along the *a* axis and (b) series 2 viewed along the *a* axis (the DMF molecules are omitted for clarity). (c) The zig-zag chain of the [NdO₉Cl] polyhedra of series 1 Ln-MOFs viewed along the *c* axis. (d)–(e) Secondary building units (SBUs) of series 2 Ln-MOFs. Colour code: O: red, C: grey, N: blue, Cl: green, Nd: lilac, Er: magenta, and Yb: purple.



Powder X-ray diffraction

The phase purity of the Ln-MOFs was analyzed by powder X-ray diffraction (PXRD) measurements (Fig. S4 and S5, ESI†). It was demonstrated that 1-Ln_xANDC_y MOFs retained their phase purity up to doping concentrations of $x \leq 0.2$ and $y \leq 0.5$ (Fig. S4, ESI†). Furthermore, when the value of y is more than 0.3, formation of the $P\bar{1}$ phase is preferred. The PXRD patterns of 2-Ln_xANDC_y and 2-LnDANDC are in a good agreement with their respective simulated diagrams. In addition, there is a shift of the diffraction peak appearing at the lowest angle for 2-LnANDC and 2-LnDANDC MOFs that depends on the nature of Ln³⁺ ion. More specifically, the observed peak positions range from 2-theta values equal to 6.85° for Ln = Nd³⁺ to 7.00° for Ln = Er³⁺ and 7.03° for Ln = Yb³⁺. This observation reflects a slight shrinkage of the MOF, which can be attributed to lanthanide(III) contraction (Fig. S5, ESI†).⁶¹

Photophysical studies

In order to assess the ability of our new MOFs to sensitize Nd³⁺, Yb³⁺ and Er³⁺, steady-state and time-resolved luminescence spectra were collected on microcrystalline samples of Ln-MOFs at room temperature. The absorption spectra of the studied bridging ligands in MeOH solution are given in Fig. 2. H₂NDC displays two structured π - π^* absorption bands in the UV range centred at 280 and 340 nm with the absorption tailing off at *ca.* 360 nm. H₂ANDC and H₂DANDC show similar absorption features in the UV range in addition to two substantially red-shifted absorption bands centred at 385 and 412 nm, respectively. The latter features, that tail off at *ca.* 450 (H₂ANDC) and 460 nm (H₂DANDC), can be assigned to charge transfer transitions involving an electron from an orbital localized mainly on the amino group to a vacant π^* orbital of the aromatic core.⁶²

Representative optical spectra of Ln-MOFs are collected in Fig. 3. Excitation spectra of 2-Nd₁ANDC₁, 2-Er₁ANDC₁ and 2-Yb₁ANDC₁ MOFs upon monitoring of the emission of the

respective Ln³⁺ (Nd³⁺: $^4F_{3/2} \rightarrow ^4I_{11/2}$ at 1064 nm, Er³⁺: $^4I_{13/2} \rightarrow ^4I_{15/2}$ at 1525 nm and Yb³⁺: $^2F_{5/2} \rightarrow ^2F_{7/2}$ at 980 nm) exhibit similar broad bands in the range of 250–470 nm (Fig. 3f). This observation confirms that for all three 2-LnANDC MOFs the sensitization of different NIR-emitting Ln³⁺ ions is occurring through the electronic states of the chromophoric ligands. The excitation spectra of 1-Ln_xANDC_y MOFs in the range up to 380 nm are dominated by the bands associated with the 2,6-naphthalene dicarboxylate chromophoric group. The ANDC contribution appears as a shoulder located around 400 nm, which becomes more pronounced as the ANDC content increases (Fig. 3c, d and Fig. S6, ESI†). Additionally, we observed narrow bands that correspond to direct excitation occurring through the f-f transitions of the respective Ln³⁺. For example, 2-Nd₁ANDC₁ exhibits these narrow excitation bands at 524, 582, 683, 746, 803 and 867 nm that can be attributed to the Nd³⁺ $^4I_{9/2} \rightarrow ^4G_{7/2}$, $^4G_{5/2}$, $^2H_{11/2}$, $^4F_{9/2}$, $^4F_{7/2} + ^4S_{3/2}$ and $^4F_{5/2} + ^2H_{9/2}$ transitions, respectively. 2-ErANDC shows the Er³⁺-based narrow excitation bands at 520 and 656 nm arising from the $^4I_{15/2} \rightarrow ^2H_{11/2}$, $^4F_{9/2}$ transitions, respectively. On the other hand, no such sharp features are observed in the excitation spectrum of 2-YbANDC since Yb³⁺ does not possess any electronic levels in the UV-visible range.

Upon excitation of the chromophoric ligands at 400 nm, 2-LnANDC MOFs (Ln = Nd, Er, Yb) exhibit characteristic Ln³⁺ luminescence signals indicating that the ligand sensitizes all three NIR-emitting Ln³⁺ ions *via* the antenna effect (Fig. 3g). In particular, the emission spectrum of 2-NdANDC revealed the presence of the Nd³⁺ bands centred at 873, 1056 and 1330 nm, attributed to the $^4F_{3/2} \rightarrow ^4I_{9/2}$, $^4I_{11/2}$, and $^4F_{3/2} \rightarrow ^4I_{13/2}$ electronic transitions, while those of 2-ErANDC and 2-YbANDC exhibit the Er³⁺ (centred at 1514 nm) and Yb³⁺ (centred at 978 nm) emission bands attributed to the $^4I_{13/2} \rightarrow ^4I_{15/2}$ and $^2F_{5/2} \rightarrow ^2F_{7/2}$ transitions, respectively. The 2-LnDANDC MOFs show similar behaviour (Fig. 3h–i) with respect to the sensitization of NIR-emitting Ln³⁺. In addition, both series of Ln-MOFs containing DANDC or ANDC ligands exhibit residual broad-band ligand-centred emission in the visible range centred at 480 and 500 nm, respectively (Fig. 3e), indicating that the energy transfer from the MOF sensitizers to the Ln³⁺ is incomplete for all three Ln³⁺ cations investigated in this work.

We further investigated the contribution of the NDC ligand on the sensitization of NIR-emitting Ln³⁺ by analyzing the photophysical properties of the mixed-ligand Ln-MOFs, *i.e.* 1-Nd_xANDC_y ($y = 0-0.2$) and 1-Yb_xANDC_{0.2} ($x = 0.005-0.5$). We observed that, in both cases, upon excitation either into the NDC-centred bands at 350 nm or into the ANDC-centred bands at 415 nm, the characteristic emissions of Nd³⁺ and Yb³⁺ in the NIR range are generated (Fig. 3a and b and Fig. S6, S7, ESI†). This observation indicates that NDC contributes to the sensitization of the NIR-emitting Ln³⁺ either directly or indirectly *via* an NDC-to-ANDC energy transfer.⁶⁰ Excitation spectra of 1-Ln_xANDC_y MOFs (Fig. 3c and d) exhibit similar ligand-centred spectral features as the ones observed in the diffuse reflectance spectra further confirming Ln³⁺ sensitization through the Ln-MOF ligands. Moreover, in the case of 1-Nd_xANDC_y MOFs,

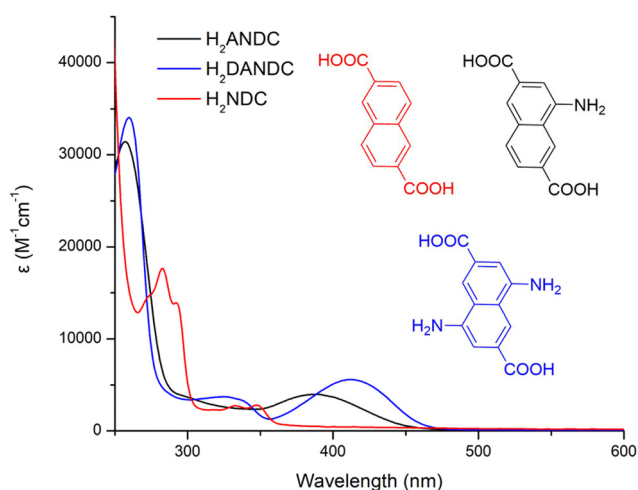


Fig. 2 Absorption spectra of the ligands H₂NDC, H₂ANDC and H₂DANDC in MeOH (10⁻⁵ M).



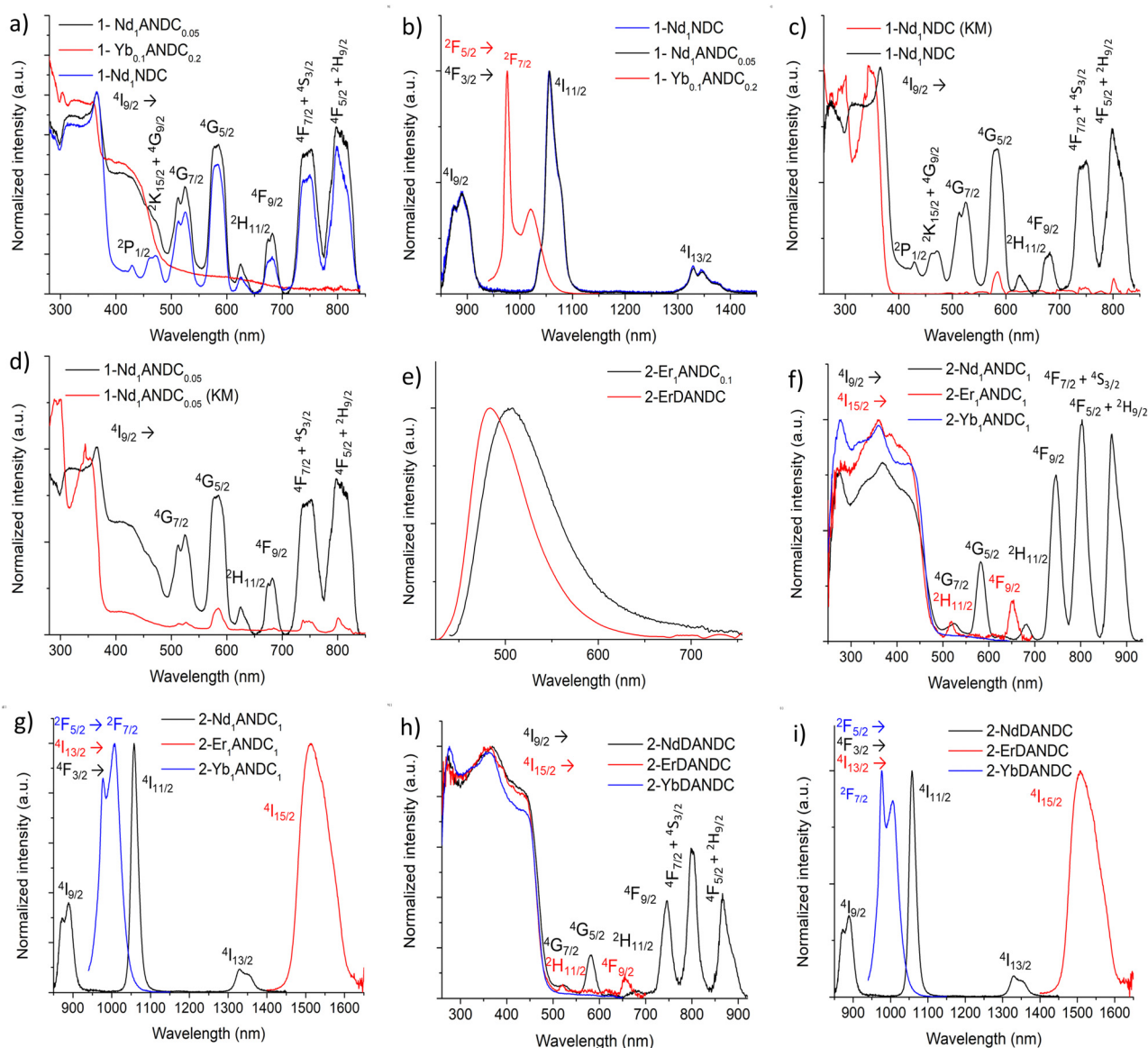


Fig. 3 (a) Excitation spectra of 1-Nd₁ANDC_{0.05}, 1-Yb_{0.1}ANDC_{0.2} and 1-Nd₁NDC. (b) Emission spectra of 1-Nd₁NDC, 1-Nd₁ANDC_{0.05} and 1-Yb_{0.1}ANDC_{0.2} ($\lambda_{\text{exc}} = 350$ nm). Overlay of excitation and diffuse reflectance spectra, after conversion to Kubelka–Munk (KM) function, for (c) 1-Nd₁NDC and (d) 1-Nd₁ANDC_{0.05}. (e) Representative ligand-centred emission spectra of 2-Er₁ANDC_{0.1} and 2-ErDANDC ($\lambda_{\text{exc}} = 400$ nm). (f) Excitation spectra of 2-Nd₁ANDC₁, 2-Er₁ANDC₁ and 2-Yb₁ANDC₁ and (g) emission spectra of 2-Nd₁ANDC₁, 2-Er₁ANDC₁ and 2-Yb₁ANDC₁ ($\lambda_{\text{exc}} = 400$ nm). (h) Excitation spectra of 2-NdDANDC, 2-ErDANDC and 2-YbDANDC and (i) emission spectra of 2-NdDANDC, 2-ErDANDC and 2-YbDANDC ($\lambda_{\text{exc}} = 400$ nm). The excitation spectra were recorded upon monitoring the emission maximum of the corresponding Ln³⁺. All spectra were corrected for detector response and normalized.

sharp features corresponding to the Nd³⁺ transitions are observed. It should be noted that the band located in the range 400–450 nm, attributed to the presence of the ANDC ligand, becomes more pronounced upon increasing the ANDC content in the 1-Ln_xANDC_y MOFs (Fig. S6 and S7, ESI[†]). The absolute quantum yields of Ln³⁺ emission under excitation into the electronic levels of the chromophoric ANDC and DANDC ligands ($\Phi_{\text{Ln}}^{\text{L}}$) were collected for all relevant Ln-MOFs (Table 2). It was found that for 1-NdANDC_y ($y = 0.05$ – 0.2) MOFs the ANDC loading level has no effect on the Nd³⁺ emission quantum yields even when the excitation takes place selectively

on the ANDC chromophore. This observation potentially reflects the presence of a mechanism in which initial ANDC excitation is followed by the thermal equilibration of the NDC and ANDC triplet (T_1) levels ($\Delta E \approx 1600$ cm⁻¹, *vide infra*), with both acting as sensitizers for Nd³⁺. In contrast to 1-NdANDC_y MOFs, the quantum yields of the Yb³⁺-based emission in mixed-metal Yb/La-MOFs with $x = 0.1$ show significant dependence on the ANDC loading. Notably, there is a ten-fold increase of the Yb³⁺ emission quantum yield, upon excitation at 410 nm, between 1-Yb_{0.1}ANDC_{0.2} and 2-Yb_{0.1}ANDC₁ where the ANDC content increases from 20 to 100 mol%. This can be



Table 2 Ln³⁺-centred luminescence quantum yields ($\Phi_{\text{Ln}}^{\text{L}}$) and observed luminescence lifetimes (τ_{obs}) of the studied Ln-MOFs in the solid state^a

Ln-MOF	$\Phi_{\text{Ln}}^{\text{L}}(10^{-2}\%)^b$	τ_{obs} (μs) ^c
1-NdNDC	13.03(6) ^c	1.035(4)
1-Nd ₁ ANDC _{0.01}	12.2(4) ^c	1.032(3)
1-Nd ₁ ANDC _{0.02}	10.5(1) ^c	1.06(1)
1-Nd ₁ ANDC _{0.05}	9.52(7)	1.08(1)
1-Nd ₁ ANDC _{0.1}	10(7)	1.072(8)
1-Nd ₁ ANDC _{0.2}	10.8(1)	1.03(1)
2-Nd ₁ ANDC	11.5(2)	0.96(1)
2-Nd ₁ ANDC _{0.45}	17.3(2)	0.957(6)
2-Nd ₁ ANDC _{0.1}	8.1(2)	1.018(5)
2-Nd ₁ DANDC	7.2(5)	0.81(4)
1-Yb _{0.005} ANDC _{0.2}	0.10(1)	N/A
1-Yb _{0.01} ANDC _{0.2}	0.116(9)	1.349(8)
1-Yb _{0.02} ANDC _{0.2}	0.24(1)	14.9(1)
1-Yb _{0.05} ANDC _{0.2}	0.58(2)	14.3(1)
1-Yb _{0.1} ANDC _{0.2}	1.56(2)	14.8(4)
1-Yb _{0.15} ANDC _{0.2}	2.46(5)	15.63(6)
1-Yb _{0.2} ANDC _{0.2}	5.41(5)	15.09(8)
1-Yb _{0.3} ANDC _{0.2}	0.880(1)	14.4(1)
1-Yb _{0.5} ANDC _{0.2}	1.93(2)	12.1(4)
2-Yb _{0.1} ANDC ₁	13.6(3)	11.35(2)
2-Yb ₁ ANDC ₁	14.7(3)	9.68(2)
2-YbDANDC	21.4(5)	8.11(4)
2-ErANDC	0.73(4)	1.68(2)
2-ErDANDC	0.48(5)	1.27(1)

^a Data obtained at room temperature. Standard deviation (2σ) is between parentheses; estimated relative errors: τ_{obs} , $\pm 2\%$; $\Phi_{\text{Ln}}^{\text{L}}$, $\pm 10\%$.

^b Under excitation at 400–410 nm. ^c Under excitation at 355 nm.

mainly attributed to the more efficient ANDC-Yb³⁺ interaction in 2-Yb_{0.1}ANDC₁ since effectively all Yb³⁺ ions are directly connected to an ANDC sensitizer. On the other hand, the Yb³⁺ emission quantum yields and lifetimes are comparable for 2-Yb₁ANDC₁ and 2-Yb_{0.1}ANDC₁ MOFs (Table 2) indicating that there is no significant Yb–Yb quenching pathway.

In addition, for 2-Nd₁ANDC₁, 2-Nd₁ANDC_{0.1} and 2-NdDANDC MOFs, the intrinsic quantum yields of the Nd³⁺ emission under direct excitation into its f–f transition at 750 nm ($\Phi_{\text{Ln}}^{\text{Ln}}$) could be acquired and were found to be 0.307(7) %, 0.293(1) % and 0.201(1) %, respectively. Using the values of $\Phi_{\text{Ln}}^{\text{Ln}}$ and $\Phi_{\text{Ln}}^{\text{L}}$, we calculated the sensitization efficiencies of Nd³⁺ (η_{sens}) according to the formula:

$$\Phi_{\text{Ln}}^{\text{L}} = \eta_{\text{sens}} \times \Phi_{\text{Ln}}^{\text{Ln}}$$

We found that the values of η_{sens} are similar for 2-Nd₁ANDC₁ and 2-NdDANDC ($38 \pm 8\%$ vs. $35 \pm 5\%$) and slightly lower for 2-NdANDC_{0.1} ($27 \pm 1\%$).

Experimental luminescence decay curves of Ln-MOFs (Ln = Yb³⁺, Nd³⁺, Er³⁺) measured under excitation of the chromophoric sensitizers ($\lambda_{\text{ex}} = 355$ nm) and upon monitoring the NIR emission signals arising from the main transitions of the corresponding Ln³⁺ (Nd³⁺: ⁴F_{3/2} → ⁴I_{1/2} at 1064 nm, Er³⁺: ⁴I_{13/2} → ⁴I_{15/2} at 1525 nm and Yb³⁺: ²F_{5/2} → ²F_{7/2} at 980 nm) are best fitted with a monoexponential function. Such behaviour indicates that Ln³⁺ ions are located in a single well-defined coordination environment within the MOF scaffold. The measured values of the observed luminescence lifetimes (τ_{obs}) are in the order of 1 μs for Nd³⁺ MOFs and range between

10 and 15 μs for the MOFs incorporating Yb³⁺ (Table 2). The similarity of the values of τ_{obs} for a particular Ln³⁺ within the studied series of MOFs suggests that they are insignificantly affected by the nature of the ligands and their content in Ln-MOFs. The values of τ_{obs} are comparable with the ones reported for similar NIR-emitting Ln-MOFs.^{63,64}

Ligand-centred photophysical properties and the Ln³⁺ sensitization mechanism

To further investigate the energy transfer processes that take place between NDC, its amino- derivatives, and the Ln³⁺ ions within the Ln-MOF matrices, we conducted steady-state and time-gated luminescence analysis at a low temperature (77 K) of the 1-Nd₁NDC, 1-LaANDC_{0.02} and 1-LaDANDC_{0.02} MOFs. The energy of an excited state is defined as the energy gap between the lowest vibrational level of the excited state and that of the corresponding ground state (the so-called 0–0 energy gap). Spectroscopically, this energy is quantified by identifying the 0–0 transition at low temperature as the highest-energy vibrational feature of the corresponding emission signal. Consequently, the energy of the lowest spin-allowed (S₁) excited state is determined by deconvoluting the fluorescence spectrum measured at 77 K. Similarly, the energy of the long-lived lowest spin-forbidden state (T₁) can be quantified from the corresponding phosphorescence spectrum which is measured in a time-gating mode by applying a delay of *ca.* 1 μs to remove short-lived fluorescence signals. Using the described method, we determined the energies of the ligands' excited states by performing Gaussian deconvolution of the emission spectra of the mixed-ligand Ln-MOFs acquired at 77 K in the steady-state mode for the S₁ excited states and in the time-gated mode for the T₁ excited states (Fig. S11, ESI[†]).⁶⁵ The energy diagram of Fig. 4 shows that the excited levels, including the long-lived T₁ states of NDC, ANDC and DANDC, lie above the main emissive levels of Nd³⁺ (⁴F_{3/2}), Er³⁺ (⁴I_{13/2}) and Yb³⁺ (²F_{5/2}), therefore, are in favourable energy positions to populate them.⁶⁶

In general, the sensitization of the Ln³⁺ in the studied series of MOFs involves the excitation of the electronic states of the NDC chromophore, since it is a dominant one in the matrix, thus creating a ligand excited state (exciton). Then, this excited state migrates through the matrix by a series of energy transfer steps (exciton hopping) until it reaches an energy acceptor site (either an amino-substituted ligand, *i.e.* ANDC or DANDC, or a low-lying excited state of Ln³⁺), or it eventually decays through radiative or non-radiative (thermal) processes.⁶⁷ To gain insight into the process of energy transfer from NDC to its amino derivatives, ANDC and DANDC, we measured and analysed the NDC-centred quantum yields (Φ^{L}) of 1-LaANDC_y and 1-LaDANDC_y with low doping levels of amino-substituted derivatives ($y = 0-0.002$) under excitation at 330 nm (experimental details in the ESI[†] and results in Table S3). This process allowed the achievement of partial NDC-to-(D)ANDC energy transfer, which is readily quantifiable with our experimental setup, instead of the almost complete quenching of NDC-centred emission which is observed in Ln-MOFs with higher acceptor loadings. We plotted the ratios of the quantum yield



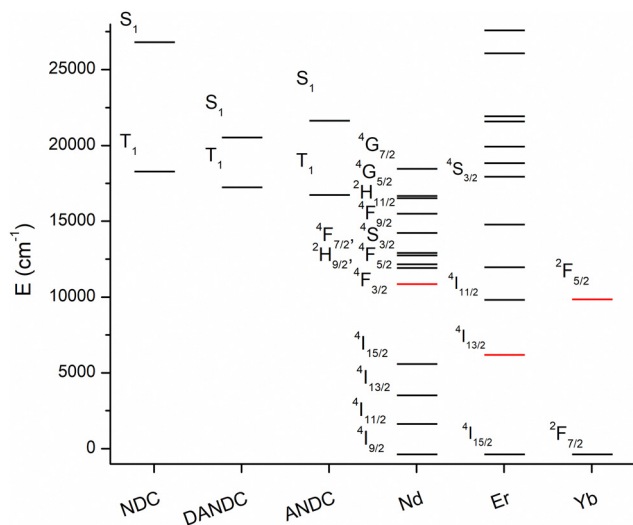


Fig. 4 Energy diagram of the singlet (S_1) and triplet (T_1) excited states of the NDC, ANDC and DANDC ligands and the electronic levels of the Ln^{3+} ions (red: emissive levels).

values obtained for the 1-LaNDC MOF (ϕ_0^L) over the ones acquired for 1-LaANDC_y and 1-LaDANDC_y (ϕ^L) MOFs against the molar fraction of ANDC and DANDC within the MOF in order to construct a Stern–Volmer type plot (Fig. 5).⁶⁷ In both cases, the quantum yield of the NDC-centred emission decreases progressively upon addition of a quencher (ANDC or DANDC) due to the incorporation of more low-energy acceptor sites into the MOF. We found that the slopes observed in the Stern–Volmer plots for 1-LaANDC_y and 1-LaDANDC_y are in the same order of magnitude. The steeper slope for the latter possibly indicates that the energy transfer efficiency from the NDC ligand is more efficient for the DANDC ligands than for the ANDC ones.^{68,69}

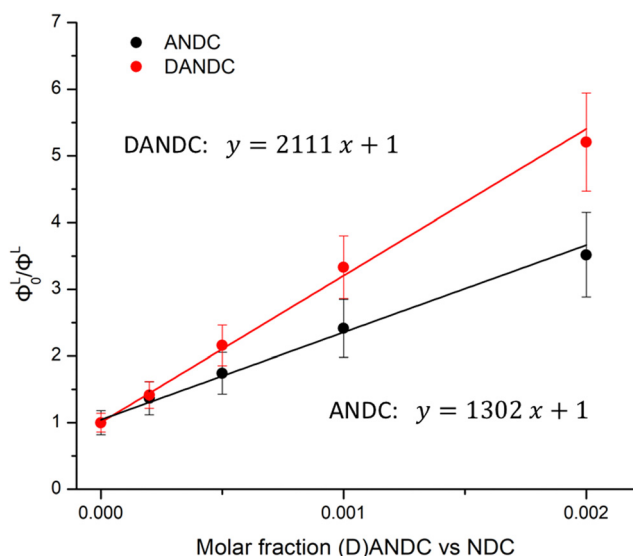


Fig. 5 Stern–Volmer-type plots of ϕ_0^L/ϕ^L for LaANDC_y and LaDANDC_y over the molar fractions of ANDC (red) and DANDC (black) ($y = 0-0.002$) in the MOF ($R^2 > 0.996$). The values of the corresponding quantum yields are given in Table S3 (ESI†).

Conclusions

In summary, we have synthesized two new series of Ln-MOFs ($\text{Ln} = \text{La}^{3+}, \text{Nd}^{3+}, \text{Er}^{3+}, \text{Yb}^{3+}, \text{Lu}^{3+}$) featuring H₂NDC and its amino derivatives, H₂ANDC and H₂DANDC, with the general formulae $[\text{La}_{1-x}\text{Ln}_x(\text{NDC})_{1-y}(\text{ANDC})_y]\text{Cl}(\text{DMF})$ ($\text{Ln} = \text{Nd}^{3+}, \text{Yb}^{3+}$, $x = 0.005-1$, $y = 0-0.2$) (series 1) and $[\text{La}_{1-x}\text{Ln}_x((\text{NDC})_{1-y}(\text{ANDC})_y)_{1.5}(\text{DMF})_2]\cdot\text{DMF}$ ($\text{Ln} = \text{Nd}^{3+}, \text{Er}^{3+}, \text{Yb}^{3+}, \text{Lu}^{3+}$, $x = 0.1$ or 1 , $y = 0.45$ or 1), $[\text{Ln}(\text{DANDC})_{1.5}(\text{DMF})_2]\cdot\text{DMF}$ ($\text{Ln} = \text{Nd}^{3+}, \text{Er}^{3+}, \text{Yb}^{3+}$) (series 2). Series 1 comprises Ln-MOFs with an orthorhombic structure, while the Ln-MOFs of series 2 adopt a triclinic structure. We have demonstrated that series 1 retains its structural integrity upon incorporation of amino-substituted ligands for $y \leq 0.3$, while for higher loadings formation of the triclinic phase is preferred. In addition, Ln-MOFs constructed using 100 mol% H₂ANDC and H₂DANDC ligands crystallize in the triclinic phase with small structural variations pertinent to the lanthanide(III) contraction.⁶¹ In particular, Nd³⁺ and Er³⁺ in the aforementioned Ln-MOFs are nine-coordinated while Yb³⁺ and Lu³⁺ are eight-coordinated. Thus, the secondary building units of the former Ln³⁺ consist of edge sharing $[\text{Ln}_2\text{O}_{16}]$ dimers while Yb³⁺ and Lu³⁺ form two discrete carboxy-bridged coordination polyhedra. The photophysical analysis of the Ln-MOFs of series 1 revealed that each of the three chromophoric ligands studied in this work is efficient for the sensitization of three different lanthanide(III) cations emitting in the NIR domain, Nd³⁺, Yb³⁺ and Er³⁺. This is a major result as the chromophoric ligands that combine absorption in the visible region with the ability to sensitize NIR-emitting Ln³⁺ within a MOF scaffold remain scarce. In mixed-ligand Ln-MOFs, both NDC and ANDC contribute to the sensitization of NIR-emitting Ln³⁺. An increase of the ANDC content in these Ln-MOFs has a beneficial effect on the Nd³⁺ and Yb³⁺ emission quantum yields. The Nd³⁺ sensitization efficiencies of the series 2 Ln-MOFs were found to be similar for 2-Nd₁ANDC₁ and 2-NdDANDC, ranging within 35–38%. Overall, this work demonstrates that the use of amino-derivatives of 2,6-naphthalene dicarboxylic acid (H₂NDC) in the construction of Ln-MOFs results in shifting the absorption wavelength to the visible region while achieving efficient sensitization of NIR-emitting Ln³⁺ ions.

Experimental

All experimental information (synthesis, characterization data, X-ray crystallography, instrumentation and software used) are provided in the ESI.†

Conflicts of interest

There are no conflicts to declare.

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