

REVIEW

View Article Online
View Journal | View IssueCite this: *Energy Environ. Sci.*,
2025, 18, 2720

Revisiting the kinetics enhancement strategies of Si anodes through deconstructing particle–interface–electrode integration

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The successive introduction of silicon (Si) graphite composite anodes into the global market highlights the tremendous commercial potential of Si anodes. Good kinetic performance related to fast charging capability is the central topic of next-generation Si anodes. However, there is a lack of critical reviews exploring the fundamental limiting factors affecting the kinetics of Si and evaluating the effectiveness of the current strategies. In this review, we deconstruct the particle–interface–electrode integration to analyze key limiting factors of kinetics from a practical application perspective for the first time, including long Li⁺ diffusion distance and poor conductivity for particles, high Li⁺ migration impedance at the interface, and insufficient or even interrupted Li⁺ diffusion paths inside the electrodes. Then, the kinetics enhancement strategies for progressively addressing the above issues are systematically investigated and the quantitative relationships between kinetics and these strategies are deeply discussed. Accordingly, the challenges in quantification and balance for fast-charging Si anodes are identified as the remaining issues, and potential solutions are provided. This review provides valuable guidance on fast-charging Si anodes and suggests promising directions in commercial-oriented Si anode studies.

Received 27th November 2024,
Accepted 3rd February 2025

DOI: 10.1039/d4ee05595k

rsc.li/ees

Broader context

The pursuit of fast charging is a theme of the contemporary era in the development of high energy density lithium-ion batteries (LIBs). Currently, silicon (Si)–graphite composite anodes are recognized as the most promising choice for high energy density LIBs. However, Si anodes suffer from poor kinetics in multiple dimensions from the perspective of practical applications, resulting in performance that is far behind the targeted fast charging performance. In this context, we deconstruct the particle–interface–electrode integration to analyze key limiting factors of kinetics from a practical application perspective and evaluate effective strategies of progressively enhancing the kinetics of Si anodes, as well as deeply discuss the quantitative relationship between the kinetics and strategies of multiple levels. This review aims to provide valuable insights from an overall perspective of Si-based batteries, and we expect a significant breakthrough in the research and development of high energy density and fast-charging power LIBs based on Si anodes.

1. Introduction

Promoting electric vehicles (EVs) is an important process toward reaching the carbon neutrality target. Over 26 million electric vehicles (EVs) were on the road in 2022, indicating vigorous development.¹ However, EVs still hold a small share of the global market when compared with the dominant traditional fuel vehicles.² One major bottleneck is their long charging time, which negatively impacts user experience. Consequently, fast charging performance is becoming a critical benchmark for lithium-ion batteries (LIBs) to enhance their competitiveness and cope with more complex driving conditions, thereby

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expanding the application scenarios.^{3–11} According to the Goals for Advanced Batteries proposed by the U.S. Advanced Battery Consortium (USABC), power batteries need to obtain a fast charging performance of 80% initial capacity for 15 min charging.¹² Currently, the Tesla model S can be recharged up to the 320 km range within 15 min at supercharger locations,¹³ and BYD can shorten the charging time to 33 minutes,¹⁴ which still falls far from the targeted fast charging performance.

For the power LIB system, fast charging requires the anode to possess good kinetic performance for fast Li storage.^{15–19} Based on the properties of the intercalation mechanism, the lithiation of graphite anodes mainly occurs from the sheet edge, which forms the only path for Li⁺ to intercalate into the graphite layers, resulting in the slow kinetic performance.^{6,20–24} The limited theoretical capacity of graphite requires the adoption of thicker electrodes to achieve sufficient energy density, which prolongs the Li⁺ diffusion distance and hampers fast charging performance. Si is regarded as a promising candidate for graphite due to its high capacity ($\sim 3580 \text{ mA h g}^{-1}$, $\text{Li}_{15}\text{Si}_4$) and low potential ($< 0.3 \text{ V Li}^+/\text{Li}$).²⁵ Particularly, notable battery manufacturers such as CATL and Tesla have successfully brought to market the practice of Si-graphite composite anodes to increase the energy density, thereby extending the driving range of EVs.²⁶ Studying the kinetics behavior of Si anodes could facilitate the development of fast-charging Si-graphite composite anodes which are highly anticipated in the market.

Numerous research studies have been conducted on Si anodes from the perspective of alleviating volume expansion.^{27–30} The volume expansion of Si materials in the composite anode could be mitigated by graphite,³¹ while the poor kinetic performance of Si materials profoundly hinders the fast charging of these anodes, resulting in them significantly lagging behind the target for fast charging performance. A comprehensive review of the kinetic behavior and performance is crucial to advancing the understanding of practically-oriented Si anode iterations.

On this basis, we deconstruct the particle–interface–electrode integration and propose these limiting factors in three dimensions as follows (Fig. 1): (i) the low intrinsic conductivity and long Li⁺ diffusion distance cause slow lithiation inside Si

materials, resulting in slow kinetics at the particle level;^{32–35} (ii) the complex solid electrolyte interface (SEI) and the exposed surface of Si materials with poor conductivity result in significant obstruction on kinetics at the interface level;^{36,37} (iii) the initial electrode porosity and the electrode cracking during operation can affect the Li⁺ diffusion in the electrode, and the insufficient contact between current collector and electrode also directly impacts the integrity of electrodes, which limits the kinetic performance at the electrode level.^{38,39}

Herein, the systematic kinetics enhancement strategies of Si anodes from particle, interface, and electrode levels are categorized, as shown in Fig. 2. First, we discuss the adjustments in particle size and porosity of Si materials to balance the shortened Li⁺ diffusion distance and the increased interface impedance to obtain the optimal Li⁺ diffusion. Additionally, the element doping and particle compositing strategies are summarized to address poor conductivity at the particle level. Next, the surface coating, SEI optimization strategies, and the quantitative relationship among them are analyzed to reduce the Li⁺ migration impedance of Si materials at the interface level. Later, the strategies of decreasing electrode porosity and cracks to stabilize the electrical contact of Si materials during operation, as well as the strategy of enhancing the electrode-current collector contact for electrode integrity are reviewed to improve the kinetic performance at the electrode level. In the conclusion, we emphasize quantification and balance as the kinetics enhancement challenges of Si anodes and propose several potential strategies to shed light on future studies.

2. Particle level

Poor Li⁺ diffusion coefficients (10^{-14} – $10^{-13} \text{ cm}^2 \text{ s}^{-1}$) caused by low intrinsic conductivity of Si leads to their poor kinetic performance.⁴⁴ Numerous studies are dedicated to revealing the intrinsic mechanism.⁴⁵ Lee and co-authors investigated the mechanical interaction effects of lithiation Si clusters through *in situ* TEM.⁴⁶ As shown in Fig. 3a–c, the unconstrained Si column largely swells along the $\langle 110 \rangle$ orientation. As a

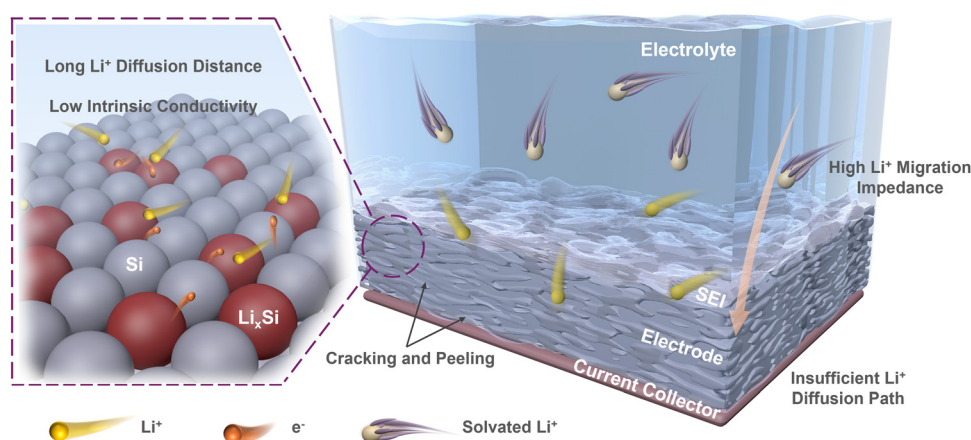


Fig. 1 Limiting factors causing poor kinetic performances of Si anodes at particle, interface, and electrode levels.



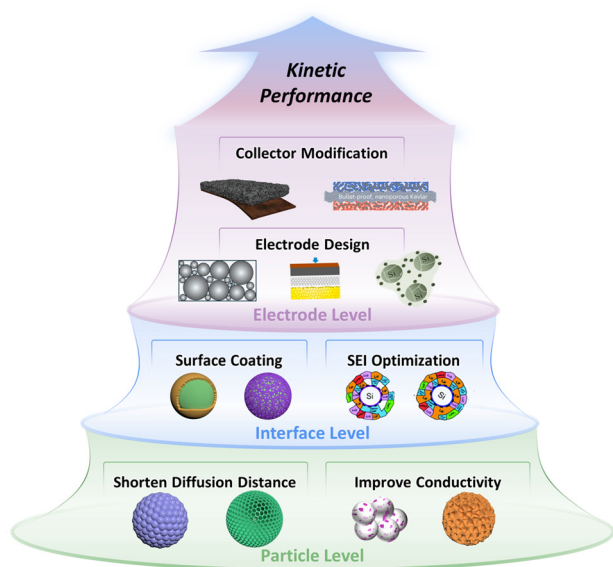


Fig. 2 Overview of kinetics enhancement strategies of particle–interface–electrode integration for Si anodes. (Partial images are cited.^{40–43} Copyright © 2019, American Chemical Society, copyright © 2020 Published by Elsevier Ltd, copyright © 2020 Published by Elsevier Ltd, copyright © 2024, The Author(s), under exclusive license to Springer Nature Limited.)

comparison, the constrained Si column first swells along the $\langle 110 \rangle$ orientation and gradually touches the two walls. And then, it expands along the $\langle 100 \rangle$ orientation as the compressive stress accumulates. The expansion orientation indicates that the lithiation kinetics of Si along different orientations are different. Yang and co-authors further developed a kinetics model to reveal the lithiation orientation of crystalline Si.⁴⁷ According to the Li flux profile shown in Fig. 3d, the lithiation orientation dependence in the $\langle 110 \rangle$ orientation is more significant than in other directions. The anisotropic lithiation of Si is controlled by the orientation mobility of crystalline Si and amorphous Li_xSi . More in-depth study at the atomic scale would be beneficial to further reveal the orientation dependence of Si.

Liu and co-authors revealed the causes of orientation mobility by observing the dynamic lithiation interface of single crystalline Si at the atomic scale.⁴⁸ As shown in Fig. 3e, the Si atom tends to dissociate due to the weakened Si–Si covalent bonds under a high-rate lithiation process.⁵¹ They provided a grain boundary ledge flow fact to explain the lithiation kinetics properties of Si anodes. It can be observed that the ledge flow causes the phase transformation from crystalline Si to amorphous Li_xSi , and the flow rate represents the lithiation kinetics (Fig. 3f). To sum up, this dynamic atomic resolution research on the interface motion and phase transformation reveals the anisotropic lithiation mechanism of Si anodes. It also demonstrates that the high hydrostatic pressure under increasing particle depth will restrict the corresponding lithiation kinetics, indicating the advantages of shortening the Li^+ diffusion distance in improving the kinetic performance of Si anodes.^{52,53}

The lithiation of amorphous Si exhibits differences from crystalline Si.⁴⁵ McDowell and co-authors proposed two-phase lithiation of amorphous Si through *in situ* observation of the lithiation process.⁴⁹ As presented in Fig. 3g, an obvious boundary can be found between the Si-rich core in the dark area and the Li-rich shell in the light region, indicating a two-phase lithiation like crystalline Si. Although the amorphous Si lacks long-range ordered structures, as shown in Fig. 3h, its lithiation still involves the breaking of Si–Si bonds.⁵⁰ Specifically, the Li-rich clusters at the interface between two phases promote the dissociation of Si covalent bonds. In contrast, the amorphous Si presents less hydrostatic pressure and its increased lithiation thickness grows approximately linearly over time, illustrating an isotropic lithiation process (Fig. 3i).⁵⁴ The kinetics process of amorphous Si is still limited by diffusion distance, as its lithiation velocity is approximately constant.

Hence, the overlong Li^+ diffusion distance and intrinsic properties restrict the kinetic performance enhancement of Si anodes. It was reported that the capacity retention of micron-sized bulk Si anodes at a current density of 1C is below 20%,^{55,56} far below the set fast-charging target of 80% at 4C. Therefore, how to shorten the Li^+ diffusion distance and increase the intrinsic conductivity of Si materials at the particle level is significant to improve their kinetic performance, and detailed targeted strategies are introduced in this section.

2.1 Shorten the Li^+ diffusion distance

It is well known that reducing Si particle size is the most fundamental way to shorten the Li^+ diffusion distance.⁵⁷ In particular, Si particles with a mean diameter of 150 nm have attracted widespread attention due to an optimized “size effect”.^{58,59} Therefore, various nanostructures such as nanoparticles, nanowires, hollow nanotubes, and nanofilms/nanosheets have been developed.^{60–63} It was reported that Si nanoparticles with a particle size smaller than 100 nm, obtained by ball milling, delivered an approximately 30% capacity retention at 2.5 A g^{-1} , an increase compared to micron-sized bulk Si.⁶⁴ As shown in Fig. 4a, the Si nanowires were initially intended to arrange nanoarrays to buffer volume expansion.⁶⁵ Importantly, Liu and co-authors found a fast lithiation process of Si nanowires through *in situ* TEM.⁶⁶ The Si nanowires possess high lithiation rates of $60\text{--}200 \text{ nm s}^{-1}$ in liquid cells (Fig. 4b), which may benefit from the advantage of one-dimensional Li^+ transport. Due to the fast kinetics, Si nanowires exhibit good rate performance over 60% capacity retention at 1C.⁶⁵ Similarly, Si hollow nanotubes also demonstrate the advantage of one-dimensional rapid Li^+ diffusion and possess the capability of fast lithiation in the interior of the nanotubes.^{67,68} Park and co-authors reported a hollow Si nanotube with carbon coatings, delivering over 88% capacity retention at a current density of 5C.⁶⁸ For the two-dimensional Si materials, Si nanofilms/nanosheets present distinctive lithiation properties.⁶³ As shown in Fig. 4c, Li^+ can rapidly diffuse in a two-dimensional orientation parallel to the current collector, which facilitates rapid kinetics compared to Si nanoparticles. The representative Si nanostructures are summarized in Table 1. However, nanostructures imply more contact with electrolytes, which is accompanied by the



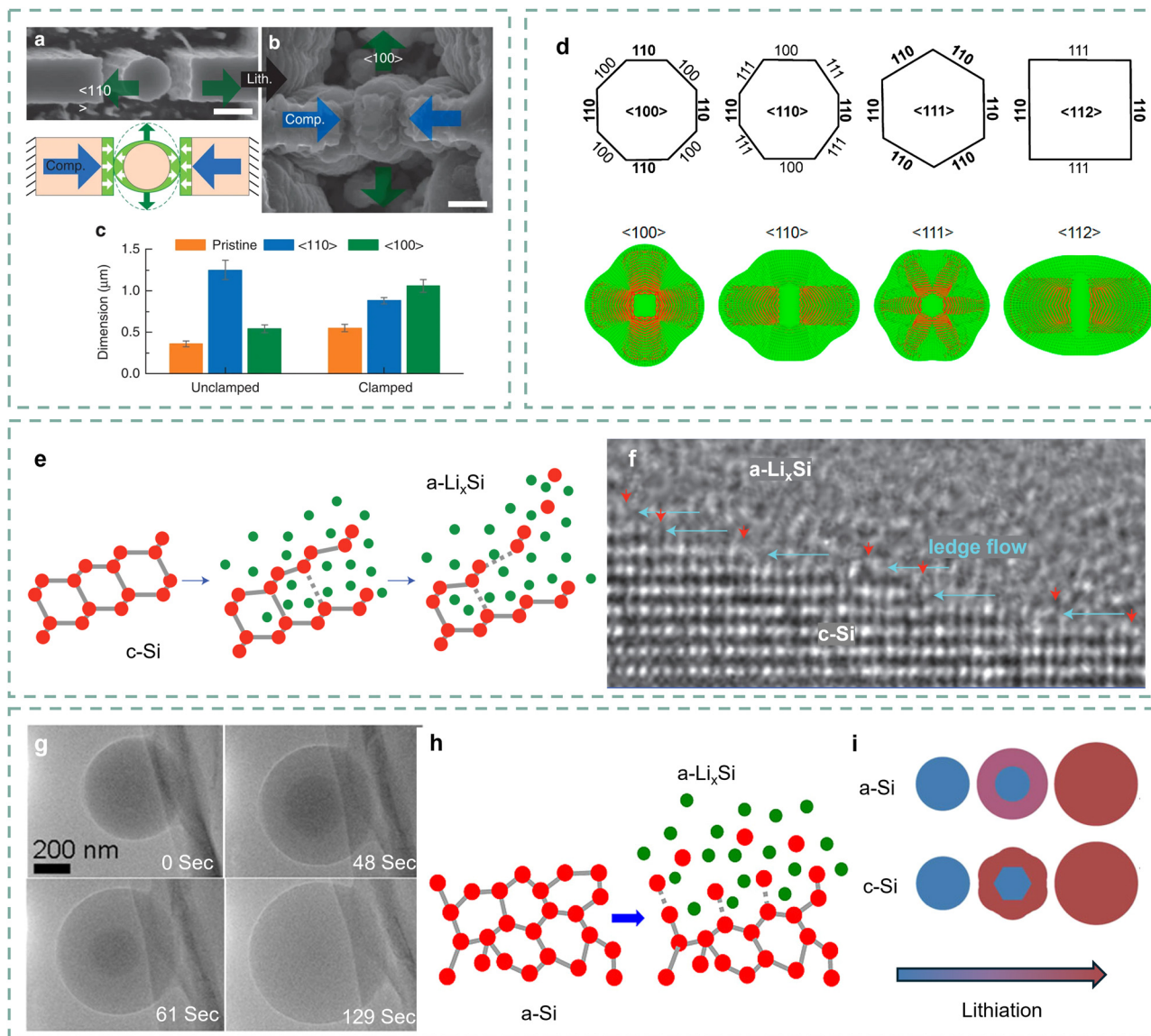


Fig. 3 (a) and (b) SEM images of the Si column positioned inside the fixture before (a) and after (b) lithiation, (c) dimension change of the Si column along the <110> (blue) and <100> (green) orientations after lithiation.⁴⁶ Copyright © 2015, The Author(s). (d) The characteristic crystallographic orientations of Si nanowires, and the corresponding Li flux profiles at a representative lithiation snapshot ($t = 0.4$).⁴⁷ Copyright © 2012, American Chemical Society. (e) Proposed amorphization mechanism of crystalline Si based on TEM observations, and (f) the ledge-flow dominated lithiation along the (111) plane. Red arrows indicate ledges and cyan arrows ledge-flow directions.⁴⁸ Copyright © 2012, Springer Nature Limited. (g) Time series of lithiation of the amorphous Si microsphere, and (i) schematic of the lithiation for amorphous Si and crystalline Si.⁴⁹ Copyright © 2013, American Chemical Society. (h) Mechanism of dissolution for Si atoms from amorphous Si to Li_xSi .⁵⁰ Copyright © 2013, American Chemical Society.

formation of more SEI, resulting in a potential increase in interface impedance.

Wu and co-authors explored the rate performance differences of Si particles with D_{50} ranging from 50 nm to 5 μm .⁷⁰ Si anodes exhibit gradually increased reversible capacity as their particle size decreases from 5 μm to 50 nm, suggesting that a short Li^+ diffusion distance can reduce the polarization of Si anodes, thereby increasing their reversible capacity. However, Si anodes with a particle diameter of 1 μm instead of 50 nm present the highest capacity retention at 5C. This confirms that too-small particle sizes promote the electrode-electrolyte contact and SEI formation, which increases the interface

impedance, thus obstructing the Li^+ migration before reaching particles. Fig. 4d reflects the fact that the kinetic performance of Si anodes has a “volcanic” relationship with particle size, which delivers the highest kinetic performance corresponding to appropriate particle size. The SEI chemical difference of Si anodes with nanometer and micron-sized particle sizes are described in further detail in Fig. 4e. It can be observed that the Si electrode with nanometer particle size forms thicker SEIs than the micron-sized Si electrode due to high surface area and electrode-electrolyte contact. It increases Li^+ migration difficulty and distance, as well as polarization during fast lithiation, resulting in poor rate performance. Therefore, balancing the



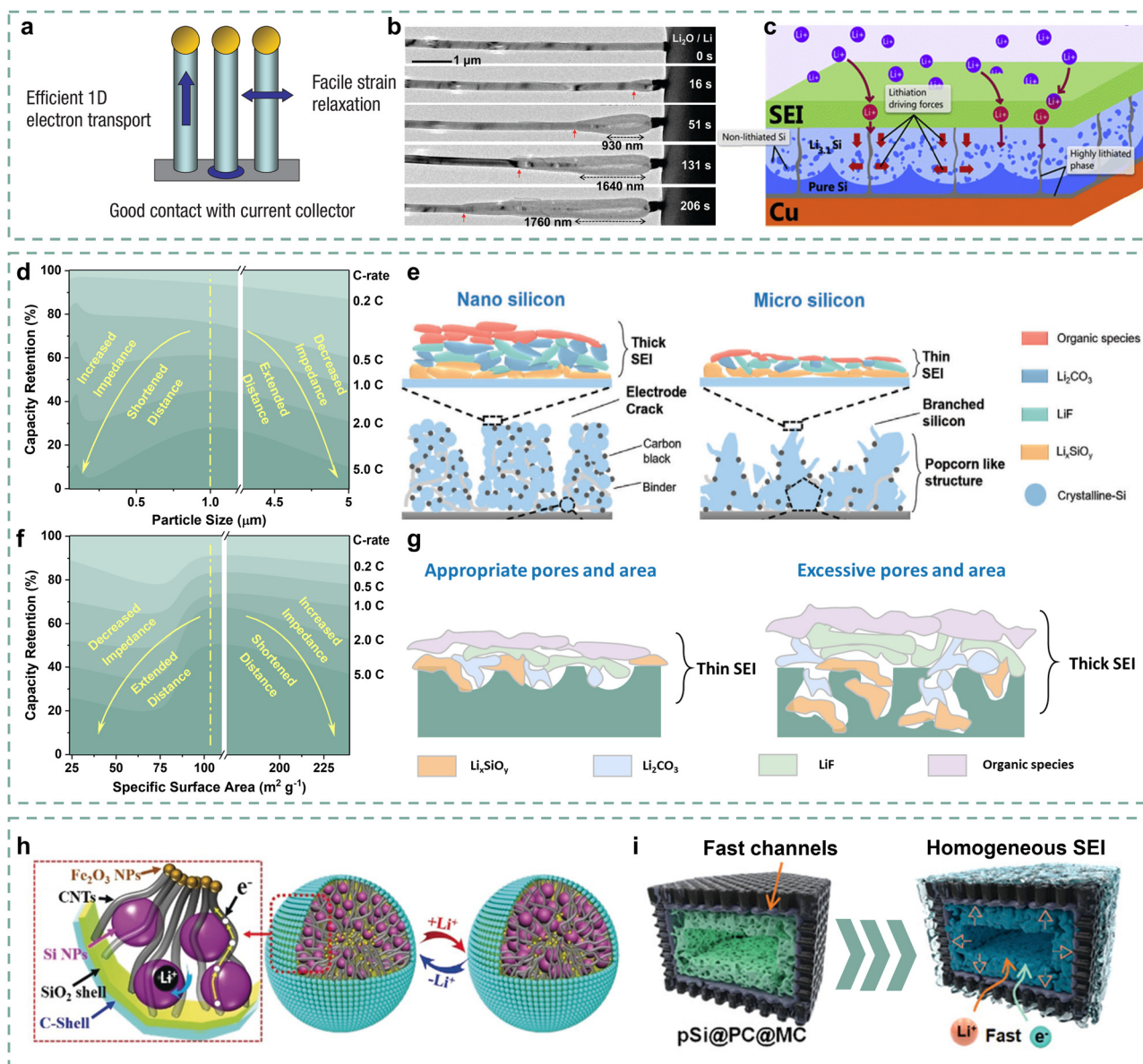


Fig. 4 (a) Schematic of lithiation for Si nanowires.⁶⁵ Copyright © 2007, Springer Nature Limited. (b) Morphology evolution of the Si nanowires during lithiation.⁶⁶ Copyright © 2011, American Chemical Society. (c) Schematic of the lithiation for Si nanofilms.⁶⁹ Copyright © 2013, Trans Tech Publications Ltd. (d) The relationship between capacity retention and particle size ranging from 50 nm to 5 μm at different rates and (e) an SEI chemical schematic of nano Si and micron-sized Si anodes.⁷⁰ Copyright © 2023 Wiley-VCH GmbH. (f) The relationship between capacity retention and surface area at different rates.⁷¹ Copyright © 2016 American Chemical Society. (g) SEI chemical schematic of porous Si with different pores and areas. (h) Structural schematic of yolk-shell structured Si/C anodes.⁷² Copyright © 2019 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim. (i) Schematic of the structural properties of pSi@PC@MC anodes.⁷³ Copyright © 2024 Wiley-VCH GmbH.

particle size and the interface impedance of the Si anodes is crucial for improving its kinetic performance.

From this perspective, nano-micro-particle design is constantly emerging.^{72,79–81} Hou and co-authors proposed the 3D plum-pudding-like Si/C anode to balance particle size and interface impedance, ensuring the good fast-charging performance of Si/C electrodes.⁷⁹ The Si/C material formed by the secondary stacking of nano Si particles delivers a particle size of 2–5 μm . Importantly, the micron-sized Si/C anode delivers a high initial Coulombic efficiency (ICE) of 88%, significantly reducing side reactions such as SEI formation. Hence, this

unique structure can shorten the Li^+ diffusion distance while reducing SEI formation and interface impedance, thereby accelerating kinetic performance. Furthermore, more functional networks have also been introduced into the constructed micron-sized Si composite particles. Zhang and co-authors developed a yolk-shell Si anode (YS-Si/C) with secondary stacking of Si, CNTs, and Fe_2O_3 particles and double-layer encapsulation.⁷² As shown in Fig. 4h, the introduced CNTs and Fe_2O_3 can serve as the conductive highway among Si particles as well as bridging particles and the shell. Benefiting from the structural design, the YS-Si/C anode delivers a high



Table 1 Typical Si nanostructures for improving the kinetics of Si anodes

Sample	Nanostructure	ICE (%)	Rate performance	Ref.
BM-Si	Nanoparticle	74	~ 30% at 2.5 A g ⁻¹	64
Si NP	Nanoparticle	~ 45	~ 35% at 1 A g ⁻¹	74
Si NW	Nanowire	73	~ 60% at 1C	65
Si NWs	Nanowire	69.5	~ 72% at 0.5C	75
Si nanotubes	Nanotube	89	~ 88% at 5C	68
DWSiNT	Nanotube	76	~ 34% at 12C	67
Si nanofilm	Nanofilm	85.1	79.7% at 1C	76
a-C/Si thin film	Nanofilm	80	~ 70% at 3.57C	77
C-SiNS	Nanosheet	79.4	~ 50% at 2C	78

ICE of 86.9% and 67% capacity retention at a current density of 2 A g⁻¹, indicating that the appropriate balance between diffusion distance and interface impedance has been obtained through this structure.

Porous Si can also greatly shorten the Li⁺ diffusion distance, thereby ensuring the Si anodes with more active sites and high fast-charging performance.⁸² Liang and co-authors designed porous Si anodes with different specific surface areas and pore size distributions to investigate the porosity effect on the corresponding Si anode fast-charging performance.⁷¹ The synthetic A-Si, C-Si, S-Si, and D-Si materials deliver specific surface areas of 239.3, 23.9, 102.6, and 74.13 m² g⁻¹ due to their different pore structures, which results in different Li⁺ diffusion distances. As described in Fig. 4f, it can be observed that the S-Si electrode with mean surface area (102.6 m² g⁻¹) and appropriate pore size delivers the highest capacity retention compared with other Si anodes, indicating that simply increasing the pore structure/surface area of Si materials cannot be beneficial to improving kinetic performance. As shown in Fig. 4g, the porous structure also tends to promote similarly high interface contact between electrodes and electrolytes, leading to high interface impedance. It can be inferred that increasing active sites and shortening the Li⁺ diffusion distance of Si materials by optimizing their pore structure, and meanwhile inhibiting the formation of thicker SEIs, contribute to the kinetic performance enhancement of Si anodes.

Furthermore, numerous studies have been conducted to balance the Li⁺ diffusion distance and interface impedance for porous structures. Coating strategies can effectively inhibit the excessive formation of SEIs and the increase in interface resistance.^{73,83,84} For example, Ren and co-authors encapsulated porous Si with dense carbon layers, promoting a significant decrease of specific surface area from 350 m² g⁻¹ to 10 m² g⁻¹ and an increase of ICE from 68% to 78%.⁸⁵ Reducing the specific surface area helps to minimize the interface contact with electrolytes, thereby suppressing the formation of SEIs. And the porous Si/C anode also exhibits lower charge transfer impedance than that of porous Si, which further confirms the inhibition of coatings on thicker SEI formation, thus contributing to the kinetics enhancement. Nevertheless, the dense coating can also cover the pore structures, which weakens the active sites. It is important to perform effective strategies to control the appropriate specific surface area and pore structures. Cheng and co-authors proposed a carbon-coated porous Si anode (pSi@PC@MC) to promote appropriate pore structures for kinetics enhancement of typical

porous Si anodes.⁷³ The pitch-derived carbon (PC) coating is relatively dense, which helps to avoid direct contact between porous Si and electrolytes. Meanwhile, the dopamine-derived mesoporous carbon (MC) introduces mesopores externally to construct more active sites, thereby promoting the widespread diffusion of Li⁺. As shown in Fig. 4i, the double coatings of PC and MC largely avoid the contact between porous Si and electrolytes, based on the decrease in specific surface area (222.8 m² g⁻¹ to 164.1 m² g⁻¹) and the leap of ICE value (65.92% to 75.6%), which contributes to the thin and homogeneous SEI. Meanwhile, the external mesoporous structures promote electrolyte infiltration and the introduction of more active sites, further providing good kinetic performance. The pSi@PC@MC anode delivers better rate performances of 57.9% and 53.6% capacity retention at current densities of 4 and 6 A g⁻¹ compared with the porous Si anode. Besides, comparison of the pseudocapacitance of porous Si, pSi@MC, and pSi@PC@MC anodes is further persuasive evidence to further demonstrate the importance of controlling appropriate surface areas and pore structures in enhancing the kinetic performance of porous Si anodes.

Therefore, regulating the structure of Si materials through optimizing the particle size of Si particles and pore size distribution of porous Si to balance the Li⁺ diffusion distance and the interface impedance is necessary to inhibit the excessive formation of SEIs and the increase of interface impedance, thereby significantly improving the kinetic performance of Si anodes.

2.2 Improving the conductivity

The charging process involves Li⁺ diffusion within the Si anodes and synchronous electron transfer. The imbalance between Li⁺ diffusion and electron transfer can lead to polarization, which induces over potential and low lithiation depth, resulting in poor rate performance.^{86–88} Therefore, fast charging requires rapid Li⁺ diffusion inside Si anodes and the matching electron transfer.^{89–91} However, the intrinsic properties of Si materials severely hinder the Li⁺ diffusion and electron transfer within particles, manifesting as locally slow kinetics at the particle level. Therefore, it is necessary to improve the electronic and ionic conductivity of Si materials and balance them.

Element doping is a typical and effective strategy to improve the electronic conductivity of Si materials by introducing charge carriers (electrons or holes). For typical III and V groups (boron, phosphorus), as shown in Fig. 5a, the electronic conductivity enhancement mechanism can be simply described as elements incorporated into the lattice structure of Si materials creating p-type Si with three valence electrons or n-type Si with five valence electrons.⁹² For other potential elements, they can generate multiple ionizations and introduce multiple energy levels, accompanied by the introduction of non-fixed forms of carriers to enhance electronic conductivity.

Metal alloying and doping are conventional strategies for improving the conductivity of Si materials.^{103,104} Nulu and co-authors developed Mn doping SiMn05% and Ni doping SiNi05% anode materials to address the poor conductivity of Si materials.⁹³



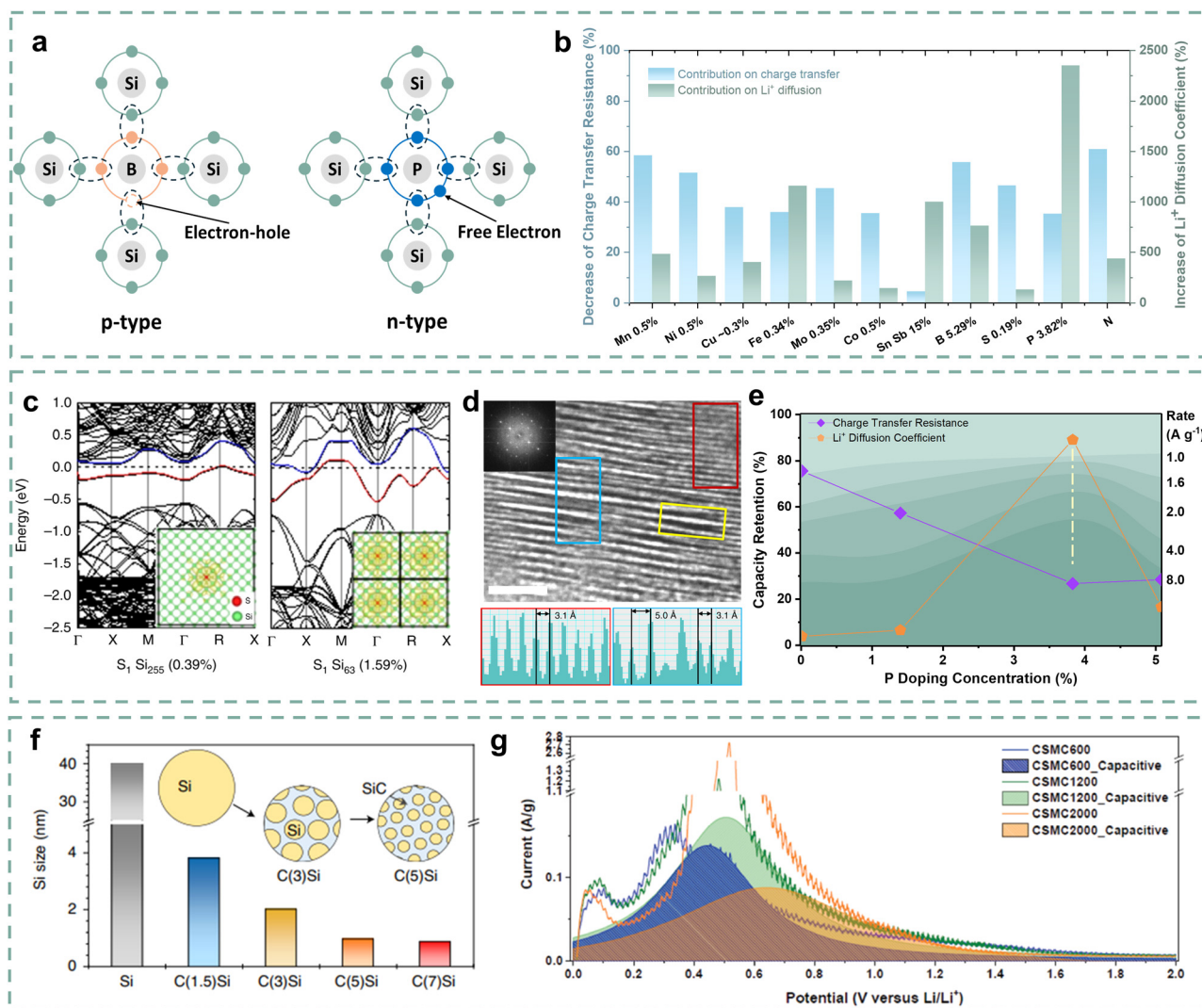


Fig. 5 (a) Schematic of the conductivity enhancement mechanism. (b) Doping effects on charge transfer and Li^+ diffusion.^{93–99} (c) Calculated band structures with the S doping in Si. (d) HR-TEM images of S doping Si and the intensity profiles of the selected areas.¹⁰⁰ Copyright © 2019, The Author(s). (e) The relationship between capacity retention and P doping concentration at different rates.⁹⁹ (f) Crystallite size of Si, C(1.5)Si, C(3)Si, C(5)Si, and C(7)Si anodes.¹⁰¹ Copyright © 2021, The Author(s), under exclusive license to Springer Nature Limited. (g) Capacitive contribution of CSMC600, CSMC1200 and CSMC2000 anodes.¹⁰² Copyright © 2019 WILEY-VCH Verlag GmbH & Co. KGaA, Weinheim.

The SiMn05% and SiNi05% anodes deliver higher Li^+ diffusion coefficients and lower charge transfer impedance than the compared Si anode. In particular, the SiMn05% anode delivers a good rate performance with 50% capacity retention at a current density of 3.2 A g^{-1} due to its good conductivity. Luo and co-authors studied the effect of molybdenum (Mo) doping on the conductivity of Si anodes.⁹⁵ First-principles calculations show that Mo doping reduces the band gap of Si from 0.61 eV to 0.45 eV and increases the nearby charge density, which means an enhancement of conductivity. Approximately 70% decrease in impedance and an 8.8 times increase in Li^+ diffusion coefficient promote the Mo-doped porous nanostructured Si (Mo-PNSi) anode to exhibit an enhanced rate performance of 38.6% capacity retention at a current density of 6.72 A g^{-1} . Various transition metal doping strategies such as Cu, Fe, and Co have also been performed to enhance the conductivity of Si anodes.^{94,96} Additionally, some

electrochemically active metals are also used as dopants to enhance the kinetics of Si anodes while ensuring high reversible capacity. Gao and co-authors proposed to introduce Sn and Sb to improve the electronic conductivity and Li^+ diffusion of Si anodes.¹⁰⁵ The as-prepared $\text{Si}_{8.5}\text{Sn}_{0.5}\text{Sb}$ anode presents 6000 times the electronic conductivity and 10 times the Li^+ diffusion coefficient higher than those of Si particles, respectively. Due to the introduction of electrochemically active Sn and Sb, the composite anode has a high reversible capacity of 1190 mA h g^{-1} at a current density of 5 A g^{-1} . Fig. 5b presents the contributions of element doping on decreasing charge transfer resistance (R_{CT}) and improving the Li^+ diffusion coefficient (D_{Li^+}), and detailed information on several element doping strategies is displayed in Table 2. It can be inferred that element doping can reduce the impedance and enhance the Li^+ diffusion ability, thereby optimizing the kinetic performance. However, some transition metal ions are easily

Table 2 Typical strategies of element doping for improving the kinetics of Si anodes

Element	Concentration (%)	R_{CT} (Ohm)	Original R_{CT} (Ohm)	D_{Li^+} ($\text{cm}^2 \text{s}^{-1}$)	Original D_{Li^+} ($\text{cm}^2 \text{s}^{-1}$)	Ref.
Mn	0.5	59	101	1.90×10^{-13}	3.94×10^{-14}	93
Ni	0.5	52	101	1.04×10^{-13}	3.94×10^{-14}	94
Cu	~0.3	40.3	106.5	1.41×10^{-11}	3.53×10^{-12}	
Fe	0.34	38.2	106.5	4.09×10^{-11}	3.53×10^{-12}	
Mn	0.32	40.7	106.5	1.07×10^{-11}	3.53×10^{-12}	
Mo	0.35	115	253.7	2.088×10^{-21}	9.488×10^{-22}	95
	0.69	78.62	253.7	8.394×10^{-21}	9.488×10^{-22}	96
Co	0.1	8.8	15.8	4.08×10^{-11}	6.92×10^{-11}	
	0.3	6.5	15.8	7.61×10^{-11}	6.92×10^{-11}	
	0.5	5.6	15.8	9.99×10^{-11}	6.92×10^{-11}	
Sn, Sb	5, 10	18.9	424.6	10 times increase		105
B	5.29	7.42	13.32	8.24×10^{-12}	1.08×10^{-12}	97
S	0.19	69.93	150.5	1.2×10^{-8}	9.26×10^{-9}	98
P	1.39	114.4	150.9	9.637×10^{-14}	5.68×10^{-14}	99
	3.82	53.1	150.9	1.335×10^{-12}	5.68×10^{-14}	107
	5.07	56.6	150.9	2.48×10^{-13}	5.68×10^{-14}	
N	—	142.7	234.3	2.82×10^{-12}	1.22×10^{-12}	

reduced to metal atoms or dendrites, which tend to cause short circuits.¹⁰⁶ Thus, the metal doping strategy of Si materials still needs to be screened and identified for commercial application.

Non-metal doping (P, S, B, N, *etc.*) for Si materials has attracted wide attention due to its potential advantages. Since the S atom generally provides two or more electrons as charge carriers,¹⁰⁸ Guo and co-authors designed a high-conductivity S-doped porous Si/SiO₂ anode for good fast-charging performance.⁹⁸ After S doping, the S-doped p-Si/SiO₂ anode delivers higher reversible capacity with lower loss and rate performance than the p-Si/SiO₂ anode. It can be inferred that the capacity and kinetic performance enhancement of the S-doped p-Si/SiO₂ anode is largely attributed to the reduced polarization due to high conductivity. Importantly, non-metal doping may alter the crystalline structure of Si, which further affects the Li⁺ diffusion within the Si particles. Ryu and co-authors proposed that the metallicity of Si will change and transit towards a quasi-metallic state, as the S doping exceeds the equilibrium solid solubility.¹⁰⁰ As shown in Fig. 5c, the spatial states of the remaining valence electrons of Si gradually approach and overlap with the increase of S doping, which enhances the band dispersion of the state at the Fermi level and forms the metallic bands. Furthermore, expanded channels (0.5–0.72 nm) have been formed, as the introduction of S-chains expanded the interlayer spacing (Fig. 5d). During the lithiation process, Li₂S structures are formed to support the expanded channels, resulting in a lower energy barrier for Li⁺ diffusion (0.32 eV). Similarly, B, P, or N element doping can also provide additional electrons or holes and impact the crystalline structure. These dopants may cause a decrease in crystal spacing due to the smaller atomic size and affect the Li⁺ diffusion inside the Si particles.^{109–111} For example, Wang and co-authors stated that the introduction of P can lead to a decrease in the Li⁺ diffusion barrier from 0.57 eV to 0.53 eV.¹¹² Je and co-authors discovered a mixed amorphous–crystalline structure with localized atomic distortion in the synthesized B-doped Si (mixed amorphous–crystalline Si, MACS) anodes.¹¹³ The localized structural distortion is mainly attributed to the heteroatom-bridge bonds (B–O–Si, B–Si) in crystalline Si, which cause the amorphous phase to wedge into the crystalline phase and reduce the crystalline size, thereby

accelerating the Li⁺ diffusion in the MACS anode. The MACS anode affords ~36% capacity retention at a current density of 5C due to the effective fast Li⁺ diffusion. Furthermore, the appropriate doping concentration is important for the high kinetic performance of Si anodes.¹¹⁴ Long and co-authors developed P doping P_xSi anode materials through a hydrolysis strategy and further adjusted the P doping concentration through the formation of P–O–Si bonds and the interaction of H–OH bonds.⁹⁹ As described in Fig. 5e, as the concentration increases, the P_xSi anodes present gradually improving rate performance. And the P_{0.125}Si anode delivers the optimal R_{CT} and D_{Li^+} , and the highest rate performance when the P doping concentration increases to 3.82% instead of 5.07%, resulting in a “volcanic” relationship with doping concentration. The P_{0.125}Si anode obtains higher capacity retention of 50% than other P_xSi anodes at a current density of 8 A g^{−1}, indicating that the appropriate P doping concentration significantly improves the kinetic performance of the Si anodes.

In addition to element doping, Si composite materials with particle compositing also have obvious structural advantages in enhancing their kinetic performance. Carbon particles have great advantages due to their high conductivity and strong mechanical structure, which mainly contains graphite,¹¹⁵ hard carbon,¹¹⁶ carbon nanotubes,¹¹⁷ and graphene.¹¹⁸ Kim and co-authors designed a Si–graphite composite anode using edge-activated graphite as the substrate.¹¹⁹ They utilized Ni to penetrate the graphite core, forming activated edges and creating numerous active sites. Si is covered on the surface of graphite, which fully fits with the shape of edge-activated graphite, thereby promoting rapid kinetics through these active sites. The Si–graphite composite anode delivers an enhanced rate performance of 20% capacity retention at a current density of 10.5 mA cm^{−2} (3C). Sung and co-authors further designed a novel C(x)Si anode by embedding nano Si particles into a SiC–C composite matrix (Fig. 5f), which obtains a high capacity retention of 79.6% at a current density of 5C due to good Li⁺/e[−] conductivity and short Li⁺ diffusion distance.¹⁰¹ Besides, adjusting the composite structure is conducive to improving the pseudocapacitive properties of Si anodes, thereby facilitating its kinetic performance. Son and co-authors



developed a nanocage-shaped Si/C anode (CSMC) by embedding nano Si particles into a 3D carbon matrix.¹⁰² Fig. 5g describes the pseudocapacitive contribution to the kinetic performance of these CSMC materials. Benefiting from the synergistic promotion effect of the carbon matrix and Si/C particles, the pseudocapacitive contribution can even reach 88%, which enhances the kinetic performance. As a result, the CSMC600 anode delivers a high capacity retention of 80% at a current density of 20C, which is far superior to the other two compared anodes.

In summary, it can be inferred that both element doping and particle compositing can improve the kinetic performance of Si anodes since they promote good conductivity and Li⁺ diffusion. The best kinetic performance enhancement can be obtained by regulating the appropriate doping or compositing concentration. Additionally, element doping can bring potential crystal structure changes, which may further improve the Li⁺ diffusion inside the Si particles. Particle compositing can promote the construction of structure with pseudocapacitive contributions, which is another important factor for developing high kinetics Si anodes at the particle level.

3. Interface level

The interface impedance is the main factor limiting the kinetic performance enhancement of Si anodes, which is also largely attributed to the surface chemical properties of Si anodes, containing material exposed surface properties and the formation of SEIs.¹²⁰ Bare Si materials generally possess high interface impedance due to poor surface conductivity. SEIs usually contain inorganic layers and organic layers with poor Li⁺ conductivity, and Li⁺ migrates cross SEI after desolvation to carry out the lithiation and de-lithiation of Si anodes.^{121,122} Thus, the chemicals, chemical distribution, and thickness of SEIs affect the Li⁺ migration and kinetic performance of the Si anodes.^{123–126} Therefore, how to optimize the surface properties and SEI structure of Si anodes for good kinetic performance is discussed in this section.

3.1 Surface coating strategy

Chen and co-authors demonstrated that Si materials naturally have a thin SiO_x coating due to the relatively high chemical activity of surface Si atoms, and both bare Si surface and thin SiO_x coating have poor Li⁺/e[−] conductivity that restricts the corresponding material kinetic performance.¹²⁷ Several researchers have attempted to adjust the SiO_x coating thickness to reduce the interface impedance.¹²⁸ However, it is a difficult technology with poor effects. Surface coatings have attracted wide attention due to their good Li⁺/e[−] conductivity in kinetic performance improvement,¹²⁹ which mainly contains carbon coatings¹³⁰ and elastic organic coatings.¹²⁷

A Si anode with carbon coating possesses good kinetic performance due to promising Li⁺/e[−] conductivity and strong structural stability.^{115,131} Xu and co-authors demonstrated the carbon coating contribution of Si anodes on kinetic performance improvement through the *in situ* TEM method.¹³²

Fig. 6a describes the structural differences and dynamic lithiation process of uncoated Si and Si/C anodes, in which the Si/C particle takes only 40 s for full lithiation while the uncoated Si particle takes longer (200 s). This illustrates that carbon coating can accelerate the kinetic performance of Si anodes since the Li⁺ diffusion of carbon coatings is far superior to Si materials ($\sim 10^{-7}$ vs. $\sim 10^{-12}$ cm² s^{−1}),^{11,133} which can be reconfirmed by the lithiation product thickness comparison between uncoated Si and Si/C particles (Fig. 6b). Numerous studies further focus on the relationship between the carbon coating thickness and the kinetic performance of Si anodes. As described in Fig. 6c, Qi and co-authors investigated the electronic conductivity and rate performance evolution of the Si@C anode with coating thickness ranging from 0 to 9.73 nm.¹³⁴ The electronic conductivity and capacity retention at high rates exhibit a similar trend, which is a “volcanic” relationship as the coating thickness increases. It can be inferred that the thickness and the accompanied defects of carbon coating can influence the surface area and electronic conductivity of Si@C anodes, thereby affecting their rate performance. The Si@C anode with appropriate carbon coating thickness has an enhanced rate performance due to the good Li⁺/e[−] conductivity of the carbon coating, while its rate performance can also be significantly inhibited as the over-thick carbon coating prolongs the Li⁺ diffusion. Since thinner coating leads to breakage while thicker coating increases Li⁺ diffusion distance, therefore, an appropriate coating thickness is of great significance in enhancing the kinetic performance of Si anodes at the interface level.

Furthermore, the structural design of carbon coatings can further improve the kinetics of Si anodes. Wang and co-authors proposed an interconnected carbon coating that encapsulated Si nanowires into a conductive network to facilitate electron transfer.¹³⁵ As shown in Fig. 6d, the interconnected carbon-coated nanowires accelerate the electron transfer, as well as the space among the nanowires and the pores inside the nanowires provide channels for Li⁺ diffusion. Conformal coating of thin carbon layers reduces the resistance in the system and ensures a thin SEI outside Si nanostructures. Benefiting from the rapid Li⁺ diffusion and the matched external electron transfer, the N-PSi@C anode (“N” stands for “network”, “P” stands for “porous”) exhibits an optimized rate performance of $\sim 48\%$ capacity retention at a current density of 7C. Son and co-authors reported a design to enhance the kinetics of Si anodes by directly growing graphene on Si as a conductive coating.¹³⁶ As shown in Fig. 6e, the layered stacking of 2D graphene on the surface of Si particles promotes an interconnected percolation network, which increases the conductivity of the Gr-Si materials to 12.8 S cm^{−1}. The Gr-Si composite anode can deliver over 90% capacity retention at a high rate of 10C. Therefore, effective structural design of carbon coatings is beneficial for Li⁺ diffusion into the Si particles, thereby promoting great kinetics of Si anodes.

Since the huge volume expansion and irreversible structural damage tend to interrupt the Li⁺ diffusion path of Si anodes, elastic organic coatings have received widespread attention in enhancing the kinetic performance at the interface level.¹⁴⁸



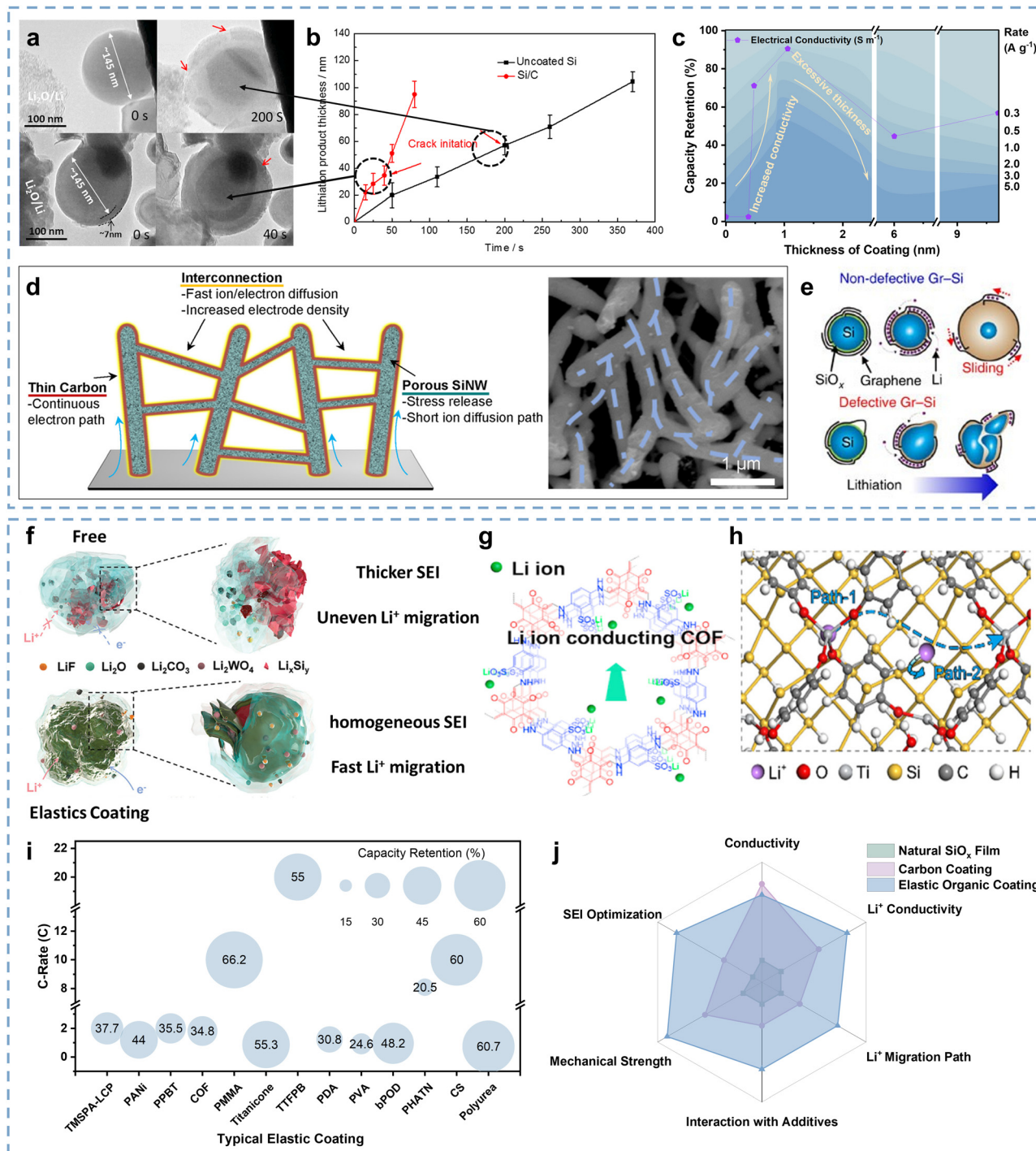


Fig. 6 (a) Dynamic lithiation process of uncoated Si and Si/C particles through *in situ* TEM. (b) Lithiation product thickness comparison of uncoated Si and Si/C anodes.¹³² Copyright © 2016 Elsevier B.V. All rights reserved. (c) The relationship between capacity retention and coating thickness at different rates.¹³⁴ Copyright © 2021 Elsevier Ltd. All rights reserved. (d) Structural schematic of the N-PSi@C anode and the corresponding SEM images.¹³⁵ Copyright © 2019, American Chemical Society. (e) Schematic of the lithiation of Gr-Si NPs with different graphene encapsulation.¹³⁶ Copyright © 2015, The Author(s). (f) Schematic of the evolution of interfacial properties and the SEI structure for Si and Si@LCP coating.¹³⁷ Copyright © 2022 Wiley-VCH GmbH. (g) Structural schematic of the COF coating.¹³⁸ Copyright © 2020 Published by Elsevier Ltd. (h) Schematic of the Li⁺ horizontal and vertical diffusion path for Si NPs@titanicore anodes.¹³⁹ Copyright © 2021 American Chemical Society. (i) Comparison of the rate performance for Si anodes with typical organic coatings.^{137–147} (j) Normalization for the functionality of natural film, carbon coating, and elastic organic coating.

Most elastic organic coatings with special chemical groups (C–O and C–F) present good conductivity and are beneficial for structural stability, which is conducive to accelerating the

corresponding Li⁺ migration at the interface.¹²⁷ For example, some electron-conductive coatings have been developed to facilitate the electron transfer near Si particles. Pan and

co-authors proposed a layered conductive polyaniline (LCP) coating for Si anodes.¹³⁷ The polyaniline chains in the LCP coating allow for electron transfer, resulting in 5 orders of magnitude higher electronic conductivity of the coated Si materials than pure Si, serving as a good electron-conductive coating. And the N-containing in LCP coatings enables dipole-dipole interaction with fluoroethylene carbonate (FEC), which induces uniform LiF formation, thereby forming a homogeneous SEI above the LCP coatings. As shown in Fig. 6f, the designed elastic LCP coating promotes the homogeneous SEI formation and avoids the fragmentation and reconstruction of the SEI, resulting in a uniform and thin SEI that is conducive to Li⁺ migration. Moreover, the high electronic conductivity of the LCP coating facilitates rapid local electron transfer, which can also enhance the kinetics near the interface. In contrast, the SEI distribution on the pure Si is uneven, and it fractures and reconstructs with volume expansion, resulting in thicker SEI and uneven Li⁺ migration. Benefiting from the enhanced electronic conductivity and uniform SEI, the composite anode delivers a good rate performance of ~46% at a current density of 5 A g⁻¹. Wu and co-authors formed a 3D interconnected electron-conductive coating for Si particles by reacting phytic acid with aniline monomers of polyaniline (PANI).¹⁴⁹ The PANi coating accelerates the local electron transfer of Si particles and simultaneously extends the electron transfer to the current collector through the interconnection. Similarly, some electron-conductive organic coatings such as poly[3-(potassium-4-butoxate)thiophene] (PPBT) and biphenyl-polyoxadiazole (bPOD) have also been reported to increase the local electronic conductivity to enhance the kinetics of Si anodes.^{144,150}

Moreover, organic coatings with high Li⁺ conductivity are also conducive to enhancing the kinetics of Si anodes. Ai and co-authors designed a novel Si@COF anode, of which the covalent-organic-framework (COF) coating has good Li⁺ conductivity (Fig. 6g).¹³⁸ It can be observed that the Li⁺ diffusion coefficient of the Si@COF anode is higher than the pure Si anode during lithiation. Additionally, the Si@COF anode delivers better rate performance than the Si anode, which is largely attributed to the high Li⁺ conductivity, strain, and strength synergy of the COF coating. A thin titanicon coating is another

promising choice for Li⁺ conductivity enhancement at the interface of the Si anodes. Fang and co-authors developed a novel SiNPs@titanicon anode with good fast-charging performance and strong mechanical strength.¹³⁹ The SiNPs@titanicon-70 anode obtains a high retention capacity of 957 mA h g⁻¹ after 450 cycles due to the low Young's modulus of the titanicon coating, which can release stress and buffer the volume expansion of the Si anodes during lithiation. It can also be inferred that the titanicon coating is beneficial for reducing the SEI thickness of the Si anodes. In contrast, a pure Si anode with thick SEIs will increase the Li⁺ migration impedance. Benefiting from the long-strand (Ti-O-benzene-O-Ti) group of the titanicon coating, the SiNPs@titanicon anode presents a lower energy barrier, good Li⁺ conductivity, and more efficient Li⁺ diffusion path, thereby facilitating the kinetic performance (Fig. 6h). Besides, the titanicon coating is also beneficial for homogenizing the stress distribution of the SiNPs@titanicon anode during lithiation. Benefiting from the lower Li⁺ migration impedance, the SiNPs@titanicon-70 anode delivers a higher capacity retention (50%) than the other two compared anodes at the current density of 2 A g⁻¹.

Furthermore, the function of coatings has been extended to optimize SEI, thereby further optimizing the surface properties of Si anodes to improve the kinetic performance at the interface. Cao and co-authors developed a chain-like framework (poly-4-trifluoromethylphenylboronic acid, PTFPBA) to form B/F-rich SEIs by polymerizing 2,4,6-tris-4-(trifluoromethylphenyl)boroxine (TTFPB) coating onto the surface of Si anodes.¹⁴¹ The PTFPBA coating exhibits a higher adsorption energy with EC, EMC, DMC, and FEC and lower adsorption energy with LiPF₆ than that of bare Si, indicating that the PTFPBA coating can inhibit the growth of SEI and take on several important SEI roles. Benefiting from good Li⁺/e⁻ conductivity and a high Young's modulus, the Si@TTFPB anode exhibits a higher Li⁺ diffusion coefficient and better rate performance than the compared Si anode.

As shown in Table 3, elastic coatings typically exhibit one or several functions including the contributions to electronic conductivity, Li⁺ migration, mechanical strength, and SEI optimization. According to the quantization of the contributions for typical elastic coatings and the comparison of rate performances

Table 3 The functions of typical elastic organic coatings and rate performances for the corresponding Si-based anode

Organic coating	Function				Rate performance	Ref.
	Electronic conductivity	Li ⁺ migration	Mechanical strength	Optimize SEI		
TMSPA-LCP	✓	✓	✓		5 A g ⁻¹ , ~46%	137
PANI	✓		✓		3 A g ⁻¹ , ~44%	149
PPBT	✓		✓		2C, 35.5%	150
COF		✓			5 A g ⁻¹ , 34.8%	138
PMMA		✓	✓		10C, 66.2%	140
Titanicon		✓	✓		2 A g ⁻¹ , 55.3%	139
TTFPB			✓	✓	20C, ~55%	141
PDA			✓		4 A g ⁻¹ , 30.8%	142
PVA			✓		3 A g ⁻¹ , 24.6%	143
bPOD	✓	✓	✓		3 A g ⁻¹ , 48.2%	144
PHATN			✓	✓	16.5 A g ⁻¹ , 20.5%	145
CS		✓		✓	10C, ~60%	146
Polyurea		✓		✓	2 A g ⁻¹ , ~60.7%	147



(Fig. 6i), coatings with higher Li^+ conductivity or inducing SEI affinity to Li^+ can exhibit better rate performance, illustrating that both coatings' properties and their impact on SEI should be considered in coating construction to enhance the kinetic performance. As shown in Fig. 6j, elastic organic coatings have a deeper potential to improve the kinetic performance of Si anodes since they possess more comprehensive functions.

As a result, the surface coating strategy of Si anodes can optimize their surface properties by improving conductivity, facilitating Li^+ migration, increasing mechanical strength, and adjusting the distributions or even components of SEIs. Carbon coating can generally improve the conductivity, and promote Li^+ migration through potential structural designs. While organic elastic coatings can obtain more functionality through potential functionalization, and better matching with Si particle interfaces and electrolyte interfaces, thereby promoting the kinetics enhancement at the interface more effectively.

3.2 SEI optimization strategy

The chemical distribution of SEIs can affect the Li^+ migration barriers. Adopting appropriate SEI optimization strategies is conducive to obtaining superior Li^+ migration ability and enhanced kinetic performance of Si anodes from another perspective at the interface level. Typical electrolytes consist of Li salts, solvents, and additives. Many studies have been focused on SEI optimization through the adjustment of electrolyte engineering including Li^+ salts, solvents, and additives, thus ensuring the good kinetic performance of Si anodes.^{151–156}

Li salt is the Li^+ migration carrier in electrolytes, and LiPF_6 is a commercial Li salt that easily tends to dissolve in carbonate solvent.¹⁵⁷ However, LiPF_6 presents a potential corrosion challenge for Si anodes due to its decomposition and HF byproduct.^{155,158} Lu and co-authors explored the kinetic performance difference among six Li salts on the Si anodes, including LiPF_6 , LiFSI, LiTFSI, LiBOB, LiDFOB, and LiClO_4 .¹⁵⁵ As shown in Fig. 7a, LiFSI delivers the highest Li^+ conductivity of 11.3 mS cm^{-1} among these individual and composite salts, which is conducive to Li^+ migration during desolvation. As a result, the Si anodes based on LiFSI obtain excellent rate performance, including over 70% capacity retention at 5C and over 45% capacity retention at 30C. Besides, the synergy of LiFSI and LiPF_6 can also facilitate kinetic performance, which delivers significant improvements in rate performance from 5C to 30C. It is better than Si anodes with pure LiFSI, even though the Li^+ conductivity is lower than LiFSI, suggesting potential contributions to SEI optimization. Fig. 7b compares the element distribution around Si particles based on LiFSI and LiPF_6 . For the Si–LiFSI electrode, Si particles are uniformly surrounded by Li and present uniform Li–O dominated inorganic layers. Since the FSI[−] possesses good hydrolytic stability and avoids HF formation,¹⁵⁹ Li_2O in SEI can be preserved while some LiOH phase may be formed.^{160,161} The Li–O dominant layer is mainly composed of Li_2O , Li_2CO_3 , and a small amount of LiF, as well as possibly LiOH, with some Li_2CO_3 attributed to brief exposure to air. In comparison, in the Si– LiPF_6 electrode, the possible oxygen-containing inorganics (Li_2O , Li_2CO_3 , etc.)

around the Si particles are largely converted into LiF after multiple cycles due to the promotion of HF by LiPF_6 , which corresponds to the isolated Li–F dominant region in Fig. 7b. Since the Li^+ conductivity of LiF is lower than that of Li_2O ($\sim 10^{-9} \text{ S cm}^{-1}$ for LiF and $\sim 10^{-7} \text{ S cm}^{-1}$ for Li_2O),^{162–164} the LiF-dominant layer derived from multiple cycles is more unfavorable for Li^+ migration. Importantly, the Li–O dominant layer can maintain uniform contact with Si particles after multiple cycles, while the Li–F dominant layer appears as isolated clusters outside the Si particles, with observed gaps between them. It can be inferred that LiFSI is beneficial for the formation of uniform SEI and promotes the smooth migration of Li^+ , while in the LiPF_6 system, the SEI distribution is uneven, which is not conducive to the Li^+ migration. Therefore, LiFSI is crucial for constructing uniform SEI to facilitate Li^+ migration, which provides more important contributions to accelerate kinetics at the interface than good Li^+ conductivity.

It is widely recognized that SEI presents various components, including Li_xPF_y , LiF, Li_2O , Li_2CO_3 , and organic Li–OR, LiO_2COR , and poly(ethylene-oxide) type polymer/oligomers, stacking on the surface of Si particles.^{123–125} Ensuring uniform distribution of each component and presenting organic–inorganic stratification to promote Li^+ migration is the key to improving kinetics at the interface level (Fig. 7c). Since SEI formation is mainly related to the reduction and decomposition of electrolyte solvents, it is necessary to regulate the solvent composition to construct SEI with optimized components and distributions.^{169–171} Si anodes largely follow carbonate solvents, which are composed of cyclic alkyl carbonates and one or several linear carbonates.^{170,172} Ethylene carbonate (EC) has been a widely used solvent component due to its high oxidation stability and high solubility to Li salts.¹⁶⁰ Nevertheless, the decomposition of EC through a ring-opening reaction will promote the formation of lithium ethylene dicarbonate (LEDC), which will further react with silicide, causing significant instability of the SEI and rapid thickening with cycling.^{173,174} Electrolyte additives can promote the SEI formation of Si anodes with thin and uniform distribution.^{151,169} Schroder and co-authors investigated the positive effect of FEC on SEI structural properties.¹⁷⁵ As the in-depth TOF-SIMS analysis described, introducing FEC can increase the inorganic chemical content (LiF and Li–X) while reducing the aliphatic carbon (C sp^3 , RCO, etc.), thus ensuring good Li^+ conductivity and the low Li^+ migration impedance of SEI at the interface.¹⁷⁶ As a result, the cycled Si–FEC cell delivers slightly higher charge transfer impedance than the compared Si cell at the initial lithiation due to its SEI thickness increase, while an obvious impedance reduction at the de-lithiation because of the stable and uniform SEI formation. It can be inferred that FEC can improve SEI structural strength during lithiation and ensure thin and uniform SEI during cycling to obtain low Li^+ migration impedance. Moreover, ether solvents have also received widespread attention due to their high ionic conductivity, low viscosity, and stability in reduction at the anode side. Ji and co-authors designed a mixed electrolyte solvent of tetrahydrofuran (THF) and 2-methyltetrahydrofuran (MTHF) to facilitate



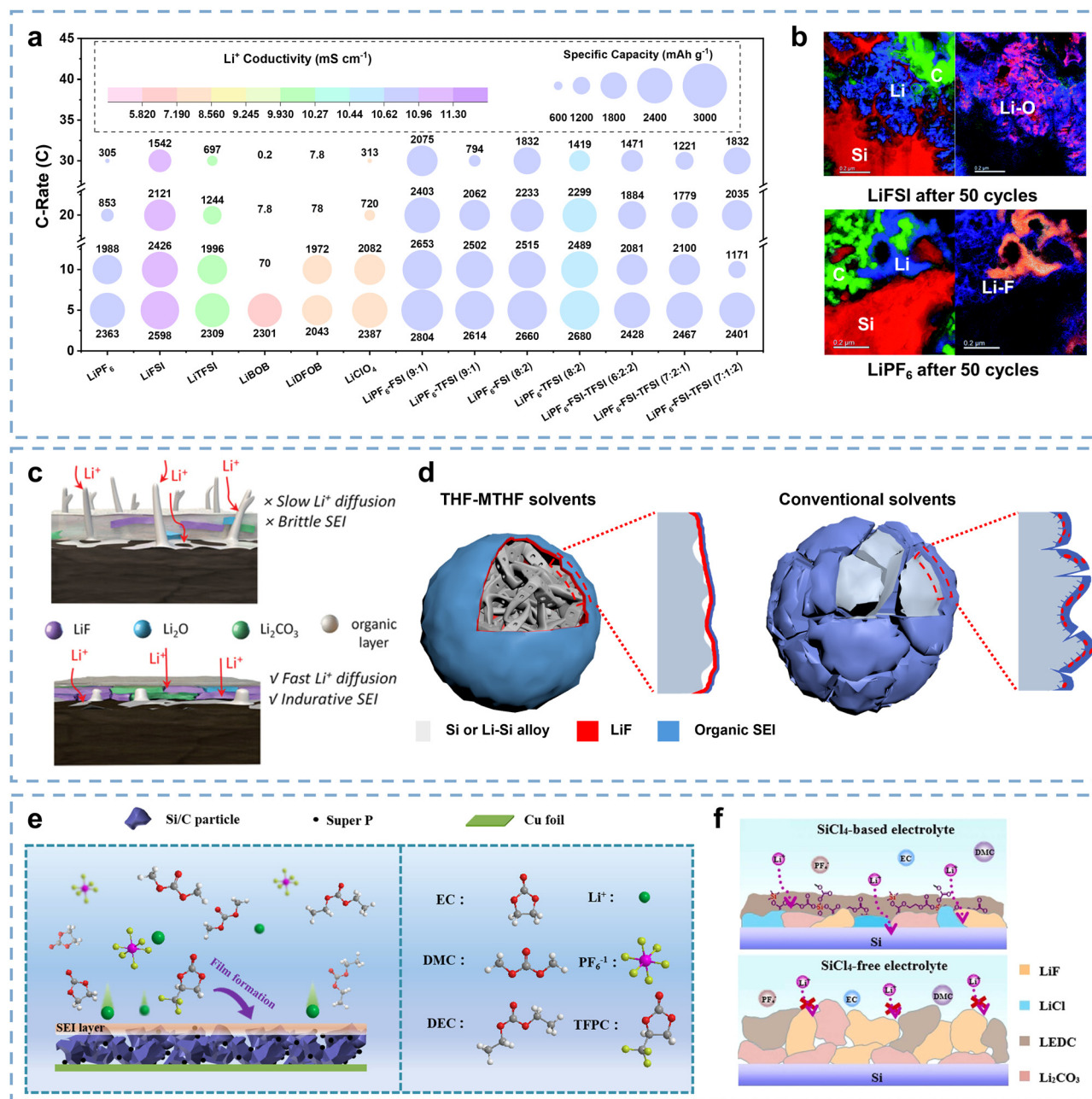


Fig. 7 (a) Comparison of rate performance and Li^+ conductivity for various Li salts.¹⁵⁵ Copyright © 2022 Elsevier Ltd. All rights reserved. (b) Element maps of Si electrodes with LiFSI and LiPF₆.¹⁶⁵ Copyright © 2022 Royal Society of Chemistry. (c) Schematic of Li^+ migration at the SEI.¹⁶⁶ Copyright © 2024 Wiley-VCH GmbH. (d) Schematic of SEI formation on Si particles in different solvents.¹²² Copyright © 2020, The Author(s), under exclusive license to Springer Nature Limited. (e) Schematic of Li^+ migration in electrolyte with TFPC.¹⁶⁷ Copyright © 2022, American Chemical Society. (f) Schematic of SEI structure with SiCl_4 .¹⁶⁸ Copyright © 2022 Elsevier Ltd. All rights reserved.

the formation of thin and uniform LiF-based SEI.¹²² The lower thermodynamic reduction sites of ether promote the preferential decomposition of LiPF₆, resulting in the preferential growth of fluoride salts on the surface of Si particles. As shown in Fig. 7d, a uniform SEI composed of internal LiF and external few organic species can form on the Si particles in the mixed THF electrolyte, avoiding the excessive thickening of SEI, which is conducive to Li^+ migration. In comparison, the SEI formed on Si particles for EC/DMC electrolytes exhibits LiF clusters

surrounded by organic components, which may thicken with cycling and hinder the Li^+ migration. Benefiting from the regulated thinner SEI, the Si electrode in mixed-THF electrolytes delivers 56.4% capacity retention at a current density of 3C and ~ 25% capacity retention in EC-DMC electrolytes. However, ether solvents present low oxidation decomposition potential, which limits their practical application. Developing multicomponent mixed solvents to fully utilize advantages is expected to be one of the necessary strategies for optimizing SEI



composition and distribution to promote kinetics at the interface.

Limited by the structural defects and challenging regulating technology of natural SEI, researchers have focused on developing artificial SEIs with enhanced uniformity and stability in the kinetic performance of Si anodes. Wen and co-authors proposed a potential trifluoropropylene carbonate (TFPC) additive to promote the formation of SEIs with good Li^+ conductivity and structural stability.¹⁶⁷ As shown in Fig. 7e, TFPC can promote the uniform formation of SEI on the surface of Si anodes, since its strong electron-withdrawing capability drives ring-opening reactions preferentially on the Si surface. After XPS-etching verification, the SEI induced by TFPC presents an organic-inorganic layered structure, which provides the possibility for the kinetics enhancement of Si anodes, delivering an improved rate performance of 33.6% retention at 5C. Yang and co-authors developed a novel artificial SEI with strong structural strength and chemical stability through the plasticizing reaction between SiCl_4 and LEDC.¹⁶⁸ As described in Fig. 7f, SiCl_4 -based electrolytes can facilitate the formation of cross-linked aliphatic carbon in the organic layer and LiCl in the inorganic layer of SEIs, thereby avoiding the decomposition of LEDC during cycling. Benefiting from unique structural advantages, the Si- SiCl_4 -based cell delivers a lower Warburg factor (σ) and shorter release time (τ_0) than the compared Si- SiCl_4 -free cell, suggesting enhanced Li^+ migration through the artificial SEI. As a result, the Si- SiCl_4 -based cell delivered a significant rate performance enhancement.

Based on uniformity and stability, optimizing the inorganic chemical component and concentration of SEIs is beneficial for improving the kinetic performance of Si anodes. The kinetics of Li^+ in different components can be estimated based on the migration barriers and conductivity of Li^+ .¹⁷⁷ The relevant information on potential SEI components is summarized in Table 4, especially Li_3N which stands out due to its extremely low migration barrier and high Li^+ conductivity, suggesting novel and efficient components urgently need to be introduced smoothly. For the typical components, more attention should be paid to the uniform distribution of components since they present no magnitude of difference among them. However, excessive inorganic chemicals will inevitably increase the density and thickness of SEIs, shielding electron transfer and hindering kinetic performance. Lin and co-authors explored

the electron tunneling-thickness effect of Li_2CO_3 , Li_3PO_4 , LiF, and Li layers as artificial SEIs.¹⁷⁸ They calculated the E_g , ΔE_t , d^* , and C of Li_2CO_3 , Li_3PO_4 , LiF, and Li layers of Si anodes. Obviously, the Li_2CO_3 layer with a critical thickness of 2 nm presents the lowest electron tunneling barrier compared to the Li_3PO_4 (1.66 nm) and LiF (1.62 nm) layers, illustrating that controlling the chemical distribution and corresponding thickness of SEIs is crucial for enhancing the kinetic performance of Si anodes from the perspective of interface design.

The components and distributions of SEIs make significant differences in the kinetics at the interface. Although the composition can be effectively adjusted by controlling Li salts, and solvents, and even introducing additives, the transfer ability of Li^+ in most SEI components is not outstanding. Therefore, the distribution of SEIs is more important to enhance kinetics at the interface. Constructing thin and uniform SEIs can facilitate the Li^+ migration at the interface more effectively and thereby promote interfacial kinetics.

4. Electrode level

During battery operation, the influence of initial structure and subsequent structural evolution of Si electrodes on their kinetic performance is generally overlooked. Although the initial kinetic performance of Si anodes can be significantly improved through particle and interface optimization, the slowly decreasing cycling performance during operation, which involves reversible capacity and cycling stability, has always been a serious challenge.^{198,199} This is largely attributed to the pore growth, electrode cracking, and material peeling caused by the huge volume expansion of the Si anodes, resulting in Li^+ diffusion interruption as another limiting factor that decreases their kinetic performance. Thus, the kinetic performance enhancement of Si anodes from the perspective of the electrode level in this section mainly contains electrode and collector structural design.

4.1 Electrode porosity optimization

Electrode pore evolutions can affect the Li^+ diffusion paths, especially in thick Si electrodes for high energy density, and the porosity directly determines the diffusion depth of Li^+ inside electrodes, greatly affecting the corresponding kinetics (Fig. 8a). Various attempts have been made to optimize the porosity of Si

Table 4 The Li^+ migration barrier and Li^+ conductivity of potential SEI components

Component	Li^+ migration barrier (eV)	Li^+ conductivity (S cm^{-1})	Ref.
Li_2O	0.28–0.33	$\sim 10^{-7}$	163, 164 and 179
Li_2CO_3	0.28–0.60	10^{-8} – 10^{-10}	180–182
LiOH	0.46	—	183
Li_3PO_4	0.3–0.5	8.62×10^{-8}	184 and 185
LiF	0.71	3×10^{-9}	183, 186 and 187
Li_3N	0.007, 0.038	6×10^{-6} (50 °C)	188
LiCl	0.268	2.085×10^{-4} , 5.767×10^{-4}	189–191
Li_2S	0.34	$\sim 10^{-8}$	192–194
Li-OR, LiO_2COR	—	$\sim 10^{-9}$	195–197



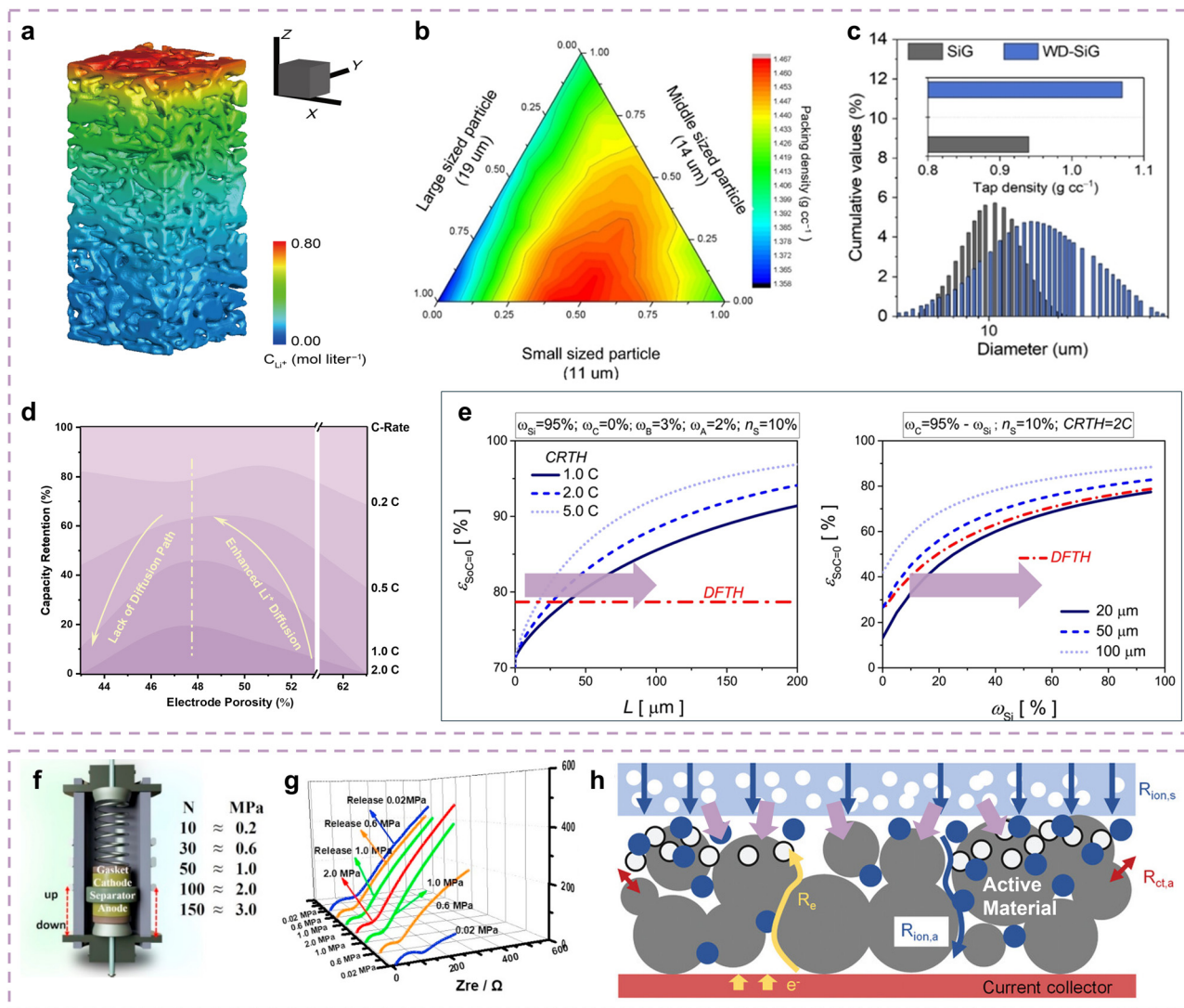


Fig. 8 (a) Schematic of Li⁺ distribution in electrode at high rate.²⁰² Copyright © 2022 The Authors, some rights reserved; exclusive licensee American Association for the Advancement of Science. (b) Particle diameter discrete element modeling simulation. (c) Statistical analysis for SiG and WD-SiG particles.²⁰¹ Copyright © 2022 Wiley-VCH GmbH. (d) The relationship between capacity retention and electrode porosity at different rates.²⁰³ Copyright © 2014 American Chemical Society. (e) Porosity kinetics contribution of Si electrodes in deformation threshold (DFT) and C-rate threshold (CRTH) conditions.²⁰⁴ Copyright © 2017 Elsevier Ltd. All rights reserved. (f) Schematic of external pressure research and (g) EIS tests of Si electrodes during pressure evolution.⁴¹ Copyright © 2020 Published by Elsevier Ltd. (h) Schematic of Li⁺ diffusion in an electrode with external pressure.²⁰⁵ Copyright © 2023 The Authors. Advanced Energy Materials published by Wiley-VCH GmbH.

electrodes for their kinetic performance enhancement. Wu and co-authors investigated the porosity effect on the kinetic performance of Si anodes by adjusting their particle size distribution.²⁰⁰ They provided four types of Si materials with D_{50} distributed at 20.1, 22.7, 15.7 and 13.5 μm. Benefiting from the Li⁺ diffusion path optimization endowed by moderate particle size distribution, the BSC2 anode delivers the lowest charge transfer impedance. It indicates that the pore structures between large particles can be filled by small particles, thereby reducing the electrode porosity and facilitating Li⁺ diffusion in the Si electrodes. Sung and co-authors conducted similar attempts on the particle size distribution of Si anodes.²⁰¹ As shown in Fig. 8b, Si electrodes with small particles generally have low packing densities, while small/large mixture particles contribute to high

packing densities. Benefiting from this design, the wide particle-size distribution Si/graphite hybrid (WD-SiG) material achieves a 14% increase in tap density when compared with the SiG material (Fig. 8c). Meanwhile, the WD-SiG anode also delivers low charge transfer impedance and rebounding tendency, indicating good adaptability during stress increases. Therefore, reducing the porosity of Si electrodes to optimize the Li⁺ diffusion paths and kinetic performance is a consensus.

In addition, appropriate porosity in thick Si electrodes is also beneficial for their kinetic performance enhancement due to the demand for good electrolyte infiltration and structural buffers. Zhao and co-authors further examined the influence of electrode porosity on the kinetic performance of Si anodes.²⁰³ As shown in Fig. 8d, the decrease in electrode porosity can

promote the Li^+ diffusion in the electrode while excessive reduction of porosity may decrease the Li^+ diffusion path, suggesting a “volcanic” relationship between the porosity and kinetic performance. Heubner and co-authors quantified the critical porosity on the kinetic performance of Si anodes depending on the deformation threshold (DFTH) and the C-rate threshold (CRTH) conditions.²⁰⁴ As described in Fig. 8e, the Si electrode with a 20 μm thickness requires 84% porosity to obtain the fast-charging performance at a current density of 5C, while the Si electrode with a 100 μm thickness requires 97% porosity.

Furthermore, appropriate external pressure on the Si electrodes also contributes to kinetic performance enhancement since it provides good Li^+/e^- contact among active materials during battery operation. As shown in Fig. 8f, Cui and co-authors determined the most appropriate external pressure of Si electrodes by investigating the stress evolution contribution to their kinetic performance.⁴¹ During the initial stage of stress increase (<1 MPa), the Si electrodes present few cracks and good Li^+/e^- conductivity. As the stress further increases (1–2 MPa), the poor electrolyte infiltration and electrode cracks cause a re-increase of charge transfer impedance (Fig. 8g). It can be observed that the Si electrode with 1 MPa external pressure exhibits the lowest impedance, demonstrating kinetic performance enhancement at the electrode level. As described in Fig. 8h, as the appropriate external pressure is applied, the Si particles are in closer contact with each other and adhere more closely to the conductive network, which undoubtedly promotes the Li^+ diffusion and electron transfer in the electrode, thereby effectively improving the kinetic performance.

Therefore, appropriate porosity is the prerequisite for enhancing the kinetics at the electrode level, as it affects electrolyte infiltration, Li^+ diffusion, and electron transfer. Specifically, loose particle distribution leads to ineffective volume occupation and prolongs the Li^+ diffusion inside the electrode, resulting in poor kinetics at the electrode level, while overly dense particle distribution may hinder electrolyte infiltration and obstruct the Li^+ diffusion path. Optimizing particle size distribution and applying external pressure are effective means to obtain appropriate gaps among particles and shorten the Li^+ diffusion distance while ensuring electrolyte infiltration.

4.2 Electrode integrity maintenance strategy

Significant volume changes of the Si anodes tend to cause electrode cracks during battery operation, which interrupts Li^+ diffusion and facilitates the formation of SEIs on the exposed crack surfaces, resulting in increased interface impedance, interrupted Li^+ diffusion paths, and kinetic performance decay.^{67,206} Numerous studies have been conducted to restrict and repair electrode cracks.^{42,207–209}

Electrode cracks can seriously affect electrolyte infiltration and Li^+ diffusion, and gas generation deteriorates the situation. Vanpeene and co-authors quantified the gas and electrolyte phase variations in electrode cracks through XRCT/XRD analysis.²¹⁰ The Si standard electrode has an electrolyte phase volume fraction of 13.3% and a gas phase volume fraction of 35% (Fig. 9a), indicating that the infiltration of electrolyte into

the electrode is incomplete, which is largely attributed to the obstruction of gas production to the electrolyte deep into the electrode. They further proposed a Si matured electrode which is obtained by placing the Si standard electrode in a humid environment to improve the microstructural environment inside the electrode for gas phase decrease.²¹¹ For the Si standard electrode, cracks start to form at 30% SOC and gradually grow with the charging depth increases, eventually connecting with each other to form the interrupted Li^+ diffusion paths. In contrast, the growth of cracks in matured electrodes is slower, which avoids the formation of interconnected cracks that penetrate the electrodes. As a result, the Si matured electrode presents a much lower crack volume fraction of 3.5% than the Si standard electrode of 11.6% and reduces the gas phase volume fraction by $\sim 75\%$ (Fig. 9b–d), indicating that cracks and gas generation can be effectively suppressed in the matured electrode. Therefore, addressing the crack and gas generation challenge to ensure electrode integrity is important for stabilizing the Li^+ diffusion paths of Si electrodes and enhancing the kinetic performance.

Binders play a crucial role in maintaining electrode integrity by fixing Si particles, conductive agents, and current collectors in the electrode, which can effectively suppress the electrode cracking, thereby ensuring stable kinetics during cycling. Importantly, binders can also be extended with good conductivity by grafting various effective e^-/Li^+ -conductive polymers to further enhance the kinetics at the electrode level, with the premise of ensuring electrode integrity.^{216–219} Endowing effective electronic conductivity to the binders can assist or replace the conductive agents to construct a complete conductive network throughout the Si electrodes, thereby accelerating electron transfer in the electrode. For example, Song and co-authors proposed a polyfluorene-type cross-linked conductive binder (CCB) serving as the secondary conductive network inside the Si electrode.²²⁰ The polyfluorene-type CCB exhibits good electronic conductivity of $5.5 \times 10^{-2} \text{ S cm}^{-1}$, promoting uniform current distribution inside the electrode and electronic depolarization effect.²²¹ The Si-CCB electrode delivers a capacity retention of $\sim 71\%$ at 1 A g^{-1} , higher than that of conventional CMC (56%) and PAA (62.5%) binders. Liu and co-authors developed a conductive polymer binder (sodium poly[9,9-bis(3-propanoate)fluorene], PF-COONa) with high conductivity provided by the main chains (n-type doped polyfluorene) and strong adhesion of the carboxyl side chains, to construct conductive agent free Si electrodes.²¹² As shown in Fig. 9e, the designed PF-COONa can maintain constant electrical connections even as the Si particles break into pieces, while the conventional conductive agents are pushed away and disconnected from Si particles due to the volume expansion of Si particles. Benefiting from good electronic conductivity and electrical connections, the Si/PF-COONa electrode delivers 41.5% capacity retention at 8.4 A g^{-1} , presenting superior kinetics at the electrode level. Moreover, the influence of ionic conductivity on kinetics is more significant, as superior ion conductive binders can provide potential Li^+ diffusion paths.^{217,219} Munaoka and co-authors designed a functional



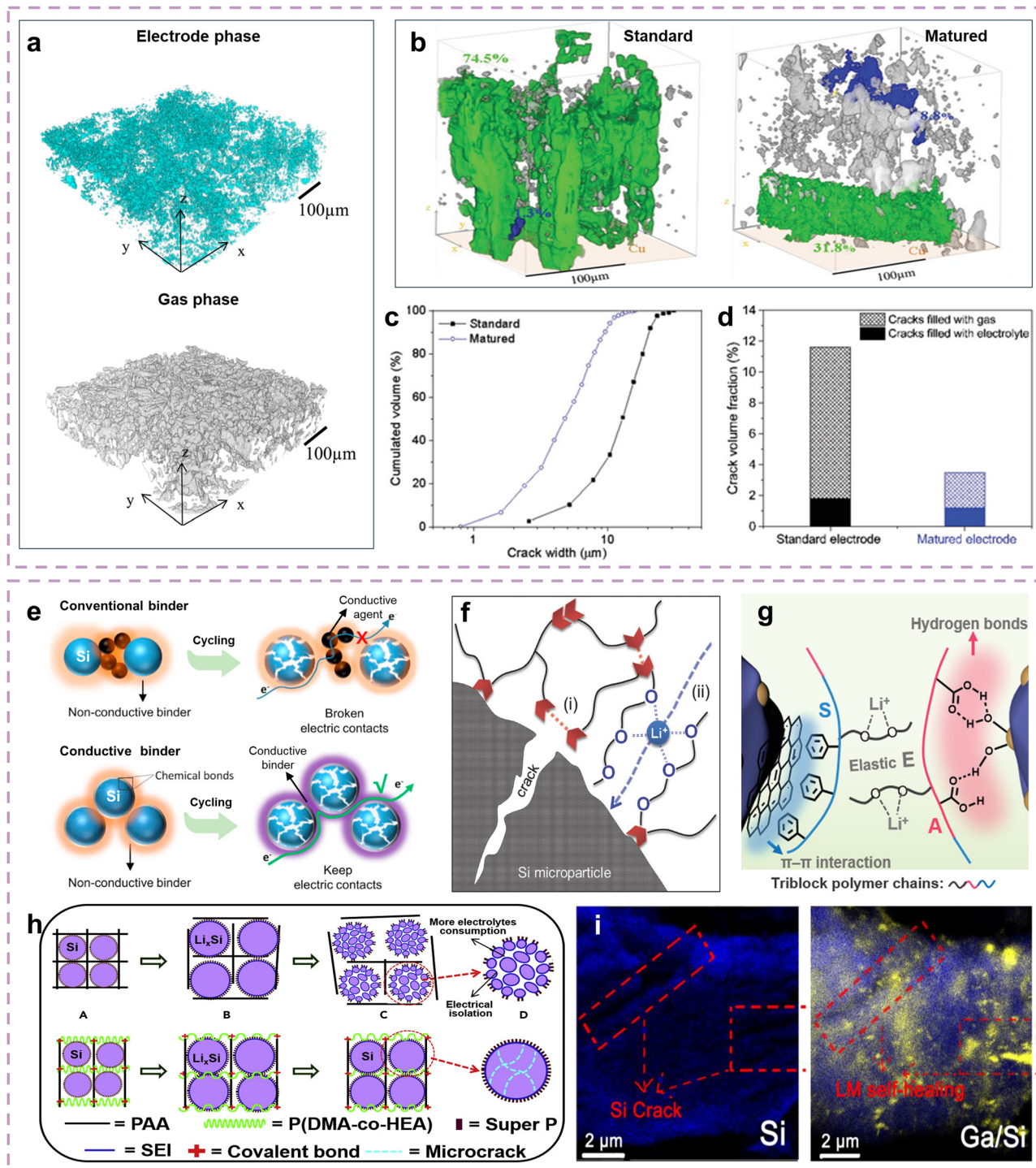


Fig. 9 (a) Electrolyte phase and gas phase distribution profiles of Si standard electrodes.²¹⁰ Copyright © 2018 Elsevier Ltd. All rights reserved. (b) Crack distribution 3D views of Si standard and matured electrodes, (c) crack volume fraction, and (d) electrolyte phase and gas phase volume fraction of Si standard and matured electrodes.²¹¹ Copyright © 2019 WILEY-VCH Verlag GmbH & Co. KGaA, Weinheim. (e) Schematic showing a comparison of electrical connections between a Si/PF-COONa electrode and conventional Si electrode.²¹² Copyright © 2017 Elsevier Ltd. All rights reserved. (f) Schematic of the self-healing and Li⁺ transfer of SHP-PEG binder.²¹³ Copyright © 2018 WILEY-VCH Verlag GmbH & Co. KGaA, Weinheim. (g) Schematic of the π-π stacking, hydrogen bonds, and the Li⁺ diffusion through electron-rich coordination sites for PSEA binder.²¹⁴ Copyright © 2022 Wiley-VCH GmbH. (h) Mechanical mechanism and schematic of PAA-P(HEA-co-DMA) binder.²⁰⁹ Copyright © 2018 Elsevier Inc. All rights reserved. (i) Element mapping of the initial cycled Si and Ga/Si anodes.²¹⁵ Copyright © 2018 Elsevier Ltd. All rights reserved.

binder (self-healing polymer incorporated with polyethylene glycol, SHP-PEG) with high Li⁺ conductivity.²¹³ As shown in

Fig. 9f, self-healing polymer (SHP) is rich in dynamic hydrogen bonds that endow the binders with significant mechanical



strength to ensure electrode integrity (Fig. 9h-i), while polyethylene glycol (PEG) units assist in the Li^+ transfer in the oxygen-rich channels of the binders (Fig. 9h-ii). Introducing 40% PEG2000 can promote an enhancement of Li^+ conductivity by 6 times. The Si/SHP-PEG electrode demonstrates 36% capacity retention at a current density of 2C. Li and co-authors proposed that oxygen heteroatoms can provide complexation sites to serve as Li^+ diffusion paths.²²² They further developed a high ionic conductivity binder (GG-g-PAM) by grafting polyacrylamide (PAM) onto the guar gum (GG) backbone, where the grafting chains buffer stress and oxygen heteroatoms promote Li^+ diffusion. Similarly, Hu and co-authors introduced side chains with oxygen atoms onto the binders to provide electron-rich coordination sites,²¹⁴ developing an interface-adaptive triblock binder (PSEA), including segments of hydrophobic polystyrene (S), elastic poly(2-(2-methoxyethoxy)ethyl acrylate) (E), and hydrophilic poly(acrylic acid) (A) (Fig. 9g). Meanwhile, the electrode integrity is double guaranteed by $\pi \cdots \pi$ stacking and abundant hydrogen bonds at the hydrophilic side. The Si/PSEA electrode presents an enhanced capacity retention of $\sim 50\%$ at a current density of 2C, while the Si/CMC-SBR electrode exhibits only $\sim 42\%$ capacity retention. Therefore, optimizing the Li^+ conductivity of the binder usually involves constructing electron-rich coordination sites to promote the Li^+ transfer, which can provide potential Li^+ diffusion paths, thereby enhancing the kinetics at the electrode level. Table 5 enumerates typical conductive binders for Si anodes, in which the increases in Li^+ conductivity often promote more significant kinetics enhancement, demonstrating the urgency of optimizing the ionic conductivity inside the Si electrodes.

Furthermore, various self-healing strategies have emerged to further maintain the integrity of Li^+ diffusion paths.^{232–234} Xu and co-authors restricted the formation of Si electrode cracks and repaired them through a novel polymer binder with high mechanical strength and good self-healing ability.²⁰⁹ The PAA-P(HEA-co-DMA) binder is synthesized through a thermal condensation reaction between PAA and P(HEA-co-DMA). Benefiting from the rigid and soft chain networks which provide good stretching and shrinking ability (Fig. 9h), the Si/PAA-P(HEA-co-DMA) electrode presents fewer cracks and better Li^+ diffusion

path integrity than the Si electrode. As a result, the Si/PAA-P(HEA-co-DMA) electrode delivers a significantly higher rate performance than the compared Si/PAA electrode. Moreover, Han and co-authors proposed a novel electrode crack repair strategy through liquid metal with high Li^+/e^- conductivity, which could fill these cracks and preserve electrode structural integrity.²¹⁵ As shown in Fig. 9i, Ga fills the Ga/Si electrode cracks, repairing the conductive network and Li^+ diffusion paths, which contributes to lower impedance. And the Si/Ga electrode delivers a higher kinetic performance of 16% capacity retention at 10C. According to Table 6, self-healing strategies can effectively promote the repair of electrode cracks, ensuring the integrity of the Li^+ diffusion path. After comparison, liquid metal can effectively maintain the Li^+ diffusion path and enhance electron transfer in the electrode due to its excellent conductivity and fluidity, thus improving the kinetics more significantly.

As a result, maintaining the integrity of Si electrodes is crucial for ensuring good kinetics during cycling. Binders should receive more attention since their potential for functionalization can satisfy multiple requirements on conductivity and self-healing. In particular, the development of binders with high conductivity helps to obtain matching rapid Li^+ diffusion and electron transfer, which can effectively improve the kinetic performance at the electrode level.

4.3 Current collector modification strategy

It is well known that the current collector plays an important role in electron collection and electrode support to maintain Li^+ diffusion path integrity and efficiency at the electrode level.²⁴⁴ Cu foil is recognized as an important commercial anode current collector due to its good electronic conductivity. However, Si materials tend to fall off the typical Cu foil since the smooth surface cannot provide enough adhesion for the electrode during battery operation. Thus, optimizing the surface characteristics of the current collector through a targeted modification strategy can also enhance and maintain the kinetic performance of the Si anodes.²⁴⁵ Xue and co-authors increased the surface roughness of Cu foil to provide an enhanced adhesion strength with Si materials, thereby ensuring the Si/m-Cu electrodes deliver higher kinetic performance than the compared Si/Cu electrode.²⁴⁶

Table 5 Typical electron/ion conductive binders for Si anodes

Binder	Conductive function	Rate performance	Cycling stability	Ref.
Polyfluorene-type CCB	Electron conductive	1 A g ⁻¹ , $\sim 71\%$	0.8 A g ⁻¹ , 250 cycles, 88%	220
PPyE	Electron conductive	2C, $\sim 45\%$	2C, 1000 cycles, $\sim 1500 \text{ mA h g}^{-1}$	223
PFPQ-COONa	Electron conductive	1C, $\sim 36\%$	0.5C, 400 cycles, 901 mA h g ⁻¹	224
Polymine	Electron conductive	2.5C, 28.2%	1C, 1000 cycles, 82.4%	225
PF-COONa	Electron conductive	8.4 A g ⁻¹ , 41.5%	4.2 A g ⁻¹ , 1000 cycles, 999 mA h g ⁻¹	212
APA/CNT	Electron conductive	0.5 A g ⁻¹ , 84.5%	0.2 A g ⁻¹ , 240 cycles, 87.7%	226
PP@CA	Electron conductive	3.2 A g ⁻¹ , 64.3%	2 A g ⁻¹ , 2000 cycles, 588 mA h g ⁻¹	227
b-POB	Electron/ion conductive	1C, 48.5%	1C, 100 cycles, $\sim 1800 \text{ mA h g}^{-1}$	228
SHP-PEG	Ion conductive	2C, $\sim 36\%$	0.5C, 150 cycles, 80%	213
GG-g-PAM	Ion conductive	8 A g ⁻¹ , $\sim 43\%$	1 A g ⁻¹ , 100 cycles, 83.9%	222
SSIP	Ion conductive	1C, $\sim 65\%$	0.5C, 400 cycles, 1620 mA h g ⁻¹	229
c-PEO-PEDOT:PSS/PEI	Ion conductive	8 A g ⁻¹ , 50.4%	1 A g ⁻¹ , 500 cycles, 2027 mA h g ⁻¹	230
PSEA	Ion conductive	2C, $\sim 50\%$	0.5C, 400 cycles, 82.1%	214
SHA	Ion conductive	1C, $\sim 44\%$	0.2C, 100 cycles, 1770.1 mA h g ⁻¹	231



Table 6 Typical self-healing strategies for Si anodes

Healing material	Rate performance	Cycling stability	Ref.
Ga/Si	10C, ~16%	4C, 1500 cycles, 81.3%	215
Ga _{12.6} Sn _{1.0}	3C, 54.6%	2C, 900 cycles, 95.7%	235
EGaIn-Mxene	2C, 30.8%	2C, 100 cycles	236
GaGeSiP ₃	20 mA cm ⁻² , 57.1%	6 mA cm ⁻² , 2000 cycles, 90%	237
PAA-P(HEA-co-DMA)	1.96C, 72.7%	5 A g ⁻¹ , 200 cycles, 90%	209
XPAA-DABBF	5C, 57.2%	0.5C, 500 cycles, 47.65%	238
PAA-DA/PVA	4 A g ⁻¹ , ~63%	4 A g ⁻¹ , 500 cycles, 50.8%	239
PAA-DA	1C, 66.8%	0.5C, 300 cycles, 1834.1 mA h g ⁻¹	240
TCB	3C, 49.8%	0.5C, 200 cycles, 621.2 mA h g ⁻¹	241
DNB	4C, 18.2%	1C, 300 cycles, 1115 mA h g ⁻¹	242
yCDp/Py-PAA	2C, 54.1%	0.5C, 300 cycles, 86.4%	243

Shen and co-authors further developed carbon-coated Cu foil through an *in situ* chemical vapor deposition strategy.²⁴⁷ The Si/

coated Cu foil pouch cell delivers lower initial impedance and operation impedance after 300 cycles than the compared Si/Cu

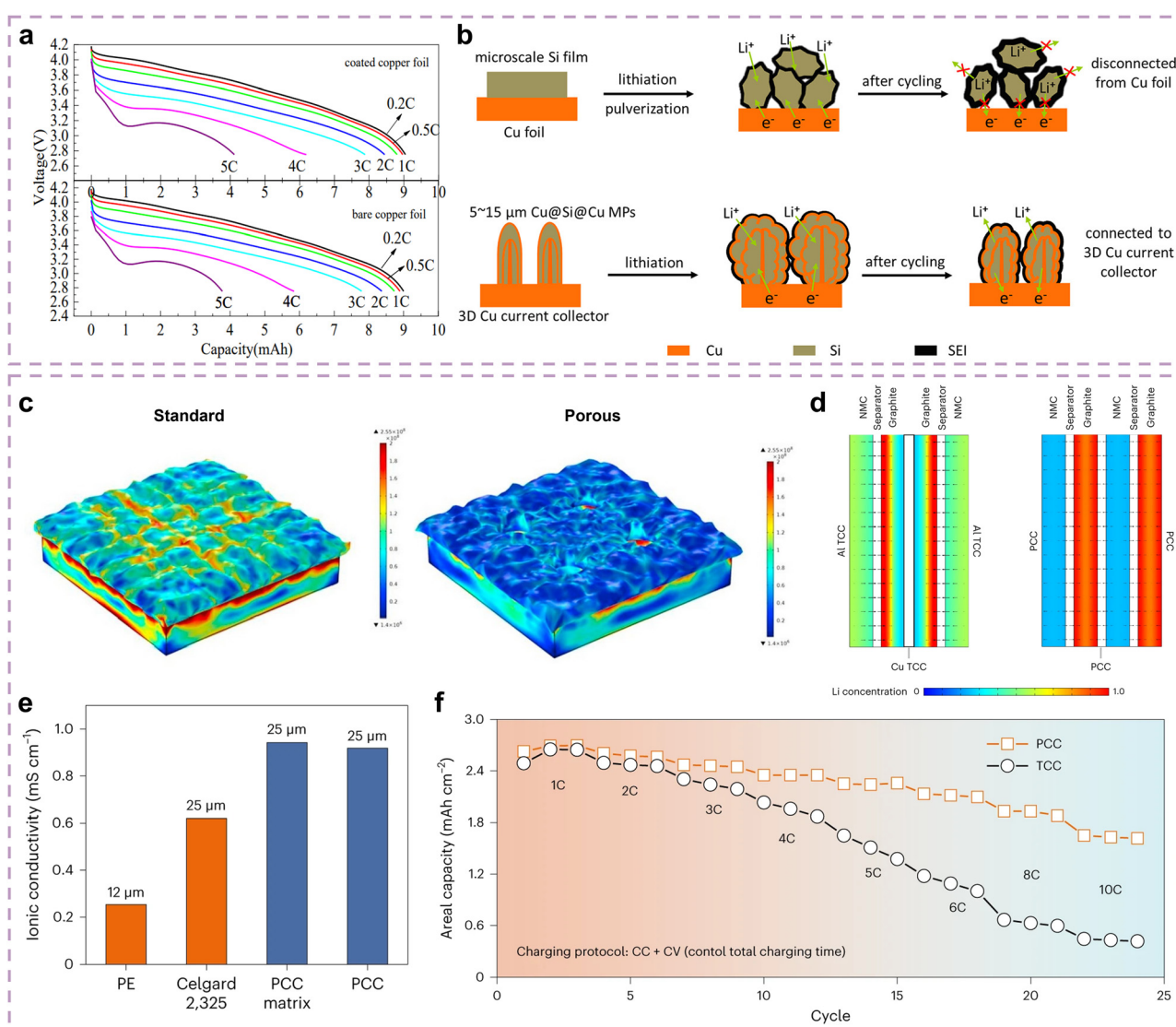


Fig. 10 (a) Rate performance of Si/bare Cu and Si/coated Cu electrodes.²⁴⁷ Copyright © 2020 The Authors. Published by ESG. Published by Elsevier B.V. (b) Schematic of kinetics enhancement for multisheathed Cu@Si@Cu microparticles coated on 3D Cu current collectors.²⁴⁸ Copyright © 2018 American Chemical Society. (c) Stress distribution of standard and porous Cu collectors.²⁴⁹ Copyright © 2018 Elsevier B.V. All rights reserved. (d) Li⁺ concentration distribution simulation of TCC and PCC electrode on a full charge, (e) ionic conductivity comparison of PE, Celgard 2325, PCC matrix and PCC, and (f) rate performance of Si/PCC and Si/TCC electrodes.⁴³ Copyright © 2024, The Author(s), under exclusive license to Springer Nature Limited.



foil pouch cell, suggesting that Cu foil with carbon coating is beneficial for maintaining the Li^+ diffusion path integrity. As a result, the Si/coated Cu foil pouch cell has a higher rate performance than the Si/Cu foil pouch cell (Fig. 10a). Furthermore, Zhang and co-authors developed multi-shelled Si@Cu microparticles supported on a 3D Cu collector to further improve the adhesion and electron transfer between Si materials and Cu collector.²⁴⁸ As shown in Fig. 10b, Si materials are confined within rigid Cu meshes, ensuring the integrity of Li^+ diffusion paths and promoting kinetic performance at the electrode level. Benefiting from the interconnected 3D Cu network with rich pore structures, the multi-shelled Cu@Si@Cu electrode has a good rate performance of 76% capacity retention at a current density of 10C.

Additionally, Moon and co-authors proposed using porous Cu collectors to release the stress of the Si anode during cycling, thereby ensuring the integrity of the electrode.²⁴⁹ As shown in Fig. 10c, the pores of the porous current collector can release the stress generated by volume expansion, thus keeping the internal stress of the electrode at a lower level compared to the standard current collector, which can improve the kinetics at the electrode level. Li^+ usually bypasses the traditional current collector (TCC) rather than passing through due to high compactness, which increases the Li^+ diffusion distance, thus restricting the kinetic performance of the Si anodes.²⁵⁰ Ye and co-authors designed a porous current collector (PCC) by modifying nanoporous Kevlar film to resolve this challenge.⁴³ As described in Fig. 10d, an even Li^+ concentration ranging from 1.0 to 0.8 along the distance from separator to PCC can be observed. As a comparison, the TCC exhibits an uneven and wider Li^+ concentration range (1.0–0.3) along the distance from the separator to the TCC. The rich porosity and high Li^+ conductivity can effectively shorten Li^+ diffusion distance and accelerate Li^+ migration, thereby ensuring the high kinetic performance of Si/PCC electrodes (Fig. 10e and f). As a result, the Si/PCC pouch cell delivers a better rate capacity retention of 78.3% at 4C (62.3% vs. Si/TCC) and 70.5% at 6C (33.4% vs. Si/TCC), respectively. Therefore, the adhesion strength and Li^+/e^- conductivity of the current collector foil in the kinetic performance enhancement of Si anodes cannot be neglected.

The ability of the current collector to collect and transfer electrons can also limit the kinetic performance at the electrode level.²⁴⁴ Although Cu foil has superior conductivity, the contact between the electrode and the current collector remains the limiting factor. For Si electrodes, volume changes make active materials tend to peel from the current collector, further exacerbating the unstable contact. Improving the adhesion between the electrode and the current collector is the most effective method to enhance the kinetic performance of Si anodes at the electrode–collector interface.

5. Summary and perspective

To conclude, the particle–interface–electrode integration of Si anodes is deconstructed in this review to identify the main

issues that limit the kinetic performance: long Li^+ diffusion distance and poor intrinsic conductivity for particles, high Li^+ migration impedance at the interface due to defective material surface and limited migration barrier, and insufficient or even interrupted Li^+ diffusion paths for initial and cycled electrodes. Through vertical combining of particle, interface, and electrode levels and fully considering their mutual influences, we have evaluated effective strategies for progressively enhancing the kinetics of Si anodes. Specifically, as follows:

(i) Si materials with appropriate particle size or pore size distribution largely possess enhanced kinetic performance due to the balance between Li^+ diffusion distance and Li^+ migration impedance.

(ii) Element doping and particle compositing strategies could effectively enhance the conductivity or promote the Li^+ diffusion inside particles. And further adjusting the doping or compositing concentration can obtain the correspondingly optimal kinetic performance.

(iii) Coating and SEI optimization strategies could improve the interface properties of Si anodes, which is conducive to facilitating Li^+ migration at the interface, thereby significantly reducing the Li^+ migration impedance of Si anodes.

(iv) Appropriate particle size distribution and external pressure, effective binder strategies, as well as modification of the roughness and pore structure of the current collector can adjust electrode porosity and suppress the occurrence of cracks and peeling for initial or cycling electrodes, thus ensuring the integrity of Si electrodes and sufficient Li^+ diffusion paths.

Although the emerging strategies promote significant academic and commercial achievements for Si anodes, quantifying the quantitative relationship between kinetics and the strategies of multiple dimensions, as well as balancing their mutual influence, remains a serious challenge of fast charging Si anodes for power batteries. For this reason, several potential solutions require further exploration.

(i) Quantifying the optimal particle size and pore size. The particle size and pore size distribution of Si materials are closely related to their Li^+ diffusion distance and involve mutual influence among particle, interface, and even electrode multiple levels. Establishing two theoretical models of “particle size – Li^+ diffusion distance” and “pore size distribution – Li^+ diffusion distance” with the optimal values is a challenging but necessary topic.

(ii) Establishing the Li^+ diffusion barrier comparison table of Si anodes. Although doping element types and the corresponding concentration of Si materials affect the conductivity, most attempts are undirected and inefficient. Therefore, it is necessary to establish the above standard comparison table based on doping elements and their concentration to overcome this challenge.

(iii) Evaluating the contributions of coatings to kinetics for Si anodes. Numerous coatings have been established for Si anodes, endowed with various functions including conductivity, Li^+ migration, mechanical strength, and SEI formation. However, their contributions to kinetic behavior and the performance of Si anodes lack quantification. Establishing a database to summarize the enhancement percentage on rate



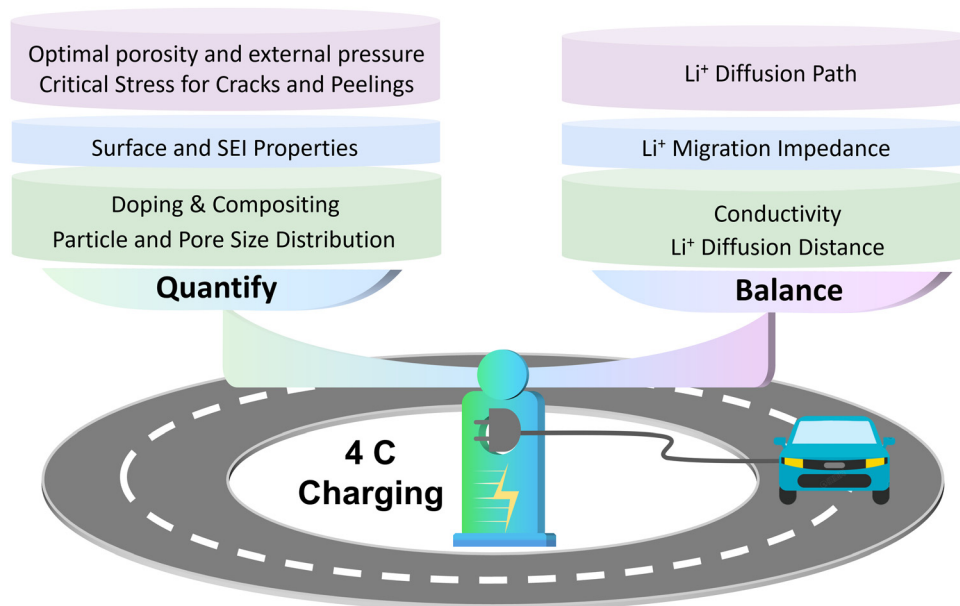


Fig. 11 Prospects for improving the kinetic performance of Si anodes.

performance through different coatings and their thickness is a significant quantitative method to select the optimal coating suitable for fast-charging Si anodes.

(iv) Constructing targeted SEI precisely. Although in-depth research has been conducted on various components of the SEI, a single component is not sufficient to support the revelation of the complex and disordered SEI integration, resulting in difficulty in controlling the uniformity, composition, thickness, *etc.* More attention should be paid to the development and introduction of SEIs dominated by components with excellent Li^+ conductivity (*e.g.* Li_3N). For typical SEIs, the optimization should be holistically discussed from multiple perspectives of composition, spatial distribution, and overall thickness, avoiding focusing solely on the effect of a single component in a single perspective.

(v) Maintaining Li^+ diffusion path integrity during cycling. The formation of porosity, cracks, and peeling of Si electrodes due to the huge volume change and uneven stress distribution tend to interrupt the Li^+ diffusion path. How to quantify their impact on the kinetic performance degradation of Si electrodes remains a challenge. The theoretical model “local stress-strain distribution” should be applied to the Si anodes that have undergone enhancement at particle and interface levels. And then calculate the critical stress of crack generation, which is beneficial for further optimizing the Si electrodes and fully exerting the enhancement strategies from multiple levels.

(vi) Promoting dry electrode technology on Si electrodes. Dry electrode technology can avoid the uncontrollable migration of non-active materials stemming from the drying process in the wet process, thereby ensuring a uniform ion/electron transfer network, which is beneficial for improving the fast-charging performance of Si anodes theoretically. However, severe volume expansion can hinder the application of the current dry process on Si anodes. Specifically, Si particles, conductive agents, and binders all suffer from agglomeration stemming from the lack of solvent dispersion,

resulting in poor conductivity and electrode instability. Furthermore, volume expansion can exacerbate this instability. It is crucial to develop materials (*e.g.* binders and conductive agents) and optimize dry processes adaptable to Si anodes. In particular, the development of dry electrode technologies suitable for Si anodes should focus on the adaptation of surface energy among binders, conductive agents, and Si particles to promote the balance between binder cohesion and the adhesion to current collector and materials, thus enhancing the kinetics at the electrode level.

In summary, the kinetic performance enhancement of Si anodes requires design strategies from multiple dimensions. Benefiting from the viewpoints of particle, interface, and electrode levels, we expect this review would contribute to the research and development of high energy density and fast-charging power LIBs based on Si anodes, leaping towards the 4C fast charging era (Fig. 11).

Author contributions

P. Jia: conceptualization, visualization, writing – original draft, J. Guo: conceptualization, writing – review & editing, Q. Li: writing – review & editing, Y. Liu: investigation, formal analysis, Y. Zheng and Y. Guo: formal analysis, validation, Y. Huang, Y. Shen, L. Long, S. Dong, and J. Jiang: validation, H. Zhang, R. Chen, C. Zhang, Z. Zhang, and J. Shen: formal analysis, M. Chang, X. Liu, and X. Wang: visualization, Y. Tang: writing – review & editing, H. Shao: writing – review & editing, supervision, funding acquisition, project administration.

Data availability

No primary research results, software or code have been included and no new data were generated or analysed as part of this review.



Conflicts of interest

The authors declare no conflicting interests regarding the content of this article.

Acknowledgements

This work was financially supported by the Science and Technology Planning Project of Shenzhen of China (Shenzhen-Hong Kong-Macao Category C) (No. SGDX20220530111004028), Macau Science and Technology Development Fund (FDCT) for funding of the Macao Centre for Research and Development in Advanced Materials (2022–2024) (No. 0026/2022/AMJ and 006/2022/ALC), the Natural Science Foundation of Guangdong Province (No. 2023A1515010765), the Science and Technology Planning Project of Guangdong Province of China (No. 2023A0505030001), and a Multi-Year Research Grant (MYRG) from University of Macau (No. MYRG-GRG2023-00140-IAPME-UMDF).

References

- 1 I. E. A. (2023), Global EV Outlook 2023, <https://www.iea.org/reports/global-ev-outlook-2023>, (accessed April, 2023).
- 2 Z. Cano, D. Banham, S. Ye, A. Hintennach, J. Lu, M. Fowler and Z. Chen, *Nat. Energy*, 2018, **3**, 279–289.
- 3 J. Goodenough, *Nat. Electron.*, 2018, **1**, 204.
- 4 D. Meng, Z. Xue, G. Chen, D. Zhou, Y. He, Z. Ma, Y. Liu and L. Li, *Energy Environ. Sci.*, 2024, **17**, 4658–4669.
- 5 M. M. Yuan, H. J. Liu and F. Ran, *Mater. Today*, 2023, **63**, 360–379.
- 6 M. Weiss, R. Ruess, J. Kasnatscheew, Y. Levartovsky, N. R. Levy, P. Minnmann, L. Stolz, T. Waldmann, M. Wohlfahrt-Mehrens, D. Aurbach, M. Winter, Y. Ein-Eli and J. Janek, *Adv. Energy Mater.*, 2021, **11**, 2101126.
- 7 A. Zhang, Z. Bi, G. Wang, S. Liao, P. Das, H. Lin, M. Li, Y. Yu, X. Feng, X. Bao and Z. Wu, *Energy Environ. Sci.*, 2024, **17**, 3021–3031.
- 8 C. Y. Wang, T. Liu, X. G. Yang, S. H. Ge, N. V. Stanley, E. S. Rountree, Y. J. Leng and B. D. McCarthy, *Nature*, 2022, **611**, 485–490.
- 9 S. H. Park, P. J. King, R. Y. Tian, C. S. Boland, J. Coelho, C. F. Zhang, P. McBean, N. McEvoy, M. P. Kremer, D. Daly, J. N. Coleman and V. Nicolosi, *Nat. Energy*, 2019, **4**, 560–567.
- 10 C. F. Zhang, S. N. Park, A. Seral-Ascaso, S. Barwich, N. McEvoy, C. S. Boland, J. N. Coleman, Y. Gogotsi and V. Nicolosi, *Nat. Commun.*, 2019, **10**, 849.
- 11 X. Su, Q. L. Wu, J. C. Li, X. C. Xiao, A. Lott, W. Q. Lu, B. W. Sheldon and J. Wu, *Adv. Energy Mater.*, 2014, **4**, 1300882.
- 12 U. S. A. B. Consortium, Low-Cost/Fast-Charge EV Goals, <https://uscar.org/usabc/>.
- 13 Tesla Model S, <https://www.tesla.com/models>.
- 14 B.Y.D., <https://www.bydauto.com.cn/pc/brandRelated?id=3>.
- 15 M. Armand and J. Tarascon, *Nature*, 2008, **451**, 652–657.
- 16 Y. Liu, Y. Zhu and Y. Cui, *Nat. Energy*, 2019, **4**, 540–550.
- 17 E. Pomerantseva, F. Bonaccorso, X. Feng, Y. Cui and Y. Gogotsi, *Science*, 2019, **366**, eaan8285.
- 18 V. Etacheri, R. Marom, R. Elazari, G. Salitra and D. Aurbach, *Energy Environ. Sci.*, 2011, **4**, 3243–3262.
- 19 M. Li, J. Lu, Z. Chen and K. Amine, *Adv. Mater.*, 2018, **30**, 1800561.
- 20 W. X. Huang, Y. S. Ye, H. Chen, R. A. Vilá, A. Xiang, H. X. Wang, F. Liu, Z. A. Yu, J. W. Xu, Z. W. Zhang, R. Xu, Y. C. Wu, L. Y. Chou, H. S. Wang, D. T. Boyle, Y. Z. Li and Y. Cui, *Nat. Commun.*, 2022, **13**, 7091.
- 21 D. P. Finegan, A. Quinn, D. S. Wragg, A. M. Colclasure, X. K. Lu, C. Tan, T. M. M. Heenan, R. Jervis, D. J. L. Brett, S. Das, T. Gao, D. A. Cogswell, M. Z. Bazant, M. Di Michiel, S. Checchia, P. R. Shearing and K. Smith, *Energy Environ. Sci.*, 2020, **13**, 2570–2584.
- 22 M. Salah, P. Murphy, C. Hall, C. Francis, R. Kerr and M. Fabretto, *J. Power Sources*, 2019, **414**, 48–67.
- 23 H. Jin, Y. Huang, C. Wang and H. Ji, *Small Sci.*, 2022, **2**, 2200015.
- 24 Y. Liu, H. Shi and Z. Wu, *Energy Environ. Sci.*, 2023, **16**, 4834–4871.
- 25 H. Zhang, L. Wang, H. Li and X. M. He, *ACS Energy Lett.*, 2021, **6**, 3719–3724.
- 26 A. Zulke, Y. Li, P. Keil, R. Burrell, S. Belaisch, M. Nagarathinam, M. P. Mercer and H. E. Hoster, *Batteries Supercaps*, 2021, **4**, 934–947.
- 27 Y. F. Li, Q. M. Li, J. L. Chai, Y. T. Wang, J. K. Du, Z. Y. Chen, Y. C. Rui, L. Jiang and B. H. J. Tang, *ACS Mater. Lett.*, 2023, **5**, 2948–2970.
- 28 X. Xia, X. Y. Qian, C. Chen, W. Y. Li, D. F. He, G. Y. He and H. Q. Chen, *J. Energy Storage*, 2023, **72**, 108715.
- 29 M. Jiang, J. Chen, Y. Zhang, N. Song, W. Jiang and J. Yang, *Adv. Sci.*, 2022, **9**, 2203162.
- 30 M. Peng, K. Shin, L. Jiang, Y. Jin, K. Zeng, X. Zhou and Y. Tang, *Angew. Chem., Int. Ed.*, 2022, **61**, e202206770.
- 31 S. Chae, S. Choi, N. Kim, J. Sung and J. Cho, *Angew. Chem., Int. Ed.*, 2020, **59**, 110–135.
- 32 C. Y. Chou and G. S. Hwang, *Surf. Sci.*, 2013, **612**, 16–23.
- 33 B. Peng, F. Y. Cheng, Z. L. Tao and J. Chen, *J. Chem. Phys.*, 2010, **133**, 034701.
- 34 W. H. Wan, Q. F. Zhang, Y. Cui and E. G. Wang, *J. Phys.:Condens. Matter*, 2010, **22**, 415501.
- 35 Q. F. Zhang, Y. Cui and E. G. Wang, *J. Phys. Chem. C*, 2011, **115**, 9376–9381.
- 36 L. Wang, J. Yu, S. Y. Li, F. S. Xi, W. H. Ma, K. X. Wei, J. J. Lu, Z. Q. Tong, B. Liu and B. Luo, *Energy Storage Mater.*, 2024, **66**, 103243.
- 37 Q. Y. Man, Y. L. An, C. K. Liu, H. T. Shen, S. L. Xiong and J. K. Feng, *J. Energy Chem.*, 2023, **76**, 576–600.
- 38 H. H. Zheng, L. Tan, G. Liu, X. Y. Song and V. S. Battaglia, *J. Power Sources*, 2012, **208**, 52–57.
- 39 I. V. Thorat, D. E. Stephenson, N. A. Zacharias, K. Zaghib, J. N. Harb and D. R. Wheeler, *J. Power Sources*, 2009, **188**, 592–600.
- 40 Q. Li, X. S. Liu, X. Han, Y. X. Xiang, G. M. Zhong, J. Wang, B. Z. Zheng, J. G. Zhou and Y. Yang, *ACS Appl. Mater. Interfaces*, 2019, **11**, 14066–14075.



- 41 J. Cui, X. Chen, Z. Zhou, M. Zuo, Y. Xiao, N. Zhao, C. Shi and X. Guo, *Mater. Today Energy*, 2021, **20**, 100632.
- 42 H. Chen, Z. Z. Wu, Z. Su, S. Chen, C. Yan, M. Al-Mamun, Y. B. Tang and S. Q. Zhang, *Nano Energy*, 2021, **81**, 105654.
- 43 Y. S. Ye, R. Xu, W. X. Huang, H. Y. Ai, W. B. Zhang, J. O. Affeld, A. Cui, F. Liu, X. Gao, Z. Y. Chen, T. Y. Li, X. Xiao, Z. W. Zhang, Y. C. Peng, R. A. Vila, Y. C. Wu, S. T. Oyakhire, H. Kuwajima, Y. Suzuki, R. Matsumoto, Y. Masuda, T. Yuuki, Y. Nakayama and Y. Cui, *Nat. Energy*, 2024, **9**, 11.
- 44 F. H. Du, K. X. Wang and J. S. Chen, *J. Mater. Chem. A*, 2016, **4**, 32–50.
- 45 M. T. McDowell, S. W. Lee, W. D. Nix and Y. Cui, *Adv. Mater.*, 2013, **25**, 4966–4984.
- 46 S. W. Lee, H. W. Lee, I. Ryu, W. D. Nix, H. J. Gao and Y. Cui, *Nat. Commun.*, 2015, **6**, 7533.
- 47 H. Yang, S. Huang, X. Huang, F. F. Fan, W. T. Liang, X. H. Liu, L. Q. Chen, J. Y. Huang, J. Li, T. Zhu and S. L. Zhang, *Nano Lett.*, 2012, **12**, 1953–1958.
- 48 X. H. Liu, J. W. Wang, S. Huang, F. F. Fan, X. Huang, Y. Liu, S. Krylyuk, J. Yoo, S. A. Dayeh, A. V. Davydov, S. X. Mao, S. T. Picraux, S. L. Zhang, J. Li, T. Zhu and J. Y. Huang, *Nat. Nanotechnol.*, 2012, **7**, 749–756.
- 49 M. McDowell, S. Lee, J. Harris, B. Korgel, C. Wang, W. Nix and Y. Cui, *Nano Lett.*, 2013, **13**, 758–764.
- 50 J. Wang, Y. He, F. Fan, X. Liu, S. Xia, Y. Liu, C. Harris, H. Li, J. Huang, S. Mao and T. Zhu, *Nano Lett.*, 2013, **13**, 709–715.
- 51 B. Key, R. Bhattacharyya, M. Morcrette, V. Seznéc, J. M. Tarascon and C. P. Grey, *J. Am. Chem. Soc.*, 2009, **131**, 9239–9249.
- 52 M. McDowell, I. Ryu, S. Lee, C. Wang, W. Nix and Y. Cui, *Adv. Mater.*, 2012, **24**, 6034–6041.
- 53 Z. Jia and T. Li, *J. Mech. Phys. Solids*, 2016, **91**, 278–290.
- 54 M. T. McDowell, S. W. Lee, J. T. Harris, B. A. Korgel, C. M. Wang, W. D. Nix and Y. Cui, *Nano Lett.*, 2013, **13**, 758–764.
- 55 Q. Y. Liu, C. L. Pang, W. L. Chen, Z. X. Rao, H. Q. Lu, L. H. Xue and W. X. Zhang, *Energy Technol.*, 2019, **7**, 1900487.
- 56 F. F. Zhao, M. Zhao, Y. R. Dong, L. Ma, Y. Zhang, S. L. Niu and L. M. Wei, *Powder Technol.*, 2022, **404**, 117455.
- 57 J. Szczech and S. Jin, *Energy Environ. Sci.*, 2011, **4**, 56–72.
- 58 X. H. Liu, L. Zhong, S. Huang, S. X. Mao, T. Zhu and J. Y. Huang, *ACS Nano*, 2012, **6**, 1522–1531.
- 59 Z. Bitew, M. Tesemma, Y. Beyene and M. Amare, *Sustainable Energy Fuels*, 2022, **6**, 1014–1050.
- 60 H. Li, H. Li, Y. Lai, Z. Yang, Q. Yang, Y. Liu, Z. Zheng, Y. Liu, Y. Sun, B. Zhong, Z. Wu and X. Guo, *Adv. Energy Mater.*, 2022, **12**, 2102181.
- 61 H. Wu and Y. Cui, *Nano Today*, 2012, **7**, 414–429.
- 62 A. Casimir, H. Zhang, O. Ogoke, J. Amine, J. Lu and G. Wu, *Nano Energy*, 2016, **27**, 359–376.
- 63 A. Mukanova, A. Jetybayeva, S. T. Myung, S. S. Kim and Z. Bakenov, *Mater. Today Energy*, 2018, **9**, 49–66.
- 64 T. Kasukabe, H. Nishihara, S. Iwamura and T. Kyotani, *J. Power Sources*, 2016, **319**, 99–103.
- 65 C. K. Chan, H. L. Peng, G. Liu, K. McIlwrath, X. F. Zhang, R. A. Huggins and Y. Cui, *Nat. Nanotechnol.*, 2008, **3**, 31–35.
- 66 X. H. Liu, H. Zheng, L. Zhong, S. Huan, K. Karki, L. Q. Zhang, Y. Liu, A. Kushima, W. T. Liang, J. W. Wang, J. H. Cho, E. Epstein, S. A. Dayeh, S. T. Picraux, T. Zhu, J. Li, J. P. Sullivan, J. Cumings, C. S. Wang, S. X. Mao, Z. Z. Ye, S. L. Zhang and J. Y. Huang, *Nano Lett.*, 2011, **11**, 3312–3318.
- 67 H. Wu, G. Chan, J. W. Choi, I. Ryu, Y. Yao, M. T. McDowell, S. W. Lee, A. Jackson, Y. Yang, L. B. Hu and Y. Cui, *Nat. Nanotechnol.*, 2012, **7**, 309–314.
- 68 M. Park, M. Kim, J. Joo, K. Kim, J. Kim, S. Ahn, Y. Cui and J. Cho, *Nano Lett.*, 2009, **9**, 3844–3847.
- 69 A. Arie and J. Lee, *Nanotechnol. Appl. Energy Environ.*, 2013, **737**, 80–84.
- 70 F. Wu, Y. Dong, Y. F. Su, C. X. Wei, T. R. Chen, W. A. Yan, S. Y. Ma, L. Ma, B. Wang, L. Chen, Q. Huang, D. Y. Cao, Y. Lu, M. Wang, L. Wang, G. Q. Tan, J. H. Wang and N. Li, *Small*, 2023, **19**, 2301301.
- 71 J. Liang, X. N. Li, Z. G. Hou, W. Q. Zhang, Y. C. Zhu and Y. T. Qian, *ACS Nano*, 2016, **10**, 2295–2304.
- 72 L. Zhang, C. Wang, Y. Dou, N. Cheng, D. Cui, Y. Du, P. Liu, M. Al-Mamun, S. Zhang and H. Zhao, *Angew. Chem., Int. Ed.*, 2019, **58**, 8824–8828.
- 73 Z. Cheng, H. Lin, Y. Liu, J. Li, H. Jiang and H. Zhang, *Small*, 2024, **20**, 2407560.
- 74 Y. Xu, E. Swaans, S. Chen, S. Basak, P. Harks, B. Peng, H. Zandbergen, D. Borsa and F. Mulder, *Nano Energy*, 2017, **38**, 477–485.
- 75 H. Chen, Y. Xiao, L. Wang and Y. Yang, *J. Power Sources*, 2011, **196**, 6657–6662.
- 76 Y. Cheng, C. Chen, S. Wang, Y. Li, B. Peng, J. Huang and C. Liu, *Nano Energy*, 2022, **102**, 107688.
- 77 M. Datta, J. Maranchi, S. Chung, R. Epur, K. Kadakia, P. Jampani and P. Kumta, *Electrochim. Acta*, 2011, **56**, 4717–4723.
- 78 J. Ryu, D. Hong, S. Choi and S. Park, *ACS Nano*, 2016, **10**, 2843–2851.
- 79 G. L. Hou, B. L. Cheng, Y. B. Cao, M. S. Yao, B. Q. Li, C. Zhang, Q. H. Weng, X. Wang, Y. Bando, D. Golberg and F. L. Yuan, *Nano Energy*, 2016, **24**, 111–120.
- 80 N. Liu, Z. Lu, J. Zhao, M. McDowell, H. Lee, W. Zhao and Y. Cui, *Nat. Nanotechnol.*, 2014, **9**, 187–192.
- 81 S. Chae, Y. Xu, R. Yi, H. Lim, D. Velickovic, X. Li, Q. Li, C. Wang and J. Zhang, *Adv. Mater.*, 2021, **33**, 2103095.
- 82 J. Entwistle, A. Rennie and S. Patwardhan, *J. Mater. Chem. A*, 2018, **6**, 18344–18356.
- 83 W. Liu, J. Z. Wang, J. T. Wang, X. Z. Guo and H. Yang, *J. Alloys Compd.*, 2021, **874**, 159921.
- 84 Y. C. Zhao, C. G. Liu, Y. Sun, R. W. Yi, Y. T. Cai, Y. Q. Li, I. Mitrovic, S. Taylor, P. Chalker, L. Yang and C. Z. Zhao, *J. Alloys Compd.*, 2019, **803**, 505–513.
- 85 W. Ren, Y. Wang, Z. Zhang, Q. Tan, Z. Zhong and F. Su, *J. Mater. Chem. A*, 2016, **4**, 552–560.
- 86 D. Kim, S. Hwang, J. Cho, S. Yu, S. Kim, J. Jeon, K. Ahn, C. Lee, H. Song and H. Lee, *ACS Energy Lett.*, 2019, **4**, 1265–1270.



- 87 M. Liu, W. Xu, S. Liu, B. Liu, Y. Gao and B. Wang, *Adv. Sci.*, 2024, **11**, 2402915.
- 88 Y. Chen, F. Guo, L. Yang, J. Lu, D. Liu, H. Wang, J. Zheng, X. Yu and H. Li, *Chin. Phys. B*, 2022, **31**, 078201.
- 89 Y. Tang, Y. Zhang, W. Li, B. Ma and X. Chen, *Chem. Soc. Rev.*, 2015, **44**, 5926–5940.
- 90 K. J. Griffith, K. M. Wiaderek, G. Cibin, L. E. Marbella and C. P. Grey, *Nature*, 2018, **559**, 556–563.
- 91 H. R. Xia, W. Zhang, S. K. Cao and X. D. Chen, *ACS Nano*, 2022, **16**, 8525–8530.
- 92 G. L. Hornyak, J. Dutta, H. F. Tibbals and A. Rao, *Introduction to nanoscience*, CRC press, 2008.
- 93 A. Nulu, V. Nulu and K. Y. Sohn, *J. Alloys Compd.*, 2022, **911**, 164976.
- 94 A. Nulu, Y. G. Hwang, V. Nulu and K. Y. Sohn, *Nanomaterials*, 2022, **12**, 3004.
- 95 C. Luo, X. Zhou, J. Ding, H. Xu, X. Wang, H. Yao, J. Tang and J. Yang, *Compos. Commun.*, 2021, **28**, 100941.
- 96 A. Nulu, V. Nulu and K. Sohn, *ChemElectroChem*, 2021, **8**, 1259–1269.
- 97 X. Han, Z. Zhang, H. Chen, L. Luo, Q. Zhang, J. Chen, S. Chen and Y. Yang, *J. Mater. Chem. A*, 2021, **9**, 3628–3636.
- 98 D. Guo, J. Wang, Y. Mai, P. Yang, J. W. Zhou, X. J. Xu, Y. Cheng, X. Y. Dai, Y. J. Gu and F. Z. Wu, *Ceram. Int.*, 2023, **49**, 5799–5807.
- 99 B. Long, Y. L. Zou, Z. Y. Li, Z. S. Ma, W. J. Jiang, H. Y. Zou and H. Y. Chen, *ACS Appl. Energy Mater.*, 2020, **3**, 5572–5580.
- 100 J. Ryu, J. Seo, G. Song, K. Choi, D. Hong, C. Wang, H. Lee, J. Lee and S. Park, *Nat. Commun.*, 2019, **10**, 2351.
- 101 J. Sung, N. Kim, J. Ma, J. H. Lee, S. H. Joo, T. Lee, S. Chae, M. Yoon, Y. Lee, J. Hwang, S. K. Kwak and J. Cho, *Nat. Energy*, 2021, **6**, 1164–1175.
- 102 Y. Son, J. Ma, N. Kim, T. Lee, Y. Lee, J. Sung, S. H. Choi, G. Nam, H. Cho, Y. Yoo and J. Cho, *Adv. Energy Mater.*, 2019, **9**, 1803480.
- 103 N. Mahmood, J. H. Zhu, S. Rehman, Q. Li and Y. L. Hou, *Nano Energy*, 2015, **15**, 755–765.
- 104 A. Nulu, V. Nulu and K. Y. Sohn, *ChemElectroChem*, 2021, **8**, 1214.
- 105 Y. Gao, L. Fan, R. Zhou, X. Du, Z. Jiao and B. Zhang, *Nano-Micro Lett.*, 2023, **15**, 222.
- 106 Y. Y. Zhu, J. Xie, A. Pei, B. F. Liu, Y. C. Wu, D. C. Lin, J. Li, H. S. Wang, H. Chen, J. W. Xu, A. K. Yang, C. L. Wu, H. X. Wang, W. Chen and Y. Cui, *Nat. Commun.*, 2019, **10**, 2067.
- 107 Y. Dong, C. Liu, F. Li, H. Jin, B. Li, F. Ding, Y. Yang, Z. Yang and F. Yuan, *ACS Appl. Electron. Mater.*, 2023, **6**, 816–827.
- 108 S. Y. Li, Z. M. Wu, L. Y. Du, Y. L. Shi, F. Tang, R. Li and Y. D. Jiang, *J. Mater. Sci.:Mater. Electron.*, 2018, **29**, 288–293.
- 109 M. Chen, Q. Zhou, J. Zai, A. Iqbal, T. Tsega, B. Dong, X. Liu, Y. Zhang, C. Yan, L. Zhao, A. Nazakat, E. SharelPeisan, C. Low and X. Qian, *Nano Res.*, 2021, **14**, 1004–1011.
- 110 C. Hou, Y. Qu, H. Xie, H. Tian, X. Wang, H. Lu, J. Wu, Y. Ma and J. Hou, *J. Mater. Sci.:Mater. Electron.*, 2024, **35**, 132.
- 111 X. Qu, X. Zhang, Y. Wu, J. Hu, M. Gao, H. Pan and Y. Liu, *J. Power Sources*, 2019, **443**, 227265.
- 112 Z. Wang, Q. Su, H. Deng, W. He, J. Lin and Y. Fu, *J. Mater. Chem. A*, 2014, **2**, 13976–13982.
- 113 M. Je, G. Song, S. Lee, H. Park, J. Kim and S. Park, *J. Mater. Chem. A*, 2023, **11**, 1694–1703.
- 114 F. Z. Zhang, W. Luo and J. P. Yang, *Chem. – Asian J.*, 2020, **15**, 1394–1404.
- 115 M. Ko, S. Chae, J. Ma, N. Kim, H. W. Lee, Y. Cui and J. Cho, *Nat. Energy*, 2016, **1**, 16113.
- 116 S. Yi, Z. L. Yan, X. D. Li, Z. Wang, P. P. Ning, J. W. Zhang, J. L. Huang, D. R. Yang and N. Du, *Chem. Eng. J.*, 2023, **473**, 145161.
- 117 C. M. Wang, X. L. Li, Z. G. Wang, W. Xu, J. Liu, F. Gao, L. Kovarik, J. G. Zhang, J. Howe, D. J. Burton, Z. Y. Liu, X. C. Xiao, S. Thevuthasan and D. R. Baer, *Nano Lett.*, 2012, **12**, 1624–1632.
- 118 M. Ko, S. Chae, S. Jeong, P. Oh and J. Cho, *ACS Nano*, 2014, **8**, 8591–8599.
- 119 N. Kim, S. Chae, J. Ma, M. Ko and J. Cho, *Nat. Commun.*, 2017, **8**, 812.
- 120 Q. Fang, S. Xu, X. Sha, D. Liu, X. Zhang, W. Li, S. Weng, X. Li, L. Chen, H. Li, B. Wang, Z. Wang and X. Wang, *Energy Environ. Sci.*, 2024, **17**, 6368–6376.
- 121 K. Cheng, S. Tu, B. Zhang, W. Wang, X. Wang, Y. Tan, X. Chen, C. Li, C. Li, L. Wang and Y. Sun, *Energy Environ. Sci.*, 2024, **17**, 2631–2641.
- 122 J. Chen, X. Fan, Q. Li, H. Yang, M. Khoshi, Y. Xu, S. Hwang, L. Chen, X. Ji, C. Yang, H. He, C. Wang, E. Garfunkel, D. Su, O. Borodin and C. Wang, *Nat. Energy*, 2020, **5**, 386–397.
- 123 C. K. Chan, R. Ruffo, S. S. Hong and Y. Cui, *J. Power Sources*, 2009, **189**, 1132–1140.
- 124 E. Peled and S. Menkin, *J. Electrochem. Soc.*, 2017, **164**, A1703–A1719.
- 125 C. Xu, F. Lindgren, B. Philippe, M. Gorgoi, F. Björefors, K. Edström and T. Gustafsson, *Chem. Mater.*, 2015, **27**, 2591–2599.
- 126 Z. Y. He, C. X. Zhang, Y. K. Zhu and F. Wei, *Energy Environ. Sci.*, 2024, **17**, 7.
- 127 Z. Chen, A. Soltani, Y. Chen, Q. Zhang, A. Davoodi, S. Hosseinpour, W. Peukert and W. Liu, *Adv. Energy Mater.*, 2022, **12**, 2200924.
- 128 G. R. Zheng, Y. X. Xiang, L. F. Xu, H. Luo, B. L. Wang, Y. Liu, X. Han, W. M. Zhao, S. J. Chen, H. L. Chen, Q. B. Zhang, T. Zhu and Y. Yang, *Adv. Energy Mater.*, 2018, **8**, 1801718.
- 129 H. Dang, Y. Y. Peng, L. Wang, X. Y. Li and F. Ran, *J. Energy Storage*, 2023, **74**, 109526.
- 130 F. L. Wang, G. Chen, N. Zhang, X. H. Liu and R. Z. Ma, *Carbon Energy*, 2019, **1**, 219–245.
- 131 W. L. An, B. A. Gao, S. X. Mei, B. Xiang, J. J. Fu, L. Wang, Q. B. Zhang, P. K. Chu and K. F. Huo, *Nat. Commun.*, 2019, **10**, 1447.
- 132 Z. L. Xu, K. Cao, S. Abouali, M. A. Garakani, J. Q. Huang, E. K. Heidari, H. T. Wang and J. K. Kim, *Energy Storage Mater.*, 2016, **3**, 45–54.



- 133 H. Tachikawa and A. Shimizu, *J. Phys. Chem. B*, 2006, **110**, 20445–20450.
- 134 C. Qi, S. Li, Z. Yang, Z. Xiao, L. Zhao, F. Yang, G. Ning, X. Ma, C. Wang, J. Xu and J. Gao, *Carbon*, 2022, **186**, 530–538.
- 135 B. Wang, J. Ryu, S. Choi, X. Zhang, D. Pribat, X. Li, L. Zhi, S. Park and R. Ruoff, *ACS Nano*, 2019, **13**, 2307–2315.
- 136 I. Son, J. Park, S. Kwon, S. Park, M. Rummeli, A. Bachmatiuk, H. Song, J. Ku, J. Choi, J. Choi, S. Doo and H. Chang, *Nat. Commun.*, 2015, **6**, 7393.
- 137 S. Pan, J. Han, Y. Wang, Z. Li, F. Chen, Y. Guo, Z. Han, K. Xiao, Z. Yu, M. Yu, S. Wu, D. Wang and Q. Yang, *Adv. Mater.*, 2022, **34**, 2203617.
- 138 Q. Ai, Q. Y. Fang, J. Liang, X. Y. Xu, T. S. Zhai, G. H. Gao, H. Guo, G. F. Han, L. J. Ci and J. Lou, *Nano Energy*, 2020, **72**, 104657.
- 139 J. B. Fang, S. Z. Chang, Q. Ren, T. Q. Zi, D. Wu and A. D. Li, *ACS Appl. Mater. Interfaces*, 2021, **13**, 32520–32530.
- 140 W. Wang, Y. Wang, W. B. Huang, M. Zhou, L. Z. Lv, M. Shen and H. H. Zheng, *ACS Appl. Mater. Interfaces*, 2021, **13**, 6919–6929.
- 141 Z. Cao, X. Y. Zheng, Y. Wang, W. B. Huang, Y. C. Li, Y. H. Huang and H. H. Zheng, *Nano Energy*, 2022, **93**, 106811.
- 142 Y. Bie, J. Yang, X. Liu, J. Wang, Y. Nuli and W. Lu, *ACS Appl. Mater. Interfaces*, 2016, **8**, 2899–2904.
- 143 W. Shi, H. Wu, J. Baucom, X. Li, S. Ma, G. Chen and Y. Lu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 39127–39134.
- 144 Y. Yu, C. Yang, Y. Jiang, J. Zhu, Y. Zhao, S. Liang, K. Wang, Y. Zhou, Y. Liu, J. Zhang and M. Jiang, *Small*, 2023, **19**, 2303779.
- 145 Q. Wang, M. Zhu, G. Chen, N. Dudko, Y. Li, H. Liu, L. Shi, G. Wu and D. Zhang, *Adv. Mater.*, 2022, **34**, 2109658.
- 146 J. Sun, Y. Li, L. Lv, L. Wang, W. Xiong, L. Huang, Q. Qu, Y. Wang, M. Shen and H. Zheng, *Adv. Funct. Mater.*, 2024, **34**, 2410693.
- 147 T. Mu, Y. Sun, C. Wang, Y. Zhao, K. Doyle-Davis, J. Liang, X. Sui, R. Li, C. Du, P. Zuo, G. Yin and X. Sun, *Nano Energy*, 2022, **103**, 107829.
- 148 T. W. Kwon, J. W. Choi and A. Coskun, *Chem. Soc. Rev.*, 2018, **47**, 2145–2164.
- 149 H. Wu, G. H. Yu, L. J. Pan, N. A. Liu, M. T. McDowell, Z. A. Bao and Y. Cui, *Nat. Commun.*, 2013, **4**, 1943.
- 150 R. Na, K. Minnici, G. Zhang, N. Lu, M. González, G. Wang and E. Reichmanis, *ACS Appl. Mater. Interfaces*, 2019, **11**, 40034–40042.
- 151 Y. P. Wang, A. Attam, H. G. Fan, W. S. Zheng and W. Liu, *Small*, 2023, **19**, 2303804.
- 152 A. M. Haregewoin, A. S. Wotango and B. J. Hwang, *Energy Environ. Sci.*, 2016, **9**, 1955–1988.
- 153 G. G. Eshetu and E. Figgemeier, *ChemSusChem*, 2019, **12**, 2515–2539.
- 154 J. M. M. de la Hoz, F. A. Soto and P. B. Balbuena, *J. Phys. Chem. C*, 2015, **119**, 7060–7068.
- 155 L. Z. Lv, Y. Wang, W. B. Huang, Y. Y. Wang, G. B. Zhu and H. H. Zheng, *Electrochim. Acta*, 2022, **413**, 140159.
- 156 Y. Yu, H. Koh, Z. Zhang, Z. Yang, A. Alexandrova, M. Agarwal, E. Stach and J. Xie, *Energy Environ. Sci.*, 2023, **16**, 5904–5915.
- 157 B. Philippe, R. Dedryvère, M. Gorgoi, H. Rensmo, D. Gonbeau and K. Edström, *Chem. Mater.*, 2013, **25**, 394–404.
- 158 A. V. Plakhotnyk, L. Ernst and R. Schmutzler, *J. Fluorine Chem.*, 2005, **126**, 27–31.
- 159 B. Philippe, R. Dedryvère, M. Gorgoi, H. Rensmo, D. Gonbeau and K. Edström, *J. Am. Chem. Soc.*, 2013, **135**, 9829–9842.
- 160 J. Shin, T. Kim, Y. Lee and E. Cho, *Energy Storage Mater.*, 2020, **25**, 764–781.
- 161 K. Tasaki, A. Goldberg, J. Lian, M. Walker, A. Timmons and S. Harris, *J. Electrochem. Soc.*, 2009, **156**, A1019–A1027.
- 162 H. Chen, A. Pei, D. Lin, J. Xie, A. Yang, J. Xu, K. Lin, J. Wang, H. Wang, F. Shi, D. Boyle and Y. Cui, *Adv. Energy Mater.*, 2019, **9**, 1900858.
- 163 J. Sun, J. Yan, F. Li, J. Li, J. Ma, G. Xu, P. Han, G. Hou, Y. Tang, S. Dong, J. Huang and G. Cui, *Adv. Mater.*, 2024, **36**, 2405384.
- 164 S. Lörger, K. Narita, R. Usiskin and J. Maier, *Chem. Commun.*, 2021, **57**, 6503–6506.
- 165 K. Asheim, P. E. Vullum, N. P. Wagner, H. F. Andersen, J. P. Maehlen and A. M. Svensson, *RSC Adv.*, 2022, **12**, 12517–12530.
- 166 X. Li, Z. Chen, X. Liu, L. Guo, A. Li, X. Chen and H. Song, *Adv. Funct. Mater.*, 2024, **34**, 2401686.
- 167 Z. Y. Wen, F. Wu, L. Li, N. Chen, G. Q. Luo, J. G. Du, L. Y. Zhao, Y. Ma, Y. J. Li and R. J. Chen, *ACS Appl. Mater. Interfaces*, 2022, **14**, 38807–38814.
- 168 Z. Yang, M. X. Jiang, C. Cui, Y. X. Wang, J. W. Qin, J. Wang, Y. X. J. Wang, B. G. Mao and M. H. Cao, *Nano Energy*, 2023, **105**, 107993.
- 169 Y. G. Zhang, N. Du and D. R. Yang, *Nanoscale*, 2019, **11**, 19086–19104.
- 170 Z. Xu, J. Yang, H. Li, Y. Nuli and J. Wang, *J. Mater. Chem. A*, 2019, **7**, 9432–9446.
- 171 W. Cai, Y. Yao, G. Zhu, C. Yan, L. Jiang, C. He, J. Huang and Q. Zhang, *Chem. Soc. Rev.*, 2020, **49**, 3806–3833.
- 172 Y. Tian, Y. An and B. Zhang, *Adv. Energy Mater.*, 2023, **13**, 2300123.
- 173 K. Leung, F. Soto, K. Hankins, P. Balbuena and K. Harrison, *J. Phys. Chem. C*, 2016, **120**, 6302–6313.
- 174 F. Soto, Y. Ma, J. de la Hoz, J. Seminario and P. Balbuena, *Chem. Mater.*, 2015, **27**, 7990–8000.
- 175 K. Schroder, J. Avarado, T. A. Yersak, J. C. Li, N. Dudley, L. J. Webb, Y. S. Meng and K. J. Stevenson, *Chem. Mater.*, 2015, **27**, 5531–5542.
- 176 P. Lu, C. Li, E. W. Schneider and S. J. Harris, *J. Phys. Chem. C*, 2014, **118**, 896–903.
- 177 Q. Zhao, S. Stalin and L. Archer, *Joule*, 2021, **5**, 1119–1142.
- 178 Y. X. Lin, Z. Liu, K. Leung, L. Q. Chen, P. Lu and Y. Qi, *J. Power Sources*, 2016, **309**, 221–230.
- 179 M. Islam and T. Bredow, *J. Phys. Chem. C*, 2009, **113**, 672–676.



- 180 X. Yi, Y. Guo, S. Pan, Y. Wang, S. Chi, S. Wu and Q. Yang, *Trans. Tianjin Univ.*, 2023, **29**, 73–87.
- 181 H. Iddir and L. Curtiss, *J. Phys. Chem. C*, 2010, **114**, 20903–20906.
- 182 S. Shi, Y. Qi, H. Li and L. Hector, *J. Phys. Chem. C*, 2013, **117**, 8579–8593.
- 183 X. Ma, X. Shen, X. Chen, Z. Fu, N. Yao, R. Zhang and Q. Zhang, *Small Struct.*, 2022, **3**, 2200071.
- 184 Y. Du and N. Holzwarth, *Phys. Rev. B:Condens. Matter Mater. Phys.*, 2007, **76**, 174302.
- 185 H. Liu and S. Yen, *J. Power Sources*, 2006, **159**, 245–248.
- 186 C. Li, L. Gu and J. Maier, *Adv. Funct. Mater.*, 2012, **22**, 1145–1149.
- 187 D. Lin, Y. Liu, W. Chen, G. Zhou, K. Liu, B. Dunn and Y. Cui, *Nano Lett.*, 2017, **17**, 3731–3737.
- 188 W. Li, G. Wu, C. Araújo, R. Scheicher, A. Blomqvist, R. Ahuja, Z. Xiong, Y. Feng and P. Chen, *Energy Environ. Sci.*, 2010, **3**, 1524–1530.
- 189 Y. Shi, L. Hu, Q. Li, Y. Sun, Q. Duan, Y. Jiang, Y. Xu, J. Yi, B. Zhao and J. Zhang, *Energy Storage Mater.*, 2023, **63**, 103009.
- 190 H. Duan, C. Wang, X. Zhang, J. Fu, W. Li, J. Wan, R. Yu, M. Fan, F. Ren, S. Wang, M. Zheng, X. Li, J. Liang, R. Wen, S. Xin, Y. Guo and X. Sun, *J. Am. Chem. Soc.*, 2024, **146**, 29335–29343.
- 191 R. Courtcastagnet, C. Kaps, C. Cros and P. Hagenmuller, *Solid State Ionics*, 1993, **61**, 327–334.
- 192 T. Wi, S. Park, S. Yeom, M. Kim, I. Kristanto, H. Wang, S. Kwak and H. Lee, *ACS Energy Lett.*, 2023, **8**, 2193–2200.
- 193 F. Han, J. Yue, X. Fan, T. Gao, C. Luo, Z. Ma, L. Suo and C. Wang, *Nano Lett.*, 2016, **16**, 4521–4527.
- 194 F. Altorfer, W. Buhrer, I. Anderson, O. Scharpf, H. Bill, P. Carron and H. Smith, *Physica B*, 1992, **180**, 795–797.
- 195 L. Ünal, V. Maccio-Figgemeier, L. Haneke, G. Eshetu, J. Kasnatscheew, M. Winter and E. Figgemeier, *Adv. Mater. Interfaces*, 2024, **11**, 2400024.
- 196 M. Zhang, H. Wang, A. Shao, Z. Wang, X. Tang, S. Li, J. Liu and Y. Ma, *Adv. Energy Mater.*, 2024, **14**, 2303932.
- 197 K. Xu, *Chem. Rev.*, 2014, **114**, 11503–11618.
- 198 N. Kim, Y. Kim, J. Sung and J. Cho, *Nat. Energy*, 2023, **8**, 921–933.
- 199 J. Xiao, F. Shi, T. Glossmann, C. Burnett and Z. Liu, *Nat. Energy*, 2023, **8**, 329–339.
- 200 S. J. Wu, B. Yu, Z. H. Wu, S. Fang, B. M. Shi and J. Y. Yang, *RSC Adv.*, 2018, **8**, 19894.
- 201 J. Sung, N. Kim, S. P. Kim, T. Lee, M. Yoon and J. Cho, *Batteries Supercaps*, 2022, **5**, e202200136.
- 202 L. Lu, Y. Lu, Z. Zhu, J. Shao, H. Yao, S. Wang, T. Zhang, Y. Ni, X. Wang and S. Yu, *Sci. Adv.*, 2022, **8**, eabm6624.
- 203 H. Zhao, N. Yuca, Z. Zheng, Y. Fu, V. Battaglia, G. Abdelbast, K. Zaghib and G. Liu, *ACS Appl. Mater. Interfaces*, 2015, **7**, 862–866.
- 204 C. Heubner, U. Langklotz and A. Michaelis, *J. Energy Storage*, 2018, **15**, 181–190.
- 205 H. Yu, L. Wang, Z. Zhang, Y. Li, S. Yang and X. He, *Adv. Funct. Mater.*, 2024, **34**, 2406966.
- 206 B. Lu, C. Q. Ning, D. X. Shi, Y. F. Zhao and J. Q. Zhang, *Chin. Phys. B*, 2020, **29**, 026201.
- 207 N. Yuca, I. Kalafat, E. Guney, B. Cetin and O. S. Taskin, *Materials*, 2022, **15**, 2392.
- 208 C. Wang, H. Wu, Z. Chen, M. T. McDowell, Y. Cui and Z. A. Bao, *Nat. Chem.*, 2013, **5**, 1042–1048.
- 209 Z. X. Xu, J. Yang, T. Zhang, Y. N. Nuli, J. L. Wang and S. I. Hirano, *Joule*, 2018, **2**, 950–961.
- 210 V. Vanpeene, A. King, E. Maire and L. Roué, *Nano Energy*, 2019, **56**, 799–812.
- 211 V. Vanpeene, J. Villanova, A. King, B. Lestriez, E. Maire and L. Roué, *Adv. Energy Mater.*, 2019, **9**, 1803947.
- 212 D. Liu, Y. Zhao, R. Tan, L. Tian, Y. Liu, H. Chen and F. Pan, *Nano Energy*, 2017, **36**, 206–212.
- 213 T. Munaoka, X. Yan, J. Lopez, J. To, J. Park, J. Tok, Y. Cui and Z. Bao, *Adv. Energy Mater.*, 2018, **8**, 1703138.
- 214 L. Hu, M. Jin, Z. Zhang, H. Chen, F. Ajdari and J. Song, *Adv. Funct. Mater.*, 2022, **32**, 2111560.
- 215 B. Han, Y. Yang, X. B. Shi, G. Z. Zhang, L. Gong, D. W. Xu, H. B. Zeng, C. Y. Wang, M. Gu and Y. H. Deng, *Nano Energy*, 2018, **50**, 359–366.
- 216 Y. M. Zhao, F. S. Yue, S. C. Li, Y. Zhang, Z. R. Tian, Q. Xu, S. Xin and Y. G. Guo, *InfoMat*, 2021, **3**, 460–501.
- 217 H. Chen, M. Ling, L. Hencz, H. Ling, G. Li, Z. Lin, G. Liu and S. Zhang, *Chem. Rev.*, 2018, **118**, 8936–8982.
- 218 H. Zhao, J. Li, Q. Zhao, X. Huang, S. Jia, J. Ma and Y. Ren, *Electrochem. Energy Rev.*, 2024, **7**, 11.
- 219 Q. He, J. Ning, H. Chen, Z. Jiang, J. Wang, D. Chen, C. Zhao, Z. Liu, I. Perepichka, H. Meng and W. Huang, *Chem. Soc. Rev.*, 2024, **53**, 7091–7157.
- 220 Z. Song, S. Chen, Y. Zhao, S. Xue, G. Qian, J. Fang, T. Zhang, C. Long, L. Yang and F. Pan, *Small*, 2021, **17**, 2102256.
- 221 W. Ren, K. Wang, J. Yang, R. Tan, J. Hu, H. Guo, Y. Duan, J. Zheng, Y. Lin and F. Pan, *J. Power Sources*, 2016, **331**, 232–239.
- 222 Z. Li, G. Wu, Y. Yang, Z. Wan, X. Zeng, L. Yan, S. Wu, M. Ling, C. Liang, K. Hui and Z. Lin, *Adv. Energy Mater.*, 2022, **12**, 2201197.
- 223 S. Park, H. Zhao, G. Ai, C. Wang, X. Song, N. Yuca, V. Battaglia, W. Yang and G. Liu, *J. Am. Chem. Soc.*, 2015, **137**, 2565–2571.
- 224 Y. Zhao, L. Yang, Y. Zuo, Z. Song, F. Liu, K. Li and F. Pan, *ACS Appl. Mater. Interfaces*, 2018, **10**, 27795–27800.
- 225 S. Gao, F. Sun, A. Brady, Y. Pan, A. Erwin, D. Yang, V. Tsukruk, A. Stack, T. Saito, H. Yang and P. Cao, *Nano Energy*, 2020, **73**, 104804.
- 226 F. Zhang, H. Xia, T. Wei, H. Li, M. Yang and A. Cao, *Energy Environ. Sci.*, 2024, **17**, 238–248.
- 227 B. Chen, D. Xu, S. Chai, Z. Chang and A. Pan, *Adv. Funct. Mater.*, 2024, **34**, 2401794.
- 228 Y. Yu, J. Zhu, K. Zeng and M. Jiang, *J. Mater. Chem. A*, 2021, **9**, 3472–3481.
- 229 Y. Cai, C. Liu, Z. Yu, W. Ma, Q. Jin, R. Du, B. Qian, X. Jin, H. Wu, Q. Zhang and X. Jia, *Adv. Sci.*, 2023, **10**, 2205590.
- 230 W. Zeng, L. Wang, X. Peng, T. Liu, Y. Jiang, F. Qin, L. Hu, P. Chu, K. Huo and Y. Zhou, *Adv. Energy Mater.*, 2018, **8**, 1702314.



- 231 Z. Sun, J. Zhu, C. Yang, Q. Xie, Y. Jiang, K. Wang and M. Jiang, *ACS Appl. Mater. Interfaces*, 2023, **15**, 12946–12956.
- 232 W. Ma, S. Wan, X. Cui, G. Hou, Y. Xiao, J. Rong and S. Chen, *Adv. Funct. Mater.*, 2023, **33**, 2212821.
- 233 S. Wang and M. Urban, *Nat. Rev. Mater.*, 2020, **5**, 562–583.
- 234 R. Narayan, C. Laberty-Robert, J. Pelta, J. Tarascon and R. Dominko, *Adv. Energy Mater.*, 2022, **12**, 2102652.
- 235 J. Yang, J. Li, Z. Yang, J. Liu, Y. Xiang and F. Wu, *ACS Appl. Energy Mater.*, 2022, **5**, 12945–12952.
- 236 Z. Ju, B. Zhang, T. Zheng, A. Marschilok, E. Takeuchi, K. Takeuchi and G. Yu, *Nano Lett.*, 2024, **24**, 6610–6616.
- 237 W. Li, Y. Li, J. Wang, S. Huang, A. Chen, L. Yang, J. Chen, L. He, W. Pang, L. Thomsen, B. Cowie, P. Xiong, Y. Zhou, G. Jang, D. Min, J. Byun, L. Xu, J. Huang, K. Roh, S. Kang, M. Liu, X. Duan and H. Park, *Energy Environ. Sci.*, 2024, **17**, 5387–5398.
- 238 W. Jang, S. Kim, Y. Kang, T. Yim and T. Kim, *Chem. Eng. J.*, 2023, **469**, 143949.
- 239 X. Wan, C. Kang, T. Mu, J. Zhu, P. Zuo, C. Du and G. Yin, *ACS Energy Lett.*, 2022, **7**, 3572–3580.
- 240 X. Wan, T. Mu, B. Shen, Q. Meng, G. Lu, S. Lou, P. Zuo, Y. Ma, C. Du and G. Yin, *Nano Energy*, 2022, **99**, 107334.
- 241 X. Lin, Y. Wen, J. Wang, S. Wang, X. Sun, H. Liu and X. Xu, *Green Chem.*, 2024, **26**, 2078–2086.
- 242 M. Jiang, P. Mu, H. Zhang, T. Dong, B. Tang, H. Qiu, Z. Chen and G. Cui, *Nano-Micro Lett.*, 2022, **14**, 87.
- 243 J. Kim, J. Choi, K. Park, S. Kim, K. Nam, K. Char and J. Choi, *Adv. Energy Mater.*, 2022, **12**, 2103718.
- 244 D. Rehnlund, F. Lindgren, S. Böhme, T. Nordh, Y. M. Zou, J. Pettersson, U. Bexell, M. Boman, K. Edström and L. Nyholm, *Energy Environ. Sci.*, 2017, **10**, 1350–1357.
- 245 Y. Yang, W. Yuan, X. Q. Zhang, Y. Z. Ke, Z. Q. Qiu, J. Luo, Y. Tang, C. Wang, Y. H. Yuan and Y. Huang, *Appl. Energy*, 2020, **276**, 115464.
- 246 W. D. Xue, X. Tian, Y. W. Xue, T. Peng, Y. Li and J. L. Sun, *ChemElectroChem*, 2019, **6**, 3039–3042.
- 247 X. H. Shen, L. Shao, Z. Y. Tian, Z. W. Hu and G. L. Cao, *Int. J. Electrochem. Sci.*, 2020, **15**, 9013–9023.
- 248 Z. Zhang, Z. Wang and X. Lu, *ACS Nano*, 2018, **12**, 3587–3599.
- 249 S. H. Moon, S. J. Kim, M. C. Kim, J. Y. So, J. E. Lee, Y. K. Shin, W. G. Bae and K. W. Park, *Mater. Chem. Phys.*, 2019, **223**, 152–156.
- 250 P. C. Zhu, D. Gastol, J. Marshall, R. Sommerville, V. Goodship and E. Kendrick, *J. Power Sources*, 2021, **485**, 229321.

