

Sustainable Energy & Fuels

Interdisciplinary research for the development of sustainable energy technologies

rsc.li/sustainable-energy



ISSN 2398-4902

PAPER

View Article Online
View Journal | View IssueCite this: *Sustainable Energy Fuels*,
2025, 9, 1165Thermogalvanic bricks: optimising large dimension
thermocells for air and water valorisation†Rebecca Houghton-James,^{‡a} Sireenya Mesawang,^{ID ‡a} Mark A. Buckingham,^{ID ab}
Robert Taylor,^{ID c} Patrick E. Phelan^d and Leigh Aldous^{ID *ae}

Thermogalvanic cells can potentially valorise the huge quantity of energy available as waste heat; using entropy-driven thermoelectrochemistry they can convert a thermal gradient into electricity. Most investigations exploit a thermal source (e.g. hot water, the human body, sunlight, electronics) *via* a heat exchanger (metal pipe, skin, housing, etc), combined with an unlimited heat sink (e.g. pumped cold water). Limited studies have used ambient air as the heat sink. This study is believed to be the first to explore using air as both the thermal source and heat sink. It compares thermogalvanic cell performance when using water–water and air–air as the thermal energy sources and sinks, respectively, for devices with relatively large physical dimensions (25 to 100 mm wide). Gelation improved power output under both scenarios, due to enhanced thermal isolation of the electrodes; power decreased with increasing width in the water–water setup, but power increased with increasing width for air–air harvesting. Water–water yielded higher power overall, yet the air–air system operated passively and could be further optimised for real-world applications, *i.e.* as thermogalvanic bricks or panels in building materials.

Received 27th October 2024
Accepted 9th December 2024

DOI: 10.1039/d4se01498g

rsc.li/sustainable-energy

Introduction

Almost every activity undertaken by humanity produces low-grade (<100 °C) thermal energy. Industrial processes, refrigerators, air conditioners, data centres, and buildings create a vast source of continuous—yet unexploited—waste energy.¹ Thermogalvanic cells (also known as thermocells) are two-electrode electrochemical devices which exploit entropy differences between two redox states in a redox couple (ΔS_{rc}). These devices valorise a temperature gradient (ΔT e.g. from a waste heat source and a heat sink, T_{hot} and T_{cold} , respectively) into electrical energy *via* a generated potential difference (ΔV , or V_{OC}), similar to the classic solid-state Seebeck effect.² A cartoon of the process is shown in Fig. 1a. Fundamentally, these properties are intrinsically linked:

$$V_{ocp} = S_e \Delta T = \frac{\Delta T \Delta S_{rc}}{nF} \quad (1)$$

^aDepartment of Chemistry, King's College London, London, UK. E-mail: leighaldous@ntu.edu.tw^bDepartment of Materials, The University of Manchester, Manchester, UK^cMechanical and Manufacturing Engineering, UNSW Sydney, Sydney, Australia^dSchool for Engineering of Matter, Transport and Energy, Arizona State University, Tempe, USA^eDepartment of Chemical Engineering, National Taiwan University, Taipei, Taiwan† Electronic supplementary information (ESI) available: Exemplar raw data; temperature study; and tabulated results. See DOI: <https://doi.org/10.1039/d4se01498g>

‡ These authors contributed equally to this work.

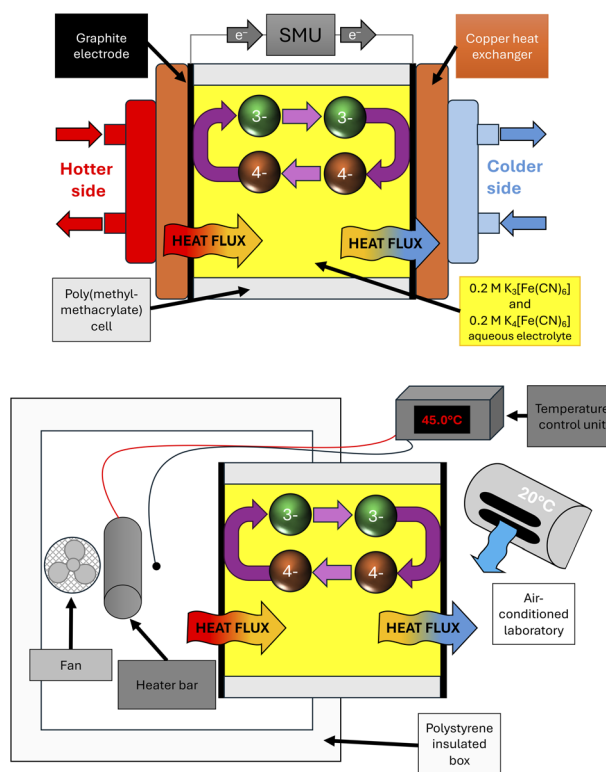


Fig. 1 Annotated cartoons of the (top) water–water thermogalvanic set-up and (bottom) air–air thermogalvanic set-up. Photographs of the setup can be found in the ESI.†



where S_e is the thermogalvanic Seebeck coefficient ($V K^{-1}$), n is the number of electrons transferred and F the Faraday constant. The S_e value is typically constant, and is typically dominated by the (de)solvation changes that occurs when the redox couple changes oxidation state.³

The S_e directly correlates with the V_{ocp} , and the V_{ocp} represents the driving force behind the flow of current between the two electrodes. The highest voltage, V_{ocp} , and the highest current generated (the short circuit current density, j_{sc}) can be modelled via Butler-Volmer kinetics, as summarised by eqn (2);⁴

$$j_{sc} = Fk_{agg} \left(C_{ox}^{\alpha_a} C_{red}^{\alpha_c} \right) \cdot \left\{ \exp \left[\frac{0.5 V_{ocp} \alpha_c F}{RT_{cold}} \right] - \exp \left[- \frac{0.5 V_{ocp} \alpha_c F}{RT_{cold}} \right] \right\} \quad (2)$$

where C represents the concentration, α represents the charge transfer coefficients (for the anodic, a , and cathodic, c , processes), R is the universal gas constant, and k_{agg} is an aggregated, limiting kinetic constant (typically represented by the limiting rate of mass transport and/or electron transfer).

The maximum electrical power, P_{max} , produced by thermocells typically obeys the relationship $P_{max} = 0.25 V_{ocp} j_{sc}$.² Under typical conditions, F , C , α and R can all be taken as constants. Removing these constants from eqn (2) (discussed in more detail in the ESI† section), then we can see that;

$$P_{max} \propto k_{agg} \cdot \frac{\Delta T^2}{T_{cold}} \propto k_{agg} \cdot \frac{(T_{hot} - T_{cold})^2}{T_{cold}} \quad (3)$$

Here, the ΔT between the two electrodes is therefore the most critical parameter in achieving higher power; secondary is reducing the mass transfer resistance and electron transfer resistance (to increase k_{agg}).

For these reasons, during investigations ‘semi-infinite’ heat sources and heat sinks are typically used to heat exchange with the two thermocell electrodes. These precisely control the applied T_{hot} and T_{cold} and therefore the ΔT , by using *e.g.* Peltier modules⁵ and water circulation.⁶ These mirror some limited ‘real life’ applications, such as in data centre cooling.⁷ However, unlimited quantities of pumped ‘hot’ and ‘cold’ water is only encountered in limited areas.

There are several examples where air forms one side of the heat exchange of a thermocell, *e.g.* during body heat harvesting^{8,9} (air = cold side), hot water pipe valorisation¹⁰ (cold side), solar irradiation^{11,12} (hot side), battery heat harvesting¹³ (cold side), and nighttime radiative cooling¹⁴ (cold side). Other cooling mechanisms can be applied, such as evaporation.¹³ However, we are unaware of prior reports of thermogalvanic cell heat exchange performed with both sides exposed to air.

Large scale (electro)chemical devices are being considered for incorporation into building structures, such as phase change materials¹⁵ and batteries.¹⁶ Other technologies have significant promise, such as electrochemical air conditioning,¹⁷ and electrochromic materials controlling UV/vis/IR-interactions.¹⁸ Building structures typically have a ΔT across dividers, reaching as high as $\Delta T = 70 K$ in loft compartments.¹⁹

Here we set out to compare a ‘model’ heat source/sink (*i.e.* water bath-regulated copper heat exchangers) *vs.* a ‘real’

scenario using air as both heat source and sink. This was undertaken using relatively large thermocells containing liquid or gelled electrolyte comprised of 0.2 M $K_3[Fe(CN)_6]$ and 0.2 M $K_4[Fe(CN)_6]$. The cell width ranged from 25 mm to 100 mm. These were chosen because ceiling panels are typically 8–25 mm thick, pre-fab wall panels 25–100 mm thick, and the standard UK metric house brick is 102.5 mm wide; all represent real-life scenarios where two air bodies are separated by bricks or panels into different temperature areas, and thus could potentially be used for ambient temperature difference valorisation in to electricity.

Experimental

Electrolyte and gelation

The electrolyte consisted of 0.2 M $K_3[Fe(CN)_6]$ and 0.2 M $K_4[Fe(CN)_6]$ dissolved in ultrapure water. Gelation was achieved by adding sodium poly(acrylate) powder (supplied as ‘Instant Artificial Snow Magic Powder’, Amazon, UK) at fixed weight per volume values (wt/v%). Typically, only *ca.* 5 minutes was required for the system to equilibrate. Because the sodium poly(acrylate) powder was slightly acidic, for safety reasons²⁰ 0.05 M potassium carbonate (K_2CO_3) was also added to the electrolyte. This was not found to influence the performance of the electrolyte but was added as a precaution due to the risks associated with potential HCN evolution from heated, acidified ferri/ferrocyanide solutions.²⁰

The sodium poly(acrylate) dry powder rapidly expands and gels liquid media (within seconds for pure water, and within minutes for electrolyte solutions). Fig. S1† shows photographs of solutions containing 0.2 M $K_3[Fe(CN)_6]$, 0.2 M $K_4[Fe(CN)_6]$ and 0.05 M K_2CO_3 , as a function of wt/v% of added sodium poly(acrylate) powder. After 5 minutes, the solutions were inverted, to determine whether they passed the inversion test.²¹ For 0.5 and 1 wt/v%, the swollen gel material tended to sink to the bottom of the electrolyte, causing a different phase. For 1.5 to 2.5 wt/v%, the swollen gel filled the electrolyte (*e.g.* was relatively homogeneously dispersed throughout the phase, but as a heterogeneous *quasi-solid* phase) but the mixture could still flow, making it a slurry. For 3.0 wt/v% or above, the mixture formed a soft material that could resist flow due to gravity, making it a genuinely gelled $K_{3/4}[Fe(CN)_6]$ phase.

Thermogalvanic bricks

Four thermocells with electrode separations of 25, 50, 75 and 100 mm were custom cut and built from extruded poly(methyl methacrylate) (PMMA, PerspexSheet.uk, UK). Example photographs are shown in Fig. S2.† They comprised 4 pieces, with dimensions 5 mm × 50 mm × y mm (where $y = 25$ –100 mm). Three pieces were secured into a u-shape with adhesive (EVO-STIK Tensol 70 Two Component Cement, RS Components, UK), with a fourth piece placed on top as a detachable lid. This fourth piece had holes drilled into it for introducing thermistors, and if not used were sealed with Blu-Tack. This arrangement resulted in an external cross-sectional area of 55 × 55 mm, with an internal area of 45 × 45 mm. Previously



characterised⁴ graphite electrodes (amorphous graphite – 99.5% pure graphite gasket foil, 1 mm thick, Xiaochengshop, China) were cut to shape, and secured to the two faces of the thermocell using the same adhesive. These cells were reused for multiple measurements; between measurements the graphite electrode surfaces were resurfaced by rubbing with low-lint tissues (Kimtech Kimwipes).

When filling the thermocell, some irreproducible results were initially observed due to capillary forces sucking up the electrolyte between the lid and the electrode, resulting in variable electrode surface areas. For this reason, the cells were filled with electrolyte which took the depth up to *ca.* 85% of the maximum depth. Due to slight differences in the cells, different filling proportions were required, but the volume added each time was kept constant. This then resulted in consistent electrode surface areas being wetted by the electrolyte, which was physically measured by callipers for the specified volumes of electrolyte. The volume used and resulting geometric surface area of the electrodes are summarised in the ESI in Table S1.†

Thermoelectrochemical setup using the water–water temperature-controlled apparatus

Water heating and cooling of the thermocells was undertaken using the setup shown by the schematic in Fig. 1a (photographs are shown in Fig. S3†). The graphite electrodes were placed in thermal contact with copper-based CPU heat exchangers (Floratek CPU Water Cooling Block with 50 mm Copper Base, Amazon, UK) using thermal paste (multicomp Silicone Heat Transfer Compound, Farnell, UK). The copper blocks had their temperature controlled by water flow from two thermostatic circulator baths (a Grant Optima refrigerated TX150-R2 circulating bath, and Grant Optima TX150-ST5 Heated Circulating Bath, respectively, both from Grant Instruments (Cambridge) Ltd, UK, both operating with <0.1 °C accuracy) with an applied temperature difference, ΔT , of 20 K (T of cold electrode = 20 °C).

Thermogalvanic measurement parameters for the water–water apparatus

Thermogalvanic measurements were recorded using a source measure unit (SMU, Keysight B2901A Precision Source/Measure Unit, Keysight, UK). Measurements initially followed the sequence of potentials technique described in detail elsewhere² to generate a 5-point power curve, with later measurements using 3-point (V_{ocp} , current at $0.5V_{\text{ocp}}$ for P_{max} , and j_{sc}). Each value was generated by continuously measuring for 600 s (0.1 s per data point), then averaging the 3000 data points from 300.1 to 600.0 seconds.

Thermoelectrochemical setup using the air–air temperature-controlled apparatus

The air–air apparatus was constructed around an insulated box (34L HoCo Polystyrene Box, I.D. 420 mm × 320 mm × 250 mm, wall thickness 24 mm, Amazon, UK). A schematic is included above (Fig. 1b), and photographs are included in the ESI (Fig. S2†). A square hole was cut in the centre of the largest wall; the thermocell was balanced half-in and half-out of the box. The

inside of the insulated box was heated to 45 °C using a heating element (Dimplex 1 foot Tubular Heater with Built in Thermostat, Argos, UK); because the hysteresis in the built-in thermostat was too large, the power to this heater was controlled *via* a temperature controller (Elitech Temperature Controller Thermostat with an NTC Probe, ebay.co.uk, UK). An electric fan (Pro-Elec 6" USB Mini Desk Fan, Farnell, UK) was used within the box for heat dispersal. The fan was positioned at the back wall, facing towards the insulated box wall holding the thermocell, with the face of the fan at a distance of 23.3 cm away from the thermocell wall. The tubular heater was positioned in between the fan and the thermocell, at an average distance of 11.1 cm from the wall housing the thermocell. The temperature controller thermistor was positioned equidistant between the heater and the thermocell.

The 'cold' side in the air–air system was the ambient air in an air-conditioned laboratory. The air conditioning was set to 20 °C. For preliminary experiments no forced convection was applied this side of the thermocell; later experiments aimed to match the convection occurring inside the box by positioning a second, identical electric fan (Pro-Elec 6" USB Mini Desk Fan, Farnell, UK) facing towards the thermocell cold electrode. The face of this fan was at a distance of 20.0 cm away from the insulated wall holding the thermocell.

Thermogalvanic measurement parameters for the air–air apparatus

Thermogalvanic measurements for the air heating set-up were performed using a potentiostat (Autolab PGSTAT101, Metrohm, UK) in two-electrode mode, and using potentiostatic and amperometric procedures (using NOVA 2.0, Metrohm, UK). The instrument was set to measure a data point every 0.1 s, to match the measurements made on the water–water setup using a SMU. Because equilibration was relatively slow, and there were some repeating fluctuations in the temperature (due to a hysteresis in the temperature control unit) and some non-repeating fluctuations (*e.g.* laboratory door opening), the two-point sequence of potentials technique described in detail elsewhere² was used, where the circuit potential (V_{ocp}) and short circuit current density (j_{sc}) were initially recorded for roughly 43 000 s and 40 500 s respectively (total run time *ca.* 23 h). This allowed averaging of data over extended periods of time (*ca.* 400 000 data point each) and ensured that genuine 'steady state' measurements were performed.² The maximum power (P_{max}) was given by $0.25V_{\text{ocp}}j_{\text{sc}}$. During subsequent repeat measurements, it was determined that shorter time period measurements (*ca.* 6000 seconds steady state data, or 60 000 data points) produced data with equivalent quality. Exemplar raw data is shown in Fig. S4.†

Temperature monitoring and temperature profile generation using thermistors

Temperature was recorded during every air–air experiment using a battery-powered data logger (GP2 Data Logger and Controller, Delta-T Devices Ltd, UK) with six thermistors (MT3 sealed catheter-style 2K thermistor probes in flexible nylon tubes, Delta-T Devices Ltd, UK). These were used to record the



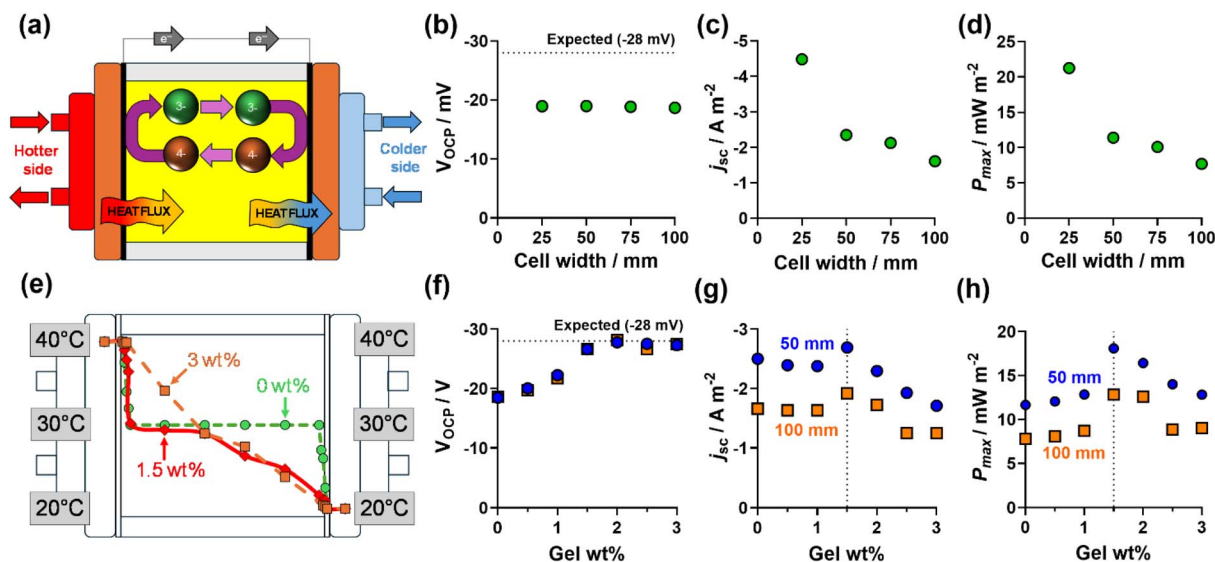


Fig. 2 The top row shows (a) a cartoon of the water–water heat exchanged thermogalvanic cell; and the measured (b) open-circuit voltage (V_{OC}), (c) short-circuit current density (j_{SC}) and (d) the maximum power density (P_{max}) measured for $\text{K}_3/\text{K}_4[\text{Fe}(\text{CN})_6]$ as a function of cell width when exposed to $\Delta T = 20 \text{ K}$ ($T_{\text{cold}} = 20^\circ \text{C}$, $T_{\text{hot}} = 40^\circ \text{C}$). The bottom row summarises the results from varying the wt/v% of polyacrylate gel added to $\text{K}_3/\text{K}_4[\text{Fe}(\text{CN})_6]$, showing (e) the temperature gradient measured across the 100 mm cell for 0, 1.5, and 3 wt% (data tabulated in the ESI†), and the measured (f) V_{OC} , (g) j_{SC} and (h) P_{max} for 50 mm (blue circles) and 100 mm (orange squares) wide thermocells.

air temperature inside and outside of the box, the air temperature *ca.* 1 mm away from the hot and cold electrodes, and the electrolyte temperature *ca.* 1 mm away from the hot and cold electrodes. They were also used to generate the temperature profiles across the cells shown later in Fig. 2e and 4e.

For standard temperature monitoring, data logging was started for *ca.* 20 minutes before the heating was switched on, and left measuring for 1 h after the heating was switched off. The thermistors were connected to channels 1–6, with channel 6 = suspended in the centre of the thermostatic heated box; 5 = suspended in the ‘hot’ air *ca.* 1 mm away from the hot electrode; 4 = immersed in the electrolyte and *ca.* 1 mm away from the hot electrode; 3 = immersed in the electrolyte and *ca.* 1 mm away from the cold electrode; 2 = suspended in the ‘cold’ air and *ca.* 1 mm away from the hot electrode; and 1 = suspended in the air of the air conditioned laboratory, *ca.* 20 cm away from the heated box and thermocell. Exemplar data showing *ca.* 19 h recordings for two thermocells is shown in the ESI, Fig. S5.†

Results and discussion

The full experimental details of this study are included above in the Experimental section, but briefly the relatively large electrodes (*ca.* 30 cm^2 each) were comprised of previously characterised amorphous graphite electrodes.⁴ The water heating/cooling apparatus has also been previously characterised^{22,23} and a cross-section schematic is shown in Fig. 2a. The air apparatus is novel to this study and comprised a tubular resistive heater and fan-convection, while the heat sink was the ambient air in an air-conditioned laboratory (schematic shown in Fig. 4a).

First the water heat source/water heat sink system (henceforth referred to as water–water) was explored by; (i)

characterising the effect of cell width; (ii) the effect of increasing gelation of the electrolyte, at various cell widths; and (iii) investigating the temperature gradient across the cell.

Fig. 2b shows the V_{OC} measured *vs.* cell width; the V_{OC} was largely constant, but also consistently below the expected value of *ca.* -28 mV (for a Seebeck coefficient of $-1.4 \text{ mV K}^{-1,4,5}$) corresponding to a lower ‘observed’ temperature difference (*ca.* 13.5 K) at the interior surface of the graphite electrodes than the applied one (20 K). The significant (*ca.* 33%) loss of temperature gradient is likely caused by significant parasitic convection though the liquid electrolyte in this physically large cell;²⁴ under extreme circumstances thermal convection can exceed $1000 \text{ W m}^{-2} \text{ K}^{-1}$.²⁵ The j_{SC} decreased with increasing cell width (Fig. 2c) due to increasing mass transport resistance in the thermocell;^{4,26} this resulted in the P_{max} following the same downward trend (Fig. 2d). One simple method of reducing natural convection and improving the temperature gradient in thermocells is by gelling the liquid.^{21,27–31}

Gelation of the electrolyte was achieved by adding polyacrylate powder. The effect of polyacrylate powder on the electrolyte is shown visually in Fig. S1;† typically $>0 \text{ wt\%}$ to $1.5 \text{ wt\%}$ (w/v) formed a heterogeneous mixture, $1.5\text{--}2.5 \text{ wt\%}$ formed a homogeneous slurry, and $3 \text{ wt\%}$ or more resulted in it passing the ‘inversion test’,²¹ confirming the electrolyte was fully gelled.

The V_{OC} was measured in the cells as a function of gel wt/v% (Fig. 2f) and increasing gelation increased V_{OC} (by suppressing the parasitic convection), such that $\geq 1.5 \text{ wt\%}$ resulted in the ΔT experienced between the two electrodes being equal to the applied ΔT . However, just as increasing cell width decreased j_{SC} due to increased mass transport resistance, increasing gelation is known to do the same.³² Therefore, as shown in Fig. 2g, increased thermal resistance increased j_{SC} to a maximum at



1.5 wt%, but beyond this point thermal resistance was already sufficient and increasing gel content only hindered mass transport, thus having a negative effect upon the current generated. As shown in Fig. 2h, $P_{\max}$ clearly peaked at 1.5 wt% sodium poly(acrylate), and this was observed for all 4 cell widths. A similar optimum was previously observed when measuring the genuine thermogalvanic conversion efficiency (in a much smaller thermocell).³³

Thermistors were next used to probe the temperature gradient across the 100 mm cell containing 0, 1.5, and 3.0 wt% polyacrylate, with the results plotted in Fig. 2e. For the liquid electrolyte, 'hot' and 'cold' electrolyte was only found immediately adjacent to the electrolyte surface, and 94% of the width of the electrolyte was homogeneously 30 °C. For 3.0 wt%, the gelled electrolyte had a linear temperature gradient across the entire width of the cell. For the 'slurry' caused by 1.5 wt%, the 'hot' electrolyte was only within 2.5 mm of the electrode surface before dropping to *ca.* 30 °C (similar to the un-gelled electrolyte), while the colder half of the cell had a clear temperature gradient (similar to the fully gelled system), suggesting the majority of the polymeric slurry accumulated in the colder electrolyte and frustrated heat transfer in primarily just this half of the cell.

The above discussion related to thermistors submerged to half the depth. However, some temperature variation as a function of depth was observed (due to convection and radiation), so 2D IR imaging was utilised to investigate this. Fig. 3 shows photographs (top) and IR camera images (bottom) for side-on and top-down views of cells containing 0 and 3 wt% polyacrylate (see Fig. S7† for 1.5 wt% photographs). The liquid (0 wt%) electrolyte had a homogeneous temperature across the middle of the electrolyte (matching the thermistor results in Fig. 2e). Conversely, the gelled electrolyte showed a steady temperature gradient and no evidence of convection, with the side-on and top-down images having similar temperature profiles (again, matching the results in Fig. 2e).

Having established the underlying trends when semi-infinite heat sources and sinks are employed, next, a novel air-based heat source and sink system (or 'air-air') was explored. This study followed the same trend, namely (i) the effect of cell width for liquid (0 wt/v% gel) electrolyte; (ii) the effect of polyacrylate wt/v% loading, and finally (iii) temperature gradient characterisation.

Fig. 4a displays a cartoon of the setup, whereby the hot air side was created using a heater and a fan; the cold side consisted of an air-conditioned laboratory without any forced convection (full details in Experimental and ESI†). Fig. 4b compares the measured V_{OCP} for the air-air set-up (purple circles) with the water-water setup above (orange squares). The air-air system applied a greater ΔT of 25 K ($T_{\text{ambient air}} \approx 20$ °C, $T_{\text{hot air}} = 45$ °C, due to enhanced stability) hence V_{OCP} should be *ca.* −35 mV. However, due to the different thermal resistance between air and electrolyte, this could not be achieved. As cell width increased, the thermal resistance of the thermocell increased; therefore, the experienced ΔT increased, and V_{OCP} increased (Fig. 4b). Consequently, j_{sc} (Fig. 4c) and $P_{\max}$ (Fig. 4d) also increased, but showed completely opposite trends to the same cells measured in the water-water setup. Furthermore, in the air-air setup, the experienced ΔT measured at the thermogalvanic brick electrodes (derived from the S_e values) reached only *ca.* 20% of the applied ΔT , resulting in overall power *ca.* 50-fold smaller than the water-water system.

Since the air-air set-up was struggling to establish reasonable ΔT values, gelation of the electrolyte was once again explored. As shown in Fig. 4f, increasing gel content in the 100 mm cell increased V_{OCP} to a plateau of *ca.* −28 mV (NB: *ca.* −35 mV was the ideal value). As seen in Fig. 4e, once again a peak in j_{sc} is observed at the optimum compromise between frustrated heat transfer and frustrated mass transfer. However, because more gel was required before this was achieved (2.5 wt%), the system suffers from lower current by this optimum point. As seen in Fig. 4h, although $P_{\max}$ has clearly

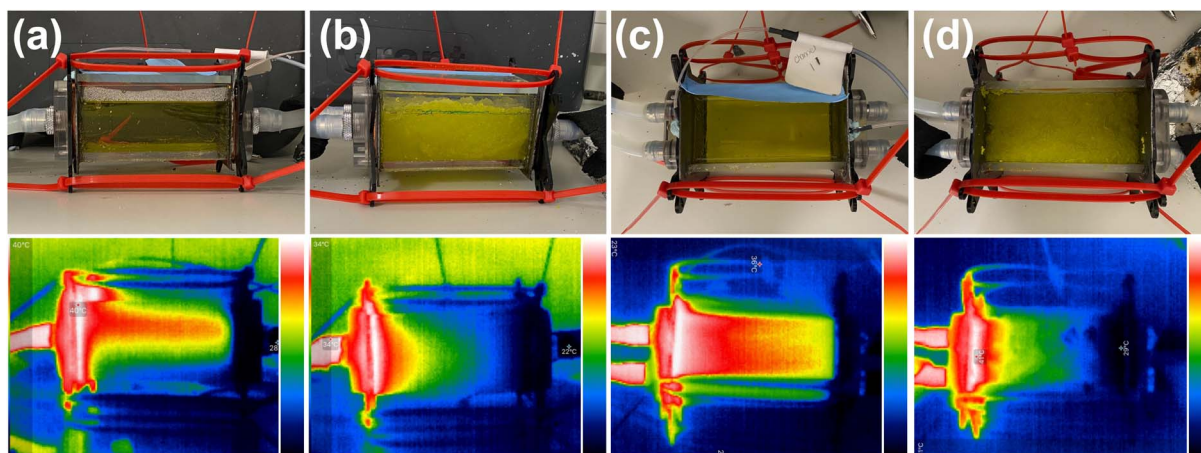


Fig. 3 Photographs (top) and IR images (bottom) of (a and c) liquid electrolyte and (b and d) 3 wt% gelled electrolyte in (a and b) side-on profile and (c and d) top-down profile. For the IR images the temperature scales are colour-coded going from white (*ca.* 40 °C) to dark blue (*ca.* 20 °C). Note that the IR images only capture the surface values and not necessarily bulk values. Furthermore the angular emissivity was not calibrated for the wide variety of different surfaces, hence the values should be taken as semi-quantitative (up to ± 5 °C). The IR images still accurately highlight relative temperature differences across surfaces composed of the same material.



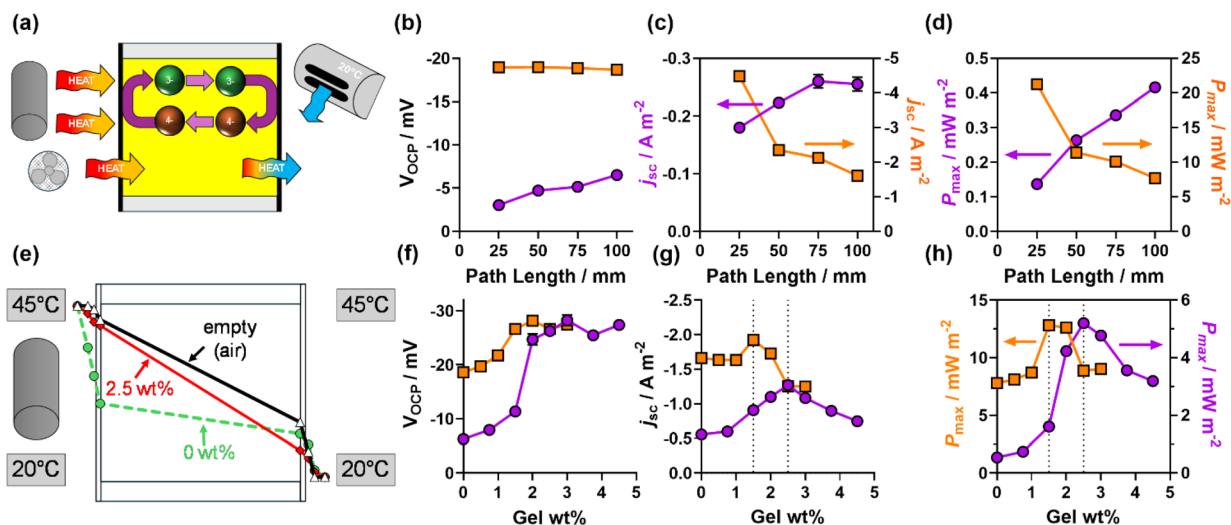


Fig. 4 The top row shows (a) a cartoon of the air–air heat exchange setup, and a comparison of the recorded (b) V_{OCP} , (c) j_{SC} and (d) P_{max} for liquid $\text{K}_3/\text{K}_4[\text{Fe}(\text{CN})_6]^{3-/4-}$ in the 100 mm cell using the water–water setup from earlier (orange squares) vs. the air heating and cooling (purple circles). The bottom row shows (e) the temperature gradient measured across the 100 mm cell for the electrolyte containing 0 wt% and 2.5 wt% polyacrylate and an empty cell (quantitative data tabulated in the ESI†), and (f–h) a comparison of the outputs for the water–water vs. air–air setup as a function of wt/v% polyacrylate gel added to $\text{K}_3/\text{K}_4[\text{Fe}(\text{CN})_6]$. Please note the water–water setup used an applied $\Delta T = 20$ K ($T_{\text{cold}} = 20^\circ\text{C}$, $T_{\text{hot}} = 40^\circ\text{C}$), while the air–air used $\Delta T = 25$ K ($T_{\text{cold}} = 20^\circ\text{C}$, $T_{\text{hot}} = 45^\circ\text{C}$). In (f) and (g) the dotted lines are used to highlight peak performance areas. Also shown in (e) are the temperature gradients measured using thermistors for an empty (air-filled) 100 mm cell, and the same cell filled with electrolyte (containing 0 wt% or 2.5 wt% polyacrylate). Full data for these and polyacrylate loadings up to 4.5 wt% are tabulated in the ESI.†

significantly increased at the optimum gelation point, the highest air–air P_{max} is only 40% of the highest in the water–water system, despite having a 20% higher applied ΔT value.

The measured temperature gradient across the cell is visualised in Fig. 4e; the electrolyte-filled cell (dashed green line) had a negligible temperature gradient across the width of the electrolyte, as expected due to extensive convective heat transfer, and slow heat exchange with the surroundings (especially at the colder electrode). As the gel wt% was increased, the steepness of this temperature gradient across the width of the electrolyte increased, reaching a maximum steepness at 2.5 wt% (shown as the red line); beyond 2.5 wt% the temperature gradient remained unchanged. As a comparison, the empty cell (*i.e.* filled with only air) was also measured. Here a temperature gradient also formed across the width of the air inside the empty cell, and interestingly the gradient was less steep than the ≥ 2.5 wt% gelled electrolyte, despite waiting for the systems to reach equilibrium. This demonstrates that the gelled electrolyte-containing ‘brick’ was a superior insulator than air.

Another factor that was clear from Fig. 4e was that the temperature gradient across the cold-side electrode was more significant than the hot-side; therefore, the colder electrode was dissipating the thermal energy to its surroundings slower than the heated side. This is to be somewhat expected, since the heated side employed a fan to improve thermal homogeneity throughout the heated box. To further probe this imbalance, a study was performed as a function of (i) cell width, and (ii) with and without forced air convection at the colder electrode. As shown in Fig. 5a, power increased as cell width increased (NB: the opposite trend of the water–water setup, showing the

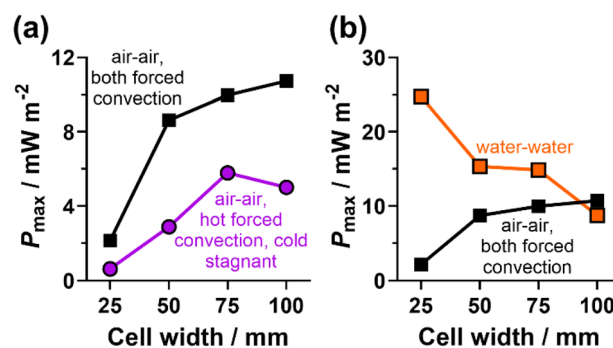


Fig. 5 Showing P_{max} for various cells widths comparing (a) air–air setup with different convection at the colder electrode, and (b) comparing the air–air setup vs. water–water. All other conditions as per Fig. 2 and 4.

different cell design compromises required). Furthermore, air convection at the colder electrode successfully increased V_{OCP} (to -33 mV) and *ca.* doubled the power (full data tabulated in the ESI). Fig. 5b highlights the contradictory trend observed in the water–water and air–air setups; power decreases with width in the water–water setup because mass transfer is the critical limiting factor, whereas in the air–air scenario, lack of thermal resistance through the electrolyte is the critical limiting factor, and is only overcome at a cell width of *ca.* 100 mm. In the future, other methods of thermal resistance should be explored, beyond gelation and electrolyte width between electrodes.

The observations made in this study apply to thermogalvanic cells in horizontal arrangements, where the hot and cold



electrodes are parallel to both each other and gravity (as might be found in a wall). Although it was not investigated herein, natural convection will play a larger role in vertical arrangements than horizontal arrangements due to temperature-induced density differences.^{34,35} For example, it is expected that if the cold electrode sits above the hot electrode, the gelation-based enhancements observed here should be more pronounced. Conversely, in a vertical arrangement where the hot electrode sits physically above the cold electrode, buoyancy should restrict flow with stagnant layers of higher temperature forming at the top.³⁴ Thus, in a hot-over-cold arrangement, such gelation effects are expected to be less significant.

This study demonstrates generally higher power for the water–water setup and could seem negative towards the air–air system of thermogalvanic energy harvesting. However, the water–water system represents a far less common real-life scenario (especially the concept of having unlimited flowing ‘cold’ water), it requires heat exchangers, and it consumes energy *via* the water-pumping process. Conversely, the air–air system operates entirely passively, and represents a scenario found all over the world, such as the temperature difference across ceiling tiles and building walls. It is also still unoptimised, *e.g.* fins could be applied to the colder electrode surface to aid in heat dissipation.

The results demonstrate that for air–air ambient valorisation by thermogalvanic devices, large dimensions and gelation are essential. The application of optimised gels (*e.g.* directed ion transfer channels³⁶) is expected to boost power further. Furthermore, these large-dimension insulating thermocells could be applied to air–water arrangements, such as on the outside of pipes which are passing through an environment with a dissimilar temperature to the interior. Sustaining a large temperature gradient across the thermogalvanic cell will ensure that the pipe contents are insulated, and electricity is generated *via* this temperature difference.

A final caveat is that more appropriate electrolytes are also required before application. Highly concentrated $K_{3/4}[Fe(CN)_6]$ devices need to be handled with care under research conditions due to the risk of toxic HCN evolution.²⁰ In large scale applications, fires, UV exposure, natural disasters *etc.* would increase the risk. This also extends to electrolytes based upon toxic and/or scarce metals (such as Co), and polyhalides (such as I^-/I_3^-). Substitution with either abundant natural elements (*e.g.* Fe) and sustainable ligands,²² or sustainable organic redox couples³⁷ may be required, but at present the performance metrics for these have not reached those of the ferri/ferrocyanide-based systems.

Conclusions

This study has demonstrated that thermogalvanic cells designed for ambient temperature valorisation, using air as both the thermal source and sink, can feasibly convert low-grade thermal energy into electricity. By comparing air–air and water–water setups, we confirmed that the water–water configuration generates higher power due to more efficient thermal transfer at the electrode-thermal source interfaces,

resulting in a larger sustained temperature difference (ΔT). However, the air–air system operates passively, requiring no external pumping. Importantly, partial gelation (and not complete gelation) was found to be crucial for both configurations; partial gelation was able to minimise parasitic convective heat transfer within the electrolyte, thus improving the temperature gradient across electrodes and optimising power output, while minimising the negative mass transfer effects typical for fully gelled electrolytes.

For real-world applications, building materials typically require good thermal resistance. For air–air thermogalvanic cells, increased cell dimensions and optimised gelation was also required to optimise the thermal resistance, and thus optimise valorising two temperature-differentiated air bodies into electricity. A thermogalvanic cell with the width of a house brick (100 mm) and partially gelled electrolyte (using 2.5 wt% sodium poly(acrylate)) demonstrated superior thermal resistance than air. Future studies should explore alternative or additional methods of increasing internal thermal resistance, combined with enhanced heat exploitation (*e.g.* solar irradiation) and especially heat dissipation (*e.g.* fins, radiative cooling) at the two electrodes.

Data availability

The data supporting this article have been included as part of the ESI.†

Author contributions

Conceptualisation (RT, PEP, LA); funding acquisition (RT, PEP, LA); methodology (RHJ, SM, MAB, LA); investigation (RHJ, SM); validation (SM); supervision (MAB, LA); writing – original draft (RHJ, SM, MAB); writing – review & editing (all).

Conflicts of interest

Three authors (R. T., P. E. P., L. A.) hold a patent related to earlier prototypes of this air–air cell.³⁸ There are no other conflicts of interest to declare.

Acknowledgements

S. M. and L. A. acknowledge the Royal Society of Chemistry for an Undergraduate Research Bursary (U21-3197306810). M. A. B. and L. A. acknowledge the EPSRC for a Standard Research Studentship (DTP; EP/N509498/1). R. T., P. E. P. and L. A. acknowledge the PLuS Alliance for a Seed Grant (PA17046).

Notes and references

- 1 R. O. Fitriani, B. D. Long, M. C. Barma, M. Riaz, M. F. M. Sabri, S. M. Said and R. Saidur, *Renewable Sustainable Energy Rev.*, 2016, **64**, 635–659.
- 2 M. A. Buckingham and L. Aldous, *J. Electroanal. Chem.*, 2020, **872**, 114280.



- 3 J. T. Hupp and M. J. Weaver, *Inorg. Chem.*, 1984, **23**, 3639–3644.
- 4 M. A. Buckingham, S. Hammoud, H. Li, C. J. Beale, J. T. Sengel and L. Aldous, *Sustainable Energy Fuels*, 2020, **4**, 3388–3399.
- 5 L. Zhang, T. Kim, N. Li, T. J. Kang, J. Chen, J. M. Pringle, M. Zhang, A. H. Kazim, S. Fang, C. Haines, D. Al-Masri, B. A. Cola, J. M. Razal, J. Di, S. Beirne, D. R. MacFarlane, A. Gonzalez-Martin, S. Mathew, Y. H. Kim, G. Wallace and R. H. Baughman, *Adv. Mater.*, 2017, **29**, 1605652.
- 6 K. Laws, M. A. Buckingham and L. Aldous, *Chem. Sci.*, 2024, **15**, 6958–6964.
- 7 A. H. Kazim, A. S. Boeshaghi, S. T. Stephens and B. A. Cola, *Sustainable Energy Fuels*, 2017, **1**, 1381–1389.
- 8 Y. Liu, S. Zhang, Y. Zhou, M. A. Buckingham, L. Aldous, P. C. Sherrel, G. G. Wallace, G. Ryder, S. Faisal, D. L. Officer, S. Beirne and J. Chen, *Adv. Energy Mater.*, 2020, **10**, 2002539.
- 9 P. Yang, K. Liu, Q. Chen, X. Mo, Y. Zhou, S. Li, G. Feng and J. Zhou, *Angew. Chem., Int. Ed.*, 2016, **55**, 12050–12053.
- 10 R. Hu, B. A. Cola, N. Haram, J. N. Barisci, S. Lee, S. Stoughton, G. Wallace, C. Too, M. Thomas, A. Gestos, M. E. Dela Cruz, J. P. Ferraris, A. A. Zakhidov and R. H. Baughman, *Nano Lett.*, 2010, **10**, 838–846.
- 11 B. Yu, J. Duan, J. Li, W. Xie, H. Jin, R. Liu, H. Wang, L. Huang, B. Hu and J. Zhou, *Research*, 2019, **1**, 1–10.
- 12 J. Zhao, X. Wu, H. Yu, Y. Wang, P. Wu, X. Yang, D. Chu, G. Owens and H. Xu, *EcoMat*, 2023, **5**, e12302.
- 13 S. Pu, Y. Liao, K. Chen, J. Fu, S. Zhang, L. Ge, G. Conta, S. Bouzarif, T. Cheng, X. Hu, K. Liu and J. Chen, *Nano Lett.*, 2020, **20**, 3791–3797.
- 14 M. Wang, J. Ruan and X. Zhang, *Therm. Sci. Eng. Prog.*, 2023, **45**, 102083.
- 15 Z. Zhang, N. Zhang, Y. Yuan, P. E. Phelan and S. Attia, *J. Energy Storage*, 2023, **62**, 106912.
- 16 H. Wang, Y. Diao, Y. Lu, H. Yang, Q. Zhou, K. Chrulski and J. M. D'Arcy, *Nat. Commun.*, 2020, **11**, 3882.
- 17 H. Zhou, F. Matoba, R. Matsuno, Y. Wakayama and T. Yamada, *Adv. Mater.*, 2023, **35**, 2303341.
- 18 Z. Jia, Y. Sui, L. Qian, *et al.*, *Nat. Commun.*, 2024, **15**, 6110.
- 19 G. Mensah, A. Andrews, F. Davis, K. Mensah-Darkwa and P. Phelan, *Therm. Sci. Eng. Prog.*, 2022, **34**, 101448.
- 20 M. A. Buckingham, K. Laws, J. T. Sengel and L. Aldous, *Green Chem.*, 2020, **22**, 6062–6074.
- 21 J. Wu, J. J. Black and L. Aldous, *Electrochim. Acta*, 2017, **225**, 482–492.
- 22 M. A. Buckingham, K. Laws, E. Cross, A. J. Surman and L. Aldous, *Green Chem.*, 2021, **23**, 8901–8915.
- 23 M. A. Buckingham, F. Marken and L. Aldous, *Sustainable Energy Fuels*, 2018, **2**, 2717–2726.
- 24 S. W. Hasan, S. M. Said, M. F. M. Sabri, A. S. A. Bakar, N. A. Hashim, M. M. I. M. Hasnan, J. M. Pringle and D. R. Macfarlane, *Sci. Rep.*, 2016, **6**, 29328–29339.
- 25 Y. Ikeda, Y. Cho and Y. Murakami, *Sustainable Energy Fuels*, 2021, **5**, 5967–5974.
- 26 Y. Mua and T. I. Quickenden, *J. Electrochem. Soc.*, 1996, **143**, 2558–2564.
- 27 H. A. H. Alzahrani, M. A. Buckingham, F. Marken and L. Aldous, *Electrochem. Commun.*, 2019, **102**, 41.
- 28 L. Jin, G. W. Greene, D. R. MacFarlane and J. M. Pringle, *ACS Energy Lett.*, 2016, **1**, 654–658.
- 29 Y. Zhou, S. Zhang, M. A. Buckingham, L. Aldous, S. Beirne, C. Wu, Y. Liu, G. Wallace and J. Chen, *Chem. Eng. J.*, 2022, **449**, 137775.
- 30 C. Han, X. Qian, Q. Li, B. Deng, Y. Zhu, Z. Han, W. Zhang, W. Wang, S. Feng, G. Chen and W. Liu, *Science*, 2020, **368**, 1091–1098.
- 31 Z. Wu, B. Wang, J. Li, Y. Jia, S. Chen, H. Wang, L. Chen and L. Shuai, *Nano Lett.*, 2023, **23**, 10297–10304.
- 32 M. A. Buckingham, S. Zhang, Y. Liu, J. Chen, F. Marken and L. Aldous, *ACS Appl. Energy Mater.*, 2021, **4**, 11204–11214.
- 33 M. A. Trosheva, M. A. Buckingham and L. Aldous, *Chem. Sci.*, 2022, **13**, 4984–4998.
- 34 A. Gunawan, H. Li, C. H. Lin, D. A. Buttry, V. Mujica, R. A. Taylor, R. S. Prasher and P. E. Phelan, *Int. J. Heat Mass Transfer*, 2014, **78**, 423–434.
- 35 Y. Cengel and A. Ghajar, *Heat and Mass Transfer: Fundamentals and Applications*, McGraw-Hill, 5th edn, 2018.
- 36 S.-H. Hong, C.-C. Hsu, T.-H. Liu, T.-C. Lee, S.-H. Tung, H.-L. Chen, J. Yu and C.-L. Liu, *Mater. Today Energy*, 2024, **42**, 101546.
- 37 B. Guo, Y. Hoshino, F. Gao, K. Hayashi, Y. Miura, N. Kimizuka and T. Yamada, *J. Am. Chem. Soc.*, 2020, **142**, 17318–17322.
- 38 P. Phelan, L. Aldous, R. Taylor and B. Obeng, *WO Pat.*, 2020082028A1, 2020.

