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Cite this: DOI: 10.1039/c0xx00000x

www.rsc.org/xxxxxx

ARTICLE TYPE

## A Review on Materials for Light Scattering in Dye-sensitized Solar Cells

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**Keywords:** DSCs, Metal Oxides, Active Layer, Scattering Layer, Mie theory.

Dye-sensitized solar cells (DSCs) offer interesting possibilities in photovoltaics which is the technology of harvesting solar photons to generate electricity. Improving the charge transport through the metal oxide film, finding dyes with better absorption both in the visible and near IR regions of the solar spectrum and fabricating innovative materials for scattering layer are the proposed way forward for improving the efficiency of DSCs. Light scattering is employed in dye-sensitized solar cells to improve the optical absorption of the incident light. The conventional method of light scattering in DSCs is by using a separate scattering layer consisting of large particles with diameters comparable to the wavelength of the incident light. An additional over-layer on nanocrystalline TiO<sub>2</sub> photoanode will encourage light scattering in DSCs especially in the red part of the solar spectrum. Different nanostructures with good dye adsorption and light scattering properties were tried as light scattering layers in DSC. Of late, scientists attempted the use of functional materials having enhanced light scattering properties and high internal surface area as dual function materials (that is a single layer of material capable of both light absorption and scattering). This review explores theoretical aspects and materials innovation for light scattering and their application in DSCs.

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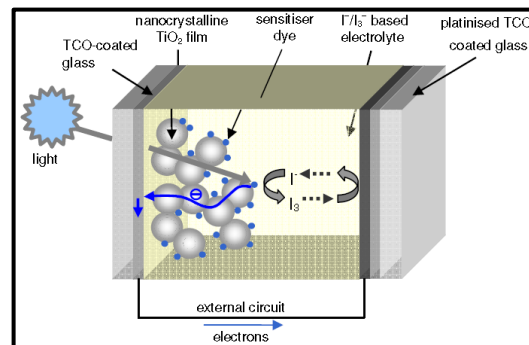
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## 1. Introduction

The dye-sensitized solar cells (DSCs) have become one of the promising solar cells to convert solar energy to electricity with good efficiency at considerably low fabrication cost.<sup>1</sup> A dye-sensitized semiconductor coated on a glass substrate is the heart of a DSC system. The TiO<sub>2</sub> semiconductor film with 10–15 μm thickness is used as the active layer in a DSC photoanode. A platinum (Pt)-coated counter electrode is sandwiched with photoanode and the space between the electrodes is filled with a liquid electrolyte that contains I<sup>-</sup>/I<sub>3</sub><sup>-</sup> redox couple. Solar radiation generates electrons from dye molecules and transfer them to the external circuit through the semiconductor oxide network.

The research on DSCs for higher conversion efficiency is intensively pursued by different research groups around the globe. The current focus is mainly on the modification of dyes, electrolyte and the semiconductor oxides for a better performance. The conventionally used dyes absorb only in the visible region of the solar spectrum. The dyes with broader absorption of solar spectrum such as Black dye are preferred in DSCs. The iodide/triiodide electrolyte offers excessive driving force of 0.6 eV for dye regeneration which leads to a potential drop of about 600 mV. This is one of major drawbacks of iodide/triiodide electrolyte and the limiting factor for increasing the conversion efficiency further from the current level. The organic solid hole-transport material (HTM) is considered as best alternative to the liquid electrolytes to obtain a high V<sub>oc</sub> of 986 mV.<sup>2</sup> But penetration of this material into nanoporous film is one of the major drawbacks of the HTMs. The morphology of TiO<sub>2</sub> in the photoanode film plays a vital role in determining the conversion efficiency.<sup>3–6</sup> One can employ one-dimensional (1D) nanostructures to provide direct pathway for electron diffusion. Using nanostructures promoting enhanced light scattering is yet another possible way to enhance the conversion efficiency.

The idea of light scattering in DSCs is to extend travelling distance by confining light propagation within the photoanode. This strategy will promote interaction of incident photons with the dye molecules. As working electrode is normally made up of TiO<sub>2</sub> nanoparticles of 10–25 nm size, it remains transparent to the visible region of solar spectrum. In this scenario, solar energy is only partially harvested by the active layer. Inclusion of large particles of size comparable to the wavelength of incident light encourages better light scattering in DSCs especially from the red part of the solar spectrum. The efficiency decrement of DSCs when scaling up to larger modules is a notorious problem. Light scattering layer approach is an effective and practical method to enhance efficiency in larger DSC modules and solar panels, though this is technologically challenging and expensive. Grätzel et al. have reported a conversion highest efficiency of 12.3% with porphyrin-sensitized solar cells using cobalt (II/III)-



**Fig.1** Schematic of liquid electrolyte Dye sensitized solar cell.<sup>8</sup>

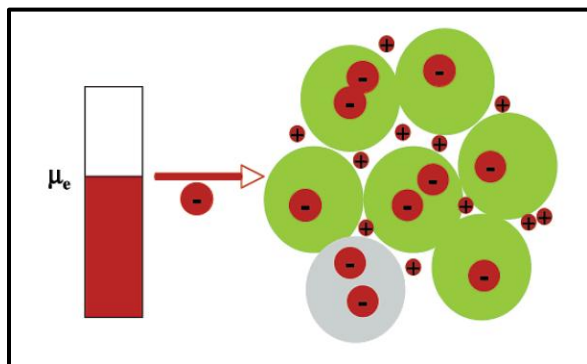
based electrolyte.<sup>2a</sup> Here they had incorporated a 5 μm thick light scattering layer (~ 400 nm sized TiO<sub>2</sub> particles) over 5 μm thick active layer. Recently the same research group has reported 13% efficiency with molecularly engineered porphyrin dye.<sup>2b</sup> Chiba et al. had achieved a highest certified conversion efficiency of 11.1 % by incorporating light scattering layer.<sup>7</sup> This is the highest conversion efficiency tested by any of public test center (National Institute of Advanced Industrial Science and Technology, AIST, Japan).

This review article presents theoretical background on light scattering and outlines recent developments in the light scattering approach to enhance current density and power conversion efficiency of DSCs. We believe the review is topical and will be interesting to the scientists working in the same and related fields.

## 2. Operational principle of DSCs

A mesoscopic semiconductor oxide in contact with the redox electrolyte or hole-conductor is the major part of DSC.<sup>8</sup> In **Fig 1**, schematic of dye sensitized mesoporous nanoparticles in contact with the liquid electrolyte is shown. TiO<sub>2</sub> (anatase) has been widely employed as the semiconductor oxide although wide band-gap materials such as ZnO and Nb<sub>2</sub>O<sub>5</sub> have also been experimented. A monolayer of sensitizer is attached to the nanocrystalline particles in the film. Electrons from dye are injected into the conduction band of the semiconductor oxide by photoexcitation. Dye is regenerated by the organic electrolyte containing a redox couple (iodide/triiodide couple). Iodide is responsible for dye regeneration; in turn iodide is regenerated by reduction of triiodide at the counter electrode. The voltage generated under illumination is the difference between quasi Fermi level of the semiconductor oxide and the redox potential of the electrolyte.

The high surface area of semiconductor oxide is critical for the high performance of DSCs. Since the dye monolayer is adsorbed on the semiconductor oxide; it is obvious that its high surface area will enhance the light absorption. The function of mesoporous semiconductor oxide is similar to the thylakoid vesicles in green leaves.<sup>9</sup> The use of mesoporous film enhances the interfacial surface area over the geometrical surface area by 1000 fold



**Fig.2** Local screening of electrons into nanocrystalline film.  $\mu_e$  is Fermi level of the electrons.<sup>9</sup>

for a 10  $\mu\text{m}$  thick film. Hence, when light passes through mesoporous structure, it intersects hundreds of dye monolayers. Moreover, the semiconductor oxide also fulfills the role of electron conduction.

A wide band-gap semiconductor oxide film acts as an insulator in dark. A single electron injection in a 20 nm-sized particle generates an electron concentration of  $2.4 \times 10^{17} \text{ cm}^{-3}$ .<sup>8</sup> Fermi level of semiconductor oxide should be close to the conduction band for the undeterred flow of electrons. Otherwise there is a certain probability of trapping of electrons which affects the diffusion coefficient. As shown in **Fig 2**; the injected electron can be screened by the cations in surrounding electrolyte which results in electron percolation.<sup>10</sup> Consider a 10  $\mu\text{m}$  thick mesoscopic film composed of 20 nm sized particles. Beer-Lambert's law explains light absorption in terms of the reciprocal length.

$$\alpha = \sigma c \quad (1)$$

where  $\alpha$  is the optical absorption cross-section of the sensitizer,  $\sigma$  is the sensitizer's concentration in the mesoporous film and  $c$  is the velocity of the light.

The light harvesting efficiency (LHE) given by

$$\text{LHE}(\lambda) = 1 - 10^{-\alpha d}, \quad (2)$$

where  $d$  is the thickness of the film.

Incorporation of 100-400 nm sized anatase particles enhances the light absorption in the near-infrared (N-IR) region of the solar spectrum.

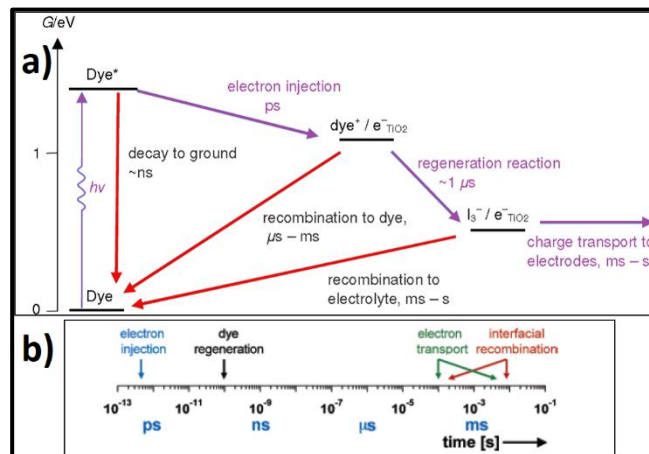
### 2.1 Dynamic competition in DSC

Dynamic competition between the electron transfer, charge-transport processes and several loss pathways determine the quantum yield of DSCs.

Different loss pathways are listed below:

- Decay of the excited dye.
- Charge recombination of injected electrons with dye cations and with the redox electrolyte.

Dynamic competition between forward processes and loss pathways decides the efficiency of DSCs.



**Fig.3a** Forward processes (light absorption, electron injection, dye regeneration and charge transport) are indicated by violet arrows and loss paths (excited-state decay to ground and electron recombination with dye cations and oxidized redox couple) ways indicated by red arrows.<sup>8</sup> **b**) Schematic of time scale involved in different processes in DSC.<sup>9</sup>

**Fig 3a&b** show the dynamic competition processes in DSCs. Photoexcited electron has to be rapidly (in picoseconds) injected to the semiconductor oxide so that it cannot fall back into the ground state. The regeneration of dye (in  $\sim \mu\text{s}$ ) should be faster than the feasible recombination with electrons in the oxide (passage time of electrons inside the oxide film must be lower than the exemplary time needed for them to recombine either with an oxidized dye molecule or electrolyte). The thicker nanoporous films cannot be used to increase the dye absorption because electrons are limited by diffusion path.

$$L_n = \sqrt{D_e \tau} \quad (4)$$

where  $D_e$  and  $\tau$  are diffusion coefficient and lifetime of electrons, respectively. For the successful collection, electron diffusion length must be greater than film thickness ( $d$ ).

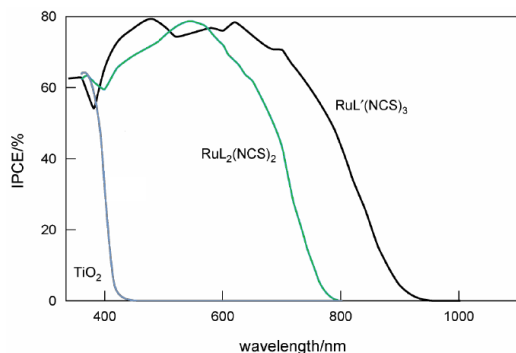
$$L_n > d \quad (5)$$

Hence, an optimum thickness should be adopted for an efficient collection of electrons.

### 2.2 Incident Photon-to-Current-Conversion Efficiency

Incident photon-to-current-conversion efficiency (IPCE) or External Quantum Efficiency is described as the number of electrons (photocurrent) in external circuit divided by incident monochromatic photons.

$$\text{IPCE}(\lambda) = \text{LHE}(\lambda) \eta_{\text{inj}} \eta_{\text{coll}} \quad (6)$$



**Fig.4** Spectral response of photocurrent for solar cell using different dyes. Violet color curve represents bare  $\text{TiO}_2$ ; green and black represent N3 and Black dye, respectively. Here,  $L=4, 4'$ -COOH-2,2'-bipyridine and  $L'=4,4',4''$ -COOH-2,2': 6',2''-terpyridine<sup>8</sup>

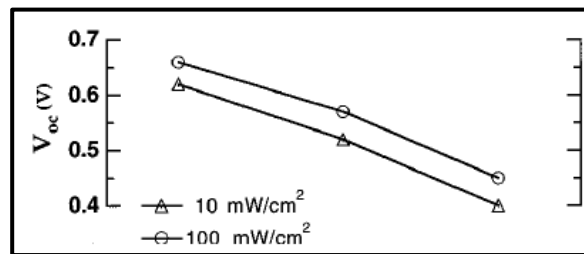
LHE ( $\lambda$ ) is the light harvesting efficiency for photons of wavelength  $\lambda$ ;  $\eta_{\text{inj}}$  is quantum yield for electron injection from excited sensitizer into oxide semiconductor and  $\eta_{\text{coll}}$  is the collection efficiency of electrons to semiconductor oxide film.

### 2.3 Charge Injection

Sensitizer absorbs an energy which is greater than the HOMO-LUMO gap. An ideal sensitizer for photovoltaic cell should absorb all photons below a threshold wavelength of 920 nm. Functional groups such as carboxylate, hydroxamate, or phosphonate moieties must be incorporated to attach it to the nanoporous semiconductor film.<sup>11</sup> The energy level of the excited state should be matched with the lower edge of conduction band. It should be stable enough to run about 100 million turnover cycles (life time about 20 years in exposure to natural light). Ruthenium complexes known for their stability are commonly used as dyes in DSCs. *Cis*-Di (thiocyanato) bis (2,2'-bipyridyl) - 4,4'- dicarboxylate ruthenium (II), known as N3 or N-719 dye depending on whether it carries four or two protons, was found to be a good solar absorber and charge-transfer sensitizer.<sup>12</sup> In 2001, Nazeeruddin et al. discovered a new dye, tri(cyanato)-2,2',2''-terpyridyl-4,4',4''(tricarboxylate) Ru(II) ('black dye') competent with N3.<sup>13</sup> The spectral response of different dyes are shown in **Fig. 4**, where black dye shows a better performance over the others.

The photoexcitation ensures a shift in electron density from the ruthenium metal to the carboxylated bipyridyl ligand. The charge separation occurs from the  $\pi^*$  orbitals of the ligand of ruthenium. Molecularly engineered sensitizers having injection times in the order of pico- or femtosecond range are also developed.<sup>14-19</sup> The injection is too rapid that it cannot redistribute its vibrational excitation energy by exchange with phonon in solids.

The quantum yield of charge injection ( $\eta_{\text{inj}}$ ) indicates the fraction of photons absorbed by dye which are converted into electrons.



**Fig.5** A graph plotted  $V_{\text{oc}}$  against  $\ln [I_3^-]$  at different radiant power.<sup>38</sup>

$$\eta_{\text{inj}} = k_{\text{inj}} / (k_{\text{deact}} + k_{\text{inj}}) \quad (7)$$

$k_{\text{inj}}$  is rate constant of injection of electron into oxide film, where  $k_{\text{deact}}$  represents rate constant of deactivation of the electronically excited state. To attain best  $\eta_{\text{inj}}$ ,  $k_{\text{inj}}$  must be at least 100 times higher than  $k_{\text{deact}}$ .

### 2.4 Charge Separation and Recombination

A parameter  $k_b$  is introduced to characterize recombination rate. It is desirable to develop a sensitizer with high  $k_{\text{inj}}$  value and low  $k_b$ . The back reaction of electrons with oxidized ruthenium complexes involves a localized d orbital in the ruthenium metal. This energy loss channel is negligible because of the rapid reduction of ionized dye molecules by highly concentrated  $\Gamma^-$  ions.<sup>20,21</sup> An important parameter in limiting the conversion efficiency of DSCs is low photovoltage.<sup>22</sup> Charge recombination in nanocrystallite/redox electrolyte interface plays a significant role in reducing the photovoltage. Voltage determining recombination process is represented by the reaction,



The above reaction takes place at the entire  $\text{TiO}_2$  / redox electrolyte solution interface. In **Fig. 5**, a graph plotted  $V_{\text{oc}}$  vs.  $\ln [I_3^-]$  shows the dependence of open-circuit voltage on recombination reaction.

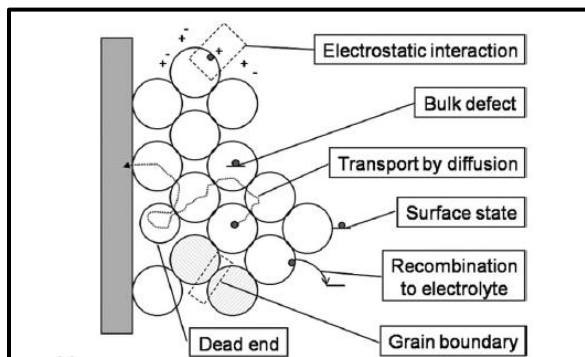
### 2.5 Electron transport in Mesoporous Semiconductor Oxide

The mesoporous semiconductor film is formed by interconnecting numerous nanocrystals. The main driving force for electron transport in mesoporous  $\text{TiO}_2$  has been attributed to the gradient in electron concentration and electronic transport occurs by diffusion.<sup>23-25</sup> The diffusion process is coupled with ions from electrolyte through a weak electric field. This is described by an ambipolar diffusion model.<sup>26,27</sup>

$$D_{\text{amp}} = n+p / (n/D_p) + (p/D_n) \quad (9)$$

$D_{\text{amb}}$  is the ambipolar diffusion coefficient and negative and positive charge carrier concentrations are represented by  $n$  and  $p$ , respectively. The  $D_n$  and  $D_p$  are, respectively, the diffusion concentrations of negative and positive charge





**Fig.6** Overview of Electron diffusion in DSC. Traps can be located at TiO<sub>2</sub>/electrolyte interface, grain boundaries or in bulk of TiO<sub>2</sub>.<sup>6</sup>

carriers. It is also observed that electron transport depends on the incident light intensity, *i.e.*, the higher the intensity the faster the transport.<sup>28,29</sup> At lower intensities, there are chances for the electrons to get trapped in the localized states below the conduction band. The different traps located at TiO<sub>2</sub>/electrolyte interface, grain boundaries or in bulk of TiO<sub>2</sub> have been depicted in **Fig. 6**.

### 2.6 Reactions at Counter Electrode

Transport of redox mediator in the electrolyte is driven by a diffusion process. This leads to diffusion impedance which acts as series resistance in the solar cell. The diffusion impedance depends on the diffusion coefficient, concentration of redox mediator and the distance between two electrodes.<sup>30</sup> Triiodide is reduced to iodide at the counter electrode of DSCs.



Pt coated on the counter electrode acts as catalyst for the above reaction. It ensures rapid reaction and low overpotential at the counter electrode.

The alternatives to I<sup>-</sup>/I<sub>3</sub><sup>-</sup> have been widely investigated and proved successful in DSCs. The good solubility and redox properties of transition metal mediators of ferrocene/ferrocenium,<sup>31</sup> copper(I)/(II) complexes,<sup>32</sup> Ni(III)/Ni(IV) bis(dicarbollide)<sup>33</sup> make them a good alternatives for conventional redox shuttles. The Copolyridyl complexes have more positive redox potentials which should contribute higher V<sub>oc</sub> in DSCs; however these are limited by slow mass transport and quick recombination between electrons and Co redox compared with the iodide<sup>34-36</sup>.

### 3. Theoretical framework of Light Scattering

Scattering of light may be described as redirection of electromagnetic (EM) waves when it encounters an obstacle. When the EM waves interact with particles, electronic orbitals of molecules would be perturbed periodically with the same frequency as the incident light waves. The oscillation of electron cloud induces a dipole moment to the molecule. This oscillating dipole moment is

a source of EM radiation, which results in the scattering of light. In elastic scattering, the scattered light (by particles) has same frequency as incident light wave. The light scattering theory is mainly categorized into two theoretical frameworks. One is the Rayleigh scattering theory applicable to small dielectric (non-absorbing), spherical particles. Another one is proposed by German Physicist Gustav Mie, known as the Mie scattering theory which is applicable to most spherical scattering systems. Mie theory has no particle size limitations; it is applicable to even large particles. So the Rayleigh theory is just a special case of Mie theory. The Rayleigh scattering is limited by the condition that  $\alpha \ll 1$  and  $|m| \alpha \ll 1$  where  $\alpha = 2\pi a / \lambda$ , where *a* is the spherical particle radius and  $\lambda$  is the relative scattering wavelength defined as

$$\lambda = \lambda_0 / m, \quad (11)$$

where  $\lambda_0$  is incident wavelength with respect to vacuum and  $m_0$  is the refractive index of medium. Refractive index of scatterer (*m*) defined as

$$m = n - ik \quad (12)$$

where *n* represents the refraction of light and complex part is related to absorption. Absorption coefficient is related to the complex term in the above equation by:

$$\text{Absorption coefficient} = 4\pi k / \lambda \quad (13)$$

The Rayleigh scattering is only applicable to particles which are having dimensions much smaller than incident wavelength. Intensity of Rayleigh scattered light by single particle can be calculated

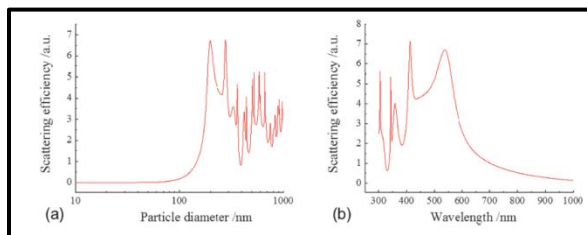
$$I = I_0 \frac{1 + \cos^2 \theta}{2d^2} \frac{2\pi}{\lambda}^4 \frac{m^2 - 1}{m^2 + 2} r^2 \quad (14)$$

Where *I<sub>o</sub>* is intensity of incident light, *d* is distance to the particle and  $\theta$  is the scattering angle. Rayleigh scattering cross section  $\sigma$  is given by the equation,

$$\sigma_{\text{scat}} = \frac{2\pi^5}{3} \frac{2r^2}{\lambda^4} \frac{m^2 - 1}{m^2 + 2}^2 \quad (15)$$

Mie theory is applied to explain scattering from spherical particles regardless of whether the material is light absorbing or non-absorbing. Mie scattering efficiency is determined by the equation<sup>37</sup>

$$Q_{\text{scat, Mie}} = \frac{\sigma_{\text{scat, Mie}}}{\pi r^2} \quad (16)$$



**Fig.7** Shows the dependence of scattering efficiency on a) particle diameter and b) incident wavelength.<sup>37</sup>

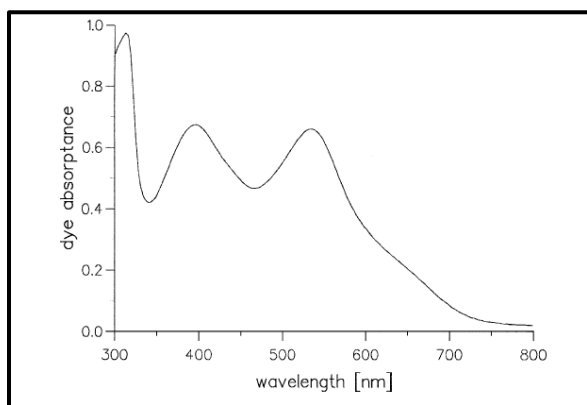
where  $\sigma_{\text{scat, Mie}}$  is the cross-section of Mie scattering given by

$$\sigma_{\text{scat, Mie}} = \sum_{n=0}^{\infty} 2n + 1 (a_n^2 + b_n^2) \quad (15)$$

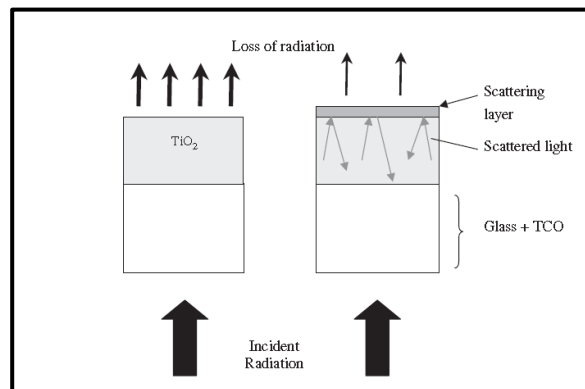
Where  $a_n$  and  $b_n$  are stated by the Riccati Bessel functions  $\psi$  and  $\zeta$ . **Fig. 7** illustrates the dependence of scattering efficiency on the incident wavelength and the particle size. The particles having size comparable to the incident wavelength will scatter effectively. In **Fig. 7a**, the incident wavelength ( $\lambda_{\text{incident}}$ ) is set to 532 nm, so the effective scattering is produced by particles having sizes above 200 nm. **Fig. 7b** shows that particles with 200 nm size scatters incident light in the region 400-550 nm wavelengths.

#### 4. Survey on Materials for Light Scattering in DSCs

Colloidal science tells us about preparation of very small semiconductor particles and how to stabilize these aggregates with respect to coagulation. Colloidal particles normally have sizes less than 1000 Å. Particles show quantum effects when they become smaller than the 'first exciton Bohr radius'. Optical properties can be changed considerably in this size domain. Nanometer sized particles arranged to thin films provide enormous inner surface area.<sup>38,39</sup> Semiconductors absorb light below a particular wavelength  $\lambda$  called the fundamental absorption edge.<sup>40</sup>



**Fig.8** Spectral response of Ru Dye in DSSC.<sup>48</sup>



**Fig.9** A model showing enhancement of scattering in DSC after introducing scattering layer.<sup>65</sup>

$$\lambda \text{ (nm)} = 1240/E_g \text{ (eV)} \quad (17)$$

Extinction of light has an exponential dependence given by:

$$I = I_0 \exp(-\alpha \ell) \quad (18)$$

where  $\ell$  is the penetration length of the light and  $\alpha$  is the reciprocal absorption length. Due to quantum effects, band edge absorption is influenced by particle sizes. TiO<sub>2</sub> is a widely used material in nanotechnology-assisted applications such as solar cells. In DSCs, TiO<sub>2</sub> film on conductive glass substrate (FTO) acts as the photoanode. Conductivity of the TiO<sub>2</sub> film is increased under UV illumination whereas it is very low in the dark at room temperature.<sup>41-44</sup> When the semiconductor is immersed into the electrolyte, the electrolyte is distributed such as semiconductor/electrolyte junction is formed at each nanocrystal. During illumination, electron-hole pair is formed and the hole-transfer to the electrolyte must be faster than the recombination of electrons. An electrochemical gradient is created by this process. Electron is transported through interconnected particles to the back contact. Charge carriers are transported *via* diffusion in the semiconductor films.

It should be noted that recombination probability is increased with increase in film thickness as electrons have to be travelled through increased number of nanoparticles and grain boundaries.<sup>45-47</sup> In addition; resistance loss is another loss mechanism as thickness of film is increased. So film thickness should be as thin as possible for high performance of DSC. Akira Usami proposed a theoretical concept of scattering layer of TiO<sub>2</sub> to improve the optical absorption of the films.<sup>46</sup> It is advisable to add larger particles to scatter light and improve the optical absorption. Nanoporous TiO<sub>2</sub> coated with dye monolayer shows poor absorption in wavelength range  $\lambda = 600-800$  nm (**Fig. 8**). Optical path length of incident light can be increased by adding large particles as scattering layer. Adding too many large particles will also reduce the internal surface area of



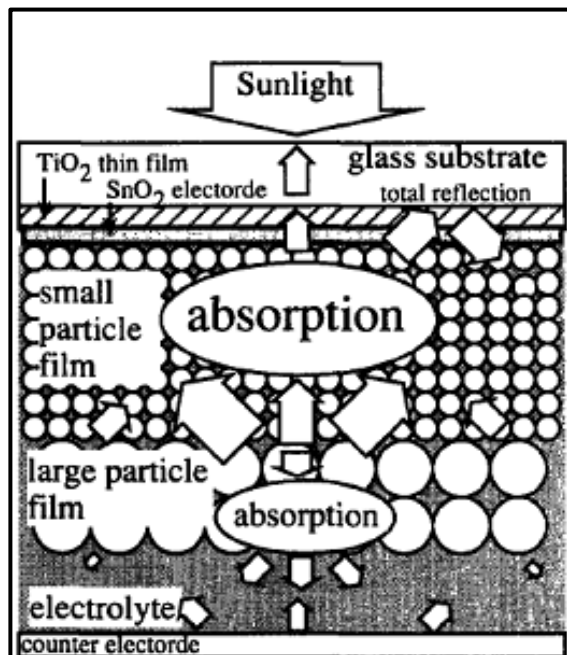


Fig.10 A newly proposed solar cell model.<sup>46</sup>

the TiO<sub>2</sub> film. The optimum thickness for the TiO<sub>2</sub> film is about 10-15  $\mu\text{m}$ . It is obvious that as the film thickness is increased, the path length of the electrons and the redox species in the electrolyte become longer. It increases the series resistance of the film. In DSC, the working electrode normally consists of  $\sim$  12-20 nm sized TiO<sub>2</sub> particles which remain transparent under visible light and eventually enhances the travelling distance of the incident light within the photoelectrode film (Fig. 9).

#### 4.1 Early studies and theoretical foundations

Barbè et al. have extensively studied the influence of processing parameters on the film microstructure at the photoanode.<sup>47</sup> The electrodes prepared by them using colloids autoclaved at  $\leq$  230  $^{\circ}\text{C}$  were fully transparent. Particle-size analysis was performed with different autoclaving temperatures.

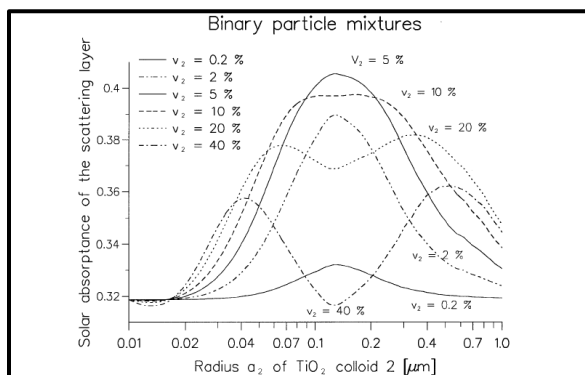


Fig. 11 The Calculated solar absorbance as function of radius of large particles.<sup>48</sup>

The electrodes prepared at 250  $^{\circ}\text{C}$  were opaque due to scattering of light by particle aggregates formed in the film. The scattering of light made a significant improvement in the photovoltaic response of the cell. The scattering of light increases the path length of photons which eventually increases its probability to interact with dye molecules. For light scattering, commercial anatase TiO<sub>2</sub> powder was mixed with the colloidal solution. Theoretical study on multiple scattering in DSC was carried out by Akira Usami.<sup>46</sup> He studied the scattering by Monte Carlo simulation method and found out scattering condition as  $kd/\pi = 1.3$ , where  $d$  is diameter of the particle and  $k$  is the wave number.

A new structure of DSC was proposed by Akira Usami including the idea of scattering within the solar cell (Fig. 10). The optimal diameter of TiO<sub>2</sub> particles was obtained as  $1.3\sim 1.4 \times \pi/k$ . Computer simulations by Ferber et al. gave a clearer picture about the particle size size needed to effectively scatter the sunlight in DSC.<sup>48</sup> The numerical solutions of radiative transport equation were used to solve multiple scattering between the TiO<sub>2</sub> particles.<sup>49,50</sup> The TiO<sub>2</sub> layer was assumed to be formed by mixing 20 nm particles with different fractions of larger particles. An optimal mixture was found to be 5% particles with radius 125-150 nm (Fig. 11). Computer simulations have shown that the fraction of large particles encourage the reflection instead of absorption and the photon absorption in near-infrared region ( $\lambda = 550\text{-}750$  nm) have significantly improved. Ferber et al. further simulated a double layer structure in which back-surface scattering layer acts as a mirror. However, the effort to put a "mirror" layer was downplayed considering the small gain and enormous extra effort in cell fabrication. Rothenberger et al. have done theoretical modeling on the optical characterization of DSCs.<sup>51</sup> They had suggested that incorporation of large particles or voids could enhance the scattering and absorption in DSCs. Two samples were prepared for optical studies of scattering layer in DSC.

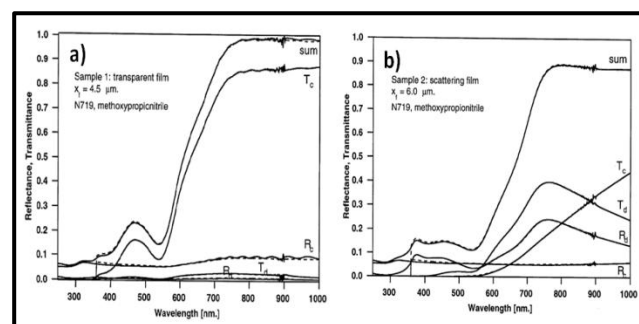
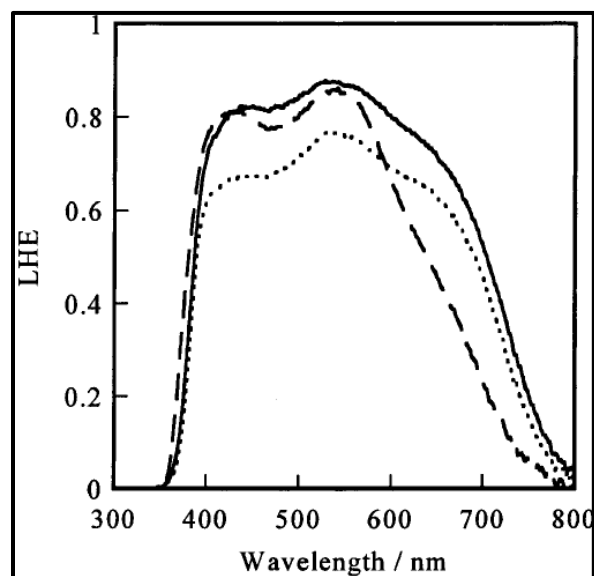
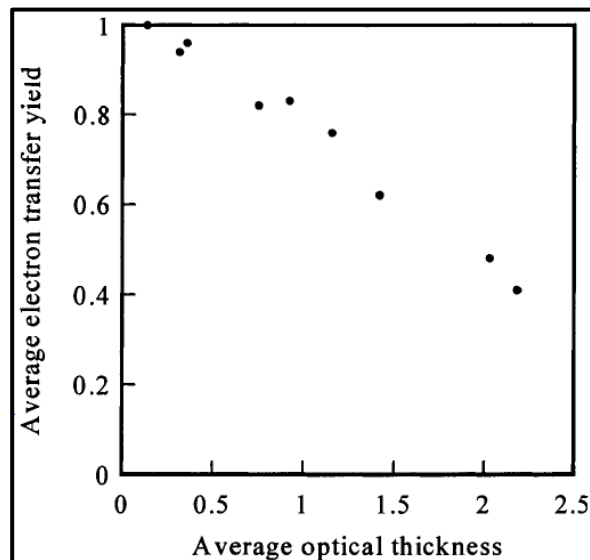


Fig.12 Transmittance and reflectance spectra for (a) sample 1 and (b) sample 2.  $T_c$ ,  $T_d$ ,  $R_c$  and  $R_d$  are the collimated transmittance, diffused transmittance, collimated reflectance and diffused reflectance respectively.<sup>51</sup>

**Table 1.** Different TiO<sub>2</sub> films with varying particle sizes.

Type of TiO <sub>2</sub>	Porosity	Roughness factor (10μm thickness)	Average diameter (nm)
S1(synthesized in acidic solution)	0.66	1450	14.0
S2(synthesized in acidic and basic solutions)	0.59	1090	22.3
S3 (commercial)	0.61	630	37.5

Sample 1 consisted of a 4 μm thick film of anatase TiO<sub>2</sub> with diameter of about 15 nm and sample 2 of 6.0 μm thick films of anatase particles with diameter of about 150 nm. The specific surface area (BET) is reduced to 35 m<sup>2</sup>/g (sample 2) from 101 m<sup>2</sup>/g (sample 1) as the size of the particle is increased. In **Fig. 12**, sample 2 shows a low diffused transmittance. As the particle size was increased, the sample 2 has shown a decrease in T<sub>c</sub> and increase in T<sub>d</sub> and R<sub>d</sub>. The major share of collinear reflectance is from the interface of substrate glass. So there is a great potential in anti-reflecting materials that can be used in substrate glass to improve solar cell efficiency. In the solar cells, the absorption of light is limited by backscattering and diffused reflection. Tachibana et al. had indirectly studied relation between the optical thickness and the electron-transfer yield in DSCs.<sup>52</sup>

**Fig.13** Harvesting Spectra for S1 (broken line), S2 (solid line) and S3 (dotted line).<sup>52</sup>**Fig.14** Relation between average electron-transfer yield and average optical thickness.<sup>52</sup>

The photocurrent in solar cell, determined by both LHE and electron-transfer yield is given by:

$$J_{sc} = \Phi_{ET} q F \lambda (1 - r \lambda) LHE \lambda d \lambda \quad (19)$$

where  $\Phi_{ET}$  is average electron-transfer yield term,  $q$  is the electron charge,  $F(\lambda)$  is the incident photon flux density at wavelength  $\lambda$  and  $r \lambda$  is the incident light loss

Tachibana et al. used differently sized TiO<sub>2</sub> particles and absorption of photons was observed by using UV-Vis steady-state absorption spectrometer. Light scattering was featured by monitoring transmittance and reflectance of unsensitized TiO<sub>2</sub> film. Light harvesting efficiency was also obtained with same configuration of particles which are sensitized by the ruthenium-based dye. The absorption ratio was calculated by comparing both sensitized and unsensitized TiO<sub>2</sub>. Then the value was corrected for incident light loss at the glass substrate because the absorption and reflection by conducting glass would affect the LHE of sensitized TiO<sub>2</sub>. Different films used by them are given in **Table 1**. In **Fig 13**, absorption spectra of different films have been depicted. The film S1 was a transparent film, which shows the characteristic absorption spectra of dye itself. Optical path has been increased by the addition of large particle into the film. So the absorption at the wavelength greater than 600 nm has been increased (i.e. Film S2). There was lowering of LHE when excessive large particles (i.e. Film S3) were added to the film. It should be noted that average electron-transfer yield was decreased with increasing average optical thickness (as shown in **Fig. 14**). In extremely light scattered condition, dye excitation becomes inhomogeneous which in turn reduce the electron-transfer yield.

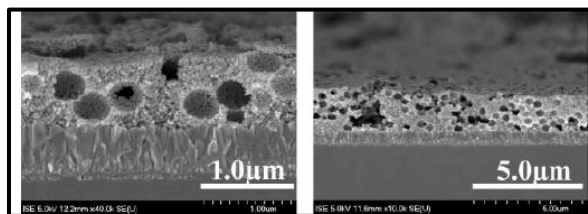


Fig.15 Spherical voids in the TiO<sub>2</sub> film.<sup>53</sup>

## 4.2 Different Nanostructures as light scatterers

### 4.2.1 Spherical voids as light scatterers

Hore et al. incorporated spherical voids as scattering centers in TiO<sub>2</sub> thin film.<sup>53</sup> Carboxyl stabilized polystyrene spheres had been used to make spherical voids since carboxyl groups are known for anchoring onto the TiO<sub>2</sub> surfaces.<sup>54,55</sup> Polystyrene spheres were added to the TiO<sub>2</sub> paste in a volume ratio 1:5 and thoroughly mixed by ultrasonic bath at room temperature. Upon sintering at 450 °C, polystyrene spheres had melted leaving voids in the film (Fig. 15). Mie theory was used to calculate scattering from spherical voids filled with the electrolyte.<sup>56</sup> Scattering cross-section ( $\sigma_s$ ) per volume is an estimate of scattering effectiveness per volume fraction. The absorption was enhanced within the film when scattering angle was very large. It is evident from Fig. 16 that spherical voids of about 50-100 nm could act as very good backscatterers. It was shown by using electrochemical impedance spectroscopy that the spherical voids facilitate the diffusion of ions. By the introduction of spherical voids in the TiO<sub>2</sub> film, the overall efficiency of device was improved by 25%.

Carbon spheres could also be employed to make pores in the TiO<sub>2</sub> film.<sup>57</sup> Yang et al. fabricated carbon spheres of different sizes (250, 500 and 700 nm, respectively) by hydrothermal method to study the influence on LHE and hence the performance of DSCs. Carbon spheres were mixed with TiO<sub>2</sub> slurries with different concentrations and different photoelectrodes were fabricated.

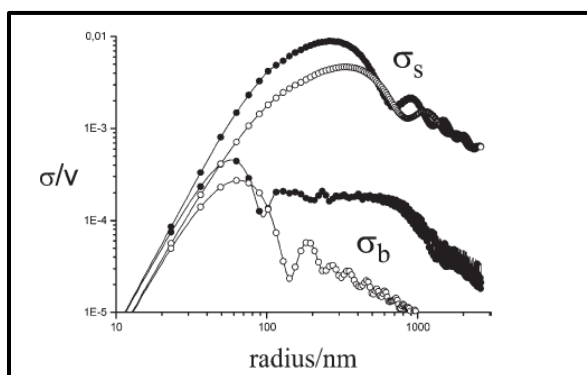


Fig.16 Total ( $\sigma_s/v$ ) and backward ( $\sigma_b/v$ ) scattering coefficient per volume as function of particle radius (white dot indicates presence of electrolyte filled voids in the film and Black dot, presence of solid particles in the film).<sup>53</sup>

It was heated at 450 °C for 2h to remove carbon spheres completely. The optimal concentration of carbon spheres were found to be 15 wt.% to the mass of TiO<sub>2</sub> nanocrystalline contained in the slurry. The best solar cell performance given by photoanode with carbon spheres of 500 nm size in TiO<sub>2</sub> film was 7.2 %.

### 4.2.2 Nano-composites for Light Scattering

#### Nanowire-nanoparticle composite

One-dimensional (1D) structures like nanowires can enhance light scattering and electron transport. Nanowire/nanoparticle composite has shown notable improvement in solar cell efficiency. The light scattering is considerably improved by the incorporation of nanowire with lengths in the range of hundreds of nanometers to micrometers which would provide the same effect as inclusion of large nanoparticles in the film (typically with a diameter of 100 nm). Mixing 5 wt. % of nanowires in nanoparticles shown an overall solar efficiency improvement from 5.84% to 6.24%.<sup>58</sup> Improvement in the light harvesting efficiency was reported by Baxter et al. by depositing ZnO nanoparticles in interstitial voids between ZnO nanowires.<sup>59</sup> Tan et al. extensively studied solar cell performance using nanowire/nanoparticle composites at the photoanode.<sup>60</sup> DSC performance was studied for different wt.% of hydrothermally grown TiO<sub>2</sub> nanowires in TiO<sub>2</sub> nanoparticles in which cell with 20 wt.% nanowires has shown an efficiency of 8.6%. The photoanodes became opaque with 77 wt. % nanowire concentration. The current density was lowered at higher concentrations of nanowires.

The DSCs showed higher open circuit voltage and fill factor because of reduced recombination and low internal resistance inside the nanowire/nanoparticle composite film. So an optimal concentration of nanowire in the film would provide critical improvement on the solar cell performance.

#### Nanofiber-nanoparticle composite

A composite of electrospun TiO<sub>2</sub> nanofibers and TiO<sub>2</sub> nanoparticles deposited on the photoanode has shown an improvement in light harvesting without much compromise on dye adsorption capability.<sup>61</sup>

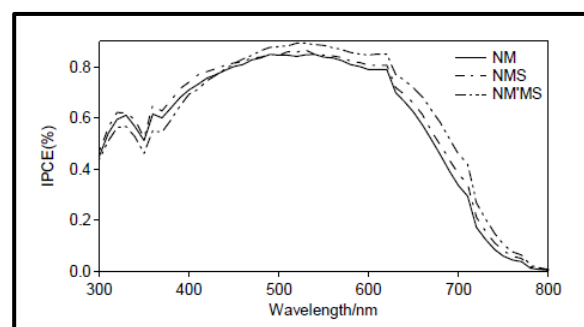


Fig.17 IPCE (%) as function of Wavelength for double layer and multi-layer photoanode.<sup>64</sup>

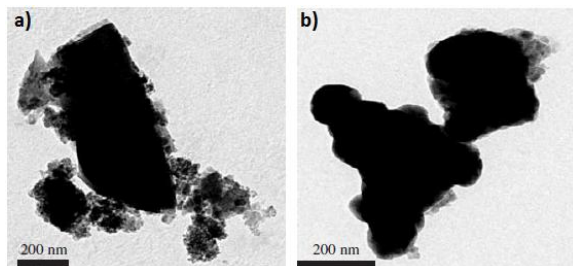


Fig.18 TEM image of rutile TiO<sub>2</sub> with small TiO<sub>2</sub> and b) ZrO<sub>2</sub>.<sup>65</sup>

The cell efficiency was improved by 44% by using nanofiber-nanowire composite than using TiO<sub>2</sub> nanoparticles alone.

The intensity of scattered light (I) by nanofibers is calculated by following equation<sup>62</sup>

$$I = \frac{I_0}{r} \pi m - 1^2 \frac{a^2}{\lambda} \frac{J_1 \frac{4\pi a}{\lambda} \sin \frac{\theta}{2}}{\sin^2 \frac{\theta}{2}}^2 \quad (20)$$

Where  $I_0$  is the incident light intensity and  $\lambda$  is the wavelength. Here  $r$  is the distance from nanofiber to the point under study, 'm' and 'a' are refractive index and radius of nanofiber respectively.  $J_1$  is the Bessel function and  $\theta$  is the angle between directions of  $r$  and incident light. From the above equation, the intensity of light scattering is to be increased substantially when nanofiber diameter is about 200 nm or above.

For the TiO<sub>2</sub> nanofiber-nanoparticle composite, electrospun nanofibers<sup>63</sup> were dispersed in anhydrous ethanol followed by adding the nanoparticle paste. Then the mixtures (with different nanofiber/nanoparticle ratios) were sonicated for uniform distribution of nanofibers and heated at 450 °C for 1 h to obtain the composites.

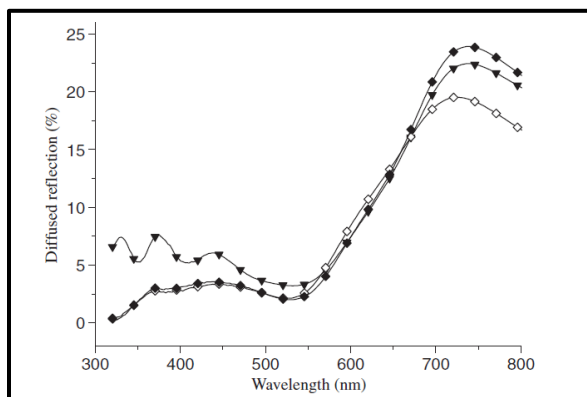


Fig.19 Diffused reflection of different scattering layers on TiO<sub>2</sub> film in a complete device. Scattering layers are mixtures of TiO<sub>2</sub>-Rutile and ZrO<sub>2</sub> in the following proportions. 25:75 (◆), 50:50 (▼) and 75:25 (◇).<sup>65</sup>

The electrospun TiO<sub>2</sub> nanofibers had diameters of about 200 to 300 nm and tens of microns length.

Randomly distributed nanofibers would scatter the light effectively to enhance light harvesting. The forward scattering was increased with increasing the diameter of the nanofibers. Dye-sensitized composite with 15% nanofibers showed a lesser transmission of light which implies a higher amount of light scattering. It has been suggested that 7.5 μm thick film of 15% nanofiber is ideal for high efficiency DSC which has shown an energy conversion efficiency and short circuit current density ( $J_{sc}$ ) of about 8.8% and 16.8 mA cm<sup>-2</sup> respectively.

#### 4.3 Light scattering layer using Nanostructures - A Double layer approach in DSCs

Wang et al. fabricated different types of photoanodes (monolayer, double layer and multilayer) using TiO<sub>2</sub> nanoparticles with various sizes of 23 nm, 50 nm and 100 nm, respectively.<sup>64</sup> In the monolayer configuration, when large particles of 100 nm were incorporated in the film, an improvement in IPCE (%) was observed due to the light scattering effect and the overall efficiency increased from 7.6% to 8.4%. A double layer structure had also been tried because large particles near the conducting glass in monolayer will cause backscattering. A separate layer for light scattering will reduce the backscattering loss. There was an overall efficiency enhancement from 8.4 to 9.2%. The multilayer film structure had shown an improved IPCE (%) compared to the monolayer and double layer photoelectrodes. (Fig. 17)

##### 4.3.1 Large particles

The inclusion of scattering layer made up of large particles above the high surface area TiO<sub>2</sub> film could enhance current density by 80%. Hore et al. tried rutile TiO<sub>2</sub> and zirconium dioxide (ZrO<sub>2</sub>) as scattering layers in varying proportions.<sup>65</sup> The light scattering depends upon change in the refractive index between the active and scattering layers. A mixture of ZrO<sub>2</sub> (with refractive index of 2.1), and rutile TiO<sub>2</sub> (refractive index of 2.8) has also been applied as scattering layer. The scattering layer was prepared from commercially available powders of ZrO<sub>2</sub> and rutile TiO<sub>2</sub> with particle sizes about 500-1000 nm. The Fig. 18 shows TEM image of rutile TiO<sub>2</sub> and ZrO<sub>2</sub>. These powders were dispersed in a 10% ethyl cellulose and  $\alpha$ -terpineol solution. Mainly four types of scattering layers were applied on nanocrystalline TiO<sub>2</sub> film such as rutile TiO<sub>2</sub>, mixture of rutile TiO<sub>2</sub>: ZrO<sub>2</sub> (25:75), mixture of rutile TiO<sub>2</sub>: ZrO<sub>2</sub> (50:50) and ZrO<sub>2</sub> alone. Reflectance and transmittance were measured for the sandwiched configurations. DSC with mixture of rutile TiO<sub>2</sub>: ZrO<sub>2</sub> (25:75) as a scattering layer has shown a higher current density amongst all the cells (Table 2). As shown in Fig. 19, the diffused reflectance is maximal for rutile TiO<sub>2</sub>:ZrO<sub>2</sub> mixture.

Lee et al. prepared two types of TiO<sub>2</sub> pastes of sizes 20 nm and 123 nm, respectively, by sol-gel method.<sup>66</sup>

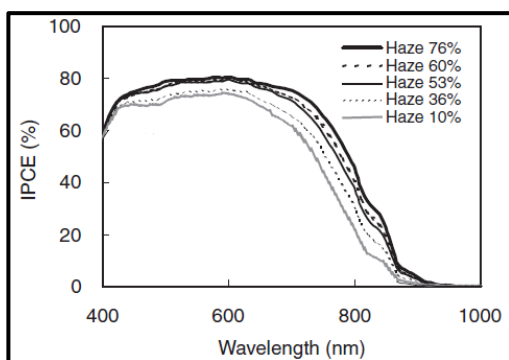


**Table 2.** Current densities shown by different cells with different scattering layers.<sup>65</sup>

DSSC	Current Density
Without scattering layer	7.6
TiO <sub>2</sub>	11.88
TiO <sub>2</sub> :ZrO <sub>2</sub> (25:75)	14.2
TiO <sub>2</sub> :ZrO <sub>2</sub> (50:50)	14.16
ZrO <sub>2</sub>	12.72

A dense layer of 9 nm particles was deposited on the FTO glass substrate. Large particles (*i.e.*, 20 nm and 123 nm) were used as scattering layer in DSC. Lee et al. observed a lesser transparency for multi-layered TiO<sub>2</sub> film with the scattering layer than the mono-layered TiO<sub>2</sub> films. An efficiency of 6.03% was achieved with PEG electrolyte-based DSC when 123 nm sized TiO<sub>2</sub> was used as light-scattering layer. Chou et al. tried TiO<sub>2</sub> with different particle sizes for light-scattering layer varying from 99.8 to 460.9 nm.<sup>67</sup> The short-circuit current density increased from 16.16 mA/cm<sup>2</sup> to 16.71 mA/cm<sup>2</sup>. Cauliflower-like TiO<sub>2</sub> rough spheres synthesized by hydrolysis of Ti(OBu)<sub>4</sub> by using tri-block copolymer as the template was used to make photoelectrode and it has shown an efficiency up to 7.36%, which is higher than that from smooth ones.<sup>68</sup>

Chiba et al. have investigated the dependence of haze factor of TiO<sub>2</sub> photoelectrode on power conversion efficiency of the DSCs.<sup>7</sup> Haze factor is defined as the ratio of diffused transmittance to the total transmittance. The light scattering by TiO<sub>2</sub> electrode depends not only on the particle size of TiO<sub>2</sub> but also on preparation methods (like casting methods, sintering temperature) and atmosphere.

**Fig.20** Dependence of IPCE (%) on the haze of TiO<sub>2</sub> film. (Adapted from *Jpn. J. Appl. Phys.*, 2006, **45**, L638.)<sup>7</sup>

So the haze is considered as the index of light scattering for obtaining better power conversion efficiencies and reproducibility. The haze was controlled by addition of sub-micron TiO<sub>2</sub> particles (~ 400 nm) to the electrodes. **Fig. 20** shows the dependence of IPCE (%) on haze. The high IPCE at infrared region demands high haze TiO<sub>2</sub> where the extinction coefficient of dye is very small (200 mol<sup>-1</sup> cm<sup>-1</sup>) (where as in the visible region molar extinction coefficient of the commercial dyes is very high, so medium haze is only required at this region). The 76% haze in the TiO<sub>2</sub> electrode shows maximum IPCE at 800 nm wavelength and with this high haze factor they have achieved a record efficiency of 11.1% with highest J<sub>sc</sub> of 21 mAcm<sup>-2</sup>. It has evolved that haze of TiO<sub>2</sub> photoelectrode has important role in deciding the device performance.

#### 4.3.2 TiO<sub>2</sub> Nanotubes

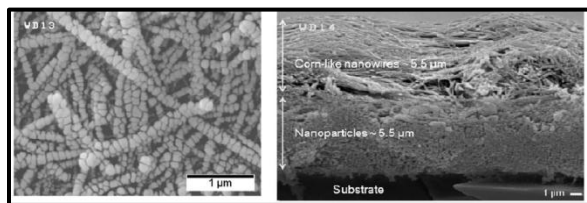
Nakayama et al. studied nanotubes (NTs) as scattering layer in DSC.<sup>69</sup> Nanotubes having dimensions that are not comparable with the wavelengths of incident light; do not generally scatter the light. The light scattering happens due to the presence of grain boundaries and voids among the nanotubes which are having sizes comparable to the wavelength of the incident light. Nanotubes prepared by the anodization method were dispersed in *tert*-butanol to obtain TiO<sub>2</sub> nanotube suspension. Light scattering layer was deposited on nanocrystalline TiO<sub>2</sub> by electrophoretic deposition (EPD). The overall efficiency of DSC using nanotubes as the scattering layer has significantly been improved to 7.53% from 6.91% when large particles were used as the scattering layer.

The DSCs with nanotubes film shows a higher current density as compared to nanoparticle based DSCs. This enhancement in J<sub>sc</sub> is mainly due to light scattering by the nanotubes. But the conversion efficiencies of both different morphologies were obtained as more or less same. This discrepancy is attributed to lesser fill factor reported by nanotube based DSCs. The insulating oxide layer formed between the NTs and the substrate during anodization process lowers the FF of the NT DSCs. So the nanotube applied as light scattering layer in DSCs is more favorable which would avoid the possibility of formation of the insulating oxide layer at the substrate/material interface.

#### 4.3.3 TiO<sub>2</sub> Nanowires

Corn like TiO<sub>2</sub> nanowires were used as scattering layer in DSCs which possess high surface area and light scattering capability (**Fig. 21**).<sup>70,71</sup> Introduction of corn like 1D nanowires in DSC contributed to a significant increase in efficiency by 12% (7.11% for corn like TiO<sub>2</sub> nanowires against 6.34% for regular TiO<sub>2</sub> nanowire cell). Corn like nanowires and regular nanowires were prepared by two hydrothermal operations. The regular nanowires were prepared as per following procedure. One gram of TiO<sub>2</sub> powder in 10 M aqueous NaOH chemically treated in autoclave at 180 °C for 48h to obtain hydrogen titanate. This was further washed with dil.HCl and dried at 25 °C.



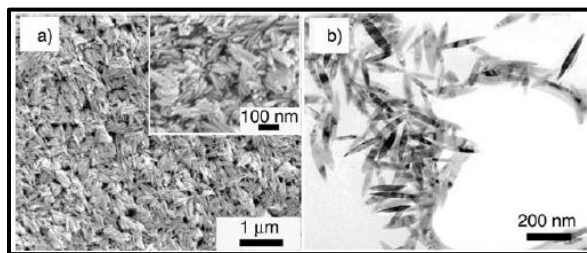


**Fig.21 a)** FE-SEM image of surface morphology of anatase TiO<sub>2</sub> corn-like nanowire, **b)** FE-SEM image of cross-sectional view of photoanode TiO<sub>2</sub> corn-like nanowire as scattering layer.<sup>70</sup>

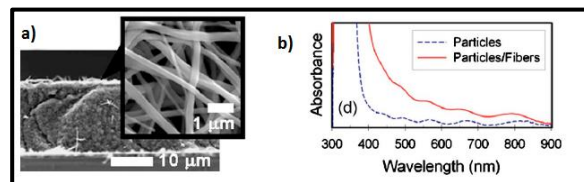
To obtain regular nanowire, the hydrogen titanate was treated in autoclave with deionized water at 180°C for 12h. The corn like nanowires were grown also by consecutive two hydrothermal operations as carried out to obtain regular nanowire except the process temperature and time were controlled to result hydrostatic stress for crack formation. The temperature and time for the both hydrothermal processes were set at 180°C for 72h. The following mechanisms have been attributed as the reasons for the increase in DSC efficiency: a) corn-like TiO<sub>2</sub> nanowires with high surface area provided an enhanced dye sensitization due to which photon absorption is increased, and b) a lower recombination loss when corn-like nanowires were used due to their one dimensionality.

#### 4.3.4 TiO<sub>2</sub> nanospindles

Qiu et al. executed a double layer structure consisting of larger anatase TiO<sub>2</sub> nanospindles as light scattering layer and smaller anatase TiO<sub>2</sub> nanospindles as underlayer for DSCs (**Fig. 22**).<sup>72</sup> The hydrothermal treatment of hydrogen titanate in water and different alcohols resulted in formation of the different nanospindles. The large nanospindles were obtained in water/ethanol solution whereas small nanospindles in water/ethylene glycol. In such a configuration, the small nanospindles facilitate an improved dye loading whereas large ones help in light scattering in the visible and N-IR region. In addition, recombination with I<sub>3</sub><sup>-</sup> in electrolyte is also limited in nanospindles photoanode due to fewer traps in the photoanode. It has shown efficiency up to 8.3% with J<sub>sc</sub> and V<sub>oc</sub> of 16.4 mA cm<sup>-2</sup> and 0.76 V, respectively.



**Fig.22 a)** SEM **b)** TEM image of large nanospindles.<sup>72</sup>

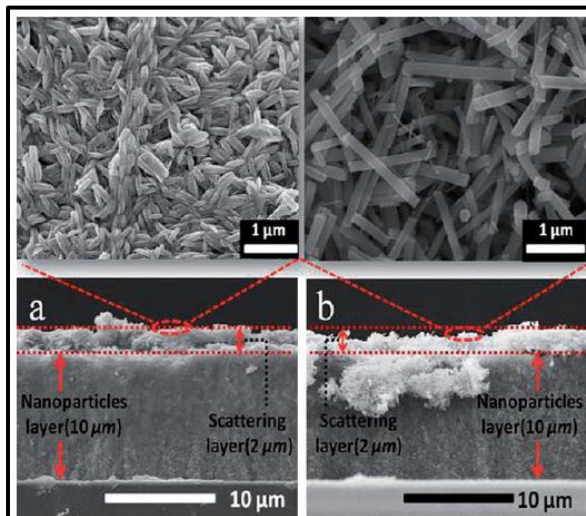


**Fig.23 a)** SEM image of nanoparticle/nanofiber electrode **b)** UV-visible spectra of nanoparticle film electrode and nanoparticle/nanofiber film.<sup>73</sup>

#### 4.3.5 Electrospun materials

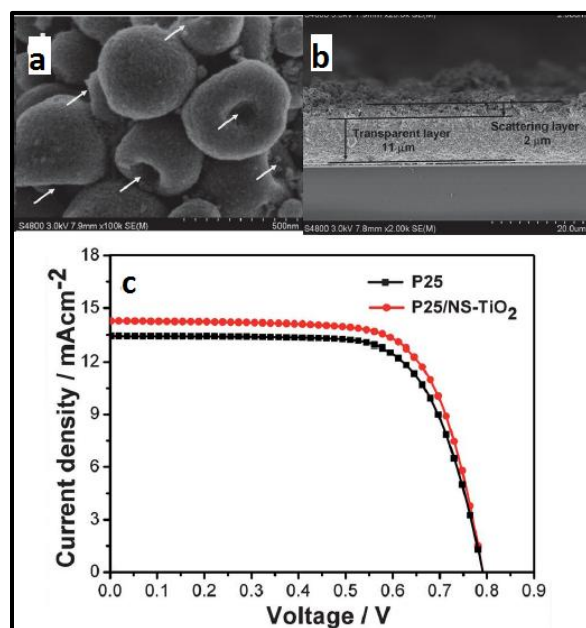
Yoshikawa et al. used 1-D TiO<sub>2</sub> nanofibers as scattering layer in DSCs.<sup>73</sup> The TiO<sub>2</sub> nanofibers were directly electrospun on TiO<sub>2</sub> paste-coated FTO (**Fig. 23a**). The solution used for electrospinning was made of 0.8g PVP, 4g of titanium (IV) butoxide and 1.175g of acetylacetone in 10 ml methanol. The nanofibers (diameter ~ 250 nm) were composed of bundle of nanofibrils (diameter ~ 20 nm). The thickness of light scattering layer deposited on nanocrystalline film was about 1 μm. An improved absorbance of light was observed for nanofibers over the nanoparticle (**Fig. 23b**). The IPCE was shot up to 85% at 540 nm when the nanofibers were used as light scattering layer which was ascribed to enhanced light scattering by TiO<sub>2</sub> nanofibers. An efficiency of 8.14 % was reported with a current density of 17 mA/cm<sup>2</sup> for an exposed area of 0.25 cm<sup>2</sup>. Zhu et al. have done a comparative study by using electrospun rice- and nanofiber-shaped TiO<sub>2</sub>, respectively, as scattering layers in DSC.<sup>74</sup> The active layer was made by TiO<sub>2</sub> commercial paste by doctor blading method on FTO. Scattering layer was fabricated by direct electrospinning of TiO<sub>2</sub> nanofibers on the coated working electrodes (**Fig. 23b**). The electrospinning solution was composed of 1.2g polyvinyl acetate, 10 ml N, N-dimethyl acetamide, 2 ml acetic acid and 1ml titanium (IV) isopropoxide. The average dimension for the rice grain shaped TiO<sub>2</sub> was about 150 nm in diameter and 400 nm in length with a high surface area of ~ 60 m<sup>2</sup> g<sup>-1</sup>.<sup>75,76</sup> The fibers used as scattering layer were having an average diameter of 150 nm with a surface area of 52 m<sup>2</sup> g<sup>-1</sup>.<sup>77</sup> When nanofibers were used as the scattering layer in DSC, efficiency had gone up to 7.06% against 6.44% using TiO<sub>2</sub> nanoparticle film alone and current density had improved from 13.6 mAcm<sup>-2</sup> to 14.9 mAcm<sup>-2</sup>. With the rice-shaped TiO<sub>2</sub> scattering layer, the overall efficiency and short circuit current density were further improved to 7.45% and 15.7 mA cm<sup>-2</sup> respectively (**Fig. 24a**). The IPCE maximum was about 74 % with the rice-shaped TiO<sub>2</sub> and 71% with nanofibers. The improved photovoltaic performance with the rice-shaped structures was attributed to its high surface area and single crystallinity in comparison to the TiO<sub>2</sub> nanofibers (less surface area and polycrystalline).

Nest-shaped TiO<sub>2</sub> (NS-TiO<sub>2</sub>) nanostructures synthesized by electrospinning route is another candidate for scattering layer in DSCs.<sup>78</sup> The precursor solution for electrospinning was composed of 3g of PVP, 14g of titanium (IV) butoxide

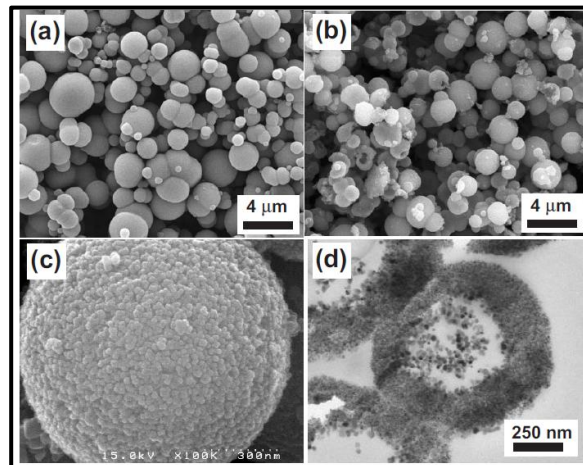


**Fig.24** SEM images of Rice grain shaped TiO<sub>2</sub> and cross-sectional image of TiO<sub>2</sub> nanoparticle – Rice shaped TiO<sub>2</sub> nanostructure working electrode **b)** SEM images of TiO<sub>2</sub> nanofiber and cross-sectional image of TiO<sub>2</sub> nanoparticle – TiO<sub>2</sub> nanofiber nanostructure working electrode.<sup>74</sup>

and 12 mL of glacial acetic acid in 48 mL of ethanol. NS-TiO<sub>2</sub> nanostructures are of diameter ~ 200-500 nm with a cavity at the center (**Fig. 25a**). It was observed that the NS-TiO<sub>2</sub> nanostructures have a higher reflectance than P-25 which implies higher light scattering capability. The photoanode composed of 11 μm thick P25 film and 2 μm thick scattering layer on the FTO showed an efficiency of 8.02% which was about 7.1% higher compared to the ones without the scattering layer (**Fig. 25b** and **Fig. 25c**).



**Fig.25** a) FESEM image of NS-TiO<sub>2</sub> nanoparticle **b)** FESEM image of P25/NS-TiO<sub>2</sub> electrode **c)** J-V curve of DSC with P25/NS-TiO<sub>2</sub> electrode.<sup>78</sup>

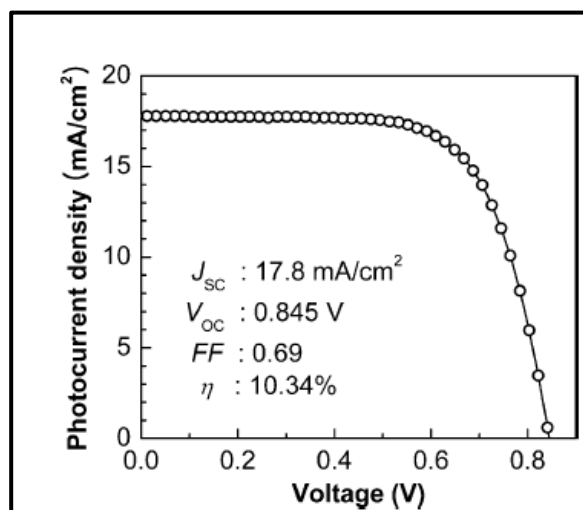


**Fig.26** SEM images of **a)** synthesized NeHS TiO<sub>2</sub> **b)** calcined NeHS TiO<sub>2</sub> at 450°C for 2h **c)** highly magnified image of calcined NeHS TiO<sub>2</sub> at 450°C **d)** TEM image of a sliced NeHS TiO<sub>2</sub>.<sup>79</sup>

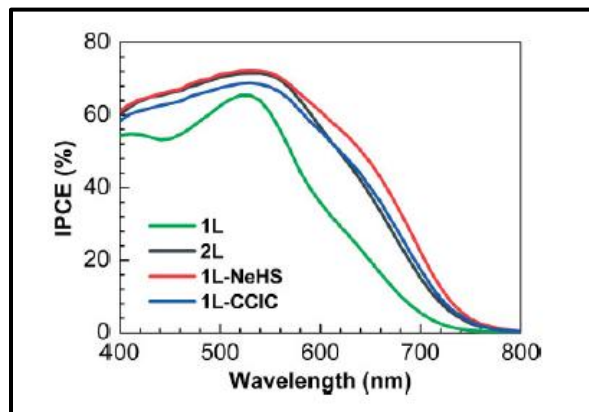
#### 4.3.6 Nano-embossed hollow spherical TiO<sub>2</sub>

Nano-embossed hollow spherical TiO<sub>2</sub> (NeHS TiO<sub>2</sub>) was used as bifunctional material for high-efficiency DSC.<sup>79</sup> The diameter of synthesized NeHS TiO<sub>2</sub> spheres was in the range of 1-3 μm (**Fig. 26**). The average pore size was about 10 nm and surface area obtained was 58 m<sup>2</sup>g<sup>-1</sup>. The photovoltaic characteristics of photoelectrodes, one with nanocrystalline TiO<sub>2</sub> film with NeHS and other with commercially available TiO<sub>2</sub> particles with diameter of 400 nm as the over layer were considered for the study.

The incorporation of NeHS as light-scattering over layer in the DSC resulted in an improvement of the overall efficiency from 7.79 to 9.43% (**Fig. 27**). The NeHS TiO<sub>2</sub> light-scattering included DSC has shown a higher photocurrent density ( $J_{sc}$ ) *i.e.* 15.8 mA cm<sup>-1</sup> than the other one.



**Fig.27** J-V curve of DSC based on the NeHS TiO<sub>2</sub> as light scattering layer.<sup>79</sup>

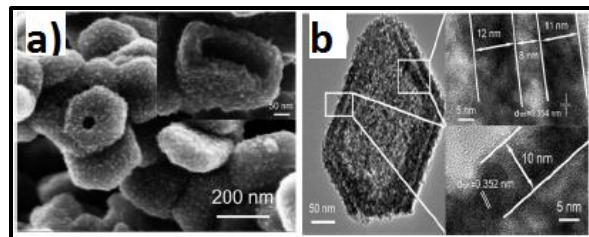


**Fig.28** IPCE spectra of DSCs with 6  $\mu\text{m}$  thick nanocrystalline  $\text{TiO}_2$  film on photoanode (1L) and 12  $\mu\text{m}$  thick film (2L) and light-scattering over layer of NeHS  $\text{TiO}_2$  (1L-NeHS) and over layer of submicrometer-sized  $\text{TiO}_2$  particles.<sup>79</sup>

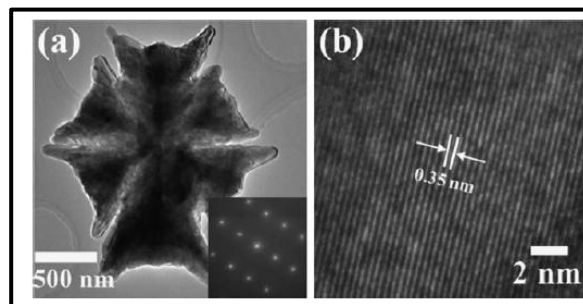
The improvement in  $J_{\text{sc}}$  has been attributed to the enhanced dye loading capability and light scattering effect of NeHS  $\text{TiO}_2$ . The absolute IPCE was higher in the NeHS  $\text{TiO}_2$  sample than the conventional larger  $\text{TiO}_2$  particles (**Fig. 28**).

#### 4.3.7 Hexagonal $\text{TiO}_2$ plates

The packing density is very low for microspherical particles which are conventionally used in DSCs as the scattering layer. This problem was solved by using novel hierarchical  $\text{TiO}_2$  nanoplates with high surface area and sufficient attachment between the particles (**Fig. 29**).<sup>80</sup> By controlling hydrolysis and deposition process of titanium precursors, hierarchical  $\text{TiO}_2$  was prepared by using  $\text{Cd}(\text{OH})_2$  nanoplates.<sup>81-83</sup> The nanoplates<sup>84</sup> were used as light scattering layer in a bilayer DSC model with a thickness of 4  $\mu\text{m}$  and P25 as bottom layer with a thickness of 8  $\mu\text{m}$ . A better reflection was obtained with increased number of mesoshells of nanoplates because of the confinement of light inside the nanoplates. The DSC fabricated with the  $\text{TiO}_2$  nanoplates as scattering layer have shown short-circuit current density of 13.52  $\text{mAcm}^{-2}$  and an open-circuit voltage of 0.79 V. The overall conversion efficiency was about 6.53%.



**Fig.29** a) SEM b) HRTEM image of hexagonal  $\text{TiO}_2$  plate.<sup>80</sup>

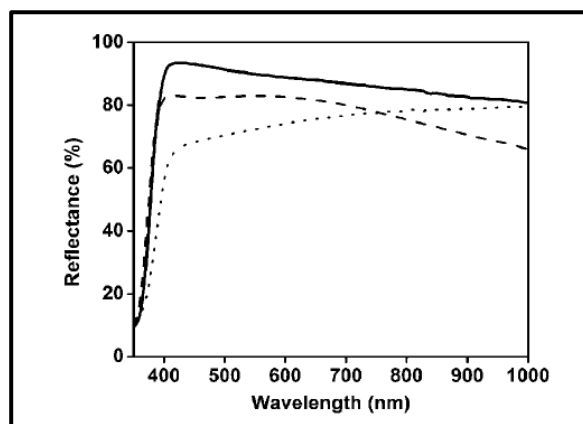


**Fig.30** a) Low magnification and b) high resolution TEM images of a microplate. Inset of (a) is selected-area electron diffraction (SAED).<sup>85</sup>

The surface area of nanoplates was more or less same as P-25 and the efficiency improvement has been attributed to the enhancement of light harvesting by reflection from nanoplates. The single crystalline hexagonal  $\text{TiO}_2$  microplates with good reflectance at the wavelength region 400-900 nm was also used as light scattering layer (**Fig. 30** and **Fig. 31**).<sup>85</sup> A paste of hexagonal  $\text{TiO}_2$  was made and screen-printed over 9.5  $\mu\text{m}$  thick nanocrystalline  $\text{TiO}_2$  layer on FTO glass substrate. The thickness of the scattering layer was about 4  $\mu\text{m}$ . The planar structure of  $\text{TiO}_2$  microplates helps to reflect incident light back to the nanocrystalline  $\text{TiO}_2$ . The photoanode made of nanocrystalline  $\text{TiO}_2$  have shown an overall efficiency ( $\eta$ ) of 6.80 % and short-circuit current density ( $J_{\text{sc}}$ ) of 12.4  $\text{mAcm}^{-2}$ . The introduction of commercial  $\text{TiO}_2$  scattering layer has made an increment in the overall efficiency to 7.45%. Further superior performances were obtained with  $\text{TiO}_2$  microplates as scattering layer with an overall efficiency of 7.91% and short-circuit current density of 14.2  $\text{mAcm}^{-2}$ .

#### 4.3.8 $\text{TiO}_2$ Photonic crystals

The scattering layer with larger particles in DSCs makes it opaque which will disqualify the DSCs as window panes in building integrated photovoltaics.



**Fig.31** Reflectance spectra of  $\text{TiO}_2$  microplate film (solid line), commercial  $\text{TiO}_2$  scattering film (dash line), and  $\text{TiO}_2$  microrod film (dotted line).<sup>85</sup>



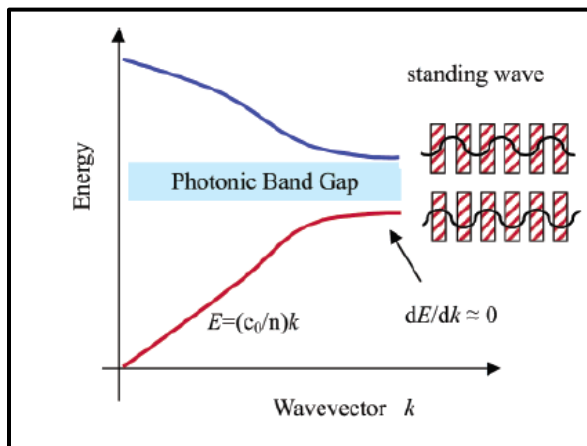


Fig.32 Optical band structure photonic crystal.<sup>95</sup>

The inclusion of photonic crystals (PC) as scattering layer will help to retain transparency of DSCs to some extent. The photonic gap materials can be employed in the form of 3D inverted TiO<sub>2</sub> opal or porous Bragg stacks to enhance light harvesting by retaining transparency.<sup>86-89</sup> Different theoretical possibilities have been put forward for enhancing light absorption such as localization of heavy photons near the edges of photonic bandgaps,<sup>90</sup> Bragg diffraction in periodic lattice,<sup>91</sup> multiple scattering at disordered regions,<sup>92</sup> and via multiple resonant modes.<sup>93</sup> The coupling of photonic crystals in DSCs can be employed as an effective way to improve the conversion efficiency of DSCs. Light propagates in ordinary materials at  $(1-3) \times 10^8$  m/s which will take about one hundred femtoseconds to travel through 10  $\mu\text{m}$  thickness. Confinement of light in the same medium reduces the velocity. Asher et al. first observed Bragg scattering with the three dimensional polystyrene colloidal crystal arrays.<sup>94</sup> They observed stop band in polystyrene colloids, which could be employed to reject light of specific wavelengths. Photonic crystals have three major properties which make them attractive.

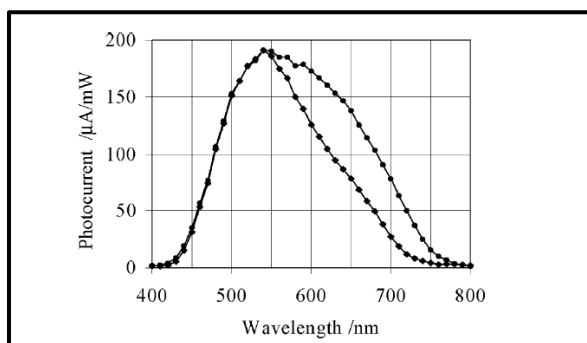


Fig.33 Wavelength dependence of the short-circuit photocurrent. Upper curve shows response of bilayer electrode and lower curve is of conventional nanocrystalline TiO<sub>2</sub>.<sup>95</sup>

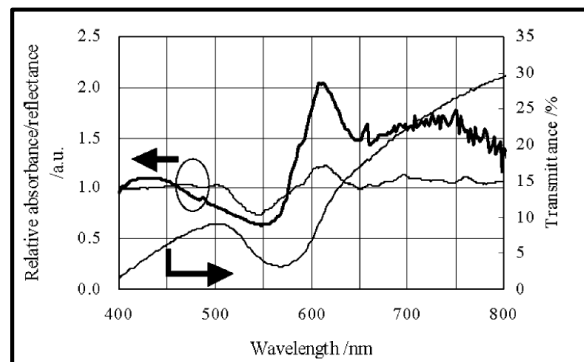
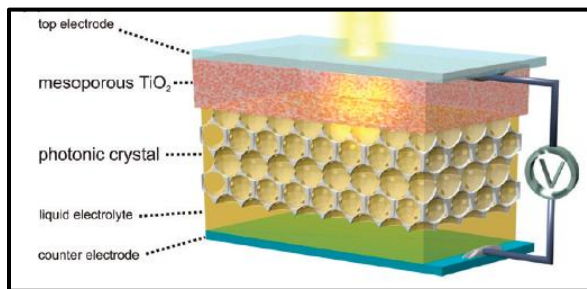


Fig.34 Relative absorption and reflectance spectral ratios. The normalized absorption spectrum of dye-sensitized photonic crystal by the dissolved dye (thick line). The reflectance spectrum of the dye adsorbed photonic crystal was normalized to that before dye absorption (thin line).<sup>95</sup>

- Prohibition of the propagation of particular wavelengths.
- Birefringent and anisotropic dispersion.
- Nonlinear dispersion.

Nishimura et al. studied red absorbance and photocurrent enhancement in DSCs where TiO<sub>2</sub> is coupled to photonic crystals.<sup>95</sup> Dispersion curve ( $E$  vs  $k$ ) shows the group velocity of light decrement near to the wavelength of stop band<sup>96</sup> (see Fig. 32). Depending on energy, light waves would be localized in different sites of PC structure and absorber in high dielectric medium will interact more with red part of the stop band. Photocurrent response has been considerably improved at longer wavelength region over using only nanocrystalline TiO<sub>2</sub> film (Fig. 33). As in Borrmann effect in X-ray scattering, red part of the spectrum absorber will “borrow” intensity from blue part (Fig. 34).<sup>97</sup> In the study reported by Nishimura et al. TiO<sub>2</sub> photonic crystals (inverse opal) were coated on the FTO. Replication of colloidal crystal templates was the method adopted by them.<sup>98-104</sup> The suspended monodisperse carboxylate-modified polystyrene latex spheres were immersed vertically in the colloidal solution which was kept at 55 °C for 1 to 2 days to evaporate the suspension completely. The colloidal crystal film was formed on the surface with a thickness of about 215 nm.<sup>105</sup> Liquid-phase deposition (LPD) method was used to infiltrate titanium dioxide into voids retaining its periodicity.<sup>106-111</sup> The latex spheres were removed by calcinations at 400 °C for 8 h, resulted in mesoporous TiO<sub>2</sub> with air holes. Thickness of titanium walls of the inverse opals can be controlled in second LPD to tune stop band precisely.<sup>112</sup> The cell was fabricated with a conventional TiO<sub>2</sub> coating over a 3  $\mu\text{m}$  thick photonic crystal. The absorption spectra of dye-sensitized photonic crystals and the same dye in ethanol solution were compared. Absorbance was increased at red edge of the stop band of dye sensitized photonic crystals as compared to the dye in solution. A suppression of absorbance for shorter wavelength was also

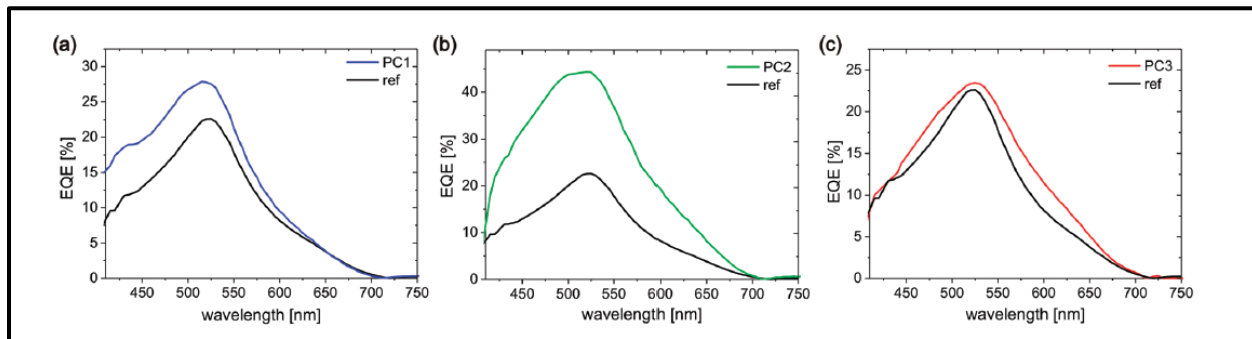


**Fig.35** Schematic of 3D photonic crystals as DSCs overlayer.<sup>114</sup>

observed (see **Fig 34**). In this study, the bilayer electrode with nanocrystalline layer in contact with FTO was found to be unsuccessful.

Miguez et al. postulated that major part of the spectrum can only be tapped when the nanocrystalline TiO<sub>2</sub> is placed in between the photonic crystal and the transparent conductive oxide substrate.<sup>113</sup> The connectivity between the nanocrystalline TiO<sub>2</sub> and the photonic crystal was one of the problems in implementing photonic crystal as the over layer. Guldin et al. constructed optically active 3D photonic crystal as overlayer (**Fig. 35**).<sup>114</sup> Copolymer-derived mesoporous underlayer was prepared on conductive transparent substrate. The TiO<sub>2</sub> inverse opals was formed as the over layer with different pore sizes. To analyze the photocurrent generation, IPCE or quantum efficiency was plotted against the wavelength of both bilayered and single layered DSCs. Here for the reference, a relatively thin layer of TiO<sub>2</sub> with 465 nm thickness was chosen to avoid cracks on the film. A double layer DSC was fabricated with 1.4 μm thick PC overlayer on the top of 465 nm thick TiO<sub>2</sub> underlayer. The external quantum efficiency (IPCE) of DSCs with three different PC over layer is shown in **Fig. 36**.

The one-dimensional photonic crystals (1D PCs) are preferred over 3D PC because of their easiness in assembling a thin layer on nc TiO<sub>2</sub> active layer. The thick 3D PCs (5-10 microns) have adverse effect on conversion efficiency as charge transport would be limited by increased recombination in the cell.



**Fig.36** External Quantum Efficiency (%) of DSCs with different PC as overlayer.

PC1:3D inverse opal with diameter 240nm. (b) PC2: 3D inverse opal with diameter 260nm. (c) PC3: 3D inverse opal with diameter 350nm.<sup>11</sup>

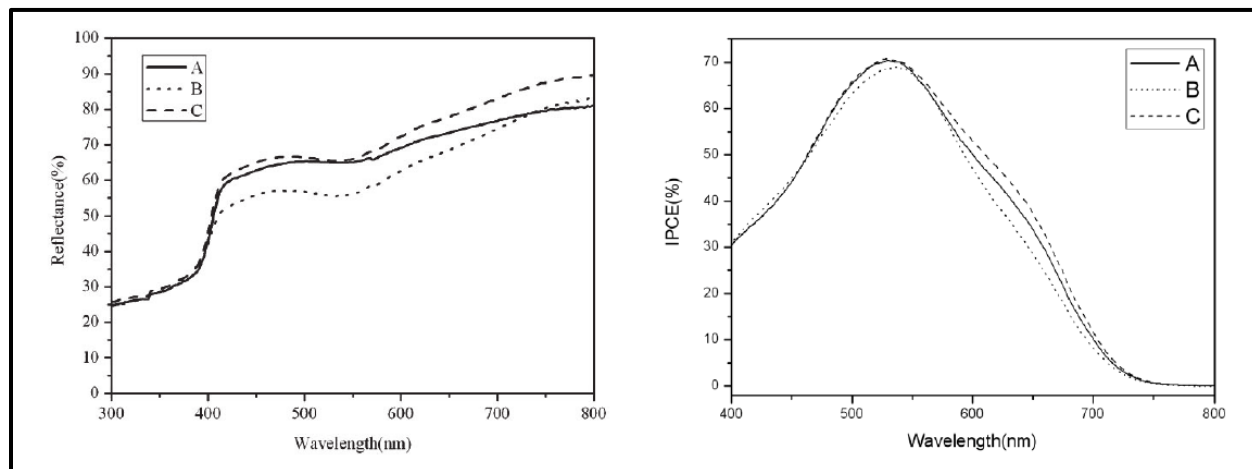
#### 4.3.9 Cubic CeO<sub>2</sub> nanoparticles

Cerium dioxide (CeO<sub>2</sub>) cubic nanostructure with high refractive index is one of the potential candidates for enhancing light harvesting efficiency by light scattering. The mirror-like facets of CeO<sub>2</sub> are responsible for the scattering effect. Yu et al. have used CeO<sub>2</sub> nanoparticles of about 400 nm as overlayer in DSCs.<sup>115</sup> With a 1.5 μm thick top layer on mesoporous TiO<sub>2</sub>, conversion efficiency has been improved to 8.6% from 7.3%. However, a reduction in dye loading was observed when CeO<sub>2</sub> was introduced as the top layer.<sup>116</sup> The diffused reflectance was very high for the bilayer photoanodes compared to pure TiO<sub>2</sub> nanocrystalline film. The CeO<sub>2</sub> overlayer enhanced the light reflectance in the wavelength region 400-800 nm (light reflectance increased over 70%). The highest current density ( $J_{sc}$ ) with CeO<sub>2</sub> as overlayer was 19.3 mAcm<sup>-2</sup>. The highest efficiency reported with CeO<sub>2</sub> as top layer was 8.6% and  $V_{oc}$  observed was about 0.69 V.

#### 4.3.10 Surface modified particles in scattering layer

The surface of particles used as scattering layer can be modified to obtain better conversion efficiency for DSCs. Choi et al. coated Al<sub>2</sub>O<sub>3</sub> on large TiO<sub>2</sub> particles.<sup>117</sup> Sol-gel method was used to modify the surface of large TiO<sub>2</sub> scattering particles.<sup>118</sup> Colloidal Al<sub>2</sub>O<sub>3</sub> of various concentrations with TiO<sub>2</sub> powder was stirred to obtain TiO<sub>2</sub> suspension. The coating of Al<sub>2</sub>O<sub>3</sub> on TiO<sub>2</sub> depends on pH of TiO<sub>2</sub> suspension and concentration of Al<sub>2</sub>O<sub>3</sub>. A few nanometer thick Al<sub>2</sub>O<sub>3</sub> nanoparticles were coated over the 300-800 nm sized TiO<sub>2</sub> particles. The photoelectrode was made of commercial TiO<sub>2</sub> paste as underlayer and large Al<sub>2</sub>O<sub>3</sub> coated TiO<sub>2</sub> as the scattering layer. The reflectance spectrum revealed that the total reflectance of TiO<sub>2</sub> film was considerably increased in the region 380-800 nm when a light scattering layer was introduced (**Fig. 37a**). The IPCE showed an increased photocurrent response for good oxide surface coated TiO<sub>2</sub> scattering layer (sample C) (**Fig. 37b**).





**Fig.37 a)** Reflectance spectrum of dye sensitized TiO<sub>2</sub> film of samples with light scattering layer **b)** IPCE of DSSC with light scattering layer. A: sample with a non-coated scattering layer, B: sample with TiO<sub>2</sub> scattering layer (with 0.235 Al<sub>2</sub>O<sub>3</sub> Concentration in wt% and pH value of 12), C: sample with TiO<sub>2</sub> scattering layer (with 0.016 Al<sub>2</sub>O<sub>3</sub> Concentration in wt% and pH value of 7).<sup>117</sup>

The best efficiency (8.71%) was reported for the sample C.

## 5. Dual function materials in DSCs

In the above sections, we examined the materials which were employed for scattering light in DSCs where the scattering material is added as an additional layer or as component. But some materials provide both higher surface area as well as light scattering capability.

### 5.1 Nanocrystalline spherical aggregates

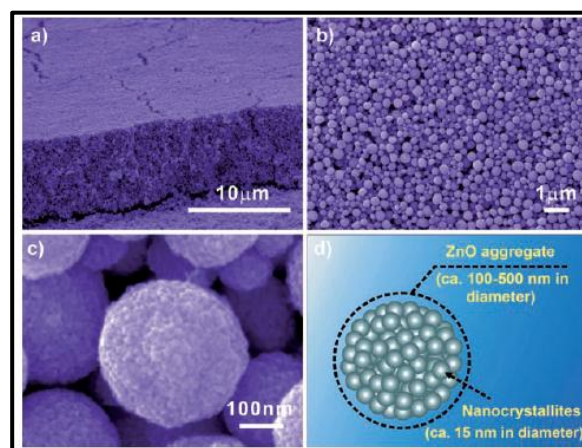
Aggregates of oxide nanocrystallites would satisfy dual requirements such as high internal surface area and light scattering ability. Nanocrystallite spherical aggregates consist of spherical assembly of nano-sized oxide crystallites with submicron diameter. The nano-sized crystallites *i.e.* building blocks of aggregates provide higher internal surface area nearly equal to the nanocrystalline film. Light scattering capability is derived from the submicron size of the aggregates. Scattering by aggregates would increase travelling distance of light within the photoanode and interaction of light with the dye is enhanced. The light scattering by nanocrystalline spherical aggregates can be analytically described by the Lorentz-Mie-Debye resonance theory in which light scattering efficiency  $Q_{sca}$  given by

$$Q_{sca} = \frac{2}{x^2} \sum_{n=1}^{\infty} 2n + 1 (a_n^2 + b_n^2) \quad (21)$$

where  $x$  is defined as size parameter,  $x = 2\pi r / \lambda$ ,  $r$  is the radius of sphere,  $\lambda$  is the incident wavelength and  $a_n$ ,  $b_n$  are the complex Mie scattering coefficients. One can see that sample consisting large aggregates of diameter  $\sim 350$  nm gives maximum scattering efficiency.

### 5.1.1 ZnO aggregates

Over the years TiO<sub>2</sub> has been the material of choice for DSCs. Recently, ZnO as a material alternative to TiO<sub>2</sub> has been widely investigated. The following reasons make ZnO a potential candidate for DSC: a) band-gap of ZnO is similar to TiO<sub>2</sub> (3.2eV), b) higher electron mobility of ZnO ( $\sim 115$ - $155$  cm<sup>2</sup>V<sup>-1</sup>s<sup>-1</sup>,<sup>119</sup> whereas that of TiO<sub>2</sub> is only  $\sim 10^{-5}$  cm<sup>2</sup> V<sup>-1</sup>s<sup>-1</sup>).<sup>120</sup> The hierarchically-structured ZnO film consisting of secondary ZnO colloids of  $\sim 300$  nm in diameter that are made up of primary ZnO nanoparticles of  $\sim 20$  nm in diameter would encourage light scattering in DSC (**Fig. 38**). The first DSC application of ZnO aggregate was introduced by Chou et al.<sup>121</sup> The ZnO aggregates were synthesized by hydrolysis of zinc salt in polyol medium at 160 °C.<sup>122,123</sup>



**Fig.38 a)** Cross sectional SEM of ZnO aggregate film **b)** SEM image of top view of polydisperse ZnO aggregate film **c)** A magnified SEM image of ZnO aggregate **d)** Schematic of ZnO aggregate comprising of closely packed nanocrystallites.<sup>125</sup>

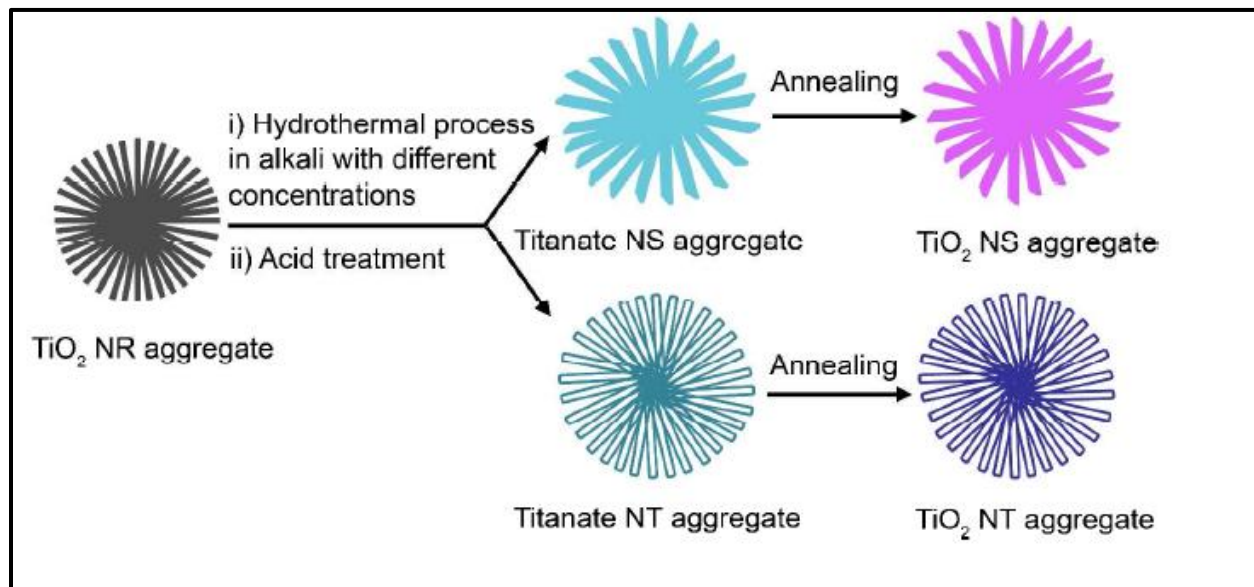


Fig.39 TiO<sub>2</sub> NS aggregate and NT aggregate were prepared from TiO<sub>2</sub> NR aggregate.<sup>134</sup>

The aggregates were separated centrifugally from the solvent. The photoelectrode was made from suspension solution by drop casting method.<sup>124</sup> The specific surface area of aggregates was reported as 80 m<sup>2</sup> g<sup>-1</sup> by BET measurement.<sup>125</sup> The large surface area provides a good dye adsorption on aggregate photoelectrode film. The optical absorbance spectrum of aggregate photoelectrode film has an additional absorption in the visible region due to light scattering, other than the intrinsic absorption of ZnO. The photocurrent peak in IPCE at 360 nm was due to direct absorption of ZnO. The sample 1 with perfect aggregates has shown highest efficiency of about 5.4%. It was found that light scattering is more intensive with aggregate size closer to incident wavelength. Memarian et al. have achieved a record efficiency of 7.5% by spray pyrolysis method.<sup>126</sup>

### 5.1.2 TiO<sub>2</sub> aggregates

The nanobeads or nanoporous microspheres are considered as TiO<sub>2</sub> aggregates.<sup>127-130</sup> The TiO<sub>2</sub> aggregates are assemblage of nanocrystallites. The TiO<sub>2</sub> beads are fabricated *via* hydrothermal treatment of amorphous TiO<sub>2</sub> beads in ammonia. The TiO<sub>2</sub> beads produced had a surface area of 89 m<sup>2</sup> g<sup>-1</sup> and diameter of 830 nm. A comparison of IPCE with P25 has shown a high performance for the TiO<sub>2</sub> mesoporous nanobeads. A higher efficiency of about 9.1 % was reported with TiO<sub>2</sub> nanobeads without TiCl<sub>4</sub> treatment (after TiCl<sub>4</sub> treatment efficiency reached at 10.6%). Another advantage of TiO<sub>2</sub> aggregate is the higher internal surface area compared with P25 which resulted in more dye loading. It was also noted that lifetime of photoelectrons in photoelectrode film is longer than that in conventional photoelectrode film because of the close packing of nanocrystallites within the TiO<sub>2</sub> aggregates. Light scattering

and the above mentioned uniqueness of TiO<sub>2</sub> aggregates contribute to the efficiency enhancement.

TiO<sub>2</sub> and SiO<sub>2</sub> aggregates synthesized by Molten Salt Method (MSM) were used in DSC applications.<sup>131,132</sup> The TiO<sub>2</sub> agglomerated nanostructures produced by MSM were of diameter ~ 100-300 nm. However, TEM studies revealed that the aggregates were actually composed of smaller nanoparticles with diameter of 5 nm. These particles were responsible for higher surface area of 200 m<sup>2</sup> g<sup>-1</sup>. Using MSM TiO<sub>2</sub> a power conversion efficiency of 7.48 % was obtained. The IPCE is enhanced at small wavelength region (< 600 nm) due to high dye loading and that at longer wavelength region (> 600 nm) due to light scattering effect by aggregates of size 100-500 nm.

#### Tailored spherical TiO<sub>2</sub> aggregates

Substitution of nanorods or nanotubes as building units for the nanocrystallites in the spherical aggregates could provide excellent electron transport property without compromising on surface area and light scattering capability. Liu et al. have prepared TiO<sub>2</sub> nanosheet (NS) aggregates and TiO<sub>2</sub> nanotube (NT) aggregates from spherical nanorod (NR) aggregates.<sup>133,134</sup>

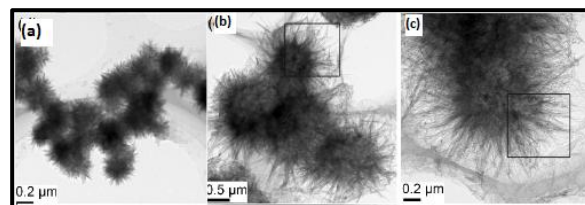
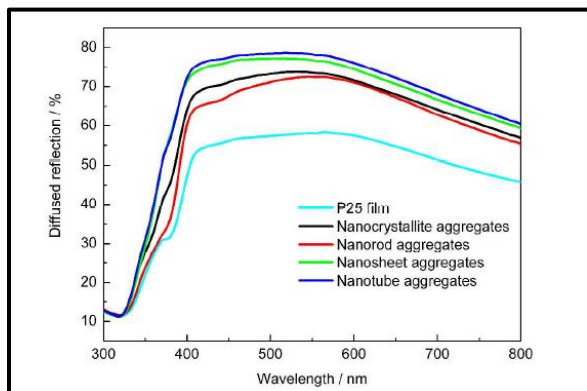
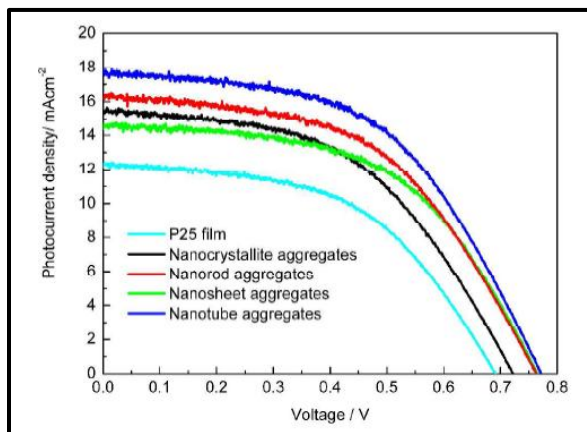


Fig.40 TEM images of a) NR aggregate b) NS aggregate c) NT aggregate.<sup>134</sup>

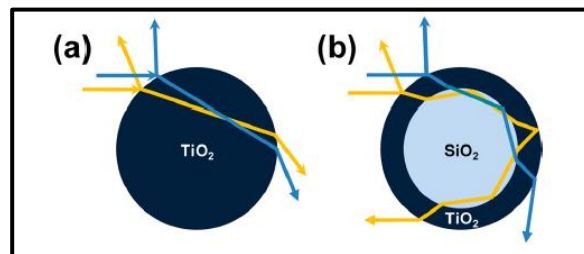


**Fig.41** Diffused reflectance spectra of different tailored TiO<sub>2</sub> aggregate films.<sup>134</sup>

The TiO<sub>2</sub> NS aggregates were prepared by following procedure (Fig. 39). The TiO<sub>2</sub> nanorod aggregates were undergone hydrothermal treatment in 5M sodium hydroxide at 140°C for 24h and further washing of sample by diluted hydrochloric acid. This was further annealed at 400°C for 30 min. The TiO<sub>2</sub> NT aggregates were also synthesized following similar route of preparing NS aggregates except changing concentration of hydrothermal reaction solution. It is interesting that dimension of NT and NS aggregate have significantly improved from 400 nm (which is the dimension of nanorod aggregate) to 1-2 μm which will ultimately increase the scattering capability of nanostructure. The TEM images have shown in Fig. 40. The diffused reflectance spectra have shown better light scattering property at the infrared region. A higher reflectance was obtained for the NT and NS aggregate compared to the NR aggregate (Fig. 41). This higher light scattering is due to the larger dimension sizes of the NT and NS aggregate. The high BET surface area have reported for NR aggregates (106.7 m<sup>2</sup>/g), NT aggregates (134.8 m<sup>2</sup>/g) and NS aggregates that resulted in better dye loading.



**Fig.42** I-V characteristics of DSCs using different tailored TiO<sub>2</sub> aggregate.<sup>134</sup>



**Fig.43** Schematic of light scattering by a) TiO<sub>2</sub> nanoparticle and b) TiO<sub>2</sub>-SiO<sub>2</sub> core/shell nanoparticle.<sup>138</sup>

The NT aggregate DSC has reported a best efficiency of 7.48% with J<sub>sc</sub> of 17.82 mA/cm<sup>2</sup> among this class of TiO<sub>2</sub> materials. The NS and NR aggregate DSC has reported an efficiency of 6.45 and 6.82 % respectively (Fig. 42).

## 6. Other innovations for light scattering in DSCs

### 6.1 Core-Shell materials for light scattering

Theoretical studies have been done by several research groups on the light scattering effect of core/shell nanoparticles.<sup>135, 136</sup> The Mie theory explains generally the light scattering by spherical particles. Arden and Kerker's scattering theory for concentric spherical particles is well suited to explain the light scattering by coated spheres (Core/shell particles). Consider a spherical, coated particle of inner radius X, outer radius Y and coated thickness Z (Y-X) which is illuminated by a monochromatic beam of wavelength λ<sub>0</sub>. The refractive indices of the core, shell, and surrounding medium were assigned as n<sub>1</sub>, n<sub>2</sub>, n<sub>0</sub> respectively. The intensity I(θ) of scattered radiation by a single particle at an angle θ for incident radiation that is polarized perpendicularly to the scattering plane given by

$$I(\theta) = I_0 \frac{\lambda_0^2}{4(\pi r)^2} S(\theta)^2 \quad (22)$$

Where r is the distance between the particle and point of interest, I<sub>0</sub> is the incident radiation and S<sub>1</sub>(θ) is infinite series.

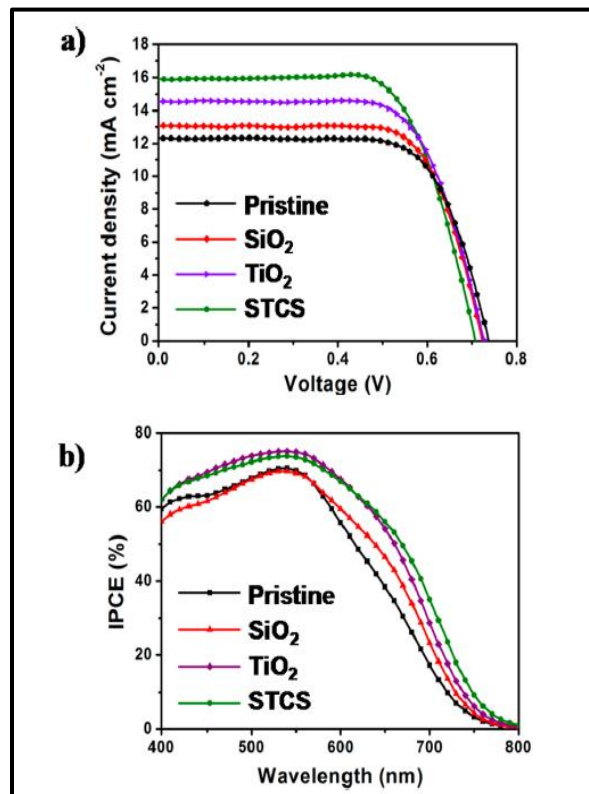
$$S_1(\theta) = \sum_{n=1}^{\infty} \frac{2n+1}{n(n+1)} [a_n \pi_n(\theta) / b_n \tau_n(\theta)] \quad (23)$$

The amplitudes a<sub>n</sub> and b<sub>n</sub> are functions of the size parameters x = kX and y = kY

$$\text{Where } k = 2\pi/\lambda = 2\pi n_0/\lambda_0.$$

High scattering efficiency was reported by Hsu et al. using core/shell nanoparticles with a diameter of 0.8-1.0 μm and shell thickness of 0.2-0.3 μm.<sup>137</sup> Son et al. prepared SiO<sub>2</sub>/TiO<sub>2</sub> core/shell nanoparticles using SiO<sub>2</sub> template by sol-gel method for DSC application.<sup>138</sup> Monodisperse SiO<sub>2</sub> nanoparticles in different diameters were synthesized by





**Fig.44 a)** Current density – voltage curve and **b)** IPCE of DSC with SiO<sub>2</sub> nanoparticle, TiO<sub>2</sub> nanoparticle and TiO<sub>2</sub>/SiO<sub>2</sub> core/shell nanoparticle as scattering layer.<sup>138</sup>

seed regrowth method.<sup>139</sup> The TiO<sub>2</sub> nanoparticles were coated over SiO<sub>2</sub> using sol-gel method. The light scattering in TiO<sub>2</sub> and SiO<sub>2</sub>/TiO<sub>2</sub> core/shell nanoparticles was explained based on geometrical optics (**Fig. 43**)<sup>140</sup>. In SiO<sub>2</sub>/TiO<sub>2</sub> core/shell nanoparticles, light undergoes back scattering unlike the forward scattering by larger TiO<sub>2</sub> nanoparticles. This will lead to optical confinement in DSCs and enhancement of the conversion efficiency. The SiO<sub>2</sub> nanoparticles with a diameter of 240 nm as core in SiO<sub>2</sub>/TiO<sub>2</sub> core/shell nanoparticles when used as scattering layer in photoelectrode has shown a better performance compared with SiO<sub>2</sub> or TiO<sub>2</sub> as scattering layer materials. The efficiency of 7.9 % was obtained with improved current density of 15.9 mAcm<sup>-2</sup> (**Fig. 44**).

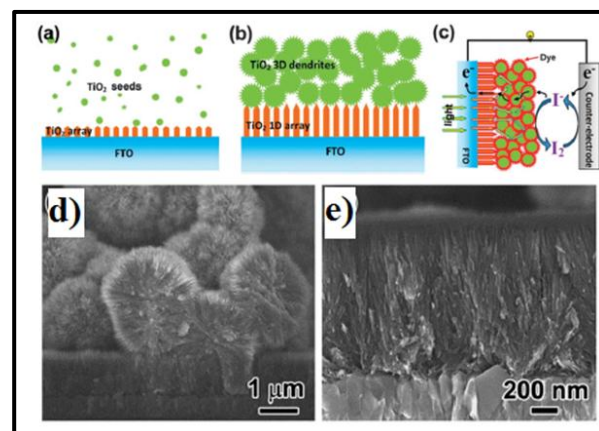
### 6.2 Double light-scattering film

Zhao et al. fabricated DSCs with double light-scattering layer consisting of TiO<sub>2</sub> nanosheets as the underlayer and TiO<sub>2</sub> hollow spheres as the overlayer.<sup>141</sup> The DSC with double light scattering layer has shown superior performance compared to TiO<sub>2</sub> hollow sphere film alone or TiO<sub>2</sub> nanosheet film. An improved efficiency of 5.03 % was reported which was 23% higher than that of TiO<sub>2</sub> hollow sphere film alone ( $\eta=3.6\%$ ) and 8.3% higher than that of TiO<sub>2</sub> nanosheet film ( $\eta=4.31$ ).

### 6.3 Surface Plasmon effects in DSCs

The optical properties of metal nanoparticles are influenced by collective motion of free electrons within a small sphere.<sup>142</sup> Electromagnetic wave induced collective motion of electrons is called as surface plasmons. The plasmonic resonance induces a dipole moment in metal nanoparticles which will eventually enhance the electric field around it and hence the optical extinctions such as absorption and scattering are also increased. Solid-state DSCs (ss-DSCs) are fabricated using hole-transport materials (HTM) replacing conventional electrolytes. It will facilitate tuning open circuit voltage by adjusting HOMO of HTM. The higher recombination<sup>143</sup> and incomplete hole filling of HTM<sup>144,145</sup> would limit thickness of active layer to 2  $\mu\text{m}$ . This thickness is much less to absorb light effectively. The plasmonic effects have been employed to improve light absorption in active layers of solar cells.<sup>146</sup> Metal nanoparticles enhance the light absorption of solar cells both by localized surface plasmon and by light scattering. Ding et al. used nanoimprint lithography (NIL) to incorporate plasmonic back reflectors into the mesoporous TiO<sub>2</sub> active layer.<sup>147</sup> The plasmonic back reflectors were designed as 2D array of metallic domes in ss-DSCs. However the challenge before the scientific community is identifying an inexpensive route to fabricate the plasmonic structures.

The metal nanoparticles (NPs) were in direct contact with dye and the electrolyte in earlier surface plasmon assisted DSCs. It normally resulted in an enhanced recombination, back reaction of photogenerated electrons and corrosion of metal NPs by the electrolyte.<sup>148-154</sup> Qi et al. has succeeded in fabricating highly efficient plasmon-enhanced DSCs using core-shell (Ag@TiO<sub>2</sub>) nanostructures.<sup>154</sup> They had used Ag as core and TiO<sub>2</sub> as thin (about 2 nm) shell.



**Fig. 45** Schematic of one step synthesis of bilayer formation (**a,b** & **c**) **a)** formation of TiO<sub>2</sub> seeds on FTO **b)** growth and assembly of 1D nanowire and 3D nanostructures **c)** bilayer photoanode of DSCs **d)** Magnified SEM image of 3D dendritic nanostructure **e)** SEM image of bilayer photoanode.<sup>148</sup>

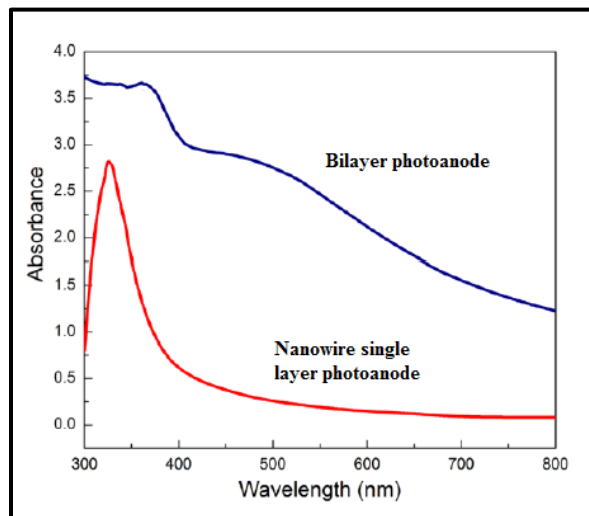


Fig.46 UV-Vis spectra of photoanodes with nanowire single layer and bilayer.<sup>148</sup>

The shell covering over metal NPs resulted in a reduced recombination and back reaction. Moreover it would also protect metal NPs from corrosion by electrolyte. The incorporation of 0.1 wt.% Ag@TiO<sub>2</sub> in thin photoanode provided increase in efficiency from 7.8% to 9.0%.

Further rod-like Ag and Au can also be utilized for developing highly efficient plasmon-enhanced DSCs. Longitudinal plasmon absorption of rod-like nanostructures can be tuned by controlling the aspect ratio of the rod.<sup>155-157</sup> Chang et al. used Ag<sub>2</sub>S-encapsulated Au nanorods (AuNR@AgS) in DSCs to enhance performance photocurrent at longer wavelength.<sup>158</sup> A 37.6% increment in photocurrent was observed at 600-720 nm range for AuNR@AgS incorporated photoanode. A power conversion efficiency about 7.1 % obtained for AuNR@AgS incorporated (11 μm thick) photoanode against 5.8% efficient TiO<sub>2</sub> alone photoanode DSCs.

#### 6.4 1D-3D nanostructured rutile TiO<sub>2</sub> bilayer photoanodes

The highest efficiency for DSCs reported for rutile TiO<sub>2</sub> photoanode with a bilayer structure consisting of one-dimensional (1D) nanowire as underlayer and three-dimensional dendritic microsphere as top layer was 7.2%.<sup>159</sup> The vertically aligned 1D nanowires as bottom layer exhibited high photoinjected collection efficiency, fast electron transport and long electron diffusion length. The top layer of TiO<sub>2</sub> dendritic nanostructure with 1-3 μm diameters provided high specific surface area and in-turn high dye loading and light scattering.

The above favorable factors result in efficient bilayer DSCs. The one-step hydrothermal synthesis was used to fabricate photoanode where a mixture of titanium tetraisopropoxide (TTIP) and cetyltrimethylammonium bromide (CTAB) together with ethylene glycol (EG) was used for hydrothermal treatment. The FTO was kept in a Teflon-lined autoclave with the conducting side facing up.

The formation of the bilayer is explained in Fig. 45. The size and morphology nanostructures can be adjusted by varying TTIP<sub>aq</sub>/EG ratio. A fill-factor of 68% and J<sub>sc</sub> of 14.51 mA cm<sup>-2</sup> were obtained for best performing cells. The UV-Vis spectra reveal the pronounced light scattering effect exhibited by the bilayer photoanode (Fig. 46).

#### 6.5 Quintuple-shelled SnO<sub>2</sub> hollow microsphere as light scattering layer

The SnO<sub>2</sub> has two major advantages over TiO<sub>2</sub> as DSC photoanode material: 1) it possesses a wider band-gap (3.8 eV), and 2) dye degradation is minimal because it creates fewer oxidative holes in its valence band under UV illumination. However the power conversion efficiency is limited for SnO<sub>2</sub>-based DSCs due to following reasons:

- 1) a 300 mV positive shift in conduction band with respect to TiO<sub>2</sub> encourages interfacial recombination.
- 2) a lesser adsorption of dyes with acidic carboxylic groups.

Dong et al. reported a power conversion efficiency of 9.53% with quintuple-shelled SnO<sub>2</sub> hollow microsphere with closed exterior double walls (5S-SnO<sub>2</sub>-HMs-CDS) as a light scattering layer.<sup>160</sup> The alkali-treated carbonaceous microspheres used as templates for synthesis of 5S-SnO<sub>2</sub>-HMs-CDS. The efficiency increased by 30.7% when 5S-SnO<sub>2</sub>-HMs-CDS with high light reflectance (%) was used as overlayer on P25 layer.

#### 6.6 Multi-stack dye solar cells

Marco et al. engineered a photoanode with a combination of three mesoporous stacks made of different sized nanorods.<sup>161</sup> In this approach, the bottom layer consisted of small nanorods capable of excellent dye adsorption; an intermediate layer of nanorods with a high aspect ratio that ensured good electron transport and an upper layer consisting of large nanorods which guaranteed superior light scattering. These favorable conditions effectively contributed to a higher efficiency of 10.23% for the DSC. The thickness of the multi-stacked film was about 17 μm. A high current density of 18.10 mA cm<sup>-2</sup> and a V<sub>oc</sub> of 0.81 V were obtained for the three-layered cell.

#### 7. Conclusions

Light scattering is an effective approach to enhance light absorption in a DSC photoelectrode. In general light scattering in DSCs is achieved in three ways: 1) mixing large particles with nanocrystalline TiO<sub>2</sub> slurry; 2) placing large particles as overlayer on TiO<sub>2</sub> photoelectrode film (double layer structure); 3) single layer of dual function materials. Theoretically increase in film thickness would enhance optical absorption of light. But the maximum allowed thickness of DSC photoanode film is limited by electron diffusion length. The first approach mentioned



above would lead to decrease in internal surface area of the photoelectrode film. The additional overlayer for light scattering would retain internal surface area; however, the light is reflected only once. Dual function materials such as hierarchical 3D nanostructures provide both internal surface area and multiple reflections of incident light. The inclusion of scattering layer in DSCs would make it opaque that restrict the use of DSCs building-integrated photovoltaic applications. However, this can be sorted out by using photonic crystal as light scattering layer in DSC which makes it transparent to some extent. Commercializing DSC at low cost is one of the challenging tasks in solar energy conversion scenario. Inclusion of light scattering layer has become a common practice in the fabrication of high performance DSCs. Minimizing the thickness of active layer is highly recommended which would reduce recombination and usage of costly commercial dyes (~ \$300/g). The diffused reflectance of NT and NS aggregate is around 80% at visible region. The microplate nanostructure has shown a better reflectance of about 80-90%. But these have a lower surface area which will limit the PV performance of DSCs. We saw that optimization of haze factor by adding light a scattering layer is also important criterion for obtaining best PV performance in DSCs.

**Table 3** summarizes effect of different nanostructures for light scattering in DSCs.

**Table 3.** Different nanostructures as scattering layer and their impact in DSCs performance.

Nanostructures as scattering layer	Current density	Efficiency ( $\eta$ )	Increase in efficiency (in percentage)
Large spherical TiO <sub>2</sub> particles (with diameter ~ 400 nm)	21 mA cm <sup>-2</sup>	11.1%	52%*
TiO <sub>2</sub> nanospindles	16.4 mA cm <sup>-2</sup>	8.3%	43%
Nano-embossed hollow spherical TiO <sub>2</sub>	15.8 mA cm <sup>-2</sup>	9.43 %	21%
CeO <sub>2</sub> nanoparticles with diameter ~ 400 nm	19.3 mA cm <sup>-2</sup>	8.6%	17.8%
Quintuple-shelled SnO <sub>2</sub> microsphere	20.1 mA cm <sup>-2</sup>	9.53%	30.7%#
Multi-stacked DSCs	18.1 mA cm <sup>-2</sup>	10.23%	40.29% <sup>s</sup>

\* # Here 7.29% is taken as efficiency of a standard P25 photoanode DSC.

It clearly depicts that spherical nanoparticles with 400 nm diameter as light scattering layer show superior performance among others. Only major nanostructures that could provide considerable increment in efficiency are only listed in **Table 3**.

Multi-stacked DSCs also show superior PV performance compared to single and double layer DSCs. But the thickness of DSCs exceed more than 15 $\mu$ m (i.e. thicker solar cells) for the best performing solar cell. The spherical TiO<sub>2</sub> aggregate with nanotubes, nanorods and nanosheets as building units exhibits better electron transport property without giving up the internal surface area and light scattering. Designing TiO<sub>2</sub> tailored aggregates such as nanorod aggregate, nanosheet aggregate and nanotube aggregate with high surface area and light scattering capability as dual functional materials in DSCs would be a viable model for obtaining highly efficient DSCs.

### Acknowledgement

The authors thank the Ministry of New and Renewable Energy, Govt. of India for financial assistance.

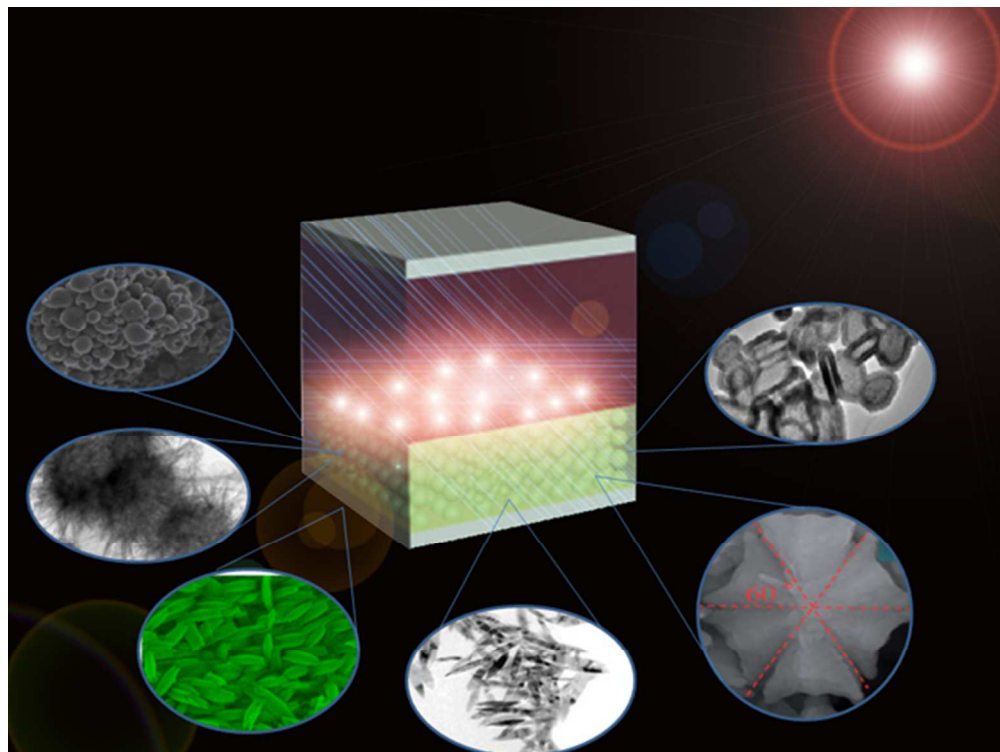
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A state-of-the-art review highlighting the theory, role and the materials of scattering layer in dye-sensitized solar cells.

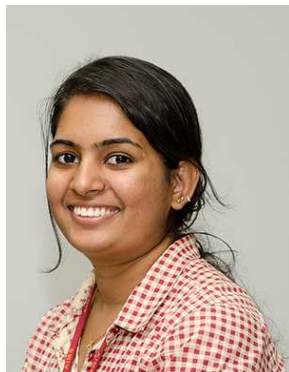
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