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Surface passivation assisted lasing emission in the quantum dots doped cholesteric liquid crystal resonating cavity with polymer template

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Copolymer have been used as template for enhance optical properties of core-shell CdS/ZnSe quantum dots doped cholesteric liquid crystal. The lasing behavior is achieved at relatively low laser excitation energy threshold at room temperature. Low-threshold value and high-optical stability make this optical system promising in visible laser emitter applications.

Cholesteric liquid crystals (CLCs) are materials with periodic dielectric structure, which leads to the formation of a photonic bandgap (PBG) [1]. CLCs phase inherently exhibits a selective reflection that is related to the Bragg optics associated with the helical structure formed from a one dimensional rotation of the birefringence axis [2]. They have attracted much interest for their promising applications in many optical fields (such as lasing emission) [3-6]. Quantum dots (QDs) are semiconductor nanocrystals whose electrons and holes are quantum-confined in all three spatial dimensions. They have attracted much attention during the past few decades due to their unique optical properties as distinguished from the corresponding bulk materials and have shown great potential in the fabrication of next generation optoelectronic devices [7, 8]. In recent years, highly efficient quantum emitters have proven to be very attractive light sources [9]. Special optical emission features from the QDs as active medium are naturally expected due to their unique quantum structure different from the conventional organic dye molecules [10]. The optical characteristics of a quantum dot doped CLCs microresonator inevitably changes with varying operation template, which greatly determines its applications in the fields of requiring specific and stable performance. However, only a few studies have been reported on the soft matter dependence of optical characteristics of QDCLCs device [10-12]. In contrast with conventional organic dyes or fluorescent proteins, QDs possess unique optical properties. They have broad absorption band, narrow emission band, size-tunable emission from visible to near infrared range, superior brightness, long fluorescent lifetime, and more importantly, they are highly resistant to photobleaching and have relatively large surface area for functionalization [13]. Another problem is that QDs tend to aggregate inside liquid crystal and are often trapped or sequestered in organelles. Owing to these cationic dendrimers and micelles are

unstable (e.g., aggregation) in soft matter or culture media and are decreasing emission performance [14,15]. Such surface passivation has been achieved using organic capping agents as well as the formation of inorganic QDs systems such as CdSe/ZnSe [16-19]. The 1-hexanethiol ($C_6H_{14}S$) surfactant also improves the dispersible and stability of the coated QDs in this case. Recently, many researchers successfully combined the liquid crystal materials with polymer template technology to fabricate optical device using template agents as structure directing agents [20, 21]. In the last years, the polymer template assisted synthesis has been proven as a promising method to obtain soft materials with defined chemical and optical properties [22, 23].

In this work, we describe the template design and procedure to control the lasing emission of QDCLCs. This report includes guidelines to comprehend the correlation between the controlled QDCLCs stability and the lasing characteristics by band edge effect, providing available route to fabricate novel active laser devices in template system. At the same time, the stable pumped energy of 11 μ J /pulse is obtained, which is 2 times better than that of reference QDCLCs band-edge laser without template. Broad gain profile and template polymerization effect in QDCLCs laser are likely playing an important role in reduction of the quantum dots aggregate sensitivity of the lasing efficiency.

Figure 1(a) and S2 shows photoluminescence (PL) and UV absorption spectra of CdSe/ZnSe QDs (before surface passivation is black line, after surface passivation is red line) dispersed in toluene. As shown, both the PL and UV spectra from the core/shell QDs essentially maintain their overall shape with absorption decay and enhanced intensity compared to those of QDs before surface passivation. The optical properties changed are an indication of the aggregation of a core/shell QDs, which has been described in detail by several authors [24-26]. Although the experimental conditions have not been optimized, the PL peak intensity increased more than 3 times after surface passivation with a core-shell of QDs, compared to that of QDs before surface passivation, without significant modifications of the absorbance and PL features. In addition, the position of the PL peak remained essentially unchanged during surface passivation process, indicating that the QDs

remained substantially optical properties. We concluded that the spectral blue shifts of QDs were not due to the presence of QDs aggregates and gradual dissociation of aggregates [27]. This was confirmed from the PL and absorption spectra of QD solutions with surface passivation: the PL intensities of QD solutions decreased without surface passivation. From these observations, we concluded not only that the aggregation of the QDs was minimized under the experimental conditions but also that increase of the PL intensity due to the presence of aggregates or changes in the physical properties such as pitch and dipole moment of the medium. However, we observed considerable decreases of the PL emission intensity of the QDs when photoactivated without surface passivation; the PL efficiency decreased about 23.1 % in toluene and 61.5 % in CLCs polymer template (Figure 1(b)). One series of ODCLCs fluorescence spectra after passivation of QDs doping and polymer template formed is shown in Figure 1(b), where text in each step (b1-b5) indicates the PL intensity of QDCLCs after each condition. It was mentioned that PL intensity of QDCLCs increase remarkable with the passivation of the QDs in the CLCs template. It could be affected by structural changes during the polymer template growth and/or doping by passivation of QDs. The result clearly demonstrates that passivation and template treated enhance the OD fluorescence, but they do not affect the wavelength of the OD fluorescence peak. Most importantly, such modification is notable; i.e., the PL intensity is totally recovered after polymer template is added to passivation of QDs in the CLC system, which could be luminescence centers as well. On the other hand, the lasing emission can only be obtained when the QDs were used by surface passivation in the template mode. The decreases of the PL or lasing efficiency in this case can be attributed to poor photoactivated surface passivation and rapid formation of aggregation in the template sample. Also, the increase of the optical efficiency is due to a decrease in the number of surface defects or an increase in the density of emissive surface states. It should be noted that, after the surface passivation process, the lasing intensity of the sample increased slightly with QDs emission matching the band edge of the CLCs polymer template. Therefore, the contribution of a template component to the interactions between CLCs and the surface of the QDs cannot be completely ruled out.



Figure 1(a) Measured fluorescence emission of the QD dissolved in toluene. The fluorescence emission spectrum of the QD is obtained by the excitation of a diode-pumped solid-state laser with wavelength of 532 nm and intensity of 100 μ W/cm². (b) Measured fluorescence emission of QDCLC cells with various conditions is obtained by the excitation of a diode-pumped solid-state laser with wavelength of 532 nm and intensity of 100 μ W/cm². (b1) with passivation of QDs in the CLCs template, (b2) without passivation of QDs in the toluene, (b5) without passivation of QDs in the toluene.

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Figure S3 indicates the reflection spectra of the sample measured at the before-curing, after-curing, after-washing, and after-refilling stages of fabrication of QDNLC (nematic liquid crystal)-refilling template sample, presented by black, green, red, and blue curves, respectively. Apparently, the spectral positions of the PBG of the sample and thus its central wavelength (λ_c) are distinct at different stages. The λ_c of the PBG of the sample shifts slightly from 640 nm to 590 nm after UV-curing. This result is attributable to the slight decrease of helical pitch of the chiral polymer following the volume shrinkage through the polymerization and crosslinking reactions [20]. The Figure S4 (a)-(d) shown the optical properties and morphologies in QDCLCs devices as observed using a polarizing optical microscope (POM). It is to be noted, however, that the almost appearance blue shift of the reflect band in the template mode is a sensitive response of the polymerization of the ODCLCs [28].

Figure 2 are optical and fluorescent microscope images of QDCLCs polymer template. Before growth of the QDCLCs polymer template, the ODCLCs device show PL emission but with quite low PL intensity due to more aggregation of QDs as shown in the Figure 2 (a). After QDCLCs polymer template growth (less aggregation of QDs, Figure 2 (b)), the lasing properties of the resulting polymer template became quite good. Without substantial optimization, the PL efficiency increased more than 1.3 times and the lasing emission appear. As mentioned above, we believe the occurrence of aggregation was originated from interparticle hydrogen bonding. The results also suggested the disaggregation was not a result from the repulsion of an electrostatic force between interparticles. The QDs of the CLC polymer template are well separated from each other, proving that QDs have been well protected by surface passivation and successfully embedded into the polymer template. From these observations, we concluded not only that the aggregation of the QDs was minimized under the polymer template conditions but also that the increase of the PL intensity were mainly due to the dissociation of aggregates or changes in the physical properties such as dipole moment of the medium. The decreases of the PL intensity in this case can be attributed to poor surface passivation and formation of the surface defects in these soft matters. In addition, the stop band of QDCLC polymer template formed with increasing static time. Over a broad range of static time (3hr), we observed complete stop band or no change in the characteristics of the band-edge for QDCLC polymer template (Figure S5). The above results indicate that static time is the crucial factor for this system.



Figure 2. Fluorescent microscopy of QDCLCs polymer template devices (a) without polymer template, (b) with polymer template.

Name

The Figure 3 (a) shows the Lasing intensity spectrum of QDCLCs polymer template structure and corresponding reflection spectra (dashed line). The raise on intensity due to selective reflection was observed around 608 nm, which corresponds to the longer wavelength edge of the CLCs. The decrease in transmittance at shorter wavelength is attributed to the absorption of the doped QDs. It should be noted that the lasing peak wavelength at 608 nm match with the band edge of the CLCs. From this result, it is found that this laser action is based on the contribution of the photonic band edge of the CLCs cell. The polarization analysis of QDCLCs polymer template lasing have been investigated for left- and right-circular polarizers (Figure S6). Figure 3 (b) shows the dependence of emission intensity on the pumping energy of Nd: YAG pulse laser and it is the spectral linewidth of the ODCLCs polymer template. This indicates the clear presence of pumping threshold approximately 11 µJ/pulse for laser action. The pumping energy dependence of emission intensity and the peak width (the full width at half maximum; FWHM) of the emission peak at 608 nm. At a low pumping energy, emission intensity increases in proportion to pumping energy. Above a threshold at the pumping energy of 11 µJ/pulse, emission intensity nonlinearly increases and a sharp emission peak appears. The FWHM of the peak is less than 2.0 nm at above the threshold excitation energy. This behavior is similar to the line narrowing expected for amplified spontaneous emission (ASE). From the results, it was confirmed that the band edge effect of QDCLC polymer template was effective to the laser action in the semiconductor compound structure with surface passivation. Figure S7 indicates that the high-energy pulses that induced the thermal effect for dyes must be very strong to instantaneously evaporate the CLC material. The DDCLC laser on the pumped region therefore becomes totally destructive at the damage threshold of 55 μ J/pulse. By contrast, the QDCLC laser was not destroyed and could be repetitively used in the pumped region even after the excitation of the high-energy pulses on the cell at $E = 100 \mu$ J/pulse. ASE emerged at the reflect band edge of photonic due to the enhanced density of the optical state near the stop band edges [29]. Experimental steady-state diffraction pattern of QDCLCs polymer template is presented in Figure S8, and the respective lasing pattern of the sample are shown in the insets. As a result, the optical properties of the ODs as a gain medium, effectively dissipates within polymer template matrix without surface passivation. In fact, adding QDs with surface passivation into the polymer template was demonstrated as a more effective way to reduce photoemission quenching than organic dye [30]. Other study reported that the PL intensity of QDs in colloidal solution can be significantly enhanced by adding proper polymers [31].



Figure 3. (a) The fluorescence emission spectrum of the QDCLC polymer template with pumped energy, and corresponding reflection spectra (dashed line) (b) Peak intensities of fluorescence emission and corresponding fullwidths at half-maximum (FWHM) of the QDCLC polymer template as functions of pumped energy.

Conclusions

We have successfully demonstrated the lasing emission in the QDCLCs polymer template. The device was fabricated by surface passivation QDs doped CLC with a polymer template. Lasing emission emerged at the reflect band edges of the QDCLCs polymer template, and the threshold was much lower with template assisted than that at the without template, which was attributed to the increase of the local stability and density of state at the reflect band edges. The lasing threshold was 11 μ J/pulse, and the fwhm was 2 nm. Our experimental results would open up a broad search for developing high efficiency QDCLCs lasers.

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Notes and references

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- L. J. Chen, J. D. Lin, S. Y. Huang, T. S. Mo and C. R. Lee, *Adv. Optical Mater.* 2013, *1*, 637–643.
- 2 H. Finkelmann, S. T. Kim and A. Munoz, *Adv. Mater.* 2001, **13**, 1069.
- 3 P. V. Shibaev, J. Madsen, and A. Z. Genack, *Chem. Mater.* 2004, *16*, 1397-1399.
- 4 P. V. Shibaev, V. I. Kopp, and A. Z. Genack, J. Phys. Chem. B 2003, 107, 6961-6964.
- 5 Y. Wang , T. Manabe , Y. Takanishi , K. Ishikawa , G. Shao , A. Orita , J. Otera , H. Takezoe, *Optics Communications* **2007**, *280*, 408–411.
- 6 L. J. Chen, J. D. Lin and C. R. Lee, J. Mater. Chem. C, 2014, 2, 4388–4394.
- 7 V. I. Klimov, S. A. Ivanov, J. Nanda, M. Achermann, I. Bezel, J. A. McGuire and A. Piryatinski, *Nature* 2007, 447 (7143), 441–446.
- 8 M. C. Schlamp, X. G. Peng and A. P. Alivisatos, J. Appl. Phys. 1997, 82 (11): 5837–5842.
- 9 S. Reitzenstein, T. Heindel, C. Kistner, A. Rahimi-Iman, C. Schneider, S. H⁻ ofling and A. Forchel, *Appl. Phys. Lett.* 2008, 93, 061104.
- 10 A. L. Rodarte, C. Gray, L. S. Hirst and S. Ghosh, *Phys. Rev. B* 2012, 85, 035430.
- 11 S. G. Lukishova, L. J. Bissell, J. Winkler and C. R. Stroud Jr, Opt. Lett., 2012, 37, 1259.
- 12 A. Bobrovsky, K. Mochalov, V. Oleinikov, A. Sukhanova, A. Prudnikau, M. Artemyev, V. Shibaev and I. Nabiev, Adv. Mater., 2012, 24, 6216.
- 13 I. L. Medintz, H. T. Uyeda, E. R. Goldman and H. Mattoussi, *Nat. Mater.* 2005, 4 (6), 435–446.

- 14 A. M. Smith, H. W. Duan, M. N. Rhyner, G. Ruan and S. M. Nie, *Phys. Chem. Chem. Phys.* 2006, *8*, 3895.
- 15 A. K. Patri, J. F. Kukowska-Latallo and J. R. Baker Jr, Adv. Drug Deliv. Rev. 2005, 57, 2203.
- 16 D. V. Talapin, A. L. Rogach, A. Kornowski, M. Haase and H. Weller, Nano Lett. 2001, 1, 207.
- 17 X. Peng, M. C. Schlamp, A. V. Kadavanich and A. P. Alivisatos, J. Am. Chem. Soc. 1997, 119, 7019.
- 18 M. A. Hines and P. Guyot-Sionnest, J. Phys. Chem. 1996, 100, 468.
- 19 M. Danek, K. F. Jensen, C. B. Murray and M. G. Bawendi, *Chem.Mater.* **1996**, *8*, 173.
- 20 M. Mitov and N. Dessaud, Nat. Mater. 2006, 5, 361-364.
- 21 I. Dierking, L. L. Kosbar, A. Afzali-Ardakani, A. C. Lowe and G. A. Held, *Appl. Phys. Lett.* **1997**, *71*, 2454–2457.
- 22 S. S. Choi, S. M. Morris, W. T. S. Huck and H. J. Coles, *Adv. Mater.* 2010, *22*, 53–56.
- 23 J. Guo, H. Cao, J. Wei, D. Zhang, F. Liu, G. Pan, D. Zhao, W. He and H. Yang, *Appl. Phys. Lett.* **2008**, *93*, 201901.
- 24 V. Biju, R. Kanemoto, Y. Matsumoto, S. Ishii, S. Nakanishi, T. Itoh, Y. Baba and M. Ishikawa, J. Phys. Chem. C 2007, 111, 7924-7932.
- 25 S. Kim and M. G. Bawendi, J. Am. Chem. Soc. 2003, 125, 14652-14653.
- 26 Z. A. Peng and X. G. Peng, J. Am. Chem. Soc. 2001, 123, 183-184.
- 27 A. M. Derfus, W. C. W. Chan and S.N. Bhatia, *Nano Lett.* **2004**, *4*(1), 11-18.
- 28 M. Mitov and N. Dessaud, Liq. Cryst. 2007, 34, 183.
- 29 R. V. Nair, A. K. Tiwari, S. Mujumdar and B. N. Jagatap, *Phys. Rev.* A 2012, 85, 023844 1–7.
- 30 S. Singamaneni, C. Jiang, E. Merrick, D. Kommireddy and V. V. Tsukruk, J. Macromol. Sci., Part B: Phys. 2007, 46, 7.
- 31 A. A. Bol and A. Meijerink, J. Phys. Chem. B 2001, 105, 10203.