

Chemical Science

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this *Accepted Manuscript* with the edited and formatted *Advance Article* as soon as it is available.

You can find more information about *Accepted Manuscripts* in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.



www.rsc.org/chemicalscience

Restructuring and Hydrogen Evolution on Pt Nanoparticle

Cite this: DOI: 10.1039/x0xx00000x

Guang-Feng Wei and Zhi-Pan Liu*

Received 00th January 2012,
Accepted 00th January 2012

DOI: 10.1039/x0xx00000x

www.rsc.org/

The restructuring of nanoparticles at the *in situ* condition is a common but complex phenomenon in nanoscience. Here, we present a first systematic survey on the structure dynamics and its catalytic consequence for hydrogen evolution reaction (HER) on Pt nanoparticle, as represented by a magic number Pt₄₄ octahedron (~1 nm size). Using first principles calculation based global structure search method, we follow stepwisely the significant nanoparticle restructuring under HER condition as driven by thermodynamics to expose {100} facets, and reveal the consequent large activity enhancement due to the marked increase of the concentration of the active site, being identified to be apex atoms. The enhanced kinetics is thus a “byproduct” of the thermodynamical restructuring. Based on the results, the best Pt catalyst for HER is predicted to be ultrasmall Pt particles without core atoms, a size below ~20 atoms.

Introduction

Nanoparticles are common forms or carriers of heterogeneous catalysts¹⁻³ and also of wide application in many other fields, e.g. as biomedical drug delivery agent⁴ and for energy conversion and storage devices⁵. Special physiochemical properties emerged at the nanoscale add new complexity in understanding and optimizing reactions on nanoparticle. Compared to chunky crystalline materials, nanoparticles are more flexible in morphology and under reaction conditions, the reshaped nanocatalyst may exhibit completely different activity, either poisoned or promoted after a so-called induction period.⁶⁻⁹ Despite the vast amount of research on nanoparticle synthesis and morphology control, major gaps in our knowledge still exist, especially with regard to our molecular level understanding on the *in situ* dynamic restructuring of nanoparticles: this is reflected in our inability to predict whether such restructuring is beneficial or detrimental to catalyst activity. Controlling the nanostructure dynamics for the desirable property, e.g. catalytic activity, is paramount for the rational catalyst design and is a general goal in nanomaterial applications.^{10, 11}

Pt is a unique metal with high catalytic performance for a wide range of reactions, and it is perhaps the most efficient HER catalyst in electro- and photocatalytic water splitting.^{12, 13} It has been constantly pursued in research to reduce the Pt usage by identifying the optimum particle size for activity. On model single crystalline surfaces, it was however found that HER is only weakly dependent on the crystal facet: the ridged Pt(110) is about two times more active than the (111) and (100) terraces¹⁴⁻¹⁷. On going to the nanoscale, there is no consensus on the particle size effect.¹⁸⁻²¹ The presence of the particle-support interaction further complexes the understanding on the particle size effect on activity.²² A very recent study by Schweinberger et al. using size-selected Pt nanoparticles supported on CdS nanorod shows that the particle of a critical particle size ~46 atoms (1 nm) can achieve the maximum H₂ production, whilst the mass activity is the highest when the particle size shrinks down to the subnanoscale with only 8 atoms (Pt₈).²³ To date, there are much uncertainties on the physical origin of the HER activity on small nanoparticles. The nature of the active site and

the dynamic structure evolution are two key issues that are required to resolve first.

Here we present the first quantum mechanics simulation on the structure dynamics of Pt nanoparticle during HER and quantify its catalytic consequence. The Pt nanoparticle considered in this work is represented by a Pt cluster of ~1 nm diameter, Pt₄₄, which is identified as a magic number size with O_h symmetry. Significant restructuring-induced promotion is revealed on the Pt₄₄ nanoparticle at the HER condition, and theory further predicts that such promotional effect due to restructuring is prominent only for nanoparticles below ~1.8 nm. In general, the restructuring as driven by the exothermicity of the adsorption of reaction intermediates may or may not increase the active site concentration that depends on the nature of the reaction and also the particle size.

As both nanoparticle restructuring and catalytic reactions are rare events with high barriers, it presents a challenge to computer simulation since the long simulation times of molecular dynamics, or even the use of enhanced sampling techniques, may not be able to capture the desired reaction patterns. For example, in HER on Pt(111), the barrier of H-H coupling to form H₂ can be as high as 0.92 eV at the working condition.²⁴ The approach we adopt here is to use the first principles density functional theory (DFT) based stochastic surface walking (SSW) global optimization method²⁵⁻²⁷, SSW-DFT, to explore the Pt nanoparticle morphology at the HER condition. The recently-developed SSW method is able to visit the minima on PES by following likely pathways, and therefore is a powerful tool for both structure prediction and pathway search²⁵. Using the new technique, we are able to follow stepwisely the particle restructuring in H₂ atmosphere and determine the HER activity.

Calculation methods

DFT Calculation. All SSW calculations and the reaction modelling were performed in combination with the DFT calculations as implemented in SIESTA package^{28,29} with Troullier-Martins norm conserving pseudopotentials³⁰. The exchange-correlation functional

utilized was at the generalized gradient approximation level, known as GGA-PBE³¹. The optimized double- ζ plus polarization (DZP) basis set with extra diffuse function was employed for metal. The orbital-confining cutoff was determined from an energy shift of 0.010 eV. The energy cutoff for the real space grid used to represent the density was set as 150 Ry. The Quasi-Newton I-BFGS method is used for geometry relaxation until the maximal force on each degree of freedom is less than 0.01 eV/Å. To correct the zero-point energy for reaction barrier, the vibrational frequency calculations were performed via the finite-difference approach. Transition states (TSs) of the catalytic reaction were searched using the Constrained-Broyden-based TS-searching methods^{32, 33}.

For all the Pt clusters from Pt₁₂ to Pt₄₆ (see Fig. S1), at least four lowest-lying isomers obtained from SSW-DFT/SIESTA search are further checked using the spin-polarized plane wave calculations with ultrasoft pseudo-potentials³⁴ or projected augmented wave^{35, 36} pseudo-potentials, as implemented in VASP³⁷. The plane-wave kinetic energy cutoff of 400 eV was used and the exchange-correlation functional utilized was at the generalized gradient approximation level, GGA-PW91³⁸ and GGA-PBE³¹. For large Pt clusters above 39 atoms, the total spin magnetic moment is zero while small Pt clusters are generally spin polarized.

SSW Calculation. In all SSW simulation, the key parameters utilized are the same with those utilized previously for exploring the PES of Carbon and Boron clusters²⁵⁻²⁷, i.e. the Gaussian width being 0.6 Å, the number of Gaussian potential being 10.

To identify the global minimum structure of Pt_x (x=12 to 46) (see Fig. S1), we set the temperature utilized in Metropolis Monte Carlo being 3000~5000 K. The higher temperature is used to verify the obtained global minimum structure. In the SSW search, we performed four to ten parallel runs and up to 300 minima are collected at the first stage, from which the most stable configuration is obtained. Next, we verified the result from the most stable configuration of the first stage and collect another 300 minima. This process is repeated until no further more stable configuration is identified at the stage of verification.

For the grand canonical Monte Carlo (GCMC) simulation for Pt_{44}H_x system, the basic procedure of the SSW simulation at each fixed H concentration is the same as that described above for pure Pt clusters. In the GCMC simulation, the major difference is that every 300 SSW steps, we will evaluate the chemical potential of adsorbed H atom with respect to that of H in the gas phase ΔG_{H} (see below in Eq. 2) based on the current most stable configuration. According to the value $\Delta G_{\text{H}} < \text{or} > 0$, we will decide to accept or refuse the current most stable configuration. To speed up the structure search for reaching the $\Delta G_{\text{H}}=0$ equilibrium, the newly-arrived H atoms will be always added to the vacant surface sites, e.g. vacant bridge site; the removal of H atoms will always chose the atop H atoms or subsurface H atoms, if present, which are calculated to have the poorest adsorption energy.

The exothermicity of restructuring under HER condition. In HER, the nanoparticle is in the H_2 atmosphere and an equilibrium of H chemical potential needs to be achieved at the steady state. The exothermicity of the H adsorption on the bare nanoparticle provides the driving force of the restructuring. This is measured by ΔG per Pt atom with reference to Pt_{44} octahedron and H_2 gas (the standard condition is utilized here), as shown in Eq. 1.

$$\Delta G = [G(\text{Pt}_{44}\text{H}_x) - G(\text{Pt}_{44}) - x/2G(\text{H}_2)]/44 \quad (1)$$

Here $G(\text{Pt}_{44}\text{H}_x)$ and $G(\text{Pt}_{44})$ can be computed from DFT directly by including the zero point energy (ZPE) correction, and $G(\text{H}_2)$ is the free energy of the gas phase H_2 that can be obtained from standard thermodynamics data³⁹.

Results and discussion

Pt_{44} octahedron. In this work, we utilize Pt_{44} as the model catalyst for investigating the HER on ~ 1 nm Pt nanoparticles. Pt_{44} is predicted to a magic number size based on the unbiased SSW-DFT global structure search (see supplementary information, SI discussion and Fig. S1), which was also suggested previously⁴⁰ by comparing with other putative high symmetry structures of Pt_{44} . From the SSW trajectories 4788 minima of Pt_{44} were collected and the GM of Pt_{44} is found to be a O_h symmetry octahedron with bulk-like face-centered cubic (fcc)

packing, exposing only $\{111\}$ facets and containing six core atoms and 38 shell atoms, see Fig. 1. It should be mentioned that Pt_{44} is the smallest octahedron of Pt nanoparticle with high stability, which exhibits a remarkable structure similarity to the bulk Pt crystal: the average Pt-Pt distances of Pt_{44} is 2.74 Å, being only 2.7% shorter than that in Pt bulk. This affords Pt_{44} to be a good model for understanding nanoparticle behaviour under reaction conditions.

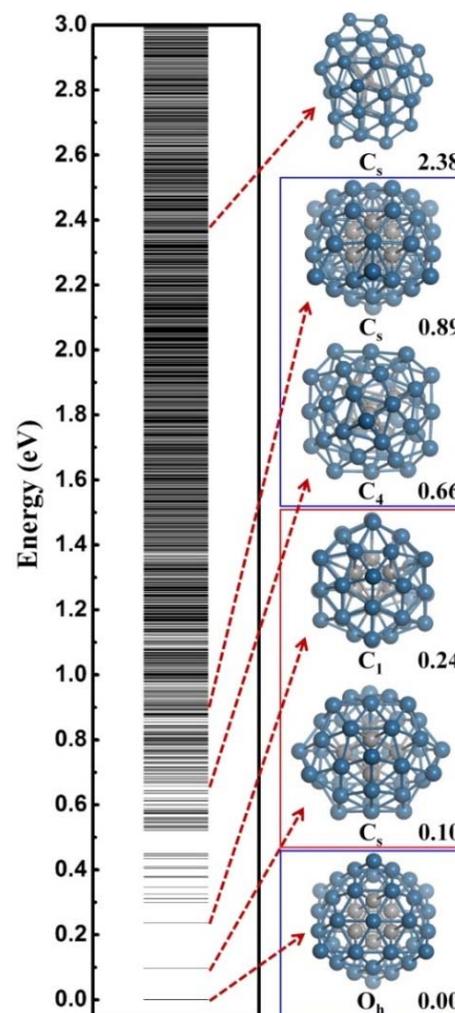


Fig. 1 The configurational spectrum of Pt_{44} obtained using SSW-DFT global structure search. Deep blue: Shell Pt atoms Grey: core Pt atoms. The PES of Pt_{44} exhibits two distinct energy funnels, a larger funnel represented by the second lowest minimum (SLM), featuring with pentastar structure units on the surface, and a smaller funnel represented by the octahedron GM with a characteristic C_4 axis. All the stable structures within 0.66 eV above GM are geometrically similar to the SLM. For the GM funnel, the next stable structure appears at 0.66 eV, which can be generated from the GM by rotating half of the shell Pt atoms around the C_4 axis.

From all the minima of Pt_{44} collected from SSW trajectories, we have constructed the configurational spectrum of Pt_{44} in Fig. 1 to

provide insights into the PES of Pt nanoparticle in general. The structures of typical less stable minima are also shown. In general, the conformation of the Pt nanoparticle is discrete at the energy window 0–0.5 eV above GM and become continuous-like above 0.5 eV. Most of the low lying structures have a common core-shell feature as the GM with 6 core atoms and 38 shell atoms. The continuous energy spectrum appears just 0.5 eV above GM, indicating that the Pt nanoparticle is highly mobile and the reconstruction of the shell is kinetically allowed even at the ambient conditions.

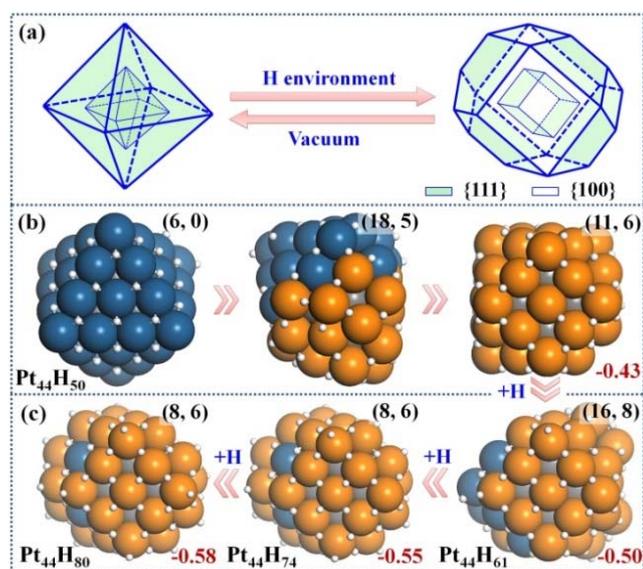


Fig. 2 Structure Evolution of Pt₄₄ under HER condition. (a) A polyhedron representation showing the nanoparticle structure evolution from octahedron to tetracapped octahedron under the HER condition. (b) The representative structures at Pt₄₄H₅₀ highlighting the initial stages of {111} to {100} reconstruction; (c) the GM structures at Pt₄₄H₆₁, Pt₄₄H₇₄ and Pt₄₄H₈₀. In the structures, the exothermicity per Pt atom due to the adsorption of H atoms (Eq. 1) is indicated at the right-bottom corner and the numbers of {111} and {100} microfacets are indicated inside the parenthesis of the right-top corner. All the apex Pt in Pt₄₄H₈₀ are coordinated with five H atoms with one atop H (the H atoms prefer the fcc hollow sites on {111} and the bridge sites on {100} terraces). Deep blue: Shell Pt atoms associated with {111} facets only; Orange: Shell Pt atoms associated with at least one {100} facet; Grey: core Pt atoms and white: H atoms.

Structure Evolution under HER condition. To simulate the structure evolution dynamics at the HER conditions, we consider the Pt₄₄/H₂ system as a grand canonical ensemble where the chemical potential of the adsorbed H atoms will eventually reach equilibrium with that in the gas phase, i.e. $\Delta G_{\text{H}} \rightarrow 0$. ΔG_{H} can be calculated as follows,

$$\Delta G_{\text{H}} = \Delta E_{\text{DFT}} + \text{ZPE} - 1/2G(\text{H}_2) \quad (2)$$

where ΔE_{DFT} and ZPE are the differential adsorption energy and the zero point energy of the newly-arrived H atom on particle; and $G(\text{H}_2)$ is the free energy of H₂ in the gas phase at the standard state. The grand canonical Monte Carlo (GCMC) simulations based on SSW/DFT global structure search (GCMC/SSW-DFT) are thus performed to investigate the structure evolution. This is done by adding/subtracting H atoms into/from the system, Pt₄₄H_x, every a few hundred (> 300) SSW steps of structure search at a fixed H concentration. In fact, at the initial stage of simulation, the addition of H atoms is always energetically preferable as new vacant sites emerge continuously due to the surface reconstruction where newly-arrived H atoms can adsorb. The simulation reaches to equilibrium at the stage of Pt₄₄H₈₀, when ΔG_{H} of newly-arrived H atoms becomes positive. From the trajectories of GCMC simulation, we selected representative structures at several Pt₄₄H_x stages, as shown in Fig. 2, and they are described as follows.

Pt₄₄H₅₀. Our GCMC/SSW-DFT simulation starts from a Pt₄₄H₅₀ octahedron, when {111} facets are fully occupied by H atoms (>1 monolayer, ML) and ΔG_{H} starts to exhibit the appreciable decrease due to the switch of the adsorption site for H atom. The simulation shows that the particle leaves the octahedron shape immediately after only 3 SSW steps. The initial reconstruction starts by the collective migration of the vertex Pt atoms of the octahedron, thereby exposing {100} facets (see SI animation-1). This is simply because {100} facet is able to adsorb more H atoms than {111} and thus thermodynamically preferred. In the restructuring, the surface atoms including H diffuse around to find the energetically favorable position and as a result, {100} facets emerge by breaking large {111} facet into small {111} microfacets (Fig. 2b). The surface becomes rough. For Pt₄₄H₅₀, additional SSW runs were performed to understand the restructuring dynamics (see SI animation-2). After a long simulation of 2021 SSW steps, Pt₄₄H₅₀ becomes significantly different from the initial octahedron, possessing 11 {111} and 6 {100} facets (Fig. 2b), although the core-shell structure still remains: Pt₄₄H₅₀ has 7 core atoms and 37 shell atoms.

Pt₄₄H₆₁₋₈₀. With the increase of H coverage, the small {100} microfacets start to merge with each other to yield large {100} facets, and simultaneously the area of {111} facets shrinks. The polyhedron shapes start to reappear as the stable forms and ΔG_H gradually approaches to zero. The identified GM of Pt₄₄H₈₀ is found to be a C_{2h} tetradecahedron of fcc packing, with 8 {111} facets, 6 {100} facets and 18 apex Pt atoms (five or six coordinated Pt atoms shared by at least three facets). ΔG_H at Pt₄₄H₈₀ is -0.07 eV.

The GCMC/SSW-DFT simulation conveys two important messages for the HER-driven nanoparticle restructuring: (i) Adsorbed H atoms are always more stable on the surface even when the equilibrium coverage is above 2 ML. Importantly, the subsurface H atoms inside Pt nanoparticle is found to be unstable (see SI Fig. S2) since the stronger Pt-Pt bond is energetically preferred compared to the Pt-H bond in forming the particle core. This implies a high stability of Pt nanoparticle under HER condition. (ii) The restructuring is driven to maximally expose {100} while the core-shell structure of Pt nanoparticles is always kept to minimize the total energy. Only {100} and {111} facets are present at the GM of Pt₄₄H₈₀. Overall, the core Pt atoms increase from 6 to 8 and the shell Pt atoms decreases from 38 to 36 (the surface density drops) after the restructuring, which is consistent with the typical surface reconstruction observed in surface science studies⁴¹.

HER activity. We are now at the position to investigate the HER on the Pt₄₄H₈₀ polyhedron. We have considered all the likely reaction patterns for the hydrogen evolution via the coupling of two adsorbed H atoms, H+H→H₂, the so-called Tafel mechanism in electrocatalytic HER that is preferable at high H coverage conditions (see our recent work on HER kinetics on surfaces²⁴ where the electrochemical potential and solvation effect have been considered; here we follow the main conclusion obtained there).

Fig. 3a shows that the calculated free energy barriers (ΔG_a) of H-H coupling on Pt₄₄H₈₀ span from 0.47 to 1.07 eV depending on the local sites. Unexpectedly, ΔG_a at the apex sites are 0.47~0.71

eV, which is much lower than that on the {111}, edge and {100} sites, 0.88, 0.88 and 1.07 eV, respectively. We also noticed that the calculated ΔG_a of the H-H coupling on the {111} facets of Pt₄₄H₈₀ (F1) is in fact similar to that on the extended Pt(111) (~0.9 eV). The optimized structures of the transition state (TS) are also similar in two cases (see SI Fig. S3). Similarly, ΔG_a at the apex sites of the unreconstructed Pt₄₄H₄₈ octahedron (AO) is also in the ΔG_a range of the apex sites on Pt₄₄H₈₀. These results indicate that the HER activity can be assessed largely by the local geometry of Pt site.

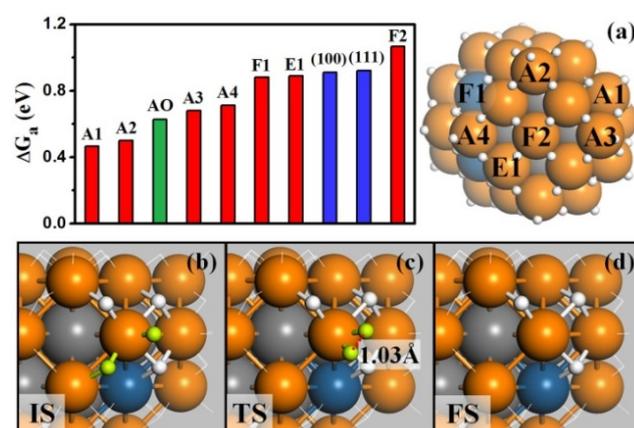


Fig. 3 (a) The free energy barriers (ΔG_a) of H-H coupling on Pt₄₄H₈₀. The sites are as indicated on the right-hand particle. Among 18 apex sites, there are ten A1, four A2, two A3 and two A4 sites. Also shown are ΔG_a on the extended (111), (100) surfaces and the apex sites of the unreconstructed Pt₄₄H₄₈ (AO); (b-d) the reaction snapshots for the lowest barrier reaction channel at the apex A1 site. The reaction features with the atop H reacting with a neighboring bridging H (both H highlighted by yellow color), where the apex Pt atom is coordinated with five H atoms. The color scheme is as that in Figure 2.

By identifying the critical role of apex sites and the local reactivity in HER, we may discuss their implication to HER catalysis. In Fig. 4, we first estimated the HER activity of Pt particles on differently sized as-synthesized nanoparticles at the equilibrium shape in solution, i.e. no restructuring due to H₂ (data taken from experiment and our recent study)^{42, 43}. We then count the apex, {111} and {100} and edge sites of the particles and sum the overall HER rate (specific activity of HER, SA, mol·cm⁻²·s⁻¹) on all the sites based on microkinetics, as shown in Eq. 3.

$$SA = 0.5 \times S^{-1} N_A^{-1} \times \sum_{site} (k_B T / h) \exp(-\Delta G_a / RT) \theta_{site} \quad (3)$$

where k_B is Boltzmann's constant and h is Planck's constant ($k_B T/h$ is 6.25×10^{12} at 298.15 K from classic TS theory); ΔG_a is the estimated free energy barrier statistically averaged according to the data in Fig. 3a, i.e. 0.48 eV for apex sites, 0.88 eV for terrace sites and 0.89 eV for edge sites; θ_{site} is the number of the active sites on the nanoparticles; S is the surface area. Here we assume the same HER activity at the same type of site based on the fact of local reactivity of HER identified above. The calculated HER rate on the nanoparticle is thus plotted in Fig. 4 black curve, showing the maximum activity around ~ 1.8 nm when the highest concentration of apex site (per unit surface area) is reached. It should be mentioned that the HER activity estimated in the black curve agrees reasonably with the kinetic data reported in experiment, validating largely the local reactivity assumption. For example, Hoshi et al. reported 1.21 mA/cm² for electrocatalytic HER on ~ 3 nm Pt nanoparticles that corresponds to $\sim 6.3 \times 10^{-9}$ mol \cdot cm⁻² \cdot s⁻¹ with the apparent ΔG_a of 0.54 eV (Eq. 3), while the estimated ΔG_a utilized in Fig. 4 is 0.48 eV and the difference in the rate is no more than 10 times.

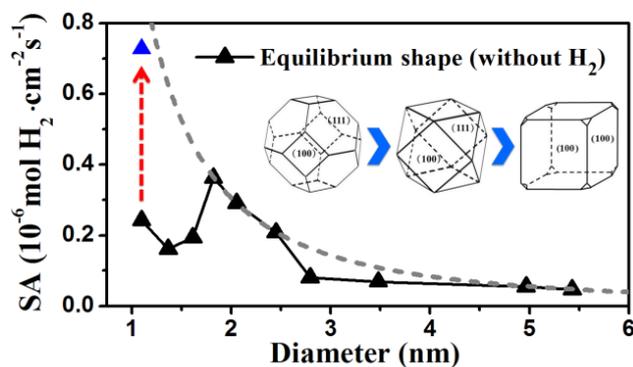


Fig. 4 Estimated HER activity (using Eq. 3) of Pt nanoparticles with the equilibrium shape in solution, i.e. no restructuring due to H₂, (black curve) and with only the {100}-dominated truncated cubic shape (grey dash curve). The nanoparticle equilibrium shape in solution is taken from experiment⁴² and our recent work⁴³, which exhibits a gradual transition from octahedron at very small particle to truncated octahedron, to cuboctahedron and to the {100}-dominated truncated cube at large particles. The grey dash curve represents the activity limit after the restructuring.

Furthermore, we may also consider the situation after the nanoparticle restructuring at the HER condition. Although we do not know the exact atomic structure, this work does show that {100} is the direction of restructuring and thus {100}-dominated truncated cube would be the preferable shape starting from small

particles, where the apex sites can reach to 24 per particle. We therefore estimate the HER rate as a possible maximum limit due to restructuring using the same approach above, as plotted in the grey dash curve in Fig. 4. Indeed, the trend for the large increase of HER activity of Pt₄₄ after restructuring is correctly reflected in the Figure (the red arrow). Interestingly, Fig. 4 predicts that for large nanoparticles above 1.8 nm, the restructuring of nanoparticle, although should occur as well, does not enhance the HER activity appreciably. The activity decays very slowly above 3 nm, when the activity can be regarded as insensitive to the particle size.

By contrast, for very small nanoparticles (e.g. 1 nm), the activity can be dramatically higher, which is caused by the dynamic restructuring at the HER condition that creates a high concentration of five or six coordinated apex sites per surface area. Along this line, we expect that ultrasmall Pt clusters without core atoms have the highest HER activity because all Pt atoms are on the surface with low coordination, where the concentration of apex sites can be maximized at the HER condition. This corresponds to a size of less than ~ 20 atoms (see ref.⁴⁰ and also SI Fig. S1), which may rationalize the highest photocatalytic HER activity of Pt₈ observed recently.²³

Conclusions

New DFT-based global optimization theoretical methods allows the observation of the dynamic catalyst structure evolution and the quantification the activity change of Pt nanoparticle for HER. Unexpectedly, we found that HER occurs preferentially on Pt low-coordinated apex Pt sites, which totally dominates the activity for Pt nanoparticles. The restructuring of nanoparticle can promote HER, but appreciably only below certain size threshold, ~ 2 nm, where the apex sites dynamically created can reach to the maximum concentration. The subnano Pt clusters without core atoms is predicted to have the highest HER activity.

Acknowledgements

We acknowledge National Science foundation of China (21173051, 21361130019), 973 program (2011CB808500, 2013CB834603), Science and Technology Commission of Shanghai Municipality (08DZ2270500), China Postdoctoral Science Foundation (2013M531112) and Shanghai Postdoctoral Scientific Program (13R21410300) for financial support.

Notes and references

Shanghai Key Laboratory of Molecular Catalysis and Innovative Materials, Department of Chemistry, Key Laboratory of Computational Physical Science (Ministry of Education), Fudan University, Shanghai 200433, China. E-mail: zpliu@fudan.edu.cn

† Electronic Supplementary Information (ESI) available: discussions on the structures of Pt clusters and the stability of the subsurface H atoms in Pt cluster, TS structure of H-H coupling on {111} facets of Pt₄₄H₈₀, XYZ coordinate of Pt₄₄ and Pt₄₄H₈₀, Movie of structure evolution at Pt₄₄H₅₀. See DOI: 10.1039/b000000x/

- H. M. T. Galvis, J. H. Bitter, C. B. Khare, M. Ruitenbeek, A. I. Dugulan and K. P. de Jong, *Science*, 2012, **335**, 835-838.
- Y. Roman-Leshkov, C. J. Barrett, Z. Y. Liu and J. A. Dumesic, *Nature*, 2007, **447**, 982-U985.
- Y. F. Li and Z. P. Liu, *J. Am. Chem. Soc.*, 2011, **133**, 15743-15752.
- R. R. Arviso, S. Bhattacharyya, R. A. Kudgus, K. Giri, R. Bhattacharya and P. Mukherjee, *Chem. Soc. Rev.*, 2012, **41**, 2943-2970.
- H. A. Atwater and A. Polman, *Nat. Mater.*, 2010, **9**, 205-213.
- F. Tao and M. Salmeron, *Science*, 2011, **331**, 171-174.
- S. Helveg, C. Lopez-Cartes, J. Sehested, P. L. Hansen, B. S. Clausen, J. R. Rostrup-Nielsen, F. Abild-Pedersen and J. K. Norskov, *Nature*, 2004, **427**, 426-429.
- P. L. Hansen, J. B. Wagner, S. Helveg, J. R. Rostrup-Nielsen, B. S. Clausen and H. Topsoe, *Science*, 2002, **295**, 2053-2055.
- C. Mager-Maury, G. Bonnard, C. Chizallet, P. Sautet and P. Raybaud, *Chemcatchem*, 2011, **3**, 200-207.
- A. F. Lee, C. V. Ellis, J. N. Naughton, M. A. Newton, C. M. A. Parlett and K. Wilson, *J. Am. Chem. Soc.*, 2011, **133**, 5724-5727.
- F. Tao, M. E. Grass, Y. W. Zhang, D. R. Butcher, J. R. Renzas, Z. Liu, J. Y. Chung, B. S. Mun, M. Salmeron and G. A. Somorjai, *Science*, 2008, **322**, 932-934.
- R. Borup, J. Meyers, B. Pivovar, Y. S. Kim, R. Mukundan, N. Garland, D. Myers, M. Wilson, F. Garzon, D. Wood, P. Zelenay, K. More, K. Stroh, T. Zawodzinski, J. Boncella, J. E. McGrath, M. Inaba, K. Miyatake, M. Hori, K. Ota, Z. Ogumi, S. Miyata, A. Nishikata, Z. Siroma, Y. Uchimoto, K. Yasuda, K. I. Kimijima and N. Iwashita, *Chem. Rev.*, 2007, **107**, 3904-3951.
- J. H. Yang, D. G. Wang, H. X. Han and C. Li, *Accounts Chem. Res.*, 2013, **46**, 1900-1909.
- N. M. Markovic, B. N. Grgur and P. N. Ross, *J. Phys. Chem. B*, 1997, **101**, 5405-5413.
- B. E. Conway and B. V. Tilak, *Electrochim. Acta*, 2002, **47**, 3571-3594.
- B. E. Conway, J. Barber and S. Morin, *Electrochim. Acta*, 1998, **44**, 1109-1125.
- J. H. Barber and B. E. Conway, *J. Electroanal. Chem.*, 1999, **461**, 80-89.
- H. Kotani, R. Hanazaki, K. Ohkubo, Y. Yamada and S. Fukuzumi, *Chem.-Eur. J.*, 2011, **17**, 2777-2785.
- N. Hoshi, Y. Asaumi, M. Nakamura, K. Mikita and R. Kajiwara, *J. Phys. Chem. C*, 2009, **113**, 16843-16846.
- M. J. Piotrowski, P. Piquini and J. L. F. Da Silva, *Phys. Rev. B*, 2010, **81**, 155446.
- B. M. Babic, L. M. Vracar, V. Radmilovic and N. V. Krstajic, *Electrochim. Acta*, 2006, **51**, 3820-3826.
- F. Beharfarid, L. K. Ono, S. Mostafa, J. R. Croy, G. Shafai, S. Hong, T. S. Rahman, S. R. Bare and B. R. Cuenya, *Phys. Chem. Chem. Phys.*, 2012, **14**, 11766-11779.
- F. F. Schweinberger, M. J. Berr, M. Doblinger, C. Wolff, K. E. Sanwald, A. S. Crampton, C. J. Ridge, F. Jackel, J. Feldmann, M. Tschurl and U. Heiz, *J. Am. Chem. Soc.*, 2013, **135**, 13262-13265.
- Y. H. Fang, G. F. Wei and Z. P. Liu, *J. Phys. Chem. C*, 2013, **117**, 7669-7680.
- C. Shang and Z. P. Liu, *J. Chem. Theory Comput.*, 2013, **9**, 1838-1845.
- X. J. Zhang, C. Shang and Z. P. Liu, *J. Chem. Theory Comput.*, 2013, **9**, 3252-3260.
- H. J. Zhai, Y. F. Zhao, W. L. Li, Q. Chen, H. Bai, H. S. Hu, Z. A. Piazza, W. J. Tian, H. G. Lu, Y. B. Wu, Y. W. Mu, G. F. Wei, Z. P. Liu, J. Li, S. D. Li and L. S. Wang, *Nat. Chem.*, 2014, **6**, 727-731.
- J. M. Soler, E. Artacho, J. D. Gale, A. Garcia, J. Junquera, P. Ordejon and D. Sanchez-Portal, *J. Phys.-Condes. Matter*, 2002, **14**, 2745-2779.
- J. Junquera, O. Paz, D. Sanchez-Portal and E. Artacho, *Physical Review B*, 2001, **64**, 235111.
- N. Troullier and J. L. Martins, *Phys. Rev. B*, 1991, **43**, 1993-2006.
- J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865-3868.
- H. F. Wang and Z. P. Liu, *J. Am. Chem. Soc.*, 2008, **130**, 10996-11004.
- C. Shang and Z. P. Liu, *J. Chem. Theory Comput.*, 2010, **6**, 1136-1144.
- G. Kresse and J. Hafner, *J Phys-condens Mater*, 1994, **6**, 8245-8257.
- P. E. Blochl, *Phys. Rev. B*, 1994, **50**, 17953-17979.
- G. Kresse and D. Joubert, *Phys. Rev. B*, 1999, **59**, 1758-1775.
- G. Kresse and J. Furthmuller, *Comput. Mater. Sci.*, 1996, **6**, 15-50.
- J. P. Perdew and Y. Wang, *Phys Rev B*, 1992, **45**, 13244-13249.
- D. R. LIDE, ed., *CRC Handbook of Chemistry and Physics*; CRC press, 2003-2004.
- V. Kumar and Y. Kawazoe, *Physical Review B*, 2008, **77**, 205418.
- K. Johnson, Q. Ge, S. Titmuss and D. A. King, *J Chem Phys*, 2000, **112**, 10460-10466.
- I. N. Leontyev, S. V. Beenov, V. E. Guterman, P. Haghi-Ashtiani, A. P. Shaganov and B. Dkhil, *J. Phys. Chem. C*, 2011, **115**, 5429-5434.
- G. F. Wei and Z. P. Liu, *Phys. Chem. Chem. Phys.*, 2013, **15**, 18555-18561.