

Journal of Materials Chemistry C

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this *Accepted Manuscript* with the edited and formatted *Advance Article* as soon as it is available.

You can find more information about *Accepted Manuscripts* in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.

Photochromic and photomechanical response in an amorphous diarylethene-based polymer. A Spectroscopic Ellipsometry investigation of ultrathin films[†]

Chiara Toccafondi,^{a‡} Luca Occhi,^b Ornella Cavalleri,^a Amanda Penco,^a Rossella Castagna,^c Andrea Bianco,^d Chiara Bertarelli,^c Davide Comoretto,^b Maurizio Canepa^{*a}

Received Xth XXXXXXXXXXXX 20XX, Accepted Xth XXXXXXXXXXXX 20XX

First published on the web Xth XXXXXXXXXXXX 200X

DOI: 10.1039/b000000x

This work deals with very thin (14–40 nm) films of a polyester containing diarylethene units in the main chain spin cast on a silicon wafer. By irradiation with UV light the colourless form turns blue due to the appearance of a strong absorption band centred at about 600 nm. The coloured state is thermally stable and the backward conversion can be triggered with visible light. Comparison of broadband (245–1700 nm) spectroscopic ellipsometry data with simulations based on a isotropic, Kramers-Kronig consistent, multiple-resonance model allowed to determine the complex index of refraction \tilde{n} of the film in its blue and colourless forms. The refractive index of the blue phase neatly exceeds the one of the transparent form for wavelengths in the NIR. In particular, out of resonance, at $\lambda \sim 1700$ nm, $\Delta\text{Re}(\tilde{n}) \sim 0.05$. Parallel to the $\Delta\text{Re}(\tilde{n})$ increase, the analysis of ellipsometry data suggests a decrease of the film thickness (about 1.5%) in the transition from the open (transparent) to the closed (coloured) form.

1 Introduction

Photonic polymers are gaining increasing interest for applications where standard inorganic materials are less competitive in terms of mechanical flexibility of devices, processing costs or properties such as color tunability and large active area.¹ In this field, photochromic materials, such as azobenzenes and diarylethenes, whose colour and other physico-chemical properties can be modulated by light stimulation, find growing applications in memories, switches, actuators, tunable filters, volume phase holographic gratings, optical nanopatterning and multiplexing/demultiplexing devices.^{2–11} Applications aiming to exploit very thin films require sensitive and non-destructive characterization techniques. Optical meth-

ods, traditionally used for the determination of optical properties of thin films^{12–14}, have been used to investigate crystallization or non-equilibrium morphologies thus flanking electron microscopy or X-ray diffraction studies.^{15,16} UV-VIS Spectroscopic Ellipsometry (SE) may couple thickness sensitivity, down to the monolayer limit^{17–20} to real-time information on broad-band spectroscopic dielectric properties, which can be possibly correlated with morphological issues²¹. Indeed, polymers (and polymerization) attracted SE studies since the early development of the technique^{22–24}; more recently, SE has found growing application to very thin films in the fields e.g. of organic electronics^{13,25}, smart surfaces^{21,26} and other stimuli responsive polymers^{20,27,28}. SE becomes particularly appealing in the investigation of ultra thin films whose optical and morphology properties may significantly differ from the bulk form ones. This is the case e.g. of conjugated polymers, where the combination of optical and morphological changes provides information on processes affecting the performances of optoelectronic devices such as the evolution of structure, relaxation of non-equilibrium morphology, crystallization, phase separation and solvent evaporation.^{28–33} Under this perspective, it is promising to exploit real-time SE in the investigation of the light-triggered conversion of films of photochromic polymers towards the ultra thin thickness limit.

Diarylethenes are well-known photochromic compounds endowed with thermal stability and high fatigue resistance^{3,5}.

[†] Electronic Supplementary Information (ESI) available: [Definition of the MSE function; UV-Vis absorption spectra of p-DTE in chloroform solution; AFM images of films; the extinction coefficient of the polymer film as function of energy; details about the comparison between experimental and simulated spectra for representative samples with different thickness;]. See DOI: 10.1039/b000000x/

^a Dipartimento di Fisica, Università di Genova and CNISM, via Dodecaneso 33, Genova, Italy.

^b Dipartimento di Chimica e Chimica Industriale, Università di Genova, via Dodecaneso 31, Genova, Italy.

^c Dipartimento di Chimica, Materiali ed Ingegneria Chimica "G. Natta", Politecnico di Milano, Piazza Leonardo da Vinci 32, 20133 Milano, Italy

^d INAF- Osservatorio Astronomico di Brera, via Bianchi 46, 23807 Merate, Italy

*E-mail: canepa@fisica.unige.it

[‡] Present address: Istituto Italiano di Tecnologia, Via Morego 30, Genova

Blended (few percent by weight) into transparent polymer matrices, their refractive index may appreciably change across the transition.^{2,5,34}

In order to overcome possible aggregation phenomena, which enhance light scattering and hamper the optical quality of molecular films, the insertion of the diarylethene moiety into a polymer chain can be of great advantage. Main-^{4,5,35,36} and side-chain^{37,38} diarylethene-based photochromic polymers combine a large content of active units in the bulk, which leads to a high optical density, and processing features which provide films of good optical quality. The material under investigation here is a polyester containing diarylethene units in the main chain (p-DTE, fig.1). The tetramethylene spacer imparts good flexibility to the molecular backbone, thus improving the solubility in common organic solvents and, in turn, material processing. By irradiation with UV light the transparent, *open* form turns *closed* and coloured as schematically indicated in fig.1. The backward conversion of the thermally stable, coloured state is obtained by exposure to red light.

The main goal of this work was the determination of the broadband (245-1700 nm) dielectric function of the photochromic polymer in each of the two forms. Significant transition-induced changes of the complex refractive index, strategic for photonic applications, were expected from results in literature^{2,4,8,34,39,40}, possibly strong for high density films⁴¹. A second intriguing objective was to exploit SE to get insight into changes of thickness of photochromic ultrathin films interesting for the development of photo-responsive actuators.^{42,43} We considered films with thickness of the order of a very few tens of nm, one order of magnitude thinner than systems hitherto investigated^{5,39,41}, at the limit of the realm of so-called ultrathin films.

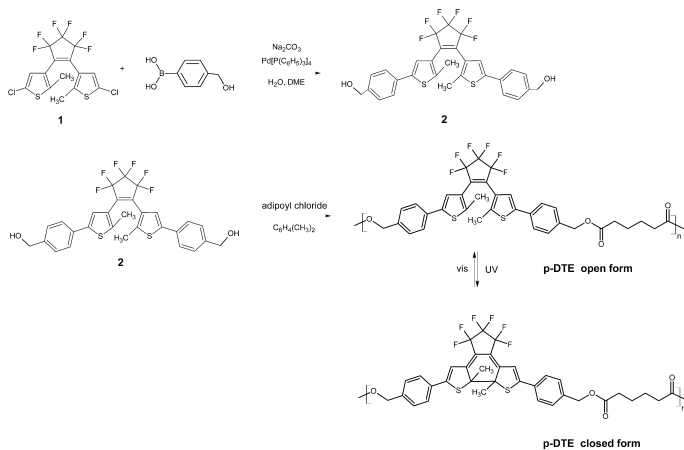


Fig. 1 Sketch of the synthetic route of p-DTE and of the photoinduced isomerization. See Sect.2.1 for details on the reagent compounds **1** and **2**.

2 Materials and methods

2.1 Materials and film preparation

Photochromic material

The diarylethene-based polyester was synthesized by step-polymerization between the 1,2-bis-(2-methyl-5-(p-hydroxymethylphenyl)-3-thienyl)perfluorocyclopentene and adipoyl chloride. This diarylethene monomer was preferred to the 1,2-bis-(2-methyl-5-(hydroxyphenyl)-3-thienyl)perfluorocyclopentene, which is difficult to purify by flash chromatography of the raw product and is rather unstable.

Reactions have been carried out under a dry, oxygen-free argon atmosphere. Unless otherwise specified, all reagents and catalysts were commercial (Aldrich, Alfa Aesar). ¹H NMR spectra of a solution in CDCl₃ were recorded on a Bruker AXR 400 spectrometer at 400 MHz.

1,2-Bis-(5-p-hydroxymethylphenyl-2-methyl-3-thienyl)perfluorocyclopentene (**2**)

Following the reaction procedure described in Ref.⁴⁴, 4-(hydroxymethyl)phenyl boronic acid (0.383 g, 2.52 mmol), 1,2-bis-(2-methyl-5-chloro-3-thienyl)perfluorocyclopentene (**1**) (0.5 g, 1.14 mmol), Na₂CO₃ · 10 H₂O (1.31 g, 4.58 mmol) and Pd(PPh₃)₄ (0.132 g, 0.11 mmol) were dissolved in previously degassed DME (20 mL) and water (5 mL). The reaction mixture was refluxed under argon for 24 hours, then extracted with diethyl ether and water and the combined organic layers were dried over Na₂SO₄. Purification by flash chromatography on silica gel (ethyl acetate) provided 570 mg of the desired product in 86% yield.

(¹H NMR (400 MHz, CDCl₃): δ 7.54 (d, 4H, J(H,H) = 8.2 Hz), 7.38 (d, 4H, J(H,H) = 8.2 Hz), 7.28 (s, 2H), 4.71 (s, 4H), 1.97 (s, 6H) ppm.)

p-DTE 1,2-Bis-(2-methyl-4-(p-hydroxymethylphenyl)-3-thienyl)perfluorocyclopentene (**2**) (0.5 g, 0.86 mmol) and adipoyl chloride (0.125 mL, in stoichiometric ratio 1:1) in anhydrous xylene (30 mL) were stirred at reflux for 120 hours. The reaction mixture was poured into petroleum ether in vigorous stirring, the white precipitate was filtered (330 mg, 56% yield) and subsequently extracted with hot diethyl ether and chloroform, obtaining 140 mg (24% yield) and 190 mg (32% yield), respectively. No insoluble residue was recovered.

(¹H NMR (400 MHz, CDCl₃): δ 7.51 (d, 4H, J(H,H) = 8.09 Hz), 7.37 (d, 4H, J(H,H) = 8.09 Hz), 5.11 (s, 4H), 2.38 (t, 4H), 1.98 (s, 6H), 1.70 (t, 4H) ppm.)

Film preparation and processing

The wafers were carefully prepared and characterised

by SE, following procedures described in a previous work⁴⁵. Films were cast by spin coating (angular speed between 1800 and 3600 rpm) on SiO₂/Si substrates using solutions with different concentration (5 mg/ml and 15 mg/ml in toluene), obtaining thickness in the 14-40 nm range. After the deposition, the polymer was isomerized into its non-cyclic form using a red laser (MRL-III, 655 nm, 400 mW, CNI Opto-electronics Tech. Co.). The cyclization reaction was obtained using a UV lamp (311 nm, PL-S 9W/01/2P ICT, Philips).

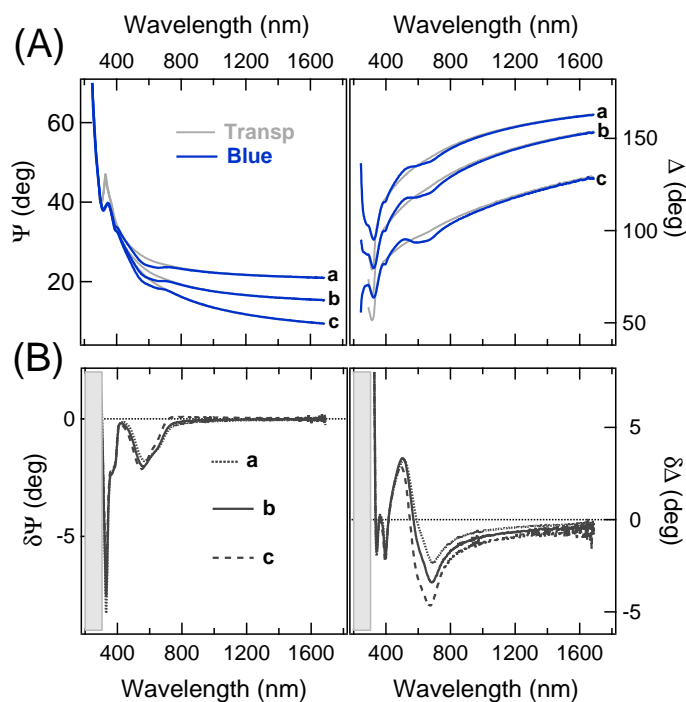


Fig. 2 (A) SE spectra of very thin films of p-DTE (thickness ~ 35 nm) in the colourless and blue form. Data are reported for three angles of incidence (**a** 60°, **b** 65° and **c** 70°) (B) Difference spectra ($\delta\Psi(\Delta) = \Psi(\Delta)_{blue} - \Psi(\Delta)_{transp}$) obtained from data of panel A.

2.2 Methods

Spectroscopic Ellipsometry

Principles and applications of SE are described at length in books^{46–48} and recent reviews^{49–51}. Specific issues pertinent to ultrathin films and monolayers are also addressed in Ref.⁵². Standard ellipsometry investigates the coefficient $\rho = \tilde{r}_p / \tilde{r}_s = \tan \Psi \exp(i\Delta)$ where \tilde{r}_p and \tilde{r}_s are the complex Fresnel reflection coefficients for p- and s-polarization, respectively. Information on the optical (e.g. the complex index of refraction) and morphological properties (such as e.g. thickness, roughness) of the system under investigation are obtained from the comparison of data with simulations. To this aim,

the system is usually modelled as a film (or a stack of layers) on a substrate. The thickness and/or some convenient quantity associated to the film optical properties are used as adjustable parameters in automated regression analysis of data looking for the best agreement between experimental Ψ and Δ spectra and simulations, exploiting the minimization of a convenient biased estimator (in this case the Mean Squared Error, MSE, function^{53,54†}).

Data reported in this paper, in the 245-1700 nm wavelength range, were obtained by a rotating compensator instrument (M-2000, J.A. Woollam Co. Inc.) equipped with a 75 W Xe lamp. The data analysis has been carried out using the software VASE supplied by the ellipsometer manufacturer. Measurements have been performed in darkened room ambient at 60°, 65° 70° angles of incidence. The spot size on the sample was of the order of a few mm². We performed *static* measurements, under steady-state conditions (transparent or coloured form of the polymer), and *dynamic* measurements, monitoring the photo-conversion process. *Static* measurements on the transparent form as well as *dynamic* measurements during the transition to the blue form were performed employing a glass sheet filtering out UV radiation from the beam ($\lambda < 305$ nm). *Static* measurements on the blue form have been taken on the full spectral range. Transition-induced spectral changes were emphasized by calculating *difference spectra* between data collected on the coloured and transparent phase ($\delta\Delta = \Delta_{blue} - \Delta_{transp}$, $\delta\Psi = \Psi_{blue} - \Psi_{transp}$), adapting approaches we recently applied to study molecular monolayers on solid substrates^{55–59}. SE data were complemented with AFM analysis of the film surface morphology †.

3 Results and discussion

Representative *static* Ψ and Δ spectra obtained on very thin films of p-DTE, in the colourless and blue form, are presented in fig.2-A (on the example of a ≈ 35 nm-thick film). The spectral differences between the two isomers are emphasized in fig.2-B. The $\delta\Psi$ pattern shows two evident negative structures, the sharpest at about 330 nm, and a broad one peaked at about 600 nm, whose position agrees with optical absorptions of the molecule measured in solution†. $\delta\Psi$ values vanish for λ in the NIR region. Regarding $\delta\Delta$ data, one can note the relatively broad, wiggling feature in the 400-800 nm range, in good correspondence with the negative feature in $\delta\Psi$ data. The $\delta\Delta$ values remain negative in the NIR range. The situation resembles the one encountered in the study of chromophores on transparent substrates⁶⁰, where $\delta\Psi$ and $\delta\Delta$ spectra were affected by adsorption-induced variation of the imaginary and real part of \tilde{n} , respectively. Such a relation between the spectral features of Fig.2-B and the blue form optical properties, intimately connected to Kramers-Kronig

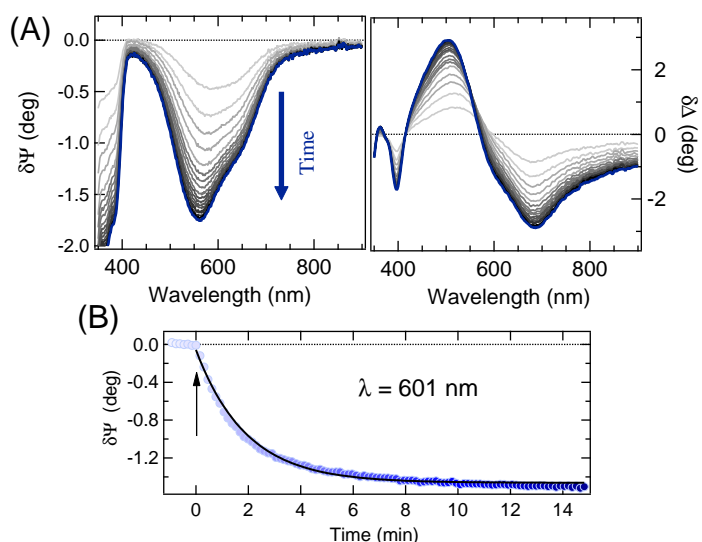


Fig. 3 (A) Selection of dynamic difference spectra (65° angle of incidence) recorded during UV irradiation of the film of Fig.2; the time separation among spectra is equal to 28 seconds. (B) Symbols: dynamic $\delta\Psi$ plot at $\lambda = 601$ nm during the photochromic transition (65° angle of incidence). The arrow marks the onset of UV irradiation. The line shows an exponential fit to data.

relations, will be more evident after the comparison with simulations.

Fig.3-A, reporting the evolution of the SE data during the irradiation at constant UV photon flux of an initially transparent film, nicely illustrates the capability of SE to monitor relatively slow transients. For a better appreciation of the transition behaviour it may be useful to consider the value of a convenient ellipsometry parameter as a function of time, as it is often done in single wavelength ellipsometry studies⁶¹. This is shown in panel (B) on the example of the $\delta\Psi$ value at 601 nm. Even though $\delta\Psi$ data alone can not be directly converted into extinction coefficient data, it is worth noting that the observed exponential behaviour is compatible with the conversion kinetics of relatively unorganized samples^{62,63}. Note that a close look at data of Fig.3-B, before switching on the UV irradiation, is useful to appreciate the overall negligible effect of the probing beam on film coloration.

Coming to the quantitative analysis of SE data, the optical properties and the thickness of the photochromic film were determined by comparison with simulations of the optical system based on suitable models.

In principle, solvent cast films exhibit preferential orientation of the polymer backbones parallel to the substrate^{64,65}. However, approaching the ultrathin thickness limit, correlations among optical fitting parameters and between optical parameters and thickness prevent intrinsically an unambigu-

ous determination of the optical anisotropy by ellipsometry methods⁶⁶⁻⁶⁸. Therefore, the simplest, isotropic two-layer model was considered, consisting of a homogeneous film on the substrate, in air environment.

As amply documented in Ref.⁶⁹, the optical constants deduced by the analysis of SE data through an isotropic model approximately correspond to the in-plane optical constants, while no information can be inferred on the extraordinary index.

We tested the introduction of the so-called surface roughness (SR) layer⁴⁶⁻⁴⁸. The SR layer thickness, treated as free parameter, was suppressed by the MSE minimization routine (\dagger). In this regard, note that AFM measurements (\dagger) indicated average roughness values ranging from 0.2 nm, on very uniform regions (typically extending for some microns squared), up to $0.5 \div 0.7$ nm on larger regions ($50 \times 50 \mu m$) including some defective areas. We didn't observe any significant change of the AFM roughness related to photo-conversion processes. Note however that the AFM laser operation may perturb the coloration of films.

The substrate pseudo-dielectric functions were predetermined from measurements carried out before film deposition. The film optical properties were described through a simple, multi-oscillator model preserving the K-K consistency of data^{24,33,58}:

$$\epsilon(h\nu) = \epsilon_\infty + \sum_i G_i(A_i, E_i, \beta_i). \quad (1)$$

ϵ_∞ is a real offset accounting for resonances located out of the measurement region. A_i , E_i and β_i are the amplitude, energy position, and broadening parameter, respectively, of the i -th oscillator, of Gaussian shape. ϵ_∞ , A_i , E_i and β_i , as well as the film thickness, were considered *free* parameters to be adjusted in the best-fit routine. Two sets of oscillators were necessary to reproduce the SE spectra of the two isomers. We opted for the criterion of keeping the number of oscillators in each set at minimum. A reliable guess of the oscillators parameters, useful to feed the fitting routine, was derived from transmission measurements in solution (\dagger) or on thick films cast on fused silica. Regarding the visible range we assumed peaks of the same width. The comparison between experimental data and best-fit simulations (multi-angle analysis) for a representative sample is presented in Fig.4 \dagger . The K-K consistent real $\text{Re}(\tilde{n})$ and imaginary $\text{Im}(\tilde{n})$ part of the complex refractive index of the blue and transparent phase derived from the best fit for the sample of Fig.4 are shown in Fig.5. The curves are fully representative of very similar results obtained on several 22-38 nm thick films (\dagger). Note that regarding $\text{Im}(\tilde{n})$, a plot as a function of photon energy is a choice more functional to emphasize the contribution of each oscillator assumed in the fit(\dagger). The attainment of the best fit of the blue-phase spectra required a minimum of two oscillators in the visible range, located by the

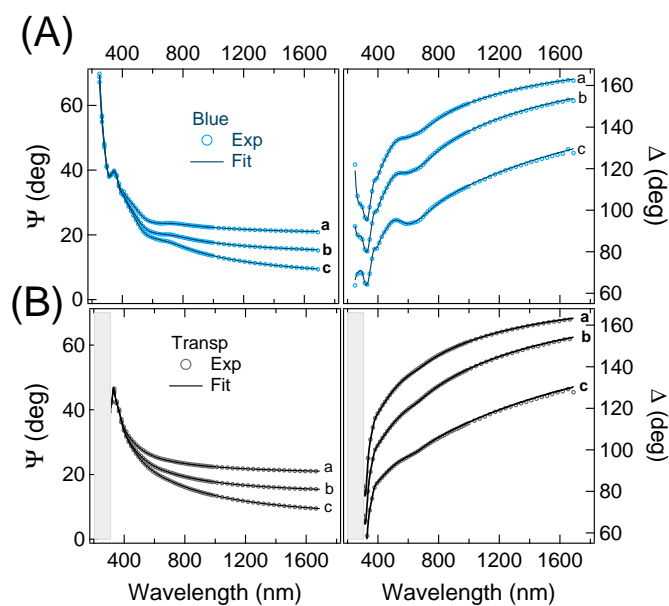


Fig. 4 Comparison between experimental and simulated (best-fit) SE spectra of a representative thin film of p-DTE in the (A) blue and (B) transparent phases (same data as in Fig. 2). The resulting film thickness was 35.4 ± 0.1 nm and 35.9 ± 0.1 for the blue and colourless phase, respectively. The resulting refractive index and extinction coefficient in the two phases are shown in Fig. 5.

fit at about 620 nm (~ 2.0 eV) and 540 nm (~ 2.3 eV), and four in the UV. The nearest UV absorption (about 385 nm, 3.2 eV) is quenched in the transparent form, in agreement with literature on related systems^{39,41}. An even better reproduction of the data in the 400-800 nm range could be obtained considering (i) three narrower peaks instead of two or (ii) independent broadening parameters for the two peaks. However, in both cases the number of free parameter increased and the improvement of MSE was overall poor.

Regarding the transparent phase, except for some blue phase residuals, one broad and intense oscillator in the UV, centred at 306 nm was enough to fit the data. Coming back to Fig. 4 and zooming in the visible range, one can observe that the experimental features related to blue-phase residuals in panel B are in practice discernible only at 70° of incidence; indeed, a reasonably good fit of the data can be obtained also for a perfectly uncoloured film. Thus, the small peak of $\text{Im}(\tilde{n})$ in Fig. 5 at about 600 nm can be considered as an upper estimate for blue phase residuals and, in a sense, may represent a figure of the overall uncertainty in the characterization of the transparent state. Finally, a low amount of residuals, possibly due to incomplete photoisomerization and imperfect UV shielding, does not invalidate the main conclusions of the analysis. In Fig. 5 note also the position of the isosbestic point, located by the fit at 315 nm, in excellent agreement with measurements in solution (\dagger). Another interesting feature in Fig. 5 is

the neat difference of $\text{Re}(\tilde{n})$ for the two phases for wavelength in the NIR. Regarding this aspect, an extended review of previous works can be found in Ref.⁸. While a close comparison with literature is not straightforward due to the diversity of materials, chromophore content and wavelengths considered in other works^{41,70}, we note that the resulting variation of the refractive index $\Delta \text{Re}(\tilde{n})$ appears higher than some values reported for other diarylethene polymers, e.g. 10^{-3} at 633 nm² and 9×10^{-3} at 830 nm⁷¹, whereas it is definitely comparable with values of about 0.04 measured, with single wavelength methods (at about 800 nm), for bulk phases of other diarylethene derivatives^{39,40}. Note that the choice of the wavelength, crucial for a reliable assessment of $\Delta \text{Re}(\tilde{n})$, appears in some papers very close to the coloured polymer resonances. Another aspect to be discussed is the way the

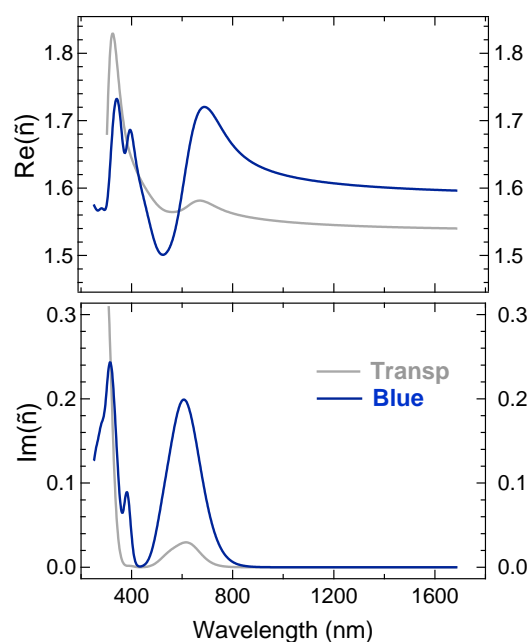


Fig. 5 Real and imaginary part of the complex refractive index of a representative film in the blue and transparent phase. The blue-form residuals in the transparent phase are due to incomplete photo-isomerization.

thickness parameter is considered in the data fitting. In our analysis it was allowed to vary upon the transition. The film thickness obtained e.g. for the represented sample of Fig. 5 showed a slight contraction, from 35.9 to 35.4 nm, from the open form to the closed form. Such a trend was observed for all films considered in our experiment (\dagger) and, to the best of our knowledge, is the first indication of a photomechanical response in diarylethene amorphous films. The thickness variation is tentatively assigned to film density changes induced by the ring-closure. Variations of density have been

observed for several diarylethenes in crystalline form, the ring closure being accompanied by either an increase^{72–77} or a decrease of density^{77–79} depending on the specific molecular structure. In the former cases, changes of density of 0.39–5.6% have been reported whereas a more limited variation ranging from 0.14 to 1.6% is shown when a density decrease results from the ring closure. If we preferably refer to the structures that are similar to the monomeric unit of p-DTE, a variation of 1.38% is reported in case of 1,2-bis(2-methyl-5-phenyl-3-thienyl)perfluorocyclopentene⁷² and 2.62% for one polymorph of the 1,2-bis(2-methyl-5-(p-methoxyphenyl)-3-thienyl)perfluorocyclopentene⁷⁷.

In a previous SE investigation of a dithienylethene embedded in PMMA matrix⁵ the thickness parameter (660 nm) was assumed constant across the photochromic transition and $\Delta\text{Re}(\tilde{n})$ was estimated to be 5×10^{-3} . Following Ref.⁵ we made a check on the correlation between thickness and refractive index, looking at the best fit on the transparent phase while keeping the thickness at the value found on the blue form; at the expense of a reduction of the fit quality (MSE=13 instead of 7), we found an increased value of $\text{Re}(\tilde{n})$ (1.57 instead of 1.54 at 1700 nm), still significantly lower than the corresponding value in the blue phase (about 1.60).

The discussion so far focused on 30–40 nm thick films, for which the relatively sharp *ellipsometric contrast* between the uncoloured and blue form (Fig.2-B) guaranteed a robust analysis of data. We were able to obtain a satisfactory reproduction of SE experimental data for film thickness down to the 20 nm limit, with relatively small adjustments of the optical parameters of Fig.5 (†). Regarding even thinner films (14–17 nm), the spectral differences between the two isomers becomes smaller and smaller (†) and we opted to analyse the data keeping the optical absorptions parameters of the blue form at the same values obtained in the analysis of the thicker films, in practice fitting only the film thickness. Within this approximation, we observe that the analysis of ultrathin films upon film colouration confirmed both a positive variation of $\text{Re}(\tilde{n})$ and a small decrease of thickness[†].

4 Concluding remarks

This work focused on the characterization of the optical properties of very thin films of a photochromic polymer spin cast on a silicon wafer, carried out by application of broad-band spectroscopic ellipsometry. The use of a rotating compensator ellipsometer allowed to perform a real time characterization of the photochromic transition during UV irradiation of the sample, clearly identifying ellipsometric features related to the polymer colour. Through the comparison of SE data with simulations based on a isotropic multi-resonance model accounting for molecular UV-VIS absorptions, we determined the complex index of refraction of the compound in both the

blue and colourless forms in a spectral region extending from the UV to the NIR range (245–1700 nm). Results obtained for films in the 14–40 nm thickness range indicate that the real part of the refractive index of the blue phase neatly exceeds the one of the transparent form for wavelengths above 600 nm: $\Delta\text{Re}(\tilde{n})$ is ~ 0.05 in the IR limit. Interestingly, the analysis of data, within the approximations involved in the proposed model, indicate a decrease (about 1.5%) of thickness in the transition from the open (transparent) to the closed (blue) form, assigned to film density changes induced by the ring-closure process. The results of this work pave the way to the exploitation of ultrathin photochromic films in stimuli responsive photonic crystals structures and related devices, as well as in photo-responsive actuators.

Acknowledgements

This work received partial support from the University of Genova (PRA grant 2011-12). Partial support from the UE FP7 Opticon Project is also acknowledged.

References

- 1 J. Clark and G. Lanzani, *Nature Photonics*, 2010, **4**, 438–446.
- 2 N. Tanio and M. Irie, *Japanese Journal of Applied Physics*, 1994, **33**, 3942–3946.
- 3 M. Irie, *Chemical Reviews*, 2000, **100**, 1685–1716.
- 4 H. Cho and E. Kim, *Macromolecules*, 2002, **35**, 8684–8687.
- 5 C. Bertarelli, A. Bianco, F. D'Amore, M. Gallazzi and G. Zerbi, *Advanced Functional Materials*, 2004, **14**, 357–363.
- 6 L. Angiolini, T. Benelli, L. Giorgini and E. Salatelli, *Macromolecules*, 2006, **39**, 3731–3737.
- 7 T. Ikeda, J.-i. Mamiya and Y. Yu, *Angewandte Chemie International Edition*, 2007, **46**, 506–528.
- 8 C. Bertarelli, A. Bianco, R. Castagna and G. Pariani, *Journal of Photochemistry and Photobiology C-Photochemistry Reviews*, 2011, **12**, 106–125.
- 9 S. Menghetti, M. Alderighi, G. Galli, F. Tantussi, M. Morandini, F. Fuso and M. Allegrini, *Journal Of Materials Chemistry*, 2012, **22**, 14510–14517.
- 10 J. Garca-Amors and D. Velasco, *Beilstein Journal of Organic Chemistry*, 2012, **8**, 1003–1017.
- 11 A. Ryabchun, A. Bobrovsky, A. Sobolewska, V. Shibaev and J. Stumpe, *J. Mater. Chem.*, 2012, **22**, 6245–6250.
- 12 D. Comoretto, G. Dellepiane, F. Marabelli, J. Cornil, D. A. dos Santos, J. L. Brédas and D. Moses, *Phys. Rev. B*, 2000, **62**, 10173–10184.
- 13 M. Campoy-Quiles, G. Heliotis, R. Xia, M. Ariu, M. Pintani, P. Etchegoin and D. Bradley, *Advanced Functional Materials*, 2005, **15**, 925–933.
- 14 C. Soci, D. Comoretto, F. Marabelli and D. Moses, *Phys. Rev. B*, 2007, **75**, 075204.
- 15 M. Brinkmann, F. Chandezon, R. B. Pansu and C. Julien-Rabant, *Advanced Functional Materials*, 2009, **19**, 2759–2766.
- 16 N. Sary, F. Richard, C. Brochon, N. Leclerc, P. Lvque, J.-N. Audinot, S. Berson, T. Heiser, G. Hadziioannou and R. Mezzenga, *Advanced Materials*, 2010, **22**, 763–768.
- 17 K. A. Bell, L. Mantese, U. Rossow and D. E. Aspnes, *The Journal of Vacuum Science and Technology B*, 1997, **15**, 1205–1211.
- 18 J. Shi, B. Hong, A. Parikh, R. Collins and D. Allara, *Chemical Physics Letters*, 1995, **246**, 90 – 94.
- 19 L. J. Richter, C. S.-C. Yang, P. T. Wilson, C. A. Hacker, R. D. van Zee, J. J.

- Stapleton, D. L. Allara, Y. Yao and J. M. Tour, *The Journal of Physical Chemistry B*, 2004, **108**, 12547–12559.
- 20 M. Prato, M. Alloisio, S. A. Jadhav, A. Chincarini, T. Svaldo-Lanero, F. Bisio, O. Cavalleri and M. Canepa, *The Journal of Physical Chemistry C*, 2009, **113**, 20683–20688.
- 21 E. Bittrich, M. Kuntzsch, K.-J. Eichhorn and P. Uhlmann, *Journal of Polymer Science part b-Polymer Physics*, 2010, **48**, 1606–1615.
- 22 Y. Kim, R. Collins, K. Vedam and D. Allara, *Journal of the Electrochemical Society*, 1991, **138**, 3266–3275.
- 23 C. Barbero and R. Kotz, *Journal of the Electrochemical Society*, 1994, **141**, 859–865.
- 24 Y.-G. Mo, R. O. Dillon and P. G. Snyder, *Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films*, 1999, **17**, 170–175.
- 25 S. Logothetidis, *Materials Science and Engineering b-Advanced Functional Solid-state Materials*, 2008, **152**, 96–104.
- 26 L. Ionov, A. Sidorenko, K. Eichhorn, M. Stamm, S. Minko and K. Hinrichs, *Langmuir*, 2005, **21**, 8711–8716.
- 27 R. Carpick, T. Mayer, D. Sasaki and A. Burns, *Langmuir*, 2000, **16**, 4639–4647.
- 28 R. W. Carpick, D. Y. Sasaki, M. S. Marcus, M. A. Eriksson and A. R. Burns, *Journal of Physics: Condensed Matter*, 2004, **16**, R679.
- 29 T. Wang, A. D. F. Dunbar, P. A. Staniec, A. J. Pearson, P. E. Hopkinson, J. E. MacDonald, S. Lilliu, C. Pizzey, N. J. Terrill, A. M. Donald, A. J. Ryan, R. A. L. Jones and D. G. Lidzey, *Soft Matter*, 2010, **6**, 4128–4134.
- 30 A. J. Pearson, T. Wang, R. A. L. Jones, D. G. Lidzey, P. A. Staniec, P. E. Hopkinson and A. M. Donald, *Macromolecules*, 2012, **45**, 1499–1508.
- 31 T. Wang, A. J. Pearson, D. G. Lidzey and R. A. L. Jones, *Advanced Functional Materials*, 2011, **21**, 1383–1390.
- 32 I. Zudans, W. Heineman and C. Seliskar, *Journal of Physical Chemistry B*, 2004, **108**, 11521–11528.
- 33 C. Akerlind, H. Arwin, F. L. E. Jakobsson, H. Kariis and K. Jarrendahl, *Thin Solid Films*, 2011, **519**, 3582–3586.
- 34 G. Callierotti, A. Bianco, C. Castiglioni, C. Bertarelli and G. Zerbi, *Journal of Physical Chemistry A*, 2008, **112**, 7473–7480.
- 35 F. Stellacci, C. Bertarelli, F. Toscano, M. C. Gallazzi, G. Zotti and G. Zerbi, *Advanced Materials*, 1999, **11**, 292–295.
- 36 G. Pariani, R. Castagna, G. Dassa, S. Hermes, C. Vailati, A. Bianco and C. Bertarelli, *Journal Of Materials Chemistry*, 2011, **21**, 13223–13231.
- 37 H. Nishi, T. Namari and S. Kobatake, *Journal of Materials Chemistry*, 2011, **21**, 17249–17258.
- 38 T. Wigglesworth, A. Myles and N. Branda, *European Journal Of Organic Chemistry*, 2005, 1233–1238.
- 39 J. Biteau, F. Chaput, K. Lahliu, J.-P. Boilot, G. M. Tsigoulis, J.-M. Lehn, B. Darracq, C. Marois and Y. Lvy, *Chemistry of Materials*, 1998, **10**, 1945–1950.
- 40 T. Kawai, N. Fukuda, D. Gröschl, S. Kobatake and M. Irie, *Japanese Journal of Applied Physics*, 1999, **38**, L1194–L1196.
- 41 M.-S. Kim, H. Maruyama, T. Kawai and M. Irie, *Chemistry of Materials*, 2003, **15**, 4539–4543.
- 42 S. Kobatake, S. Takami, H. Muto, T. Ishikawa and M. Irie, *Nature*, 2007, **446**, 778–781.
- 43 M. Morimoto and M. Irie, *Journal of the American Chemical Society*, 2010, **132**, 14172–14178.
- 44 S. Hermes, G. Dassa, G. Toso, A. Bianco, C. Bertarelli and G. Zerbi, *Tetrahedron Letters*, 2009, **50**, 1614–1617.
- 45 G. Gonella, O. Cavalleri, I. Emilianov, L. Matterna, M. Canepa and R. Rolandi, *Materials Science & Engineering C-Biomimetic And Supramolecular Systems*, 2002, **22**, 359–366.
- 46 R. M. A. Azzam and N. M. Bashara, *Ellipsometry and Polarized Light*, North-Holland, New York, 3rd edn, 1977.
- 47 *Handbook of ellipsometry*, ed. H. Tompkins and E. Irene, Andrew, Norwich, 1st edn, 2005.
- 48 H. Fujiwara, *Spectroscopic Ellipsometry: Principles and Applications*, Wiley, Chichester, United Kingdom, 1st edn, 2007.
- 49 M. Losurdo, M. Bergmair, G. Bruno, D. Cattelan, C. Cobet, A. de Martino, K. Fleischer, Z. Dohcevic-Mitrovic, N. Esser, M. Galliet, R. Gajic, D. Hemzal, K. Hingerl, J. Humlicek, R. Ossikovski, Z. V. Popovic and O. Saxl, *Journal of Nanoparticle Research*, 2009, **11**, 1521–1554.
- 50 T. W. H. Oates, H. Wormeester and H. Arwin, *Progress in Surface Science*, 2011, **86**, 328–376.
- 51 D. E. Aspnes, *Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films*, 2013, **31**, 058502.
- 52 M. Canepa, in *A Surface Scientist's View of Spectroscopic Ellipsometry*, ed. G. Bracco and B. Holst, Springer Series in Surface Sciences, Vol. 51, Berlin, Germany, 1st edn, 2013, pp. 99–135.
- 53 G. E. Jellison Jr., *Appl. Opt.*, 1991, **30**, 3354–3360.
- 54 J. N. Hilfiker, N. Singh, T. Tiwald, D. Convey, S. M. Smith, J. H. Baker and H. G. Tompkins, *Thin Solid Films*, 2008, **516**, 7979 – 7989.
- 55 F. Bordi, M. Prato, O. Cavalleri, C. Cametti, M. Canepa and A. Gliozzi, *The Journal of Physical Chemistry B*, 2004, **108**, 20263–20272.
- 56 M. Prato, R. Moroni, F. Bisio, R. Rolandi, L. Matterna, O. Cavalleri and M. Canepa, *The Journal of Physical Chemistry C*, 2008, **112**, 3899–3906.
- 57 H. Hamoudi, M. Prato, C. Dablemont, O. Cavalleri, M. Canepa and V. A. Esaulov, *Langmuir*, 2010, **26**, 7242–7247.
- 58 C. Toccafondi, M. Prato, G. Maidecchi, A. Penco, F. Bisio, O. Cavalleri and M. Canepa, *Journal of Colloid and Interface Science*, 2011, **364**, 125–132.
- 59 M. Canepa, G. Maidecchi, C. Toccafondi, O. Cavalleri, M. Prato, V. Chaudhari and V. A. Esaulov, *Phys. Chem. Chem. Phys.*, 2013, **15**, 11559–11565.
- 60 C. Toccafondi, O. Cavalleri, F. Bisio and M. Canepa, *Thin Solid Films*, 2013, **543**, 78 – 82.
- 61 H. Knobloch, S. Katholy, J. Hesse, H. Orendi, D. Prescher, R. Ruhmann and L. Brehmer, *Materials Science and Engineering: C*, 1998, **5**, 307 – 310.
- 62 J. Sworakowski, K. Janus and S. Nepurek, *Advances in Colloid and Interface Science*, 2005, **116**, 97 – 110.
- 63 G. Pariani, A. Bianco, R. Castagna and C. Bertarelli, *The Journal of Physical Chemistry A*, 2011, **115**, 12184–12193.
- 64 T. Sosnowski and H. Weber, *Applied Physics Letters*, 1972, **21**, 310–311.
- 65 W. M. Prest and D. J. Luca, *Journal of Applied Physics*, 1979, **50**, 6067–6071.
- 66 M. J. Dignam, M. Moskovits and R. W. Stobie, *Trans. Faraday Soc.*, 1971, **67**, 3306–3317.
- 67 C. Ramsdale and N. Greenham, *Advanced Materials*, 2002, **14**, 212–215.
- 68 M. Campoy-Quiles, G. Heliotis, R. Xia, M. Ariu, M. Pintani, P. Etchegoin and D. Bradley, *Advanced Functional Materials*, 2005, **15**, 925–933.
- 69 M. Campoy-Quiles, P. G. Etchegoin and D. D. C. Bradley, *Phys. Rev. B*, 2005, **72**, 045209.
- 70 E. Kim, Y.-K. Choi and M.-H. Lee, *Macromolecules*, 1999, **32**, 4855–4860.
- 71 S. Cho, M. Yoo, H.-W. Shin, K.-H. Ahn, Y.-R. Kim and E. Kim, *Optical Materials*, 2003, **21**, 279 – 284.
- 72 M. Irie, T. Lifka, S. Kobatake and N. Kato, *Journal of the American Chemical Society*, 2000, **122**, 4871–4876.
- 73 M. Irie, S. Kobatake and M. Horichi, *Science*, 2001, **291**, 1769–1772.
- 74 S. Kobatake, H. Hasegawa and K. Miyamura, *Crystal Growth & Design*, 2011, **11**, 1223–1229.
- 75 S. Kobatake, K. Shibata, K. Uchida and M. Irie, *Journal of the American Chemical Society*, 2000, **122**, 12135–12141.
- 76 K. Uchida, N. Izumi, S. Sukata, Y. Kojima, S. Nakamura and M. Irie, *Angewandte Chemie International Edition*, 2006, **45**, 6470–6473.
- 77 M. Morimoto, S. Kobatake and M. Irie, *Chemistry A European Journal*, 2003, **9**, 621–627.
- 78 S. Yamamoto, K. Matsuda and M. Irie, *Organic Letters*, 2003, **5**, 1769–

1772.

the American Chemical Society, 2000, **122**, 9631–9637.

79 T. Kodani, K. Matsuda, T. Yamada, S. Kobatake and M. Irie, *Journal of*