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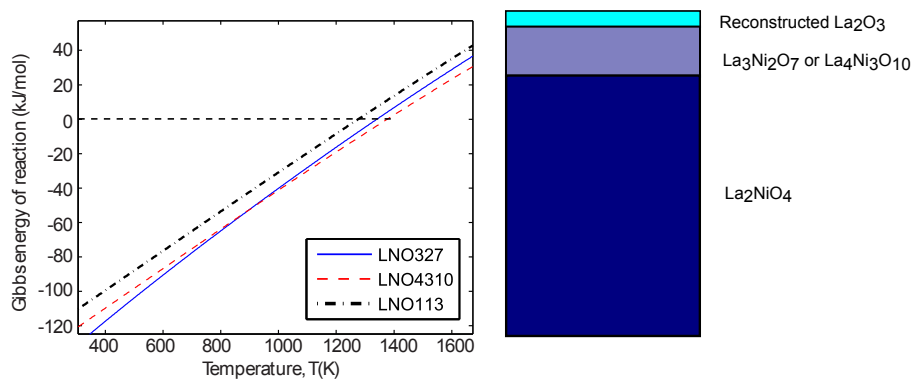


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Calculated Gibbs reaction energy for La_2NiO_4 decomposition (Left) and proposed resulting layered structure (Right) provides an explanation of the experimentally observed Ni-absent surfaces and Ni-enriched subsurface layers on La_2NiO_4 .



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Why Ni is Absent from the Surface of $\text{La}_2\text{NiO}_{4+\delta}$

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$\text{La}_2\text{NiO}_{4+\delta}$ (LNO214) is a potential intermediate temperature solid oxide fuel cell (IT-SOFC) cathode material which belongs to the Ruddlesden-Popper (RP) structure series $\text{A}_{n+1}\text{B}_n\text{O}_{3n+1}$. There is interest in this material as it offers a way to avoid Sr segregation and associated degradation, as LNO214 can take in oxygen interstitials and become catalytically active without A-site doping. While the bulk ionic conduction mechanisms are well studied, its surface structure and chemistry are still a matter of debate. Recent experimental works (both with and without dopants) reveal it has a La-terminated surface and a highly Ni deficient surface layer. These results disagree with earlier computer simulations, and undermine the conventional explanation for the oxygen reduction process at the surface. In this work we evaluate the thermodynamic stability of $\text{La}_2\text{NiO}_{4+\delta}$ at IT-SOFC operation temperatures. We find that the decomposition of $\text{La}_2\text{NiO}_{4+\delta}$ to produce La_2O_3 and higher order RP phases is indeed thermodynamically favourable. A hypothesis for the formation mechanism of the La-terminated and Ni deficient surface based on partial decomposition and surface passivation is proposed and evaluated.

1 INTRODUCTION

The Ruddlesden-Popper (RP) family material lanthanum nickelate, $\text{La}_2\text{NiO}_{4+\delta}$ (LNO214) (general formula $\text{A}_{n+1}\text{B}_n\text{O}_{3n+1}$), has attracted much interest in the past decade as a potential cathode material for intermediate temperature solid oxide fuel cells (IT-SOFC) operating at 600-800°C¹⁻⁶. It has decent mixed ionic and electronic conductivity (55-65 Scm^{-1} overall at 500-750°C)⁷ without the need of any A-site aliovalent doping due to the RP structure's ability to take in oxygen interstitials in the rock-salt like AO layers.⁴ Comparing to $\text{La}_{1-x}\text{Sr}_x\text{CoO}_{3-\delta}$ (LSC) and $\text{La}_{0.6}\text{Sr}_{0.4}\text{Fe}_{0.2}\text{Co}_{0.8}\text{O}_{3-\delta}$ (LSCF), the LNO214 has lower total conductivity and oxygen self-diffusion coefficient at 500-750°C, but it also does not suffer from the Sr-segregation problem of the Sr-doped materials.^{7,8} The Sr-segregation severely damages LSC-based materials performance over time (typically 0.05% efficiency loss per hour)^{8,9}. Therefore, LNO214 becomes a nice candidate for IT-SOFC cathode considering long-term stability and cost-effectiveness. Both the bulk structure and oxygen ion diffusion mechanism in LNO214 are well understood with insights coming from computer simulations⁶ and experiments^{4,5}. However, the $\text{La}_2\text{NiO}_{4+\delta}$ surface structure and chemistry, which strongly affects surface oxygen exchange and limits the rate of overall oxygen ion migration, have not been explored in detail.

There are only two sets of simulations of LNO214 surfaces performed so far^{10,11} to the best of the authors' knowledge. Read *et al.*¹⁰ suggested that LNO surfaces should terminate with Ni due to its lower surface energy. Zhou *et al.*¹¹ studied the oxygen adsorption mechanism based on the La_2NiO_4 surface proposed in Read's work. However, recent surface characterisation of similar RP structured materials, including $\text{La}_2\text{NiO}_{4+\delta}$, $\text{La}_{2-x}\text{Sr}_x\text{NiO}_{4+\delta}$ (LSNO) and $\text{PrLaNiO}_{4+\delta}$ (PLNO), has revealed that these surfaces terminate with purely A-site cations (La, Sr, Pr)¹²⁻¹⁴; no Ni was found in the low energy ion scattering (LEIS) spectra of the surfaces within the detection limit. The Ni:La(Sr) ratio was also found to be far from stoichiometric down to a depth of about 7 nm¹². A similar absence of Ni and Ni:La non-stoichiometry were obtained in LNO¹⁴. These results suggest the conventional oxygen reduction mechanism which requires oxygen atoms to adsorb onto surface Ni sites might not be possible.

In this study, we provide an explanation for the Ni absence at La_2NiO_4 surfaces by evaluating the thermodynamic stabilities of a set of related phases over a range of temperatures and oxygen partial pressures. The calculations are based on existing experimental thermodynamic data for relevant materials^{15,16}. Based on our calculations, a hypothesis explaining the absence of Ni on LNO surfaces is proposed and discussed. We find that it is thermodynamically advantageous for La_2NiO_4 to form a combination of La_2O_3 (containing no Ni, and presumed to appear at the surface) and higher order RP phases ($\text{La}_{n+1}\text{Ni}_n\text{O}_{3n+1}$ where $n = 2$ or 3).

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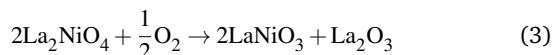
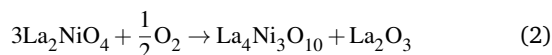
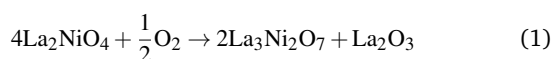
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Table 1 Thermodynamic parameters used to evaluate the change in Gibbs free energy for La_2NiO_4 decomposition. Note that for the phase transition of $\text{La}_3\text{Ni}_2\text{O}_7$ at 490K, we have $\Delta_r H = 1800\text{Jmol}^{-1}$, $\Delta_r S = 3.674\text{Jmol}^{-1}\text{K}^{-1}$.¹⁶ The values here are applicable from 305K to 1673K

	a (J mol^{-1})	b ($\text{J mol}^{-1}\text{K}^{-1}$)	c ($\text{J mol}^{-1}\text{K}^{-1}$)	d_2 ($\text{J mol}^{-1}\text{K}^{-2}$)	d_{-1} (J K mol^{-1})	Ref.
La_2NiO_4	-2095507.6	1035.849	-178.04	-0.007089	1214190	16
$\text{La}_3\text{Ni}_2\text{O}_7$	-3358980.1	1781.990	-301.86	-0.01712	1580000	16
$\text{La}_4\text{Ni}_3\text{O}_{10}$	-4599431.3	2368.174	-405.25	-0.0326	2130000	16
LaNiO_3	-1248967.7	696.185	-118.12	-0.0096	648000	16
La_2O_3	-1835600	674.72	-118	-0.008	620000	17

2 METHOD

The thermodynamic stability of LNO at any given temperature and oxygen partial pressure can be determined by considering the change in Gibbs free energy for its possible decomposition reactions. In this work, the following reactions are considered:



For ease of comparison, all reactions were normalised to consume half a mole of oxygen gas and produce one mole of La_2O_3 .

The change in Gibbs free energy for the above reactions at different temperatures and oxygen partial pressures can be written as:

$$\Delta G_{\text{reaction}}^{327}(T, p_{\text{O}_2}) = 2\Delta G_{\text{La}_3\text{Ni}_2\text{O}_7}(T) + \Delta G_{\text{La}_2\text{O}_3}(T) - 4\Delta G_{\text{La}_2\text{NiO}_4}(T) - \frac{1}{2}\Delta\mu_{\text{O}_2}(T, p_{\text{O}_2}) \quad (4)$$

$$\Delta G_{\text{reaction}}^{4310}(T, p_{\text{O}_2}) = \Delta G_{\text{La}_4\text{Ni}_3\text{O}_{10}}(T) + \Delta G_{\text{La}_2\text{O}_3}(T) - 3\Delta G_{\text{La}_2\text{NiO}_4}(T) - \frac{1}{2}\Delta\mu_{\text{O}_2}(T, p_{\text{O}_2}) \quad (5)$$

$$\Delta G_{\text{reaction}}^{113}(T, p_{\text{O}_2}) = 2\Delta G_{\text{LaNiO}_3}(T) + \Delta G_{\text{La}_2\text{O}_3}(T) - 2\Delta G_{\text{La}_2\text{NiO}_4}(T) - \frac{1}{2}\Delta\mu_{\text{O}_2}(T, p_{\text{O}_2}) \quad (6)$$

where $\Delta G_{\text{reaction}}(T, p_{\text{O}_2})$ is the change in Gibbs free energy of the respective reactions depending on temperature, T , and oxygen partial pressure p_{O_2} , $\Delta G_{\text{component}}(T)$ is the Gibbs formation energy of the component at temperature T and $\Delta\mu_{\text{O}_2}(T, p_{\text{O}_2})$ is the relative chemical potential (with respect to standard condition) of oxygen molecules at temperature T and oxygen partial pressure p_{O_2} . The standard condition are to 298.15 K and 1 bar pressure.

The standard molar Gibbs formation energy, $\Delta G^\circ(T)$, of a compound phase can be expressed as:

$$\Delta G^\circ(T) = G(T) - H^{\text{SER}} = a + bT + cT \ln(T) + \sum d_n T^n \quad (7)$$

where H^{SER} is the phases' Stable Element Reference of its constituent elements at 298.15 K and 1 bar. The coefficients a , b , c and d_n presented in Table 1 are taken from fits to experimental heat capacity measurements of the respective compounds using

Table 2 Oxygen thermodynamic constants from NIST handbook for various temperature ranges.¹⁹

Temperature (K)	100-700	700-2000	2000-6000
A ($\text{kJ mol}^{-1}\text{K}^{-1}$)	31.32234	30.03235	20.91111
B ($\text{kJ mol}^{-1}\text{K}^{-2}$)	-20.23531	8.772972	10.72071
C ($\text{kJ mol}^{-1}\text{K}^{-3}$)	57.86644	-3.988133	-2.020498
D ($\text{kJ mol}^{-1}\text{K}^{-4}$)	-36.50624	0.788313	0.146449
E (kJ K mol^{-1})	-0.007374	-0.741599	9.245722
F (kJ mol^{-1})	-8.903471	-11.32468	5.337651
G (kJ mol^{-1})	246.7945	236.1663	237.6185
H (kJ mol^{-1})	0.0	0.0	0.0

differential scanning calorimetry (DSC)^{16,17}.

The chemical potential of oxygen is given by^{18,19}

$$\Delta\mu_{\text{O}_2}(T, p_{\text{O}_2}) = H_{\text{O}_2}^0 - H_{\text{O}_2, RT}^0 - TS_{\text{O}_2}^0 + k_B T \ln\left(\frac{p_{\text{O}_2}}{p_{\text{O}_2}^0}\right) \quad (8)$$

where

$$H_{\text{O}_2}^0 - H_{\text{O}_2, RT}^0 = At + B\frac{t^2}{2} + C\frac{t^3}{3} + D\frac{t^4}{4} - \frac{E}{t} + F - H \quad (9)$$

and

$$S_{\text{O}_2}^0 = A \ln(t) + Bt + C\frac{t^2}{2} + D\frac{t^3}{3} - \frac{E}{(2t^2)} + G \quad (10)$$

The constants A, B, C, D, E, F, G, H are given in Table 2, and are obtained from the NIST chemical data handbook.¹⁹ The temperature coefficient t is $\frac{T}{1000}$, where T is temperature in kelvin. The unit of $H_{\text{O}_2}^0 - H_{\text{O}_2, RT}^0$ is kJmol^{-1} , where RT stands for room temperature (298.15K), and the unit of $S_{\text{O}_2}^0$ is $\text{Jmol}^{-1}\text{K}^{-1}$.

The Gibbs free energy of selected decomposition reactions can now be calculated at various temperatures and oxygen partial pressures, based on thermodynamic equations and associated constants given above.

It is important to note that the lanthanum nickelates are not stoichiometric when operating as a cathode under SOFC working conditions, where oxygen partial pressure is equal or more than 0.2 bar. La_2NiO_4 is usually oxygen rich, while the higher order $\text{La}_3\text{Ni}_2\text{O}_7$ and $\text{La}_4\text{Ni}_3\text{O}_{10}$ are oxygen deficient. In this study, the thermodynamic equation constants of the compounds used are fitted as stoichiometric phases based on thermal measurements on real, non-stoichiometric phases¹⁵⁻¹⁷. To provide a better prediction under experimental conditions, the error due to this difference should be estimated. For example, the Gibbs free energy

of formation of the non-stoichiometric $\text{La}_2\text{NiO}_{4+\delta}$ is given as:

$$\Delta G_{\text{La}_2\text{NiO}_{4+\delta_{214}}}(T, p_{\text{O}_2}) = \mu_{\text{La}_2\text{NiO}_{4+\delta_{214}}} - (2\mu_{\text{La}} + \mu_{\text{Ni}} + 2\mu_{\text{O}_2} + \frac{\delta_{214}}{2}\mu_{\text{O}_2}) \quad (11)$$

where $\Delta G_{\text{La}_2\text{NiO}_{4+\delta_{214}}}$ is the Gibbs free energy of formation of $\text{La}_2\text{NiO}_{4+\delta_{214}}$ at temperature, T , and oxygen partial pressure, p_{O_2} . The chemical potential of $\text{La}_2\text{NiO}_{4+\delta_{214}}$ at this temperature and pressure is $\mu_{\text{La}_2\text{NiO}_{4+\delta_{214}}}$, while μ_{La} , μ_{Ni} and μ_{O_2} are the chemical potentials of reference elemental La, Ni and oxygen gas at standard temperature and pressure.

To get the uncertainty in the experimental value, we need the stoichiometric Gibbs free energy of formation to compare with. The formation free energy of stoichiometric La_2NiO_4 is given by:

$$\Delta G_{\text{La}_2\text{NiO}_4}(T, p_{\text{O}_2}) = \mu_{\text{La}_2\text{NiO}_4} - (2\mu_{\text{La}} + \mu_{\text{Ni}} + 2\mu_{\text{O}_2}) \quad (12)$$

where $\Delta G_{\text{La}_2\text{NiO}_4}$ is the Gibbs free energy of formation of La_2NiO_4 at temperature, T , and oxygen partial pressure, p_{O_2} . The chemical potential of La_2NiO_4 at this temperature and pressure is $\mu_{\text{La}_2\text{NiO}_4}$. The definition of μ_{La} , μ_{Ni} and μ_{O_2} are the same as above.

The difference between the experimental $\Delta G_{\text{La}_2\text{NiO}_{4+\delta}}$ and ideal $\Delta G_{\text{La}_2\text{NiO}_4}$ is then:

$$\begin{aligned} \Delta G_{\text{La}_2\text{NiO}_4}^{\text{diff}} &= \Delta G_{\text{La}_2\text{NiO}_{4+\delta_{214}}} - \Delta G_{\text{La}_2\text{NiO}_4} \\ &= \mu_{\text{La}_2\text{NiO}_{4+\delta_{214}}} - \mu_{\text{La}_2\text{NiO}_4} - \frac{\delta_{214}}{2}\mu_{\text{O}_2} \\ &= \mu_{\text{La}_2\text{NiO}_4}^{\text{diff}} - \frac{\delta_{214}}{2}\mu_{\text{O}_2} \end{aligned} \quad (13)$$

which defines $\Delta G_{\text{La}_2\text{NiO}_4}^{\text{diff}}$, and $\mu_{\text{La}_2\text{NiO}_4}^{\text{diff}}$ is equal to $\mu_{\text{La}_2\text{NiO}_{4+\delta_{214}}} - \mu_{\text{La}_2\text{NiO}_4}$.

Similarly, for the higher order $\text{La}_3\text{Ni}_2\text{O}_7$ and $\text{La}_4\text{Ni}_3\text{O}_{10}$, there is:

$$\begin{aligned} \Delta G_{\text{La}_3\text{Ni}_2\text{O}_7}^{\text{diff}} &= \Delta G_{\text{La}_3\text{Ni}_2\text{O}_{7-\delta_{327}}} - \Delta G_{\text{La}_3\text{Ni}_2\text{O}_7} \\ &= \mu_{\text{La}_3\text{Ni}_2\text{O}_{7-\delta_{327}}} - \mu_{\text{La}_3\text{Ni}_2\text{O}_7} + \frac{\delta_{327}}{2}\mu_{\text{O}_2} \\ &= \mu_{\text{La}_3\text{Ni}_2\text{O}_7}^{\text{diff}} + \frac{\delta_{327}}{2}\mu_{\text{O}_2} \end{aligned} \quad (14)$$

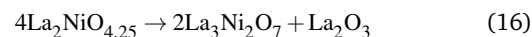
$$\begin{aligned} \Delta G_{\text{La}_4\text{Ni}_3\text{O}_{10}}^{\text{diff}} &= \Delta G_{\text{La}_4\text{Ni}_3\text{O}_{10-\delta_{4310}}} - \Delta G_{\text{La}_4\text{Ni}_3\text{O}_{10}} \\ &= \mu_{\text{La}_4\text{Ni}_3\text{O}_{10-\delta_{4310}}} - \mu_{\text{La}_4\text{Ni}_3\text{O}_{10}} + \frac{\delta_{4310}}{2}\mu_{\text{O}_2} \\ &= \mu_{\text{La}_4\text{Ni}_3\text{O}_{10}}^{\text{diff}} + \frac{\delta_{4310}}{2}\mu_{\text{O}_2} \end{aligned} \quad (15)$$

where $\mu_{\text{La}_3\text{Ni}_2\text{O}_7}^{\text{diff}}$ and $\mu_{\text{La}_4\text{Ni}_3\text{O}_{10}}^{\text{diff}}$ are defined in the same way as $\mu_{\text{La}_2\text{NiO}_4}^{\text{diff}}$.

The differences in the chemical potential of respective LNO RP compounds can be estimated using their DFT calculated ground state energies, assuming a small free energy change due to entropy. The chemical potential of oxygen gas can be calculated using data from the NIST chemical database.¹⁹ The defect structures selected for DFT calculation are $\text{La}_2\text{NiO}_{4.25}$, $\text{La}_3\text{Ni}_2\text{O}_{6.75}$ and

$\text{La}_4\text{Ni}_3\text{O}_{9.5}$. These non-stoichiometries are selected to represent extreme cases observed in experiments. The energy differences between these defect structures and their respective ideal structures are used to plot the Gibbs free energies of reaction when defect species are involved.

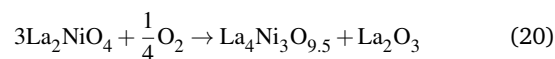
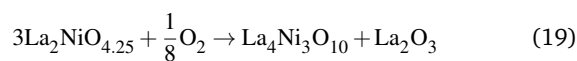
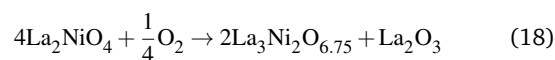
Taking the reaction between $\text{La}_2\text{NiO}_{4.25}$ and $\text{La}_3\text{Ni}_2\text{O}_7$ as an example, for the overall reaction, we have:



Its Gibbs free energy of reaction will be:

$$\begin{aligned} \Delta G_{\text{def.r.}}^{327}(T, p_{\text{O}_2}) &= 2\Delta G_{\text{La}_3\text{Ni}_2\text{O}_7}(T) + \Delta G_{\text{La}_2\text{O}_3}(T) - 4\Delta G_{\text{La}_2\text{NiO}_{4.25}}(T) \\ &= 2\Delta G_{\text{La}_3\text{Ni}_2\text{O}_7}(T) + \Delta G_{\text{La}_2\text{O}_3}(T) \\ &\quad - 4(\Delta G_{\text{La}_2\text{NiO}_4}(T) + \mu_{\text{La}_2\text{NiO}_4}^{\text{diff}} - \frac{1}{8}\mu_{\text{O}_2}) \end{aligned} \quad (17)$$

The other three defect reactions considered are:



Free energy changes for these reactions are then plotted together with their respective ideal reactions.

All density functional theory calculations in this work have been performed using the CRYSTAL09 software package^{20,21}, based on the expansion of the crystalline orbitals as a linear combination of a local basis set consisting of atom-centered Gaussian orbitals. The basis sets used for Ni and O atoms are triple-valence all-electron basis sets. A Ni atom is described by a 28-411d(41) contraction (one s, four sp, and two d shells). An O atom is described by an 8-411d(1) contraction (one s, three sp, and one d shells). The most diffuse sp(d) exponents are $\alpha_{\text{Ni}} = 0.6144(0.411)\text{bohrs}^{-2}$ and $\alpha_{\text{O}} = 0.1843(0.6)\text{bohrs}^{-2}$.²¹ The La core electrons are described with a pseudopotential, which is adapted from Dolg *et al.*'s quasirelativistic La pseudopotential²². The La valence electron basis functions are 411p(411)d(311)f(11) contracted sets (three s, three p, three d and two f shells) with the most diffuse exponent being $\alpha_{\text{La}} = 0.15\text{bohrs}^{-2}$ for each s, p, d and f shell.²³ A "ghost" atom is used to describe La vacancy by removing all but the two most diffuse contractions of each original La s, p, d and f shells. These basis sets ensure we have an accurate description of electron density at the La vacancy site.

The B3LYP hybrid-exchange electron density functional²⁴, implemented in the CRYSTAL09 package, is used to approximate electron exchange and correlation. This hybrid functional is expected to give a reliable description of solid oxide systems such as the LNO-RP family oxides.^{23,25,26} Numerical integration of the exchange and correlation potentials and energy functional are performed on an atom-centered grid of points. The XXLGRID option implemented in CRYSTAL09 was used for all calculations,

which consisted of a pruned grid consisting of 99 radial points and 5 subintervals with (146, 302, 590, 1454 and 590) angular points²¹. The cutoff threshold parameters of direct summation of the Coulomb and exchange series are set to 7, 7, 7, 7 and 14 as described in the CRYSTAL09 manual²¹. A Pack-Monkhorst mesh with a shrinking factor of 8 was chosen for reciprocal space sampling for La_2NiO_4 . For larger $\text{La}_3\text{Ni}_2\text{O}_7$ and $\text{La}_4\text{Ni}_3\text{O}_{10}$ systems, the shrinking factor used is 4. The structures selected as initial input for DFT simulation are taken from Skinner's⁴ (La_2NiO_4) and Ling's²⁷ ($\text{La}_3\text{Ni}_2\text{O}_7$ and $\text{La}_4\text{Ni}_3\text{O}_{10}$) neutron diffraction works. The supercells used in the calculations are the respective minimum conventional crystallographical cells of La_2NiO_4 (28 atoms), $\text{La}_3\text{Ni}_2\text{O}_7$ (24 atoms) and $\text{La}_4\text{Ni}_3\text{O}_{10}$ (34 atoms). We computed the La vacancy migration barrier height using a $(2 \times 1 \times 1)$ supercell of the conventional La_2NiO_4 lattice. This defective supercell contains 55 atoms. One of the La atoms is replaced by the La "ghost" atom mentioned above. Charge neutrality is maintained by promoting three Ni^{2+} atoms to Ni^{3+} . The La vacancy migration energy is evaluated using the Distinguished Reaction Coordinate (DRC) method implemented in the CRYSTAL09 package. A full description of the DRC method and its robustness has been given by Rimola *et al.* and will not be detailed here.²⁸

3 RESULTS AND DISCUSSION

3.1 Thermodynamic Calculations

To study the thermodynamic feasibility of La_2NiO_4 decomposition, the Gibbs free energy of the reactions given in Eqns. 4 to 6 were calculated for temperatures in the range of 305 K to 1673 K, and oxygen partial pressure ranging from 1×10^{-25} bar to 1 bar. Plots of the results are given in Figs. 1 to 3.

According to Fig. 1, the decomposition of La_2NiO_4 into higher order RP La-Ni-O materials is thermodynamically favourable up to about 1400 K at 0.2 bar oxygen partial pressure (atmospheric pressure of oxygen). $\text{La}_3\text{Ni}_2\text{O}_7$ is the energetically favourable phase below 1000 K, while $\text{La}_4\text{Ni}_3\text{O}_{10}$ will be more favourable above 1000K. LaNiO_3 is never the preferred phase over the entire temperature range considered. These results show that LNO would start to decompose when it is in contact with air even at room temperature. The decompositions will also take place at the designated IT-SOFC operation temperature (600-800°C). We suspect that the decomposition products La_2O_3 and $\text{La}_3\text{Ni}_2\text{O}_7$ will form a layered structure on top of the bulk La_2NiO_4 . Taking the (001) surface as an example, there are two possible stacking orders, $\text{La}_3\text{Ni}_2\text{O}_7$ - La_2O_3 - La_2NiO_4 and La_2O_3 - $\text{La}_3\text{Ni}_2\text{O}_7$ - La_2NiO_4 , with both being ordered from the surface to the bulk. The latter one is predicted to be the favourable stacking order. The first stacking order contains a large lattice mismatch between La_2O_3 and the two La-Ni-O RP phases, while the second one only contains one such large mismatch and one very small interfacial strain between the two La-Ni-O RP phases. The lattice mismatch between $\text{La}_3\text{Ni}_2\text{O}_7$ and La_2NiO_4 is only about 1% at the (001) interface^{4,27,29}. This predicted stacking order is consistent with Druce *et al.*'s secondary ion mass spectrometry (SIMS) - low energy ion scattering (LEIS) depth profiling of polycrystalline La_2NiO_4 . Druce's depth profile also showed a La enriched - Ni en-

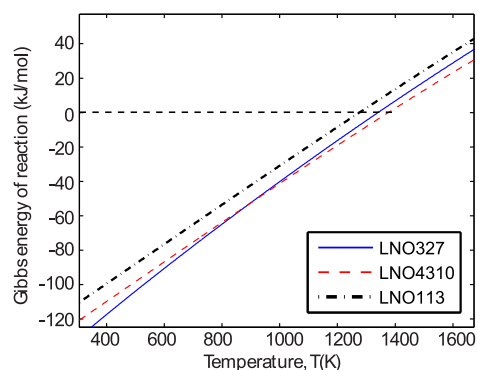


Fig. 1 Gibbs energy change for the reaction La_2NiO_4 decomposing into LaNiO_3 , $\text{La}_3\text{Ni}_2\text{O}_7$ or $\text{La}_4\text{Ni}_3\text{O}_{10}$ and La_2O_3 at 0.2 bar oxygen partial pressure over the temperature range of 305-1673 K. In the legend, the label LNO327 is for decomposition into $\text{La}_3\text{Ni}_2\text{O}_7$, LNO4310 is for $\text{La}_4\text{Ni}_3\text{O}_{10}$, and LNO113 is for LaNiO_3 .

riched - stoichiometric layered structure from the surface to the bulk of the crystal.¹⁴ The peak Ni:La ratio they observed in the Ni enriched region is about 2:3, which agrees to the Ni:La ratio in our predicted sub-surface $\text{La}_3\text{Ni}_2\text{O}_7$ layer.

These predictions agree partly with a long-term decomposition experiment on $\text{La}_2\text{NiO}_{4+\delta}$. In this experiment Gauquelin *et al.* heated $\text{La}_2\text{NiO}_{4+\delta}$ to 1273K in flowing air for two weeks³⁰. Characterisation of the heated surface with transmission electron microscopy (TEM) revealed layers of $\text{La}_3\text{Ni}_2\text{O}_7$ and $\text{La}_4\text{Ni}_3\text{O}_{10}$ with thicknesses of about 700 nm. On the (100) surface, both $\text{La}_3\text{Ni}_2\text{O}_7$ and $\text{La}_4\text{Ni}_3\text{O}_{10}$ were found, but only $\text{La}_4\text{Ni}_3\text{O}_{10}$ was present on the (001) surface. According to our calculations, $\text{La}_4\text{Ni}_3\text{O}_{10}$ should be the most stable phase at 1273K. The mixed layer of $\text{La}_4\text{Ni}_3\text{O}_{10}$ and $\text{La}_3\text{Ni}_2\text{O}_7$ may be due to incomplete conversion of $\text{La}_3\text{Ni}_2\text{O}_7$ formed at a lower temperature during the heating process and previous thermal history.

To form the transformed layers, La needs to move from the bulk to the surface. As there are different diffusion paths for La to the (001) and (100) surfaces we expect different La diffusion rates, and hence different growth rates for the transformed layers at these surfaces. One significant difference between our prediction and Gauquelin's experiment is that the other predicted product, La_2O_3 , was not observed in the experiment³⁰. A surface layer only containing La_2NiO_4 and its higher order RP counterparts cannot give the La rich region observed with XPS¹² and SIMS-LEIS depth profiling¹⁴ but rather the opposite (La deficient region³⁰). Burriel *et al.* reported that the La and Sr rich surface layer is about 7 nm thick in the related $\text{La}_{2-x}\text{Sr}_x\text{NiO}_{4+\delta}$ phase. In Druce's more recent work¹⁴, the La enriched layer is only about 1 nm thick in the depth profile of polycrystalline ceramic La_2NiO_4 . Such a thin layer may be damaged by the destructive Focused Ion Beam (FIB) technique, which was used to cut a suitable lamellar cross-section at the surface for TEM study in Gauquelin's experiment.

Despite the thermodynamically favoured decomposition process (about -100kJ/mol for La_2O_3 production at room temperature), La_2NiO_4 is stable and does not transform into its higher

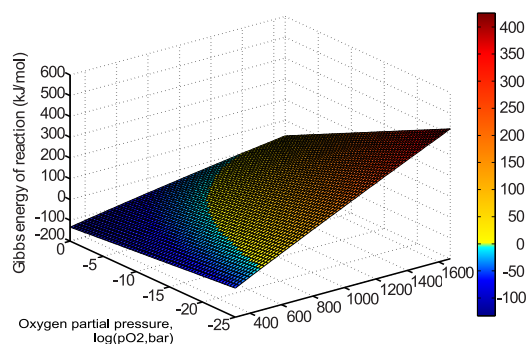


Fig. 2 Gibbs reaction energy for La_2NiO_4 into $\text{La}_3\text{Ni}_2\text{O}_7$ and La_2O_3 over the range of 1 bar to 1×10^{-25} bar oxygen partial pressure, and 305K to 1673K. The colour bar reflects the Gibbs energy of reaction in kJ/mol. The Gibbs energy of reaction is equal to zero at the edge between the blue and yellow regions.

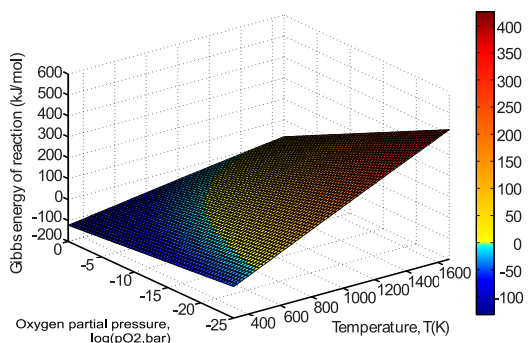


Fig. 3 Gibbs reaction energy for La_2NiO_4 into $\text{La}_4\text{Ni}_3\text{O}_{10}$ and La_2O_3 over the range of 1 bar to 1×10^{-25} bar oxygen partial pressure, and 305K to 1673K. The colour bar reflects the Gibbs energy of reaction in kJ/mol. The Gibbs energy of reaction is equal to zero at the edge between the blue and yellow regions.

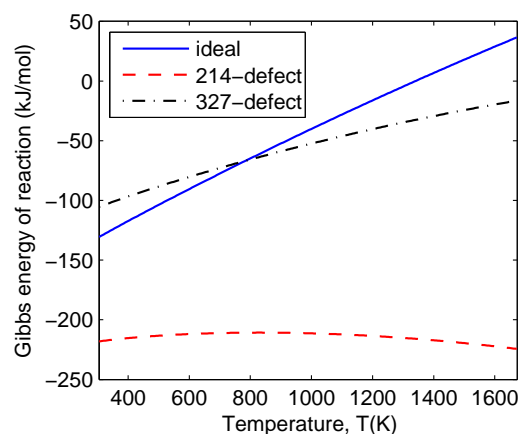


Fig. 4 Comparison of Gibbs reaction energy for La_2NiO_4 decomposition into $\text{La}_3\text{Ni}_2\text{O}_7$ (labelled as ideal), $\text{La}_2\text{NiO}_{4.25}$ decomposition into $\text{La}_3\text{Ni}_2\text{O}_7$ (214-defect) and La_2NiO_4 decomposition into $\text{La}_3\text{Ni}_2\text{O}_{6.75}$ (327-defect) at 0.2 bar oxygen partial pressure over the range of 305-1673 K.

order R-P phases completely at room temperature. In the long-term decomposition experiment, only about 700 nm of the surface La_2NiO_4 was converted into $\text{La}_3\text{Ni}_2\text{O}_7$ or $\text{La}_4\text{Ni}_3\text{O}_{10}$ after two weeks³⁰. Therefore, there must be a very slow rate determining step or a surface passivation mechanism to prevent complete decomposition of La_2NiO_4 . The three experimental works mentioned earlier treated their samples differently (1273 K for two weeks (Gauquelin), 723 K for 72 hours (Burriel) and 1273K for 12 hours (Druce)). A shorter heating time at a lower temperature gives a thinner transformed layer (only a few nanometers) than the high temperature, long-term heated sample (about 700 nm).^{12,14,30} Heating temperature and duration do have significant impact on the amount of La_2NiO_4 transformed. The reactions that result in the formation of the surface oxide layer must therefore include at least one activated process. A strong candidate is the diffusion of La, a process that is necessary for the oxide formation, and which will be activated. The diffusion barriers of several possible La diffusion paths are briefly evaluated in the next section. The large difference in the thickness of A-site enriched layers reported by Burriel and Druce, despite having a similar thermal treatments, is believed to be mainly due to Sr segregation which only occurs in Burriel's Sr-doped $\text{La}_{2-x}\text{Sr}_x\text{NiO}_{4+\delta}$. This would lead to a much thicker (7 nm compared to Druce's 1 nm) A-site (La+Sr) enriched region despite similar thermal history. In addition, Burriel's sample is a single crystal while Druce tested a polycrystalline ceramic sample. The different natures of the samples should also affect the amount of decomposed La_2NiO_4 .

Figures 2 and 3 clearly show that oxygen partial pressure has a significant impact on the Gibbs free energy of these reactions. According to the graphs, the critical oxygen partial pressure required to allow decomposition varies with temperature. Therefore, the thickness of maximum decomposition layers formed at different temperatures after long heating times will vary with critical oxygen partial pressure. In a controlled low p_{O_2} or high vacuum environment, this maximum transformation thickness can be mea-

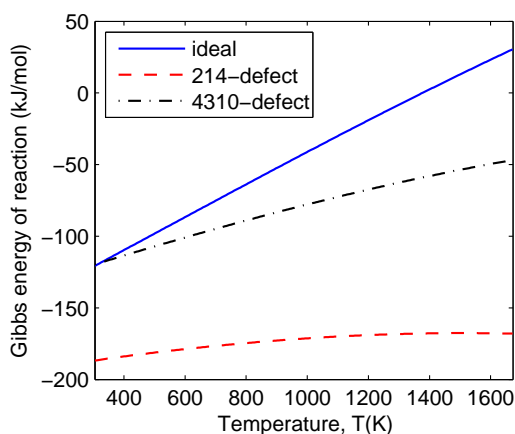


Fig. 5 Comparison of Gibbs reaction energy for La_2NiO_4 decomposition into $\text{La}_4\text{Ni}_3\text{O}_{10}$ (labelled as ideal), $\text{La}_2\text{NiO}_{4.25}$ decomposition into $\text{La}_4\text{Ni}_3\text{O}_{10}$ (214-defect) and La_2NiO_4 decomposition into $\text{La}_4\text{Ni}_3\text{O}_{9.5}$ (4310-defect) at 0.2 bar oxygen partial pressure over the range of 305-1673 K.

sured to verify the hypothesis. If this proves to be true, controlling the oxygen partial pressure and the temperature might provide a way to suppress or reverse the decomposition when desired.

3.2 Density Functional Based Calculations

To obtain a more complete picture of the decomposition thermodynamics, we use density functional theory simulations to check the effect of oxygen defects on the reactions as described in the method section. Ideal experimental structures of the La-Ni-O compounds are fully relaxed and optimised first with antiferromagnetic ordering. The results are summarised in Table 3. The differences between the experimental structures and the calculated structures are small (less than 3%) in all three cases. Therefore, we believe the DFT simulations describe these Ruddlesden-Popper compounds well. Respective defective compounds are constructed by adding oxygen interstitials or vacancies to the relaxed ideal structures. The position of interstitials or vacancies are chosen according to literature results.^{4,27} The defective compounds are then also relaxed and optimised for the evaluation of the Gibbs reaction energies' change due to the defects. The resulting Gibbs energies of reaction for the different defect cases are plotted in Figures 4 and 5.

In Figures 4 and 5 we can see that all the defective reactions still have negative reaction Gibbs energy in the IT-SOFC working temperature range (873-1073K). The defective reactions seem to be always favourable even at temperatures beyond 1673K according to the trend shown in Figs. 4 and 5. However, the assumption of small energy change due to entropy will not hold at temperatures much higher than 1673K, so the estimated trends based on ground-state DFT calculations should only be used to predict reactions at lower temperature ranges. For the temperature range we are interested in (873-1073K), these four reactions can give an idea of the order of magnitude of the decomposition Gibbs energies under experimental condition since the involved phases are some of the extreme defective cases. It is noted that the 214-

Table 3 Comparison of calculated and experimental lattice parameters of relevant La-Ni-O species. Expt. is experimental structure, and Calc. is calculated structure. All literature experimental structures are measured at 298K.

Compound		a (Å)	b (Å)	c (Å)	Ref.
La_2NiO_4	Expt.	5.459	5.464	12.685	4
	Calc.	5.538(8)	5.539(3)	12.752(2)	
	Error	1.47%	1.39%	0.52%	
$\text{La}_3\text{Ni}_2\text{O}_7$	Expt.	5.397	5.450	20.507	27
	Calc.	5.435(0)	5.586(1)	20.868(3)	
	Error	0.70%	2.49%	1.76%	
$\text{La}_4\text{Ni}_3\text{O}_{10}$	Expt.	5.413	5.462	27.960	27
	Calc.	5.448(3)	5.570(8)	28.512(2)	
	Error	0.65%	1.99%	1.97%	

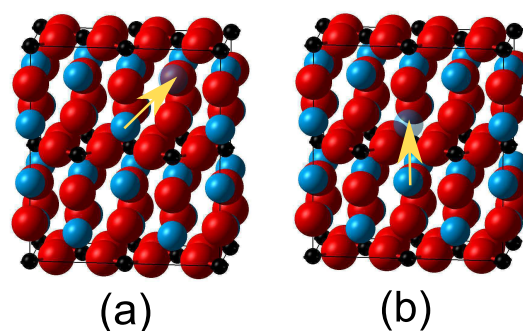


Fig. 6 Schematics of two La diffusion pathways via vacancy mechanism in a LNO_{214} $2 \times 1 \times 1$ supercell with one La vacancy. In the figure, blue is La atom, black is Ni atom and red is O atom. The semi-transparent atom represents La vacancy. (a) shows the La migration pathway in the rocksalt layer. (b) shows the La migration pathway in the perovskite layer.

defect reactions are particularly favourable as shown in the figures (about 100 kJmol^{-1} more favourable than ideal reactions). This suggests that a significant amount of energy is required to accommodate the oxygen interstitials in $\text{La}_2\text{NiO}_{4.25}$ at the ground state (DFT simulation condition). At ground state (OK), it is more difficult for the ideal lattice to take in an interstitial oxygen (like defective LNO214) than to lose an oxygen and form a vacancy (like defective LNO327 and LNO4310) due to contraction of lattice at low temperature. This is one of the sources of errors when using DFT simulation to estimate the Gibbs energies of reaction for the defective reactions. At IT-SOFC working temperatures, the energy differences between the defective LNO214 reaction and the defective LNO327 (or LNO4310) reactions should be smaller. The actual Gibbs free energy of reaction will be affected by specific experimental conditions, including temperature, oxygen partial pressure, thermal histories and the initial composition of the materials, but they will fall somewhere close to the calculated region where the Gibbs free energy of decomposition reaction is less than zero. Therefore, the simple conclusion drawn from these estimations is that the decomposition will be favoured at IT-SOFC operation temperatures even for the extremely defective compositions.

As noted earlier, for the LNO214 structure to decompose, La atoms need to diffuse to the surface. As an activated process, this diffusion has an associated activation energy. We have computed the activation energy of two possible diffusion steps (shown in Figure 6). We obtain a barrier height of 4.31 eV for the rocksalt path, and a height of 4.57 eV for the perovskite path. Although no experimental cation diffusion data for La_2NiO_4 is available, our calculated value for the perovskite layer (4.57 eV) is comparable to the La diffusion activation energy of 4.98 eV in LaCrO_3 .³¹ Comparing to the oxygen diffusion activation energy (0.88 eV)⁵ in single crystal La_2NiO_4 , the La diffusion activation energy is about five times as large. Assuming a simple Arrhenius relationship between the diffusion rate and diffusion activation energy, we can estimate that La diffusion rate is at least 3 orders of magnitudes smaller than O diffusion rate. That might be part of the reason behind the limited decomposition. A more detailed study on the La diffusion during the decomposition reaction is being carried out, and the results will be presented in a future publication.

4 CONCLUSION

The decomposition of La_2NiO_4 into LaNiO_3 , $\text{La}_3\text{Ni}_2\text{O}_7$ or $\text{La}_4\text{Ni}_3\text{O}_{10}$ and La_2O_3 has been found to be thermodynamically favourable over the temperature range of 300 K to about 1373 K under 0.2 bar oxygen partial pressure, based on calculations using literature thermodynamic equation constants. The Gibbs free energy of reactions for La_2NiO_4 decomposition into LaNiO_3 , $\text{La}_3\text{Ni}_2\text{O}_7$ or $\text{La}_4\text{Ni}_3\text{O}_{10}$ and La_2O_3 are plotted over a range of temperature and oxygen partial pressure. It is then suggested that a partial decomposition of La_2NiO_4 takes place at the surface when exposed to air. Further estimation of the reaction energetics for defective reactions based on DFT also supports the hypothesis. The decomposition leads to a stacking of (La_2O_3) - $(\text{La}_3\text{Ni}_2\text{O}_7/\text{La}_4\text{Ni}_3\text{O}_{10})$ - $(\text{La}_2\text{NiO}_4)$ (labelled from surface to bulk) depending on temperature. Such stacking could be respon-

sible for the observed Ni-absence and Ni:La non-stoichiometry at La_2NiO_4 surfaces^{12,14,30}. This hypothesis agrees particularly well with Druce's depth profile of near-surface composition of La_2NiO_4 measured with SIMS-LEIS¹⁴. While the kinetics of the cation diffusion necessary for LNO214 decomposition are not known yet, the microstructures of the actual ceramic cathode in operation will play an important role. As pointed out by Kubicek *et al.*, the grain boundaries and dislocations will act as short paths for fast cation diffusions, like Sr diffusion in $\text{La}_{0.6}\text{Sr}_{0.4}\text{CoO}_{3-\delta}$.³² The long term stability of La_2NiO_4 over 773-1023 K and its long-term performance as cathode for IT-SOFC is therefore uncertain. In addition, instead of an ideal La_2NiO_4 surface, the proposed oxidized La surface will quickly become the relevant surface for oxygen reduction and exchange under the IT-SOFC working environment. The oxygen reduction mechanism on such a surface requires investigation.

Some questions remain unclear and cannot be answered with the hypothesis. Although thermodynamically favourable, the decomposition reactions studied were very slow as shown in experiments. The diffusion barrier of La is briefly evaluated and shows a quite high diffusion activation energy. The details of the decomposition's reaction kinetics are being studied at the moment. These reaction kinetics will be critical to the evaluation of the long term stability of La_2NiO_4 . Only a few particular decompositions reactions were considered in this study. Therefore, there might be other combinations of stable species. To find out the most stable combinations of phases, a density functional theory (DFT) based full phase diagram of the La-Ni-O system is planned.

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References

- 1 V. V. Kharton, A. P. Viskup, E. N. Naumovich and F. M. B. Marques, *J. Mater. Chem.*, 1999, **9**, 2623–2629.
- 2 S. J. Skinner and J. A. Kilner, *Solid State Ionics*, 2000, **135**, 709–712.
- 3 V. Kharton, A. P. Viskup, A. Kovalevsky, E. Naumovich and F. Marques, *Solid State Ionics*, 2001, **143**, 337–353.
- 4 S. J. Skinner, *Solid State Sci.*, 2003, **5**, 419–426.
- 5 J. Bassat, P. Odier, A. Villesuzanne, C. Marin and M. Pouchard, *Solid State Ionics*, 2004, **167**, 341–347.
- 6 A. Chroneos, D. Parfitt, J. A. Kilner and R. W. Grimes, *J. Mater. Chem.*, 2010, **20**, 266.
- 7 J. A. Kilner and M. Burriel, *Annu. Rev. Mater. Res.*, 2014, **44**, 365–393.
- 8 E. Bucher and W. Sitte, *Solid State Ionics*, 2011, **192**, 480–482.

- 9 H. Ding, A. V. Virkar, M. Liu and F. Liu, *Phys. Chem. Chem. Phys.*, 2013, **15**, 489–496.
- 10 M. S. D. Read, M. S. Islam, G. W. Watson and F. E. Hancock, *J. Mater. Chem.*, 2001, **11**, 2597–2602.
- 11 J. Zhou, G. Chen, K. Wu and Y. Cheng, *J. Phys. Chem. C*, 2013, **117**, 12991–12999.
- 12 M. Burriel, S. Wilkins, J. P. Hill, M. A. Muñoz Márquez, H. H. Brongersma, J. A. Kilner, M. P. Ryan and S. J. Skinner, *Energy Environ. Sci.*, 2014, **7**, 311.
- 13 J. Druce, T. Ishihara and J. A. Kilner, *Solid State Ionics*, 2014, **262**, 893–896.
- 14 J. Druce, H. Téllez, M. Burriel, M. D. Sharp, L. J. Fawcett, S. N. Cook, D. S. McPhail, T. Ishihara, H. H. Brongersma and J. A. Kilner, *Energy Environ. Sci.*, 2014, **7**, 3593–3599.
- 15 M. Zinkevich and F. Aldinger, *J. Alloys Compd.*, 2004, **375**, 147–161.
- 16 M. Zinkevich, N. Solak, H. Nitsche, M. Ahrens and F. Aldinger, *J. Alloys Compd.*, 2007, **438**, 92–99.
- 17 A. N. Grundy, B. Hallstedt and L. J. Gauckler, *J. Phase Equilib.*, 2001, **22**, 105–113.
- 18 K. Johnston, M. Castell, A. Paxton and M. Finnis, *Phys. Rev. B*, 2004, **70**, 085415.
- 19 P. Linstrom and W. Mallard, *NIST Chemistry WebBook, NIST Standard Reference Database Number 69*, 2000.
- 20 R. Dovesi, R. Orlando, B. Civalleri, C. Roetti, V. R. Saunders and C. M. Zicovich-Wilson, *Z. Kristallogr.*, 2005, **220**, 571–573.
- 21 R. Dovesi, V. R. Saunders, C. Roetti, R. Orlando, C. M. Zicovich-Wilson, F. Pascale, B. Civalleri, K. Doll, N. M. Harrison, I. J. Bush, P. D'Arco and M. Llunell, *CRYSTAL09 Users' Manual*, Università di Torino, Torino, 2010.
- 22 M. Dolg, H. Stoll, a. Savin and H. Preuss, *Theor. Chim. Acta*, 1989, **75**, 173–194.
- 23 D. Muñoz, N. Harrison and F. Illas, *Phys. Rev. B*, 2004, **69**, 085115.
- 24 A. D. Becke, *J. Chem. Phys.*, 1993, **98**, 1372.
- 25 J. Muscat, A. Wander and N. Harrison, *Chem. Phys. Lett.*, 2001, **342**, 397–401.
- 26 G. Mallia and N. Harrison, *Phys. Rev. B*, 2007, **75**, 165201.
- 27 C. D. Ling, D. N. Argyriou, G. Wu and J. Neumeier, *J. Solid State Chem.*, 2000, **152**, 517–525.
- 28 A. Rimola, C. M. Zicovich-Wilson, R. Dovesi and P. Ugliengo, *J. Chem. Theory Comput.*, 2010, **6**, 1341–1350.
- 29 V. Voronin, I. Berger, V. Cherepanov, L. Gavrilova, A. Petrov, A. Ancharov, B. Tolochko and S. Nikitenko, *Nucl. Instr. Meth. Phys. Res. A*, 2001, **470**, 202–209.
- 30 N. Gauquelin, T. E. Weirich, M. Ceretti, W. Paulus and M. Schroeder, *Monatsh. Chem.*, 2009, **140**, 1095–1102.
- 31 T. Akashi, M. Nanko, T. Maruyama, Y. Shiraishi and J. Tanabe, *J. Electrochem. Soc.*, 1998, **145**, 2090.
- 32 M. Kubicek, G. M. Rupp, S. Huber, A. Penn, A. K. Opitz, J. Bernardi, M. Stöger-Pollach, H. Hutter and J. Fleig, *Phys. Chem. Chem. Phys.*, 2014, **16**, 2715–26.