

**High thermoelectric figure-of-merit in Sb<sub>2</sub>Te<sub>3</sub>/Ag<sub>2</sub>Te bulk composites as Pb-free p-type thermoelectric materials**

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## ARTICLE

# High thermoelectric figure-of-merit in $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$ bulk composites as Pb-free *p*-type thermoelectric materials

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We investigated the thermoelectric properties of the  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  (ST/AT) composites with molar ratios of ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1. The extrinsic composites, synthesized by wet ball milling of two separate powders of  $\text{Sb}_2\text{Te}_3$  and  $\text{Ag}_2\text{Te}$ , are differentiated with intrinsic composites by high temperature phase separation because of the low temperature synthesis process. The thermoelectric properties of the composites show systematic behaviour of decreased electrical and thermal conductivities with increasing  $\text{Ag}_2\text{Te}$  dispersion concentration. The ST/AT = 1/1 composite exhibits extremely low lattice thermal conductivity with high power factor over a wide temperature range, resulting in a high *ZT* value of 1.5 at 700 K, which is a significantly enhanced value of *ZT* compared with those of other Pb-free *p*-type chalcogenide thermoelectric materials.

## 1. Introduction

Thermoelectric (TE) materials have been attracting attention because of possible applications in fields such as solid state cooling and waste heat power generation. The efficiency of a TE generator requires a high dimensionless figure-of-merit  $ZT = S^2\sigma T/\kappa$ , where *S*,  $\sigma$ , *T*, and  $\kappa$  are the Seebeck coefficient, electrical conductivity, absolute temperature, and thermal conductivity, respectively. Recent investigations have reported progressively increased thermoelectric performance *ZT* in PbTe based materials; for example, Fermi level tuning at the resonant level in Na-doped PbTe (1.5 at 773 K),<sup>[1]</sup> convergence of band structures in PbTe<sub>1-x</sub>Se<sub>x</sub> (1.8 at 850 K),<sup>[2]</sup> and hierarchical microstructure control in an SrTe- and Na-doped PbTe system (2.2 at 915 K).<sup>[3]</sup> Exceptionally high *ZT* materials ( $ZT \geq 2.0$ ) were followed by the nano- and micro-structure modulation.<sup>[4,5]</sup>

Even though PbTe based materials are promising candidate for waste heat power generation, such as space application, the environmentally hazardous nature of Pb can cause limitations in terms of widespread application of PbTe based materials. As a Pb-free thermoelectric material, (GeTe)<sub>x</sub>(AgSbTe<sub>2</sub>)<sub>100-x</sub>, which is called TAGS-m, has been studied since the early 1960s<sup>[6-8]</sup> and (GeTe)<sub>1-x</sub>[(Ag<sub>2</sub>Te)<sub>1-y</sub>(Sb<sub>2</sub>Te<sub>3</sub>)<sub>y</sub>]<sub>x</sub>.<sup>[9]</sup> Recently, Dy-doped TAGS-85 has been found to have a high *ZT* value of 1.5 at 750 K.<sup>[10]</sup> The high thermoelectric performance of TAGS may be related with the crystallographic transformation from a high-temperature cubic to a low-temperature rhombohedral polymorph.<sup>[11]</sup>

The matrix of the TAGS, AgSbTe<sub>2</sub> compound is promising for use in *p*-type thermoelectric materials in the mid-temperature range.<sup>[12]</sup>

The ternary Ag-Sb-Te system undergoes a eutectic reaction as follows: Liquid → AgSbTe<sub>2</sub> + Ag<sub>2</sub>Te + δ-Sb<sub>2</sub>Te. Because of the eutectic decomposition, the quaternary melt of the Pb-Ag-Sb-Te system, which is known as the LAST system, exhibits a dendrite or phase segregated morphology.<sup>[13]</sup> It is well known that phase separated morphologies are beneficial in reduction of lattice thermal conductivity by interfacial phonon scattering.<sup>[14,15]</sup> By using the phase transformation  $\text{Ag}_2\text{Te} + \text{AgSbTe}_2 \leftrightarrow \text{Ag}_2\text{Te} + \text{Sb}_2\text{Te}_3$ , the composite with phase separation showed high thermoelectric performance around  $ZT \ 1.0 \sim 1.3$ .<sup>[16,17]</sup> The intrinsic thermodynamic phase separation of  $\text{Ag}_2\text{Te} + \text{Sb}_2\text{Te}_3$  at the temperature of 360 °C can cause thermal instability of thermoelectric properties in the AgSbTe<sub>2</sub> compound.

There are two different strategies for fabricating nano- or micro-structured composites. One is metallurgical phase separation [18] or precipitation;<sup>[19]</sup> the other is the synthesis of an artificial composite.<sup>[20,21]</sup> Here, we explore the thermoelectric properties of the composites of  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  (ST/AT) in systematic variations of ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1 molar ratios. The extrinsic composite fabricated by mechanical ball milling of separate mixtures of  $\text{Sb}_2\text{Te}_3$  and  $\text{Ag}_2\text{Te}$  powders can show improved thermal stability because of the low temperature synthesis of the composite. We achieved high thermoelectric figure-of-merit *ZT* of 1.5 at 700 K in the composite due to the decreased thermal conductivity caused by grain boundary phonon scattering. In spite of the *p*-type  $\text{Sb}_2\text{Te}_3$  and *n*-type  $\text{Ag}_2\text{Te}$  composites, the Seebeck coefficients were found not to be degraded by the *p*-/*n*-composite, preserving the high electrical conductivity. The high *ZT* value with Pb-free elements should have a

significant impact on the field of environmentally friendly waste heat power generation.

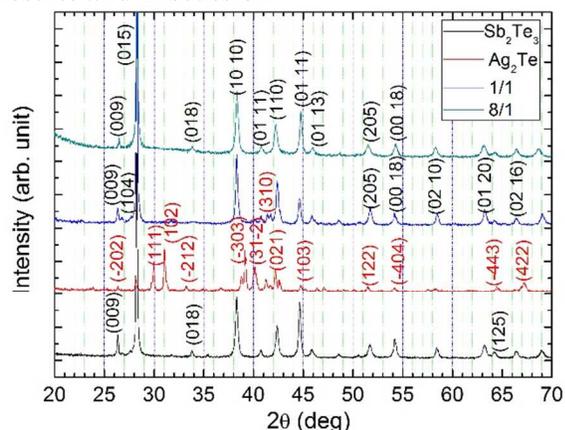
## 2. Experimental

We synthesized the  $\text{Sb}_2\text{Te}_3$  and  $\text{Ag}_2\text{Te}$  ingots by melting using high purity Ag (99.999 %), Sb (99.999 %), and Te (99.999 %) granules. For the synthesis of  $\text{Sb}_2\text{Te}_3$ , the evacuated quartz ampoules were melted at 800 °C for 24 h and slowly cooled to room temperature with a cooling rate of 10 °C/h.

The  $\text{Ag}_2\text{Te}$  ingots were melted at 1000 °C for 24 h in the evacuated quartz ampoules, slowly decreased the temperature to 500 °C, and water quenching.  $\text{Sb}_2\text{Te}_3$  and  $\text{Ag}_2\text{Te}$  ingots were pulverized by hand grinding for 6 h in an agate mortar. The pulverized powders were mixed and stirred in an *n*-hexane solution. The dried powder mixtures of  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  (ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1) were sintered by hot press sintering at 350 °C with application of uniaxial pressure of 70 MPa for 1 hour.

The phase identification and structural characterization of the composites were performed by powder X-ray diffraction (XRD) with Cu  $K\alpha$  radiation. The microstructure and relative concentration of elements were determined by scanning electron microscopy (SEM) and energy dispersive X-ray spectroscopy (EDX) measurements (JEOL, Japan). The Seebeck coefficient and electrical resistivity were determined via a four point probe method using a thermoelectric measurement system (ZEM-3 ULVAC, Japan). The thermal conductivity was obtained by  $\kappa = \rho_s C_p \lambda$ , where  $\rho_s$ ,  $C_p$ , and  $\lambda$  are the sample density, specific heat measured by the physical property measurement system (PPMS, Quantum Design, USA), and thermal diffusivity measured by a laser flash method, respectively.

## 3. Results and Discussion



**Fig. 1** Powder X-ray diffraction patterns of the  $\text{Sb}_2\text{Te}_3$ ,  $\text{Ag}_2\text{Te}$ , and  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  (ST/AT) composites with molar ratios of (ST/AT = 1/1 and 8/1). The red miller indices represent the  $\text{Ag}_2\text{Te}$ .

Figure 1 presents the powder X-ray diffraction patterns of the hot-press sintered ST/AT composite samples of ST/AT = 1/1 and 8/1 together with the major diffraction peaks of  $\text{Sb}_2\text{Te}_3$  and  $\text{Ag}_2\text{Te}$  single phase compounds. The diffraction peaks of the  $\text{Sb}_2\text{Te}_3$  (black, 1<sup>st</sup> from the bottom) and  $\text{Ag}_2\text{Te}$  (red, 2<sup>nd</sup> from the bottom) show single phases of rhombohedral ( $R\bar{3}m$ )  $\text{Sb}_2\text{Te}_3$  and monoclinic ( $P2_1/n$ )  $\alpha$ -

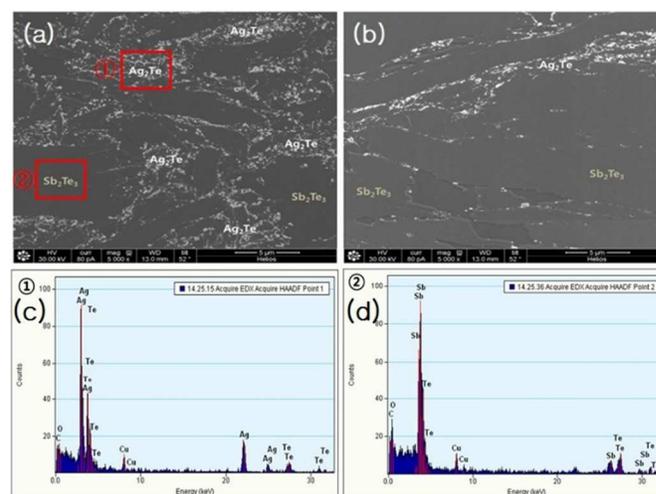
$\text{Ag}_2\text{Te}$  crystal structures, respectively. The mixed phases of  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  = 1/1 (blue, 3<sup>rd</sup> from the bottom) and ST/AT = 8/1 (sky blue, the top line) show that the major peak corresponds to  $\text{Sb}_2\text{Te}_3$  with small  $\text{Ag}_2\text{Te}$  impurity peaks. The lattice parameters of  $\text{Sb}_2\text{Te}_3$ ,  $\text{Ag}_2\text{Te}$  and their composites with molar ratios of (ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1) are presented in Table 1.

		a (Å)	b (Å)	c (Å)
Single Phases	$\text{Sb}_2\text{Te}_3$	4.279	-	30.481
	$\text{Ag}_2\text{Te}$	8.165	8.935	8.061
$\text{Sb}_2\text{Te}_3$ (ST/AT)	1/1	4.278	-	30.467
	2/1	4.278	-	30.469
	4/1	4.278	-	30.476
	8/1	4.279	-	30.475
	16/1	4.279	-	30.479
	32/1	4.279	-	30.481

**Table 1** Lattice parameters of  $\text{Sb}_2\text{Te}_3$ ,  $\text{Ag}_2\text{Te}$ , and  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  composites with molar ratios of ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1.

The lattice parameters of  $\text{Sb}_2\text{Te}_3$  and  $\alpha$ - $\text{Ag}_2\text{Te}$  correspond to those in previous reports.[22,23] The *a*-axis lattice parameters of  $\text{Sb}_2\text{Te}_3$  for the composite are not changed, while the *c*-axis lattice parameters are slightly increased with decreasing of the  $\text{Ag}_2\text{Te}$  mole fraction. The lattice parameters of the small mole fraction of  $\text{Ag}_2\text{Te}$  (32/1) are almost identical to those of the pristine  $\text{Sb}_2\text{Te}_3$  compound. The lattice parameters of  $\text{Ag}_2\text{Te}$  for the composites cannot be obtained due to the small impurity peaks of the compound.

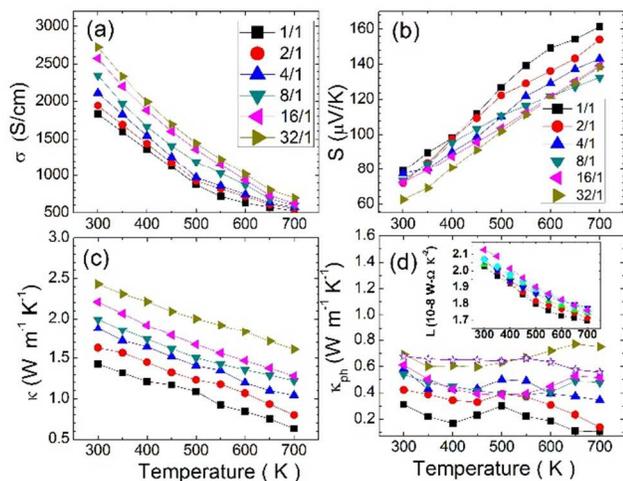
The small decrease of the *c*-axis lattice parameter, without significant change of the *a*-axis lattice parameter with increasing  $\text{Ag}_2\text{Te}$  volume fraction, implies a small solid solution reaction with  $\text{Ag}_2\text{Te}$  and  $\text{Sb}_2\text{Te}_3$  because the *c*-axis lattice parameter of  $\text{Ag}_2\text{Te}$  is smaller than that of  $\text{Sb}_2\text{Te}_3$ . If a small amount of Ag can be incorporated in the van der Waals gap of  $\text{Sb}_2\text{Te}_3$ , the *c*-axis lattice parameters should be increased. Also, if Ag atoms substitute at the Sb, Te, or interstitial sites, the *a*- and *c*-axis lattice parameters are also increased due to the larger atomic size of Ag (160 pm) than those of Sb (145 pm) and Te (140 pm).



**Fig. 2** Scanning electron microscopy images of  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  composite of ST/AT = 1/1 for perpendicular (a) and parallel direction (b) of the uniaxial pressure direction. (c) and (d) represent the energy dispersive X-ray spectroscopy data in the white  $\text{Ag}_2\text{Te}$  and grey  $\text{Sb}_2\text{Te}_3$  regions, respectively as indicated in the red boxes of (a).

Figure 2(a) and 2(b) presents scanning electron microscope (SEM) images of the ST/AT = 1/1 composite cross-sectional and parallel directions respectively, for the uniaxial hot press direction. The mole fraction of ST/AT = 1/1 corresponds to an almost 2/1 volume fraction. Energy dispersive x-ray spectroscopy (EDX) shows that the dark and white grey regions of Fig. 2(a) and 2(b) are  $\text{Sb}_2\text{Te}_3$  and  $\text{Ag}_2\text{Te}$  phases, as shown in Fig. 2(c) and 2(d), respectively. The  $\text{Ag}_2\text{Te}$  phase is well dispersed in the  $\text{Sb}_2\text{Te}_3$  matrix in the cross-sectional plane across the uniaxial pressure direction of the hot press. Along the uniaxial direction, the  $\text{Ag}_2\text{Te}$  phase is not homogeneously distributed but shows striped distribution, indicating that the  $\text{Ag}_2\text{Te}$  more homogeneously dispersed across the perpendicular plane of the uniaxial pressure direction of the hot press.

When we examine many different spots of the  $\text{Sb}_2\text{Te}_3$  region using EDX, there are no Ag peaks, as shown in Fig. 2(d). Also, the white regions in Fig. 2(a) shows only Ag and Te peaks without Sb concentration. This indicates that the  $\text{Ag}_2\text{Te}$  phase separated well from the  $\text{Sb}_2\text{Te}_3$  matrix. Near the grain boundaries between the  $\text{Sb}_2\text{Te}_3$  and  $\text{Ag}_2\text{Te}$  phases, a solid solution reaction can be possible; this is the result of the small decrease of the *c*-axis lattice parameters in the ST/AT composites in the powder X-ray diffraction. The grain sizes of  $\text{Ag}_2\text{Te}$  range from 200 nm to 5  $\mu\text{m}$  with anisotropic dispersion. The anisotropic distribution of  $\text{Ag}_2\text{Te}$  particles implies anisotropic thermoelectric properties. Here we measured the thermoelectric properties of the composites across the uniaxial pressure direction.



**Fig. 3** Temperature-dependent thermoelectric properties of  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  (ST/AT) composites with molar ratios of (ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1); Electrical conductivity (a), Seebeck coefficient (b), total thermal conductivity  $\kappa_{\text{tot}}$  (c), lattice thermal conductivity  $\kappa_{\text{ph}}$  (d), and Lorenz number  $L(T)$  (inset of (d)). Star open symbol in (d) represents the lattice thermal conductivity of  $\text{Sb}_2\text{Te}_3$ .

Figure 3 shows the temperature-dependent electrical conductivity  $\sigma(T)$  (a), Seebeck coefficient  $S(T)$  (b), thermal conductivity  $\kappa(T)$  (c), and lattice thermal conductivity  $\kappa_{\text{ph}}(T)$  (d), and Lorenz number  $L(T)$  (inset of (d)) of ST/AT composite with molar ratios of ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1, measured along the vertical direction of the uniaxial hot press direction. The electrical conductivity  $\sigma(T)$  shows typical metallic behaviour, as shown in Fig. 3(a), which is similar to but lower than those of the single crystalline  $\text{Sb}_2\text{Te}_3$  (5105 S/cm).[24] The decreased electrical conductivity is due to the  $\text{Ag}_2\text{Te}$  dispersion in the polycrystalline  $\text{Sb}_2\text{Te}_3$  matrix. The electrical transport properties of the nano-sheets of sintered  $\text{Sb}_2\text{Te}_3$  bulk samples show typical semiconductor behaviour with an electrical conductivity value of 230 S/cm at 300 K.[25,26] The electrical conductivity values of nano-sheet sintered compounds are remarkably lower by one order of magnitude than those of bulk  $\text{Sb}_2\text{Te}_3$  due to significant electron scattering at the nano-grain boundaries.

With increasing of the mole fraction of  $\text{Ag}_2\text{Te}$ , the electrical conductivity is decreased to 1832 S/cm (at 300 K) for the ST/AT = 1/1 composite. The electrical conductivity of the  $\text{Sb}_2\text{Te}_3$  rich composite (ST/AT = 32/1) has a value of 2700 S/cm at 300 K, but the  $\sigma$  values at high temperature (700 K) are not sensitive with  $\text{Ag}_2\text{Te}$  dispersion concentration. The electrical conductivity values of  $\alpha$ - and  $\beta$ - $\text{Ag}_2\text{Te}$  were 400 S/cm and 100 S/cm, respectively, at room temperature which are both lower than that of  $\text{Sb}_2\text{Te}_3$ . [27] Therefore, it is not surprising to observe a decrease of electrical conductivity with increasing  $\text{Ag}_2\text{Te}$  dispersion concentration due to the low electrical conductivity of  $\text{Ag}_2\text{Te}$  and the grain boundary scattering of the carriers.

The Seebeck coefficients of the ST/AT composites are presented in Fig. 3(b). The  $S(T)$  increases with increasing temperature, which is a typical metallic behaviour. Because of the decreased electrical conductivity, the Seebeck coefficient increases with increasing  $\text{Ag}_2\text{Te}$  dispersion concentration. The Seebeck coefficient dependence on the  $\text{Ag}_2\text{Te}$  dispersion concentration at high temperature is more significant than that of the room temperature Seebeck coefficient variation with varying of the  $\text{Ag}_2\text{Te}$  concentration. The Seebeck coefficient of the single crystalline  $\text{Sb}_2\text{Te}_3$  compound had value of 83  $\mu\text{V}/\text{K}$  along the cleavage planes and 92  $\mu\text{V}/\text{K}$  along the perpendicular direction at 300 K, which is similar to the value of ST/AT = 1/1.[28] In the case of the  $\text{Sb}_2\text{Te}_3$  nano-particle systems, the value is 210  $\mu\text{V}/\text{K}$  due to low electrical conductivity.[25,26] The Seebeck coefficients of  $\alpha$ - and  $\beta$ - $\text{Ag}_2\text{Te}$  were found to be in a range from 120~170  $\mu\text{V}/\text{K}$  due to the *p*-type conduction of the carriers.[27] Because the carriers are of the same type between  $\text{Sb}_2\text{Te}_3$  and  $\text{Ag}_2\text{Te}$ , the Seebeck coefficients are not significantly degraded with increasing  $\text{Ag}_2\text{Te}$  dispersion concentrations.

From the *ab-initio* first principle band structure calculation of  $\text{Sb}_2\text{Te}_3$ , it can be seen that this is a degenerate semiconductor in which the Fermi level ( $E_F$ ) is positioned near the top of the valence band maximum with a small band gap  $\Delta \leq 0.1$  eV at the  $\Gamma$ -point.[29,30] The original electronic band structure of  $\alpha$ - $\text{Ag}_2\text{Te}$  is that of a zero-gap semiconductor; however there is a strong Te 4*p*-Ag 4*d* hybridization shift up to the Te 4*p*-orbital and a push down to the Ag 4*d* band, resulting in the formation of a semiconducting energy band gap; the Fermi level resides at the top of the valence band maximum.[31] The electronic band structure of  $\beta$ - $\text{Ag}_2\text{Te}$  shows that the Fermi level is pinned near the conduction band minimum.[31] Even if it is well known that the band gap is overestimated in theoretical calculation, the

Fermi level shift to the conduction band minimum from  $\alpha$ - to  $\beta$ -phase indicates the lower Seebeck coefficients of  $\beta$ -Ag<sub>2</sub>Te than those of  $\alpha$ -Ag<sub>2</sub>Te.

Table 2 shows the Hall carrier concentration  $n_H$ , Hall mobility  $\mu_H$ , Seebeck coefficient  $S$ , and effective mass  $m^*$  at 300 K of the Sb<sub>2</sub>Te<sub>3</sub>/Ag<sub>2</sub>Te composites. For increased Ag<sub>2</sub>Te dispersion concentration, the Hall carrier concentrations are decreased and the Hall mobilities are systematically increased. The lower Hall mobility of Sb<sub>2</sub>Te<sub>3</sub> rich composite (32/1) than the  $\mu_H$  value of the ST/AT = 1/1 composite implies that the systematic increase of the Hall mobility with increasing Ag<sub>2</sub>Te dispersion concentration does not originate from the grain boundary scattering of carriers but rather from the decreased electron-electron scattering due to the lower carrier concentration in the 1/1 composite.

The effective mass of the carrier can be calculated by the following relation in a degenerated semiconductor.[32]

$$S = \frac{8\pi^2 k_B^2}{3e\hbar^2} m^* T \left(\frac{\pi}{3n}\right)^{2/3}$$

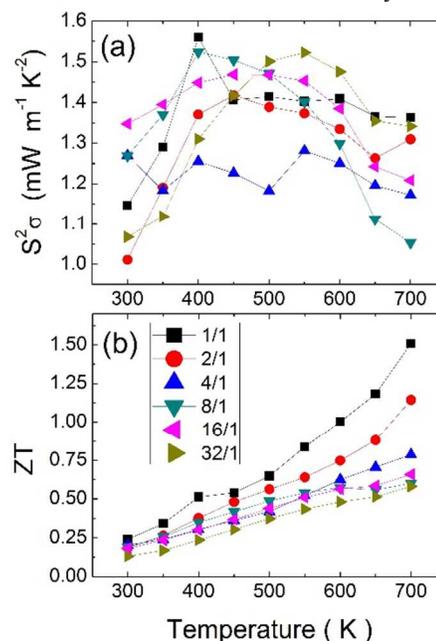
where  $m^*$  and  $n$  are the effective mass of the carrier and the carrier concentration, respectively. The effective masses of the ST/AT = 2/1, 4/1, 8/1, and 16/1 composites are almost identical. For the sample of ST/AT = 1/1, the effective mass is significantly lower than the other's effective masses due to the relatively low carrier concentration of the composite.

Sb <sub>2</sub> Te <sub>3</sub> / Ag <sub>2</sub> Te	$n_H$ ( $\times 10^{19} \text{ cm}^{-3}$ )	$\mu_H$ ( $\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ )	$S$ ( $\mu\text{V/K}$ )	$m^*$ ( $m_e$ )
1/1	2.32	493.770	69.91	0.28
2/1	4.29	283.381	65.61	0.40
4/1	5.67	232.828	58.81	0.43
8/1	5.82	252.039	52.92	0.40
16/1	6.87	234.177	50.81	0.42
32/1	8.96	189.975	48.82	0.49

**Table 2** The Hall carrier concentration  $n_H$ , Hall mobility  $\mu_H$ , Seebeck coefficient  $S$ , and effective mass  $m^*$  of the Sb<sub>2</sub>Te<sub>3</sub>/Ag<sub>2</sub>Te (ST/AT) composites with a molar ratios of (ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1) at 300 K.

The total and lattice thermal conductivity are presented in Fig. 3(c) and 3(d), respectively. The temperature-dependent total thermal conductivity  $\kappa(T)$  exhibits decreased values with increasing temperature. Also, the thermal conductivities are found to decrease with increasing Ag<sub>2</sub>Te dispersion concentration, this trend is similar to the behaviour of the electrical conductivity. For the ST/AT = 1/1 composite, the  $\kappa(T)$  ranges from 0.6~1.4 W m<sup>-1</sup> K<sup>-1</sup>, while the ST/AT = 32/1 composite has higher thermal conductivity from 1.6~2.4 W m<sup>-1</sup> K<sup>-1</sup>. It was reported that the thermal conductivity value of Ag<sub>2</sub>Te have relatively low, in a range of 0.6~0.9 W m<sup>-1</sup> K<sup>-1</sup> (Ag<sub>2</sub>Te).[27] The thermal conductivity of the single crystalline Sb<sub>2</sub>Te<sub>3</sub> compound is 5.6 W m<sup>-1</sup> K<sup>-1</sup> along the in-plane direction and 1.6 W m<sup>-1</sup> K<sup>-1</sup> in the out-of-plane direction.[34] Polycrystalline bulk Sb<sub>2</sub>Te<sub>3</sub> still has a high thermal conductivity of 4.7 W m<sup>-1</sup> K<sup>-1</sup>. [35] Accordingly, the decreased thermal

conductivity with increasing Ag<sub>2</sub>Te concentration can be understood as resulting from the lower thermal conductivity contribution of Ag<sub>2</sub>Te than those of Sb<sub>2</sub>Te<sub>3</sub> as well as from the behaviour of the electronic thermal conductivity.



**Fig. 4** Temperature-dependent power factor  $S^2\sigma$  (a) and  $ZT$  values (b) of the Sb<sub>2</sub>Te<sub>3</sub>/Ag<sub>2</sub>Te (ST/AT) composites with molar ratios of (ST/AT = 1/1, 2/1, 4/1, 8/1, 16/1, and 32/1).

Because the thermal conductivity consists of electronic  $\kappa_{el}$  and lattice thermal conductivity  $\kappa_{ph}$ ;  $\kappa = \kappa_{el} + \kappa_{ph}$ , we should subtract the electronic thermal conductivity in order to obtain the lattice thermal conductivity. The electronic thermal conductivity  $\kappa_{el}$  can be calculated by the Wiedemann-Franz law  $\kappa_{el} = L_0\sigma T$ , where  $L_0$ ,  $\sigma$ , and  $T$  are the Lorenz number, electrical conductivity, and absolute temperature, respectively. In usual cases, the Lorenz number is written as:

$$L_0 = \frac{\pi^2}{3} \left(\frac{k_B}{e}\right)^2 = 2.45 \times 10^{-8} \text{ W } \Omega \text{ K}^{-2}$$

However, the Lorenz number is incorrect in correlated metal and many degenerated semiconductors. In order to obtain a more reliable Lorenz number, we calculated the Lorenz number by using the following equation:[36]

$$L = \left(\frac{k_B}{e}\right)^2 \left( \frac{\left(r + \frac{7}{2}\right) F_{r+\frac{5}{2}}(\eta)}{\left(r + \frac{3}{2}\right) F_{r+\frac{1}{2}}(\eta)} - \left[ \frac{\left(r + \frac{5}{2}\right) F_{r+\frac{3}{2}}(\eta)}{\left(r + \frac{3}{2}\right) F_{r+\frac{1}{2}}(\eta)} \right]^2 \right)$$

where  $r$  is the scattering parameter,  $\eta = E_F/k_B T$  is the reduced Fermi energy, and  $F_n(\eta)$  is the  $n$ -th order Fermi integral given by

$$F_n(\eta) = \int_0^\infty \frac{x^n}{1 + e^{x-\eta}} dx$$

For most cases, the scattering parameter for acoustic phonon scattering is  $r = -1/2$ . The reduced Fermi energy  $\eta$  can be obtained from a fitting of the Seebeck coefficient to the following equation:

$$S = \pm \frac{k_B}{e} \left\{ \frac{(r + \frac{5}{2})F_{r+\frac{3}{2}}(\eta)}{3(r + \frac{3}{2})F_{r+\frac{1}{2}}(\eta)} - \eta \right\}$$

The calculation result for the Lorenz number by the above equations is shown in the inset of Fig. 3(d). The Lorenz number decreases with increasing temperature in a range of  $(1.7\sim 2.1) \times 10^{-8} \text{ W } \Omega \text{ K}^{-2}$  which is lower than the value of  $L_0 = 2.45 \times 10^{-8} \text{ W } \Omega \text{ K}^{-2}$ . By subtracting the electronic thermal conductivity  $\kappa_{el} = L\sigma T$ , the lattice thermal conductivities are in the range of  $\kappa_{ph} = 0.1\sim 0.7 \text{ W m}^{-1} \text{ K}^{-1}$ . The lattice thermal conductivity of the  $\text{Sb}_2\text{Te}_3$  rich composite ( $\text{ST}/\text{AT} = 32/1$ ) has a trend similar to that of the stoichiometric one (open star symbol) shown in Fig. 3(d). The lattice thermal conductivities decrease with increasing  $\text{Ag}_2\text{Te}$  dispersion concentration, presumably due to grain boundary phonon scattering. In the case of the  $\text{ST}/\text{AT} = 1/1$  composite, the extremely low lattice conductivity at high temperature,  $0.1 \text{ W m}^{-1} \text{ K}^{-1}$  should be investigated in a further research.

The power factor  $S^2\sigma$  and dimensionless figure-of-merit  $ZT$  of  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  ( $\text{ST}/\text{AT}$ ) composites are presented in Fig. 4 (a) and (b), respectively. The power factor shows high values  $1.0\sim 1.5 \text{ mW m}^{-1} \text{ K}^{-2}$  with temperature-insensitive behaviour over a wide temperature range; this is an important criterion for practical application in thermoelectric materials. While the power factors are not sensitive to the  $\text{Ag}_2\text{Te}$  concentration, the  $ZT$  values increase with increasing of the  $\text{Ag}_2\text{Te}$  dispersion concentration due to decreased thermal conductivity.

The maximum high  $ZT$  of 1.5 is reached at 700 K for the  $\text{ST}/\text{AT} = 1/1$  composite. The  $ZT = 1.5$  at this temperature range is one of the highest values ever recorded in Pb-free  $p$ -type thermoelectric chalcogenides. In addition, this value represents a more enhanced value of thermoelectric performance than the  $ZT$  value of  $\text{Ag}_2\text{Te}$  related phase separated composites of  $\text{Ag}_2\text{Te}/\text{AgSbTe}_2$ , which exhibit thermoelectric performance around  $ZT 1.0\sim 1.3$ .<sup>[16,17]</sup> It is very likely that it will be possible to increase the  $ZT$  value further by controlling the texture of the composites and the  $\text{Ag}_2\text{Te}$  dispersion concentration then these materials can be practically applied in the field of waste heat power generation as Pb-free  $p$ -type thermoelectric materials.

## Conclusions

We synthesized  $\text{Sb}_2\text{Te}_3/\text{Ag}_2\text{Te}$  ( $\text{ST}/\text{AT}$ ) composites with molar ratios of  $\text{ST}/\text{AT} = 1/1, 2/1, 4/1, 8/1, 16/1, \text{ and } 32/1$  by melting, hand grinding, wet milling, and hot press sintering processes. The low temperature wet milling process can disperse  $\text{Ag}_2\text{Te}$  particles effectively in the  $\text{Sb}_2\text{Te}_3$  matrix. Thermoelectric property investigation indicated that the  $\text{Ag}_2\text{Te}$  dispersion decreases the electrical conductivity as well as the thermal conductivity reduction. The metallic or semi-metallic behaviour of the composites, because of high electrical conductivity at room temperature and high Seebeck coefficient at high temperature, respectively, gives rise to the temperature insensitive behaviour of the power factor over a wide temperature range. The high power factor over a wide temperature range and the low thermal conductivity are beneficial to increase the  $ZT$  value. Thereby, the  $ZT$  value reached 1.5 for the  $\text{ST}/\text{AT} = 1/1$  composite, which is a more enhanced value for related Pb-free chalcogenide thermoelectric

materials. Further investigation of texture and concentration control can increase the  $ZT$  value significantly allowing it to reach a level suitable for practical application in waste heat power generation.

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## Notes and references

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