Journal of Materials Chemistry C



### Flexible Single-Crystalline GaN Substrate by Direct Deposition of III-N Thin Films on Polycrystalline Metal Tape

Journal:	Journal of Materials Chemistry C
Manuscript ID	TC-ART-09-2020-004634.R1
Article Type:	Paper
Date Submitted by the Author:	10-Dec-2020
Complete List of Authors:	Shervin, Shahab; University of Houston System Moradnia, Mina; University of Houston System Alam, Md Kamrul; University of Houston, Tong, Tian; University of Houston, Department of Electrical and Computer Engineering Ji, Mi-Hee; Oak Ridge National Laboratory Chen, Jie; University of Houston System, Mechanical Engineering Pouladi, Sara; University of Houston System Detchprohm, Theeradetch; Georgia Institute of Technology Forrest, Rebecca; University of Houston, Department of Electrical and Computer Engineering Bao, Jiming; University of Houston, Department of Electrical and Computer Engineering Dupuis, Russell; Georgia Institute of Technology Ryou, Jae Hyun; University of Houston System,



## ARTICLE

# Flexible Single-Crystalline GaN Substrate by Direct Deposition of III-N Thin Films on Polycrystalline Metal Tape<sup>+</sup>

Shahab Shervin<sup>§,a,b,c</sup>, Mina Moradnia<sup>§</sup>,<sup>a,c</sup>, Kamrul Alam<sup>d</sup>, Tain Tong<sup>d</sup>, Mi-Hee Ji<sup>e</sup>, Jie Chen<sup>a,b,c</sup>, Sara Pouladi<sup>b,c</sup>, Theeradetch Detchprohm<sup>e</sup>, Rebecca Forrest<sup>f</sup>, Jiming Bao<sup>d</sup>, Russell D. Dupuis<sup>e</sup>, and Jae-Hyun Ryou<sup>\*,a,b,c</sup>

Flexible electronics and mechanically bendable devices based on Group III-N semiconductor materials are emerging, while there are several challenges in manufacturing-cost reduction, device stability and flexibility, and device-performance improvement. To overcome the limitations, it is necessary to replace the brittle and expensive semiconductor wafers with single-crystalline flexible templates for a new-bandgap semiconductor platform. The substrates in the new-concept of semiconductor materials have a hybrid structure consisting of a single-crystalline III-N thin film on a flexible metal tape substrate which provides a convenient and scalable roll-to-roll deposition process. We present a detailed study of a unique and simple direct epitaxial growth technique for the crystallinity transformation to deliver single-crystalline GaN thin film with highly oriented grains along both a- and c- axis directions on a flexible and polycrystalline copper tape. A 2-dimensional (2D) graphene having the same atomic configuration as the (0001) basal plane of wurtzite structure is employed as a seed layer which plays a key role in following III-N epitaxy growth. DC reactive magnetron sputtering method is then applied to deposit an AlN layer in an optimized condition to achieve preferred-orientation growth. Finally, single-crystalline GaN layers (~1 µm) are epitaxially grown by metalorganic chemical vapor deposition (MOCVD) on the biaxially-textured buffer layer. The obtained flexible single-crystalline GaN film by this method suggests a new way for a wide-bandgap semiconductor platform pursuing flexible, high-performance, and versatile device technology.

#### Introduction

Group III-nitride (III-N) wide-bandgap (WBG) semiconductor has become a major material platform in the solid-state electronics and photonics because of their excellent properties including a wide range of direct bandgap energies ( $E_g = 0.7 - 6.2$  eV), high saturation velocity of carriers (~1.5–2×10<sup>7</sup> cm/s), and high breakdown field (~3–6×10<sup>6</sup> V/cm). Their devices, therefore, play a dominant role in short-wavelength optoelectronics and highfrequency and high-power electronics<sup>1–5</sup>. III-N thin-film materials are conventionally grown on single-crystalline

<sup>a.</sup> Department of Mechanical Engineering, University of Houston, Houston, Texas 77204-4006, United States.

- <sup>c</sup> Texas Center for Superconductivity at UH (TcSUH) and Advanced Manufacturing Institute (AMI), University of Houston, Houston, Texas 77204, United States.
- <sup>d</sup>. Department of Electrical and Computer Engineering, University of Houston, Houston, Texas 77204, United States.
- e. Center for Compound Semiconductors, School of Electrical and Computer Engineering, Georgia Institute of Technology, Atlanta, Georgia 30332-0250, United States.
- <sup>f.</sup> Department of Physics, University of Houston, Houston, Texas 77204, United States.
- \$ These authors contributed equally to this work.

sapphire (Al<sub>2</sub>O<sub>3</sub>), silicon carbide (SiC), and silicon (Si) (111) substrates by strained heteroepitaxy<sup>6</sup>. Sapphire is the most dominantly used substrate due to its relatively small thermal expansion coefficient (TEC) difference and in-plane lattice mismatch (LM) to achieve high-quality gallium nitride (GaN) and aluminum nitride (AIN) templates<sup>7,8</sup>. In-plane LMs and TEC differences between GaN and different substrates are shown in Fig. S1 (ESI). Despite the small LM and TEC difference between SiC and GaN for higher-quality epitaxial films, the wide adoption of SiC substrate in the commercial market has been hampered due to its high cost<sup>9</sup>. The use of Si (111) substrate generally results in the epitaxial film with a higher density of crystalline defects, even though it is the most economical approach for the manufacturing of III-N devices<sup>10,11</sup>.

Demands for wearable and flexible devices are growing in solidstate lighting (SSL) and electronics for a wide range of applications such as bendable/curved luminaire and display of flexible electronics. Recent developments in flexible electronics demonstrate that semiconductor materials can be very flexible when their thickness is small enough (on the order of micrometers), deviating from the brittle nature of bulk properties<sup>12–16</sup>. However, currently developed substrates cannot provide a platform for flexible electronic and photonic devices due to the brittle nature of single-crystalline wafers with a thickness of hundreds of micrometers. To fabricate mechanically flexible III-N devices and components, therefore, it is required to transfer the thin film for a device to a bendable secondary substrate followed by the removal of the original wafer substrate. This film-transfer technique involves additional

<sup>&</sup>lt;sup>b.</sup> Materials Science and Engineering Program, University of Houston, Houston, Texas 77204, United States.

<sup>\*</sup> Correspondence and requests for materials should be addressed to J.H.R (email: <u>iryou@uh.edu</u>).

<sup>&</sup>lt;sup>+</sup> Electronic supplementary information (ESI) available: TEC and LM calculation, AlN sputtering condition optimization, schematic of biaxial and uniaxial AlN growth, EBSD texture, and TEM analysis of GaN. See DOI:

complicated processing steps, hence, poses technical and economic challenges in the manufacturing of flexible semiconductor devices and systems, such as lower production yields, potential degradation of device performance, higher manufacturing cost, and process scalability limited by the current choices of wafer substrates. The challenges are associated with the current "top-down" process and can be resolved by the adoption of the "bottom-up" process. For the bottom-up process, it is necessary to begin with a flexible substrate instead of the brittle wafer substrates for the deposition of active thin-film device layers. However, all the current flexible substrates are non-single-crystalline materials. The deposition of thin films on such substrates results in nonsingle-crystalline materials for active semiconductor layers, which significantly degrade the performance characteristics of the devices. For the development of high-performance flexible devices, therefore, it is critical to obtain single-crystalline thin films even on non-crystalline substrates, overcoming the traditional epitaxial relationship between the substrate and film.

In the present study, we develop a direct deposition technique in which high crystalline quality of III-N thin films can be epitaxially grown on a non-single-crystalline substrate. We demonstrate a single-crystalline GaN template layer on a flexible copper (Cu) tape substrate by employing graphene as a seed layer and AIN as a buffer layer. Direct growth of the 2D material on the polycrystalline substrate followed by the deposition of the buffer layer with crystallographically-selected orientations in both the in-plane and out-of-plane directions of hexagonal lattice provides a critical solution for the crystallinity transformation from the polycrystalline substrate to single-crystalline GaN thin film. The result is the flexible single-crystalline GaN substrate consisting of a GaN template layer, crystallinity-transformational layers, and Cu tape. The use of this new flexible, inexpensive, and lengthy tape substrate, instead of brittle, expensive, and limited-area wafer substrates, will enable low-cost manufacturing of easily-scalable flexible wide-bandgap semiconductor devices compatible with the roll-to-roll process. Moreover, the flexible GaN substrate will help the easy fabrication of multifunctional III-N devices for electronic, photonic, sensing, and energy harvesting applications<sup>17-24</sup>. Among potential applications, large-scale, flexible inorganic LEDs could be applied for flexible LED TV displays, large-scale energy-efficient lighting, or window displays in the future. Furthermore, single-crystalline, GaN-based, flexible electronics, which promise higher efficiency, reliability, and a long lifetime, can replace the aim of organic flexible electronics in a range of applications, i.e., biomedical applications exploiting high electron mobility transistor-based sensors or wearable electronics.

#### Experimental

#### Graphene Growth on Cu

Chemical vapor deposition (CVD) was applied for the growth of graphene in a custom-built system, which consists of a quartz tube as a reaction chamber and a heating furnace with  $\sim$  30-cm

uniform temperature zone. Cu foil tapes with a thickness of 25 µm and a surface area of 16×21 mm<sup>2</sup> (99.8%, Alfa), sitting inside a quartz sample holder, were loaded into the reaction chamber and placed in the uniform-temperature zone of the furnace. The reaction chamber was evacuated thoroughly and then quickly filled to ambient pressure with argon (Ar) gas. Subsequently, the Cu tapes were annealed at 1050 °C in a slight reduction environment using a mixture of Ar (600 sccm [standards cubic centimeter per minute]) and hydrogen (H<sub>2</sub>, 60 sccm) during the heating and stabilizing process. A precursor with a mixture of methane (CH<sub>4</sub>, 100 sccm) and H<sub>2</sub> (30 sccm) was introduced into the reaction chamber by a carrier gas of Ar (1200 sccm) for the graphene growth for ~2 hours. Samples were fast-cooled down to room temperature after the shutoff of CH<sub>4</sub>, typically within an hour, by pushing the sample holder out of the heating zone, under the protection of Ar and H<sub>2</sub>.

#### AIN Thin-Film Growth on Graphene/Cu

AlN thin films were grown on the graphene-deposited Cu tapes using a DC reactive magnetron sputtering system (AJA International) equipped with a DC power supplier (Advanced Energy MDX 500). The base and working pressures of the sputtering chamber were 10<sup>-9</sup> Torr and 2 mTorr, respectively. The working pressure was set by adjusting a gate valve to a cryopump. The AIN films were deposited with an AI target (1.5inch diameter, 99.5%, AJA International) at 110 W of a DC cathode power in a mixture of Ar and nitrogen  $(N_2)$  gases, which were introduced into the chamber by separate mass flow controllers. Before and during the sputtering, the sample holder was rotated at 20 revolutions per minute (rpm) to ensure the uniformity of the film. Prior to each run, pre-sputtering was performed with Ar<sup>+</sup> ion bombardment to clean the Al target for 5 minutes with the movable shutter covering the substrate. Reactive sputtering was initiated by introducing N<sub>2</sub> into the chamber and opening the shutter. After the deposition, a postsputtering was performed for 10 min to avoid the target poisoning. The oxygen  $(O_2)$  and water  $(H_2O)$  partial pressures in the chamber were monitored by a residual gas analyzer (RGA) (Kurt J. Lesker) and kept as low as 10<sup>-10</sup> Torr and 10<sup>-11</sup> Torr ranges, respectively. A 200-nm-thick AIN film was deposited for 90 minutes using an optimized condition.

#### GaN Thin-Film Growth on AlN/Graphene/Cu

Epitaxial GaN layers were grown by metalorganic chemical vapor deposition (MOCVD) in a 6×2" reactor system equipped with a close-coupled showerhead chamber (Thomas Swan). Trimethylgallium (TMGa, Ga(CH<sub>3</sub>)<sub>3</sub>), and ammonia (NH<sub>3</sub>) were used as Group-III and Group-V precursors with H<sub>2</sub> carrier gas. The growth of a GaN layer was carried out at a chamber pressure of 500 Torr and a temperature of 900–950 °C for a target thickness of ~1  $\mu$ m. The growth temperature was intentionally set at lower than the typical growth temperature of GaN by MOCVD (~1050 °C) due to a concern of potential damage in Cu tape by overshoot during the heating-up.

#### Characterizations

The graphene was characterized by Raman spectroscopy using a 532-nm solid-state laser for the source of excitation in a home-built setup with a spectrometer (Horiba Jobin Yvon) at room temperature. Microscopic surface morphology of the graphene and III-N layers were characterized by atomic force microscopy (AFM) (Veeco Dimensions 3000). Microscopic surface images were taken by scanning-electron microscopy (SEM) (JEOL JSM 6400). Structural and crystalline qualities of the AlN and GaN layers were characterized by electron backscattered diffraction (EBSD) (Hikari Super EBSD camera from Ametek/EDAX attached to a Philips XL 30SFEG electron microscope) and high-resolution X-ray diffraction (HR-XRD) (Bruker D8 with the parallel beam for thin-film characterization).

#### Results

#### 2-Dimensional Material Seed Layer

A planar hexagonal atomic configuration of *sp*<sup>2</sup> bonded carbon atoms in graphene, which also creates weak van-der-Waals outof-plane binding at the cleavage planes, is expected to provide the preferential nucleation sites for the epitaxial growth of the hexagonal-lattice 3-dimensional structures. A graphene domain can grow over polycrystalline Cu grains crossing the grain boundaries of Cu due to the weak interaction between them for the surface catalyzed process, hence, no epitaxial relationship between the Cu lattice and graphene grains<sup>25</sup>. Therefore, relatively large single-crystalline graphene grain can be formed on polycrystalline Cu foil<sup>25</sup>. Figure 1a shows a microscopic surface image of a graphene film on a Cu tape substrate grown by chemical vapor deposition (CVD). The presence of wrinkles (as pointed by yellow arrows in the scanning-electron microscopy (SEM) image of Fig. 1a) confirms that the graphene film is continuous. The wrinkles are formed as the graphene expands during cool-down due to its negative TEC value. A red arrow shows a grain boundary of Cu, which also indicates the graphene domain covers the multiple grains of the Cu. The average grain size of Cu is estimated to be ~0.5 mm. Raman spectrum of the graphene shown in Fig. 1b presents a typical feature of few-layered graphene with a symmetric 2D-band peak centered at about 2690 cm<sup>-1</sup> with a full-width-at-halfmaximum (FWHM) of ~40 cm<sup>-1</sup>. An intensity ratio of 2D-band peak to G-band peak  $(I_{2D}/I_G)$  is sensitive to the layer numbers of graphene and  $I_{\rm 2D}/I_{\rm G}$  characterizes the disorder degree of graphene<sup>26,27</sup>. The  $I_{2D}/I_{G}$  intensity is ~1.6 and no D-band peak is present, confirming that high-quality defect-free graphene with less than three layers is grown on the Cu tape substrate.

#### **AIN Buffer Layer**

A proper choice of a buffer layer between the substrate and epitaxial film is critical to improve the crystalline quality of the III-N films grown by strained heteroepitaxy<sup>28–31</sup>. Especially, the growth of a GaN layer, which serves as a template for most III-N devices, on the graphene is significantly different from those on other bulk substrates. For the epitaxial growth of 3D materials on 2D substrates, it is necessary to have at least one side of the interface with dangling bonds. However, the lack of

dangling bonds of the graphene, which makes it chemically inert with very low surface energy, causes a thermodynamic problem for the nucleation of foreign atoms with the strong *sp*<sup>3</sup> bonding of GaN<sup>32</sup>. Therefore, directly-grown GaN thin films on graphene were found to contain microstructural defects<sup>32–34</sup>. For the epitaxial growth of GaN on graphene, it is necessary to mitigate the large surface free energy difference between GaN and graphene layers (1970 and 52 mJ/m<sup>2</sup> respectively) by introducing a proper nucleation layer compensating the lack of dangling bonds on the graphene surface.

We employed an AIN buffer by sputtering, which has reasonable dangling bonds and comparable surface energy to promote uniform nucleation of III-N layers. The AIN possesses the same crystal structure as GaN and its basal plane has the same atomic configuration to the underlying 2D material<sup>33,35</sup>. The sputtered AIN film does not necessarily follow the crystal structure of the underlying layer and sputtering parameters directly affect the final structure and quality of the deposited thin film. Reactive sputtering parameters of the AIN film, such as deposition rate, the target-to-substrate distance, and a ratio of reactive ion gas flow rate to total flow rate, were optimized based on the x-ray diffraction (XRD) characterization results. The Si (100) substrate was used for the calibration of AIN sputtering parameters due to its easy availability. By introducing N2 gas to the Ar plasma, the deposition rate decreases drastically, as the collision of particles in the plasma increases. A reactive gas ratio was optimized at R = ~44% with flow rates of Ar and N<sub>2</sub> set at 10 sccm and 8 sccm, respectively, based on the XRD 2theta-omega  $(2\vartheta - \omega)$  scans (Fig. S2 (ESI<sup>+</sup>)) to achieve the highly *c*-axis-textured AlN in the out-of-plane direction. A mean free path ( $\lambda$ ) of sputtered atoms should be larger than the target-to-substrate distance (d). The mean free path of Al atoms, which depends on the working pressure,  $\lambda = \sim 10$  cm at P = 2 mTorr<sup>36,37</sup>; hence, the substrate should be kept less than 10 cm away from the target. The optimum target-to-substrate distance of d = 8 cm was obtained using the reactive gas ratio of R = ~44% (Fig. S3 (ESI<sup>+</sup>)). Fig. 2a shows XRD  $2\vartheta$ - $\omega$  scans of the AIN films on the graphene/Cu substrates at various deposition temperatures. For all the deposition temperatures, the AIN films are preferably grown in the *c*-axis of wurtzite structure normal to the surface, i.e., along the out-of-plane direction. However, small peaks of (1010) and (1011) are also detected for AIN deposited at lower temperatures (400-600 °C), showing minor polycrystalline nature of the AIN films, which may originate from adatoms' energy reaching the substrate and bombardment of sputtering ions<sup>38</sup>. Fig. 2b compares the FWHM values of the peaks from AIN (0002) rocking curves ( $\omega$  scan) at various deposition temperatures. With increasing deposition temperature, the FWHM values decrease, indicating less spread in the tilt of cplane out of surface; hence, the crystalline quality of the AIN film is improved. The alignment of a c-axis of AIN in the out-ofplane direction is confirmed (refer to a uniaxially textured structure schematically shown in Fig. S4a (ESI<sup>+</sup>).

To achieve the single-crystallinity of the film, in-plane alignment (*a*-axis) as well as out-of-plane alignment of AlN needs to be confirmed. To evaluate the in-plane alignment of the deposited AlN films, XRD phi ( $\varphi$ ) scan was performed by rotating the sample at the diffraction angles of AlN {1012} asymmetric

planes, i.e.,  $\omega = 42.73^{\circ}$  and  $2\vartheta = 49.84^{\circ}$ . Fig. 3a shows the XRD  $\varphi$ scan for the peaks of {1012} planes of AIN film deposited at 650 °C. Six peaks separated by 60°, showing the six-fold rotational symmetry, are observed, which confirms that the AIN film is highly aligned in both *a*-axis and *c*-axis directions (refer to a biaxially textured structure schematically shown in Fig. S4b (ESI<sup>+</sup>). The FWHM value for one of the peaks is estimated to be  $\Delta \varphi_{AIN (10^{1}2)} = ~4^{\circ}$ . The electron backscatter diffraction (EBSD) data further confirm the single-crystalline nature of AIN. Fig. 3b shows that the surface of the AIN has a uniform (0001) plane. The pole figures in Fig. 3c also indicates intense (0001) reflections at the center, showing the growth of highly *c*-axis oriented AIN. Furthermore, a clear six-fold rotational symmetric feature is observed in the pole figures of (1120) and (1010) planes. Figure 3d and e show the crack-free and flat surface morphology of the AIN. The root-mean-square (rms) roughness of the AIN surface measured by AFM is ~ 6 nm.

#### **GaN Template Layer**

After the successful development of nearly a single-crystalline AlN buffer layer, the epitaxial growth of a GaN film becomes less challenging. Fig. 4a shows an XRD  $2\vartheta - \omega$  scan in a wide range for the GaN directly grown on the AIN/graphene/Cu substrate. The presence of only (000/) peaks such as (0002) and (0004) peaks and no other peak of GaN and AIN further confirm the epitaxial growth of III-N films along the c-axis of wurtzite structure on graphene. The polycrystalline nature of the Cu tape is clear from different peaks related to different crystallographic orientations of the Cu grains. Fig. 4b shows an XRD  $\varphi$  (rotational) scan around GaN {1012}. Again, six separate peaks with 60° apart indicate both the in-plane and out-of-plane alignments, supporting the epitaxial growth of a single-crystalline GaN film. The XRD results also show that the crystalline quality of GaN is further improved. The FWHM values of the peaks are slightly decreased for those of AIN peaks,  $\Delta \varphi_{\text{GaN}(10)}^{1} \sim 3^{\circ}$ . Considering the large size of the XRD slit and the low  $\Delta \phi$  values in GaN {101 2} peaks, it is clear that single-crystalline GaN domains are large enough in the range of more than hundreds of micrometers and millimeters.

Motivated by scientific interests and practical importance, the epitaxial growth of III-N films on graphene has been actively investigated using different approaches in recent years. Among them, Yu et al. reported a direct growth of GaN on flexible graphene-deposited Cu substrates. While the single-crystalline AIN and GaN were claimed in a selected condition, it was not fully confirmed by characterizations in a large area such as asymmetric XRD scan<sup>39</sup>. Moreover, the study employed migration-enhanced molecular beam epitaxy (ME-MBE), which is an inherently very low throughput process. Another work by Ramesh et al. by laser molecular beam epitaxy (LMBE) for the growth of GaN film on graphene/Cu foil reported a non-singlecrystalline mixed-phase of GaN grain film<sup>40</sup>. Chung et al.'s work employed a high throughput CVD process to grow a GaN film on a non-single-crystalline silica substrate<sup>41</sup>. However, the work was not a direct growth process, as the graphene layer had to be transferred to the substrate. Also, the substrate is not mechanically flexible which makes the process incompatible to

the roll-to-roll process and flexible devices. Kim et al. and Balushi et al. reported single-crystalline GaN and AlN layers on graphene/(0001) SiC wafer substrate<sup>33,42</sup>. Also, Zhang et al. reported the reduced threading dislocation density in the GaN grown on graphene/AlN on a (0001) sapphire substrate<sup>43</sup>. In those studies, it is not clear if the epitaxial growth occurs by the graphene, since the single-crystalline SiC (or AlN) can also play a role in the epitaxial growth. The studies also require an additional transfer process for flexible devices. Our study is the first demonstration of single-crystalline GaN directly grown on polycrystalline flexible metal tape using high throughput processes.

The outcome of this single-crystalline GaN epitaxy produces a flexible GaN substrate consisting of the GaN template layer on Cu tape with transfer-free seed/buffer layers as shown in the photo of Fig. 4e. The EBSD characterization further confirms the formation of a singlecrystalline GaN film (Fig. S5a and S5b (ESI<sup>+</sup>)). However, the AIN and GaN films contain threading dislocations as shown in the crosssectional transmission-electron microscopy (TEM) image of Fig. S5c (ESI<sup>+</sup>), similar to all the III-N epitaxial films grown by heteroepitaxy. The density of threading dislocations in the single-crystalline GaN layer is estimated to be ~4×10<sup>10</sup> cm<sup>-2</sup> and ~7×10<sup>10</sup> cm<sup>-2</sup> for screw and edge dislocations, respectively, from XRD rocking curves of (0002) and (1012) planes using a method suggested by Kaganer et al.<sup>44</sup> The density is higher than typical values of well-matured GaN epitaxial films on single-crystalline substrates of SiC, sapphire, and Si (e.g., 10<sup>8</sup>–10<sup>9</sup> cm<sup>-2</sup>), as expected. The density of dislocation is required to be reduced by further optimization of the deposition processes to achieve comparable performance of devices to those of wafer-based non-flexible devices. Fig. 4c and d show the surface morphology of the GaN layer. The SEM image shows a continuous film with a similar surface to that of AIN. The surface of the GaN probed by AFM is slightly rougher than that of AIN, possibly related to lower growth temperature than the typical MOCVD growth temperatures of GaN by 100-150 °C.

#### Van-der-Waals Epitaxy vs Remote Epitaxy

Single-crystalline AIN is confirmed to be epitaxially grown on the graphene layer. Van-der-Waals epitaxy (vdWE) was proposed as a possible mechanism of a thin-film growth on a 2D material. The vdWE is different from the conventional epitaxy where strong chemical bondings exist at the interface of the substrate and epitaxial materials. The vdWE makes crystalline growth possible between 2D and 3D materials by a weak van der Waals interaction at the heterojunction. The vdWE eliminates the necessity of reducing lattice mismatch between epitaxial layers and underlying 2D materials and allows the heteroepitaxial growth for various materials with different crystal structures due to having no broken bond on the vdW surface<sup>40,45</sup>. The vdWE has been experimentally observed in III-N film on graphene. Qi et al. introduced graphene as an interlayer for the growth of AIN on a sapphire substrate. Decreasing threading dislocations along with strain relaxation could come from reducing lattice mismatch between AIN and sapphire substrate by graphen<sup>46</sup>. Despite the less crystalline quality of GaN which was grown on the transferred multilayer graphene on sapphire, Li et al. found that light-emitting diodes grown on such multilayer graphene resulted in high output power<sup>47</sup>. These results could address the problem of large lattice and thermal mismatch between the III-N films and substrate by using multilayer graphene. However, in the previous reports, the growth of III-N occurred on the graphene that was transferred onto a sapphire substrate, i.e., on a graphene/sapphire substrate. Therefore, the epitaxy of III-N film was not confirmed to originate from either graphene or sapphire. The epitaxial relationship between III-N and sapphire substrate was reported, suggesting that the III-N epitaxy may still originate from the sapphire<sup>48</sup>. Kim et al. recently introduced remote epitaxy as another possible epitaxial growth mechanism of the thin film on 2D materials. The remote epitaxy occurs when the substrate directly affects the orientation of epitaxial growth and transfers its crystallographic information to the over-layers during the epitaxy. Because of sufficiently thin and electrical penetration of graphene characteristics, the over-layers can follow the orientation of the underlying substrate<sup>49</sup>.

This study provides an ideal structure to study the mechanism of III-N epitaxy on graphene, which is yet inconclusive because it uses a polycrystalline substrate underneath the graphene. In the III-N epitaxial layers on graphene/Cu of the present study, AlN and GaN are single-crystalline while Cu is polycrystalline as described earlier (Fig. 3a and Fig. 4a and b), which provides evidence to support vdWE over remote epitaxy. According to the remote epitaxy mechanism where a structure mirroring of a substrate happens in an over-layer, the graphene interlayer should transfer the crystallographic structure of the underlying Cu substrate to the AIN over-layer. If the remote epitaxy is the dominant mechanism of the III-N epitaxy on а graphene/underlying substrate, III-N films tend to be polycrystalline too. The remote epitaxy mechanism cannot explain the single-crystal growth of III-N on the polycrystalline Cu substrate. A further comprehensive study was performed on the structural properties of AIN films deposited on different substrates of poly- and single-crystals in combination with with/without the graphene layer, as shown in Fig. 5. All the AIN layers are uniaxially textured along the *c*-axis of the wurtzite structure regardless of the substrate when the optimized sputtering condition was employed (Fig. 5a). This is conventionally observed for AIN deposited by sputtering mainly due to the surface energy of its polar surface<sup>50</sup>. The crystalline quality of the films appears to be the best and worst for the AIN on graphene/single-crystalline Cu and polycrystalline Cu without graphene, respectively, from the intensity and FWHM of the AIN (0002) peaks (Figure 5a). In-plane alignment, however, cannot be obtained without the help of epitaxial relationships from the underlying layer. Figure 5b shows the XRD  $\varphi$  scan around AIN {1012} planes. All the AIN grown on graphene show the in-plane alignment of single-crystals regardless of the crystallinity of Cu, i.e., on graphene/singlecrystalline Cu (black curve of Fig. 5b) and graphene/polycrystalline Cu (red curve of Fig. 5b). In contrast, no in-plane alignment is observed for AIN on polycrystalline Cu without graphene (green curve of Fig. 5b). These results further support the vdWE for III-N on graphene. In other words, singlecrystalline III-N can only be epitaxially grown on graphene. Only polycrystalline uniaxially textured AIN can be achieved without graphene, suggesting that AIN single-crystal growth cannot occur without the effect of graphene on the polycrystalline Cu while it is independent of copper crystal orientation<sup>51</sup>. To exclude a possibility of polycrystalline AIN growth by nonoptimized deposition condition on Cu, we deposited AIN on a single-crystalline Cu substrate without graphene using the same condition. The AIN is single-crystalline on a single-crystalline Cu substrate without graphene (blue curve of Fig. 5b). Therefore, these results suggest that the epitaxy of III-N is influenced by directly underlying layer and vdWE occurs for the growth of III-N layers on graphene/Cu substrate.

GaN and Cu (111) have a significant discrepancy in LM and TEC. Graphene growth is of critical importance for reducing the effect of the lattice mismatch between AIN and copper foil, which comes from the vdWE growth for III-nitrides on graphene<sup>32,41</sup>. Thanks to the weak vdW bonding between AIN and graphene, AIN can be epitaxially grown on Cu substrate without the necessity of reducing LM and TEC mismatch between AIN and Cu (111) substrate.

#### Conclusions

We developed a unique structure material, which enables continuous roll-to-roll (R2R) thin-film epitaxial GaN growth on low-cost and bendable Cu poly-crystalline tape by employing suitable buffer layers. In this method, a 2D graphene interlayer grown by a CVD method provides an in-plane hexagonal atomic arrangement for single-crystalline AIN buffer layer growth by vdW epitaxy growth. Optimized conditions for DC reactive magnetron sputtering was applied to achieve preferredorientation AIN buffer. The XRD pole-figure of AIN (1012) shows a six-fold symmetry of AIN which is highly textured in both aand *c*-axis crystallographic directions. The single-crystalline wurtzite structure of GaN is formed on the biaxial textured AIN buffer layer by MOCVD. The XRD  $2\vartheta$ - $\omega$  scan shows that the *c*plane of GaN and AIN films are preferably grown from (0002) and (0004) peaks. The outcome is the flexible single-crystalline GaN substrate on a Cu tape. The present findings, focusing on the interface, epitaxial relationship, and vdW epitaxy growth mechanisms to achieve single-crystalline III-nitride materials on flexible/polycrystalline substrates, offer a promising solution to meet the rapidly increasing demand for bendable/curved luminaire and display of flexible electronics.

#### **Conflicts of interest**

There are no conflicts to declare.

#### Acknowledgements

The work at the University of Houston is partially supported by the National Science Foundation under Grant No. 1842299 (Electrical, Communications and Cyber Systems (ECCS)) and Grant No. 1907626 (I-CORPS of Industrial Innovation and Partnerships (IIP)) and King Abdullah University of Science and Technology (KAUST), Saudi Arabia (Contract No. OSR-2017-CRG6-3437.02). J.H.R. also acknowledges partial support from the Texas Center for Superconductivity at the University of Houston (TcSUH) and Advanced Manufacturing Institute (AMI).

#### **Notes and References**

1. M. R. Krames, O. B. Shchekin, R. Mueller-Mach, G. O. Mueller, L. Zhou, G. Harbers and M. G. Craford, *J. Display Technol.*, 2007, **3**. 160–175.

2. F. A. Ponce and D. P. Bour, Nature, 1997, 386, 351-359.

3. S. T. Sheppard, K. Doverspike, W. L. Pribble, S. T. Allen, J. W. Palmour, L. T. Kehias and T. J. Jenkins, *IEEE Electron Device Lett.*, 1999, **20**, 161–163.

4. J. B. Limb, D. Yoo, J. H. Ryou, W. Lee, S. C. Shen, R. D. Dupuis, M. L. Reed, C. J. Collins, M. Wraback, D. Hanser, E. Preble, N. M. Williams and K. Evans, *Appl. Phys. Lett.*, 2006, **89**, 011112.

5. M. Ishida, T. Ueda, T. Tanaka and D. Ueda, *IEEE Trans Electron Devices*, 2013, **60**, 3053–3059.

6. J.-H Ryou, in *Nitride Semiconductor Light-Emitting Diodes (LEDs): Materials, Technologies, and Applications,* ed. J.-J. Huang, H.-C. Kuo and S.-C. Shen, Elsevier, Amsterdam, 2nd, 2017, 3, PP. 66–98.

7. S. Nakamura, Jpn. J. Appl. Phys., 1991, 30, L1705-L1707.

8. H. Okano, N. Tanaka, Y. Takahashi, T. Tanaka, K. Shibata and S. Nakano, *Appl.Phys. Lett.*, 1994, **64**, 166–168.

9. M. A. L. Johnson, S. Fujita, W. H. Rowland, K. A. Bowers, W. C. Hughes, Y. W. He, N. A. El Masry, J. W. Cook, J. F. Schetzina, J. Ren and J. A. Edmond, *Solid-State Electron*, 1997, **41**, 213–218.

10. H. Liu, Z. Ye, H. Zhang and B. Zhao, *Mater. Res. Bull.*, 2000, **35**, 1837–1842.

11. A. M. Mizerov, S. N. Timoshnev, M. S. Sobolev, E. V. Nikitina, K. Y. Shubina, T. N. Berezovskaia, I. V. Shtrom and A. D. Bouravleuv, *Semiconductors*, 2018, **52**, 1529–1533.

12. D. Shahrjerdi and S. W. Bedell, *Nano Lett.*, 2013, **13**, 315–320. 13. S. Pouladi, M. Rathi, D. Khatiwada, M. Asadirad, S. K. Oh, P. Dutta, Y. Yao, Y. Gao, S. Sun, Y. Li, S. Shervin, K. H. Lee, V. Selvamanickam and J. H. Ryou, *Prog Photovolt*, 2019, **27**, 30–36.

14. M. Asadirad, Y. Gao, P. Dutta, S. Shervin, S. Sun, S. Ravipati, S. H. Kim, Y. Yao, K. H. Lee, A. P. Litvinchuk, V. Selvamanickam and J. H. Ryou, *Adv. Electron. Mater.*, 2016, **2**, 1600041.

15. W. Wang, J. Chen, J. S. Lundh, S. Shervin, S. K. Oh, S. Pouladi, Z. Rao, J. Y. Kim, M.-K. Kwon, X. Li, S. Choi and J.-H. Ryou, *Appl. Phys. Lett.*, 2020, **116**, 123501.

16. J. Chen, S. K. Oh, N. Nabulsi, H. Johnson, W. Wang and J. H. Ryou, *Nano Energy*, 2019, **57**, 670–679.

17. N. I. Kim, Y. L. Chang, J. Chen, T. Barbee, W. Wang, J. Y. Kim, M. K. Kwon, S. Shervin, M. Moradnia, S. Pouladi, D. Khatiwada, V. Selvamanickam and J. H. Ryou, *Sens. Actuator A Phys.*, 2020, **305**, 111940.

18. J. Chen, H. Liu, W. Wang, N. Nabulsi, W. Zhao, J. Y. Kim, M. K. Kwon and J. H. Ryou, *Adv. Funct. Mater.*, 2019, **29**, 1903162.

19. J. Chen, S. K. Oh, H. Zou, S. Shervin, W. Wang, S. Pouladi, Y. Zi, Z. L. Wang and J. H. Ryou, *ACS Appl. Mater. Interfaces*, 2018, **10**, 12839–12846.

20. W. Wang, S. Shervin, S. Member, S. K. Oh, J. Chen, Y. Huai, S. Pouladi, H. Kim, S. Lee, J. Ryou and S. Member, *IEEE Electron Device Lett.*, 2017, **38**, 1086–1089.

21. Shervin, S. *et al.* Flexible deep-ultraviolet light-emitting diodes for significant improvement of quantum efficiencies by external bending. *J. Phys. D Appl. Phys.*, 2018, **51**, 105105.

22. S. Shervin, S. H. Kim, M. Asadirad, S. Y. Karpov, D. Zimina and J. H. Ryou, *ACS Photonics*, 2016, **3**, 486–493.

23. S. Shervin, S. H. Kim, M. Asadirad, S. Ravipati, K. H. Lee, K. Bulashevich and J. H. Ryou, *Appl. Phys. Lett.*, 2015, **107**, 193504.

24. N. Palavesam, S. Marin, D. Hemmetzberger, C. Landesberger, K. Bock and C. Kutter, *Flex. Print. Electron.*, 2008, **3**, 014002.

25. Q. Yu, L. A. Jauregui, W. Wu, R. Colby, J. Tian, Z. Su, H. Cao, Z. Liu, D. Pandey, D. Wei, T. F. Chung, P. Peng, N. P. Guisinger, E. A. Stach, J. Bao, S. S. Pei and Y. P. Chen, *Nat. Mater.*, 2011, **10**, 443–449.

26. X. Li, W. Cai, J. An, S. Kim, J. Nah, D. Yang, R. Piner, A. Velamakanni, I. Jung, E. Tutuc, S. K. Banerjee, L. Colombo and R. S. Ruoff, *Science*, 2009, **324**, 1312–1314.

27. L. M. Malard, M. A. Pimenta, G. Dresselhaus and M. S. Dresselhaus, *Phys. Rep.*, 2009, **473**, 51–87.

28. T. Warren Weeks, M. D. Bremser, K. S. Ailey, E. Carlson, W. G. Perry and R. F. Davis, *Appl. Phys. Lett.*, 1995, **67**, 401–403.

29. E. Feltin, B. Beaumont, M. Laügt, P. De Mierry, P. Vennéguès, H. Lahrèche, M. Leroux and P. Gibart, *Appl. Phys. Lett.*, 2001, **79**, 3230–3232.

30. H. Amano, N. Sawaki, I. Akasaki and Y. Toyoda, *Appl. Phys. Lett.*, 1986, **48**, 353–355.

31. W. H. Sun, J. P. Zhang, J. W. Yang, H. P. Maruska, M. A. Khan, R. Liu and F. A. Ponce, *Appl. Phys. Lett.*, 2005, **87**, 211915.

32. J. H. Choi, J. Kim, H. Yoo, J. Liu, S. Kim, C. W. Baik, C. R. Cho, J. G. Kang, M. Kim, P. V. Braun, S. Hwang and T. S. Jung, *Adv. Opt. Mater.*, 2016, **4**, 505–521.

33. Z. Y. Al Balushi, T. Miyagi, Y. C. Lin, K. Wang, L. Calderin, G. Bhimanapati, J. M. Redwing and J. A. Robinson, *Surf. Sci.*, 2015, **634**, 81–88.

34. A. Ishii, T. Tatani, H. Asano and K. Nakada, *Phys Status Solidi C* 7, 2010, 7, 347–350.

35. Q. Zeng, Z. Chen, Y. Zhao, T. Wei, X. Chen, Y. Zhang, G. Yuan and J. Li, *Jpn. J. Appl. Phys.*, 2016, **55**, 085501.

36. K. Kusaka, D. Taniguchi, T. Hanabusa and K. Tominaga, *Vacuum*, 2002, **66**, 441–446.

37. T. Aubert, M. B. Assouar, O. Legrani, O. Elmazria, C. Tiusan and S. Robert, *J. Vac. Sci. Technol. A.*, 2011, **29**, 021010.

38. H. Y. Liu, G. S. Tang, F. Zeng and F. Pan, *J. Cryst. Growth*, 2013, **363**, 80–85.

39. J. Yu, Z. Hao, J. Wang, J. Deng, W. Yu, L. Wang, Y. Luo, Y. Han, C. Sun, B. Xiong and H. Li, *J. Alloys Compd.*, 2019, **783**, 633–642.

40. C. Ramesh, P. Tyagi, S. Bera, S. Gautam, K. M. Subhedar, M. S. Kumar and S. S. Kushvaha, *J. Nanosci. Nanotechnol* 2020, **20**, 3929–3934.

41. K. Chung, S. I. Park, H. Baek, J. S. Chung and G. C. Yi, *NPG Asia Mater.*, 2012, **4**, 1–5.

42. J. Kim, C. Bayram, H. Park, C. W. Cheng, C. Dimitrakopoulos, J. A. Ott, K. B. Reuter, S. W. Bedell and D. K. Sadana, *Nat. Commun.*, 2014, **5**, 1–7.

43. Y. Zhang, K. Su, R. Guo, S. Xu, D. Chen, J. Zhu, W. Bao, J. Zhang, J. Ning and Y. Hao, *Phys. Stat. Solidi RRL*, 2019, **13**, 1900167.

44. V. M. Kaganer, O. Brandt, A. Trampert and K. H. Ploog, *Phys. Rev. B*, 2008, **72**, 045423.

45. G. Li, Y. Y. Zhang, H. Guo, L. Huang, H. Lu, X. Lin, Y. L. Wang, S. Du and H. J. Gao, *Chem. Soc. Rev.*, 2018, **47**, 6073–6100.

46. Y. Qi, Y. Wang, Z. Pang, Z. Dou, T. Wei, P. Gao, S. Zhang, X. Xu, Z. Chang, B. Deng, S. Chen, Z. Chen, H. Ci, R. Wang, F. Zhao, J. Yan,

X. Yi, K. Liu, H. Peng, Z. Liu, L. Tong, J. Zhang, Y. Wei, J. Li and Z. Liu, J. Am. Chem. Soc., 2018, **140**, 11935–11941.

47. Y. Li, Y. Zhao, T. Wei, Z. Liu, R. Duan, Y. Wang, Z. Xiang, Q. Wu, J. Yan, X. Yi, G. Yuan, W. Junxi and J. Li, *Jpn. J. Appl. Phys. 2017*, **56**, 085506.

48. W. C. Ke, Z. Y. Liang, S. T. Tesfay, C. Y. Chiang, C. Y. Yang, K. J. Chang and J. C. Lin, *Appl. Surf. Sci.*, 2019, **494**, 644–650.

49. Y. Kim, S. S. Cruz, K. Lee, B. O. Alawode, C. Choi, Y. Song, J. M. Johnson, C. Heidelberger, W. Kong, S. Choi, K. Qiao, I. Almansouri, E. A. Fitzgerald, J. Kong, A. M. Kolpak, J. Hwang and J. Kim, *Nature*, 2017, **544**, 340–343.

50. B. H. Hwang, C. S. Chen, H. Y. Lu and T. C. Hsu, *Mater. Sci. Eng. A.*, 2002, **5**, 380–388.

51. K. P. Sharma, S. M. Shinde, M. S. Rosmi, S. Sharma, G. Kalita and M. Tanemura, *J. Mater. Sci.*, 2016, **51**, 7220–7228.

Figures



Fig 1 (a) SEM image and (b) Raman spectrum of a graphene film on a Cu tape substrate grown by CVD. In (a) a red arrow points to a grain boundary of Cu and yellow arrows point to graphene wrinkles. In (b) an intensity ratio of 2D-band peak to G-band peak (I2D/IG) is ~1.6 and no D-band peak is present, confirming that high-quality defect-free graphene with less than three layers is formed.



Fig 2 (a) XRD 2θ-ω scans and (b) FWHM values of AIN (0002) peak measured by XRD rocking curves (ω scans) for AIN layers on graphene/Cu tape substrates deposited at different temperatures.

8



Fig 3 (a) XRD  $\varphi$  scan of AIN {10<sup>1</sup>2} peaks, (b) color-coded orientation map and (c) (0001), (10<sup>1</sup>0), and (11<sup>2</sup>0) pole figures measured by EBSD, (d) SEM and (e) AFM (2.5×2.5  $\mu$ m<sup>2</sup> with z scale of 65 nm) images for an AIN film on a graphene/Cu tape substrate.



Fig 4 (a) XRD 20- $\omega$  scan of GaN on AlN/graphene/Cu foil, (b) XRD  $\varphi$  scan of GaN {10<sup>1</sup>2} peaks, (c) SEM and (d) AFM (4×4  $\mu$ m<sup>2</sup> with z scale of 80 nm) images for a GaN film on an AlN/graphene/Cu tape substrate. (e) A photo of a flexible GaN substrate consisting of a GaN template layer on Cu tape with graphene/AlN seed/buffer layers.



