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A large-scale deep-blue tetraphenylbenzene-bridged hybridized local and charge transfer fluorophore exhibiting small efficiency roll-off and low amplified spontaneous emission threshold†

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Efficient and stable deep-blue organic luminescent materials are important but rare for application in full-color displays due to their intrinsic wide bandgap characteristics. Here, a high-performance deep-blue hybridized local and charge transfer fluorophore **PTPC** with the aggregation enhanced emission (AEE) effect has been designed and synthesized with 1,2,4,5-tetraphenylbenzene as the bridge. High morphological stability, high photoluminescence efficiency, fast radiative rate, suitable energy levels, and bipolar-transporting properties are obtained simultaneously by incorporating a twisted molecular skeleton with weak, rigid and nearly flat donor/acceptor groups. As a result, a high forward-viewing external quantum efficiency of 6.78% as well as a deep-blue emission with the Commission Internationale de l'Éclairage coordinates of (0.156, 0.059) is achieved in a non-doped organic light-emitting diode (OLED). The EQEs still remain as high as 6.73% and 6.18% at 100 and 1000 cd m⁻², showing a small efficiency roll-off. Furthermore, **PTPC** exhibits extraordinary amplified spontaneous emissions (ASEs) with a low threshold of 1.18 μJ cm⁻² and a gain coefficient of 36.6 cm⁻¹. This work provides a new insight into creating high-performance deep-blue organic semiconductors with marked ASE performance for further practical applications.

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Introduction

With the rapid development of high and ultrahigh definition displays, highly efficient and stable blue organic optoelectronic materials have been a major requirement for applications in organic light-emitting diodes (OLEDs) and organic solid-state lasers (OSSLs).^{1–5} However, although extensive efforts have been devoted, owing to the intrinsic wide bandgap characteristics, high-performance and cost-effective deep-blue emitters

and devices meeting the requirement of Commission Internationale de l'Éclairage (CIE) y coordinate value < 0.06 are really rare.⁶ So far, blue phosphorescent and thermally activated delayed fluorescence (TADF) OLEDs have realized high external quantum efficiencies (EQEs) exceeding 25%,^{7–11} but the emitters have to be dispersed into an appropriate host matrix to alleviate the concentration quenching effect,^{12–14} which not only requires a stable host and adjacent layers with high triplet state energy (E_T) higher than that of the emitters, but is also related to host-guest doping technology, posing the phase separation risk and increasing the production costs.^{13,15} In addition, the phosphorescent and TADF OLEDs usually suffer from serious efficiency roll-off at high current density, resulting from exciton annihilation involving long-lived triplet excitons.^{16,17} Despite the ease of achieving deep-blue emitters based on triplet-triplet annihilation (TTA) and hybridized local and charge transfer (HLCT) mechanisms, achieving high efficiency OLEDs needs very delicate molecular and device engineering.^{18–22} Thus, continued effort is still urgent to pursue stable and efficient deep-blue pure organic light-emitting materials and non-doped OLEDs with CIEy < 0.06 for application in (ultra-) high definition displays.

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Designing an excellent deep-blue emitter should indeed meet several criteria with regard to color quality, efficiencies, stability, and charge-carrier injection/transportation:^{3,4,23–25} (1) satisfying the standard of CIE $y < 0.06$ with good emission color purity, which is essential to develop high-end electronic displays such as (ultra-) high definition televisions, (2) a high photoluminescence quantum yield (PLQY), generally, the higher the PLQY is, the better the electroluminescence (EL) performance of the device will be achieved, (3) high thermal and morphological stabilities, which is beneficial for extending the lifetime of OLEDs, and (4) suitable highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO) energy levels, which can ensure efficient hole and electron injection, transportation, and recombination within the emissive layer. To achieve deep-blue emission, the π -conjugation length of molecules must be limited, which in turn imposes restraint on the molecular size. Nevertheless, the limited molecular size is detrimental to forming morphologically stable and uniform amorphous films.⁴ On the other hand, deep-blue materials should have a high PLQY to realize a high EQE in the device. The PLQY at the aggregation-state can be enhanced by impeding intermolecular π - π packing interactions, and increasing the conjugated degree of the emitter that is dependent on the molecular rigidity and planarity. However, the enlargement of molecular planarity will lead to strong π - π packing and bathochromic shift of emission. Therefore, achieving a good trade-off among the necessary properties to develop high-performance deep-blue emitters and non-doped devices really represents a challenging scientific task to researchers.

1,2,4,5-Tetraphenylbenzene (1,2,4,5-TPB) featuring the typical aggregation enhanced emission (AEE) characteristic possesses moderate dihedral angles of 40.50–65.98°.²⁶ The rigid and twisted molecular structure makes it realize an efficient ultraviolet emission peak at ~ 363 nm in the aggregate state due to restriction of intramolecular motion. These properties are beneficial to construct pure blue emitters by linking appropriate electron-donating and accepting groups.^{27–29} However, it is difficult to build high-performance non-doped deep-blue OLEDs with CIE $y < 0.06$ and EQE $\geq 5\%$.³⁰ What's more, the efficiency roll-off at high luminance needs to be improved in the 1,2,4,5-TPB-based OLEDs. In this study, we designed and synthesized a large-scale donor- π -acceptor (D- π -A) type HLCT material **PTPC** with 1,2,4,5-TPB as the π -bridge linking weak, rigid, and bulky donor (D) and acceptor (A) units. This not only interrupts the π -conjugation and relieves the intramolecular charge transfer between D and A units by the 1,2,4,5-TPB spacer, realizing a deep-blue emission, but also increases the length and rigidity of the non-coplanar molecular skeleton and suppresses the rotation or vibration of groups, achieving a high PLQY and radiative rate in the pure film state. As a result, the non-doped device with the compound **PTPC** as the emitter realizes deep-blue light with the EL peak at 411 nm and a narrow full width at half maximum (FWHM) of 55 nm, corresponding to the CIE coordinates of (0.158, 0.059). It is worthwhile to note that a high forward-viewing EQE of 6.78% is achieved and the value remains as high as 6.18% at a practical

brightness of 1000 cd m⁻². Encouragingly, owing to the high PLQY and fast radiative rate ($S_1 \rightarrow S_0$ transition) exceeding 8×10^8 s⁻¹, **PTPC** exhibits excellent amplified spontaneous emission (ASE) characteristics with a low threshold of 1.18 μ J cm⁻².

Experimental methods

Synthesis and characterization of materials

Synthesis of 2',5'-dibromo-1,1':4',1''-terphenyl (TP-Br). Under an N₂ atmosphere, 1,4-dibromo-2,5-diiodobenzene (4.88 g, 10 mmol), phenylboronic acid (2.74 g, 22.5 mmol), and tetrakis(triphenylphosphine)palladium(0) (0.41 g, 0.35 mmol) were dissolved in 50 mL of toluene. And 2 M potassium carbonate (6.91 g, 50 mmol) in 25 mL of water and CH₃CH₂OH (25 mL) were added to the reaction mixture. Then the reaction mixture was stirred at 95 °C for 20 hours. After cooling to ambient temperature, the reaction mixture was extracted with dichloromethane and water and dried by anhydrous magnesium sulfate. The organic layer after filtration was evaporated under reduced pressure with a rotary evaporator. The crude product was further purified by silica-gel column chromatography using petroleum as the solvent and finally dried under vacuum to give a white solid (2.94 g, yield: 70%). ¹H NMR (400 MHz, CDCl₃, δ , ppm): 7.66 (s, 2H), 7.50–7.39 (m, 10H).

Synthesis of 9-(4'-bromo-[1,1'-biphenyl]-4-yl)-3,6-di-*tert*-butyl-9*H*-carbazole (tCZ-Br). 3,6-Bis(*tert*-butyl)-9*H*-carbazole (2.79 g, 10 mmol), 4-bromo-4'-iodobiphenyl (3.95, 11 mmol), copper(I) iodide (1.90 g, 10 mmol) and potassium carbonate (2.76 g, 20 mmol) were dissolved in 50 mL of *N,N*-dimethylformamide in a 250 mL 2-neck round bottom flask under a nitrogen atmosphere. The reaction mixture was refluxed for 12 hours with continuous stirring. Then the solvent was removed under reduced pressure. The reaction mixture was extracted with dichloromethane and water and dried using anhydrous magnesium sulfate. After filtration and evaporation under reduced pressure with a rotary evaporator, the crude product was further purified by silica-gel column chromatography using dichloromethane/petroleum (1/5, v/v) as the solvent and finally dried under vacuum to give a white solid (4.17 g, yield: 75%). ¹H NMR (400 MHz, CDCl₃, δ , ppm): 8.18 (s, 2H), 7.76 (d, $J = 8.4$ Hz, 2H), 7.67–7.60 (m, 4H), 7.58–7.52 (m, 2H), 7.49 (d, $J = 9.3$ Hz, 2H), 7.42 (d, $J = 8.6$ Hz, 2H), 1.50 (s, 18H).

Synthesis of 3,6-di-*tert*-butyl-9-(4'-(4,4,5,5-tetramethyl-1,3,2-dioxaborolan-2-yl)-[1,1'-biphenyl]-4-yl)-9*H*-carbazole (tCZ-B). Under an N₂ atmosphere, tCZ-Br (2.55 g, 5 mmol), bis(pinacolato)diboron (1.91 g, 7.5 mmol), [1,1'-bis(diphenylphosphino)ferrocene]dichloropalladium(II) Pd(dppf)Cl₂ (0.15 g, 0.2 mmol), KOAc (1.47 g, 15 mmol) and 50 mL of 1,4 dioxane were added into a 250 mL 2-neck round-bottom flask and heated at 110 °C for 15 hours. Then the reaction mixture was cooled down to room temperature, extracted with dichloromethane, and dried over anhydrous magnesium sulfate. The organic layer was filtered and purified by silica-gel column chromatography with dichloromethane/petroleum (1/3, v/v) as the solvent. Finally, the product was further dried under vacuum to give a white solid (3.83 g, yield: 90%). ¹H NMR

(500 MHz, CDCl_3 , δ , ppm): 8.16 (s, 2H), 7.96 (d, $J = 7.8$ Hz, 2H), 7.83 (d, $J = 7.6$ Hz, 2H), 7.72 (d, $J = 7.7$ Hz, 2H), 7.64 (d, $J = 8.2$ Hz, 2H), 7.48 (d, $J = 8.7$ Hz, 2H), 7.42 (d, $J = 8.6$ Hz, 2H), 1.49 (s, 18H), 1.40 (s, 12H).

Synthesis of 9-(5'-bromo-4'-phenyl-[1,1':2',1'':4'',1'''-quaterphenyl]-4''-yl)-3,6-di-*tert*-butyl-9H-carbazole (TPC-Br). The procedure was analogous to that described for TP-Br. As a result, the crude product was purified by silica-gel column chromatography using dichloromethane/petroleum (1/5, v/v) as the solvent and finally dried under vacuum to give a white solid (yield: 72%). ^1H NMR (500 MHz, CDCl_3 , δ , ppm): 8.19 (s, 2H), 7.82 (s, 1H), 7.79 (d, $J = 8.1$ Hz, 2H), 7.62 (d, $J = 7.5$ Hz, 2H), 7.57 (s, 4H), 7.52 (s, 5H), 7.45 (d, $J = 7.2$ Hz, 1H), 7.41 (d, $J = 8.6$ Hz, 2H), 7.33–7.26 (m, 7H), 1.50 (s, 18H).

Synthesis of 2-(4'''-(3,6-di-*tert*-butyl-9H-carbazol-9-yl)-2',5'-diphenyl-[1,1':4',1'':4'',1'''-quaterphenyl]-4-yl)-1-phenyl-1H-phenanthro[9,10-*d*]imidazole (PTPC). The procedure was analogous to that described for TP-Br. As a result, the crude product was purified by silica-gel column chromatography using dichloromethane/petroleum (2/1, v/v) as the solvent and finally dried under vacuum. Then the resulting compound was further purified by recrystallization with the mixed solvent ethanol and toluene to give a white solid (yield: 68%). ^1H NMR (500 MHz, CDCl_3 , δ , ppm): ^1H NMR (500 MHz, CDCl_3 , δ , ppm): 8.91 (s, 1H), 8.78 (d, $J = 8.4$ Hz, 1H), 8.72 (d, $J = 8.4$ Hz, 1H), 8.18 (s, 2H), 7.80 (d, $J = 8.3$ Hz, 2H), 7.76 (d, $J = 7.6$ Hz, 1H), 7.67 (t, $J = 7.7$ Hz, 1H), 7.65–7.45 (m, 17H), 7.41 (d, $J = 8.6$ Hz, 2H), 7.37 (d, $J = 8.2$ Hz, 2H), 7.33–7.26 (m, 10H), 7.23–7.18 (m, 3H), 1.50 (s, 18H). ^{13}C NMR (125 MHz, CDCl_3 , δ , ppm): 150.69, 142.88, 141.43, 140.85, 140.63, 140.10, 139.76, 139.34, 139.17, 139.08, 138.94, 138.67, 138.42, 137.34, 132.99, 130.43, 130.09, 129.90, 129.87, 129.73, 129.26, 129.07, 128.26, 128.13, 127.29, 127.11, 126.88, 126.79, 126.56, 126.25, 125.64, 124.89, 124.10, 123.61, 123.40, 123.10, 122.97, 120.84, 116.23, 109.23, 34.72, 32.01. TOF-MS (ESI) m/z calcd for $\text{C}_{77}\text{H}_{61}\text{N}_3$: 1028.3283; $[\text{M} + \text{H}]^+$ found: 1028.4933.

Results and discussion

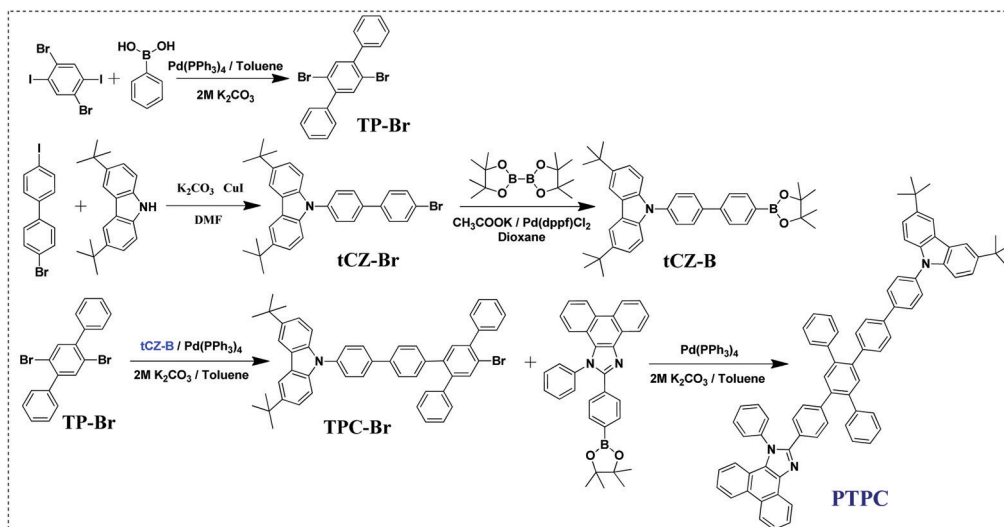
Molecular design strategy and synthesis

Donor–acceptor (D–A) type molecules consisting of electron-rich and electron-deficient moieties have been considered as an effective way to develop high-performance emitters due to their bipolar-transporting characteristics. However, these molecules generally create a strong intramolecular charge-transfer state, inducing a large red-shift of emission and a broad FWHM. To circumvent such a problem, on the one hand, rigid and bulky group phenanthrimidazole (PI) with good thermal stability was selected as the weak A unit, while rigid-flexible group 9-phenyl-3,6-di-*tert*-butyl-carbazole (tPCz) possessing a similar HOMO energy level with the PI group was used as the D unit (Fig. S1, ESI[†]). The introduction of *tert*-butyl at the 3,6-position of carbazole can not only suppress the generation of radical cations or/and anions upon electrochemical oxidation or/and reduction, stabilizing electrical operation of the device, but also decrease the intermolecular π – π interactions and improve the solubility of compound. On the other hand, introducing 1,2,4,5-TPB as the π -bridge linking D and A unit at *para*-positions can interrupt the π -conjugation of both donor and acceptor units. As a result, the rigid and twisted D– π –A type deep-blue emitter PTPC with a large non-coplanar molecular skeleton was successfully designed.

As illustrated in Scheme 1, the intermediate tCZ-Br was synthesized by the Ullmann reaction, while tCZ-Br was obtained by the Miyaura borylation reaction. The target compound PTPC was synthesized *via* three-step Suzuki cross-coupling reactions with tetrakis(triphenylphosphine)palladium as the catalyst. The chemical structures of the target compound were determined by the ^1H and ^{13}C NMR spectroscopies, and high-resolution mass spectrometry. The detailed synthetic procedures and characterization data were described in the Experimental methods.

Thermal properties and morphological stability

The thermal properties of PTPC were examined by using TGA and DSC measurements. As shown in Fig. S1a (ESI[†]), the



Scheme 1 Synthesis routes and molecular structures of intermediates and PTPC.

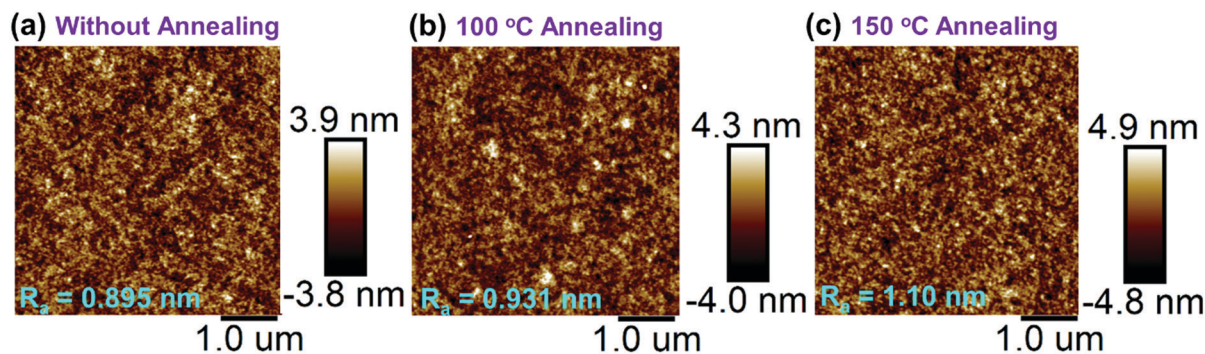


Fig. 1 AFM images of vacuum-evaporation films without annealing (a), and with annealing under 100 (b) and 150 °C (c) for 24 hours in air.

decomposition temperature (T_d) corresponding to 5% weight loss appeared at 520 °C, whereas the glass-transition temperature (T_g) was not detected in the temperature range from 20 to 300 °C, illustrating that the compound could form amorphous films in the wide temperature range. To verify this, we further investigated the morphological stability of 40 nm-thick vacuum-evaporation films of PTPC by annealing at different temperatures in the atmosphere. As displayed in Fig. 1 and Fig. S2b, c (ESI[†]), compared with the film without annealing, after annealing at 100 and 150 °C for 24 hours, the films not only show slight variations in the average surface roughness (R_a) values performed by atomic force microscopy (AFM), but also exhibit stable and unchanged emission spectra. This result proves that high-quality films with good morphological stability can be achieved based on the compound PTPC, which is conducive to the development of long-lived OLEDs.

Theoretical calculations

As shown in Fig. 2, geometry optimization in the ground state and FMO calculations were carried out *via* the DFT method at the RB3LYP/6-31G(d,p) level. The LUMO of the compound PTPC is mostly delocalized on 1,2,4,5-TPB and extended to

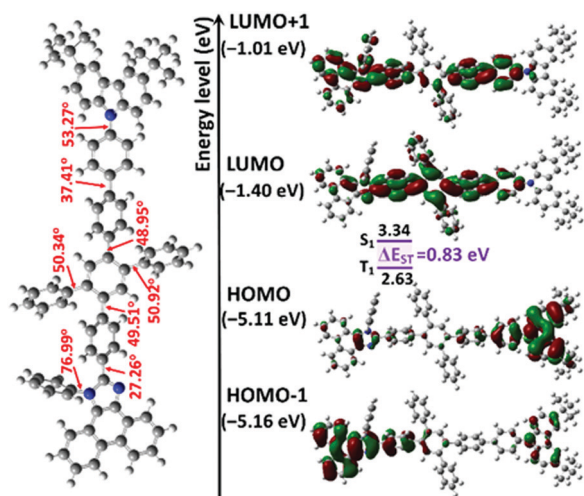


Fig. 2 The calculated optimized structures and HOMO/LUMO distributions of PTPC.

the adjacent phenyl ring and several segments in the PI unit, while the HOMO is mainly localized on the tPCz unit with a residual on the PI unit, and the HOMO-1 is mainly distributed on the PI unit the adjacent phenyl ring with a residual on the tPCz unit. This type of orbital feature containing spatial separation in conjunction with a certain degree of overlap is of benefit to the improvement of fluorescence intensity. Moreover, PTPC shows a large and highly twisted molecular skeleton with torsion angles of $\sim 48^\circ$ – 51° between the central phenyl group and external phenyl ring in the 1,2,4,5-TPB core, torsion angles of $27.26^\circ/37.41^\circ$ between 1,2,4,5-TPB and the PI/tPCz unit, and a large torsion angle of 53.27° between the phenyl ring and carbazole in the tPCz unit, respectively. The efficacious integration between the twisted molecular geometry and rigid and bulky chromophores limits the π -conjugation length, weakens the intramolecular charge transfer intensity, and overcomes the intermolecular stacking interactions in the solid-state, which can prevent aggregation-caused fluorescence quenching, and further achieve more efficient luminescence of the emitter for OLEDs. To further explore the nature of the excited state, NTO analyses were also performed (Fig. S3 and Table S1, ESI[†]). For the $S_0 \rightarrow S_1$ transition of PTPC, the hole is mostly distributed on the whole molecular backbone, while the particle is localized on 1,2,4,5-TPB and extended to the adjacent phenyl and imidazole rings, which clearly reflects the formation of the HLCT excited state. It should be noted that the $S_0 \rightarrow S_1$ excitations can mainly be attributed to the combination of HOMO and HOMO-1 to LUMO transitions (70% and 20%), which ensures the efficient overlap between hole and particle distributions, being beneficial for enhancing the PLQY, corresponding to a large oscillator strength of 1.3043. However, for the $S_0 \rightarrow T_1$ excitations, both holes and particles are mainly delocalized to phenanthrene and adjacent imidazole and biphenyl rings groups, showing an obvious locally excited (LE) state.

Photophysical properties

UV-visible absorption and PL spectra of PTPC were obtained to investigate its photophysical properties. As displayed in Fig. 3a and b, the absorption band around 260 and 294 nm can be considered as the π - π^* transition of 1,2,4,5-TPB and carbazole. And the absorption band at 305–355 nm can be assigned to the

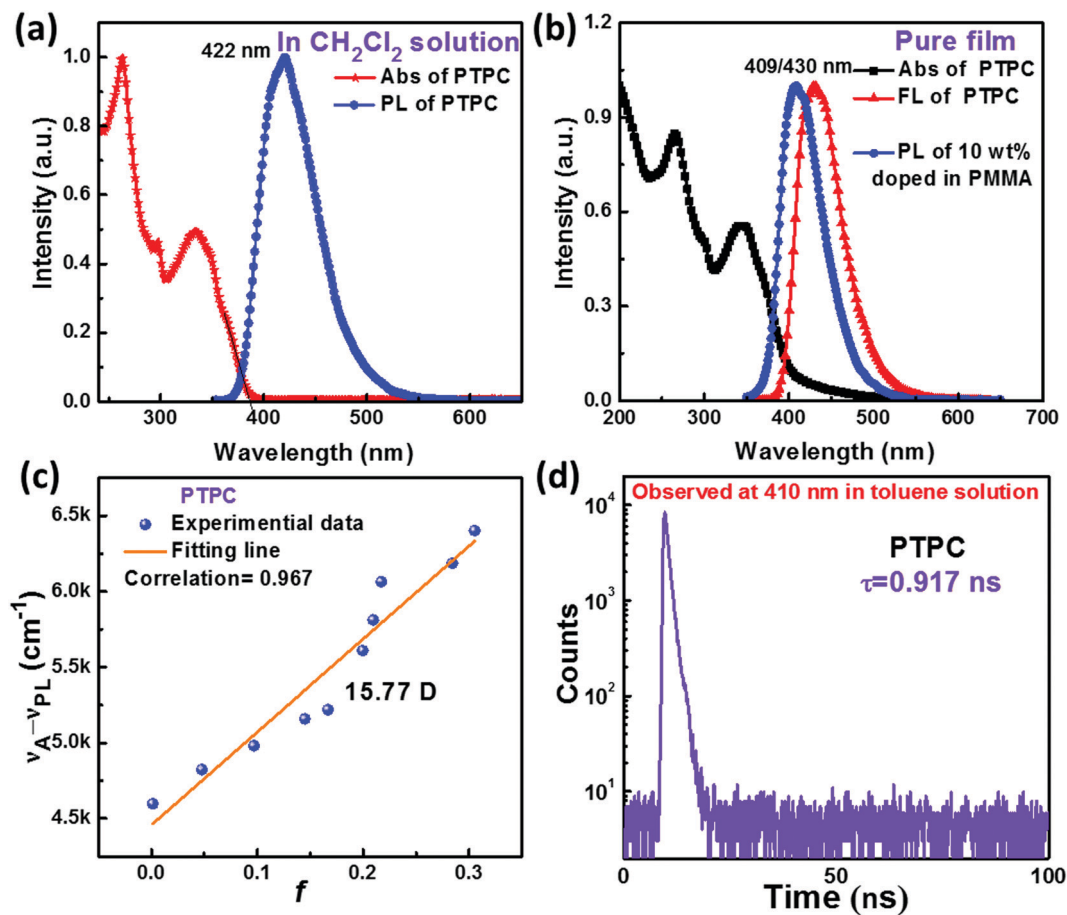


Fig. 3 (a) UV-Vis absorption and PL spectra of PTPC in CH₂Cl₂ solution (10 μM). (b) UV-Vis absorption and PL spectra of PTPC in neat films and the PL spectrum of 10 wt% PTPC doped into PMMA. (c) Solvatochromic Lippert–Mataga models of PTPC with the Stokes shift ($\nu_a - \nu_f$) versus orientational polarizability (f). (d) Transient PL decay curve in toluene solution at an excitation of 280 nm.

$\pi-\pi^*$ transition of the large and rigid molecular skeleton. The emission maxima of PTPC are at 422 nm in dilute CH₂Cl₂ solution (10 μM), while the emission in the neat vacuum-evaporation film (with 40 mm thickness) shows a slight red-shift with the peak at 430 nm due to the weak intermolecular interaction. In comparison, the PL spectrum of the doped film with 10 wt% PTPC in polymethylmethacrylate (PMMA) exhibits an obvious blue-shift with a peak at 409 nm. From the onset of UV-vis absorption in CH₂Cl₂ solution, the value of E_g was deduced to be 3.20 eV. Interestingly, PTPC exhibits the AEE feature confirmed by their PL spectra in THF/water mixtures with different water fractions (f_w) (Fig. S4, ESI[†]). When the $f_w \leq 60\%$, the reduction and red-shift of the PL spectra should be due to the increase of the solvent polarity of the THF/water mixture, which leads to the twisted intramolecular charge transfer process that quenches the emission. When the $f_w \geq 70\%$, owing to the worsened solvating ability of mixtures, the formation of aggregates can effectively activate the process of restriction of intramolecular motion, resulting in increased PL intensity. The reduction of PL intensity in the mixtures with $f_w \geq 90\%$ is probably caused by the formation of less emissive aggregates *via* random packing of

molecules. Such a phenomenon was reported frequently in D–A type molecules.^{26–29}

To further study the properties of the excited state, UV-Visible absorption and PL spectra of PTPC in various solvents are recorded and are shown in Fig. S5 (ESI[†]). When the polarity of the solvents increases (from hexane to acetonitrile), there is no obvious change of absorption spectra, while a remarkable solvatochromic effect with the red shifts of emission peaks from 396 to 426 nm is observed in the PL spectra. According to the Lippert–Mataga solvatochromic model, the Stokes shift ($\nu_a - \nu_f$) versus the solvent orientational polarizability (f) characteristic was estimated to analyze the properties of the S₁ state of PTPC. As shown in Table S2 (ESI[†]) and Fig. 3c, PTPC shows a linear relationship between $\nu_a - \nu_f$ and f with a slope of 6141 (Correlation = 0.967), corresponding to the dipole moment of S₁ of 15.77 D, demonstrating that there exists only one excited state in either low polar solvents or high polar ones. This character should be a quasi-equivalent hybridization HLCT state due to the strong interstate coupling of LE and CT (charge transfer) states,^{31–34} which coincides well with theoretical calculations. This can be further proved using the transient PL decay curves in different solvents where only nanosecond-scale

single-exponential decay lifetime (τ) exists. Owing to the quasi-equivalent hybridization between LE and CT states, the emitters in both toluene solution and neat films exhibit high PLQYs of 92.1% and 82.7%. According to the emission peaks of the low temperature fluorescence and phosphorescence spectra at 77 K (Fig. S8, ESI[†]), the S_1 and T_1 state energies can be calculated to be 3.04 and 2.48 eV, corresponding to the energy gap (ΔE_{ST}) of 0.56 eV. These results illustrate that the TADF mechanism is almost impossible for **PTPC**. As displayed in Fig. 3d, Fig. S7b and Table S3 (ESI[†]), the transient PL decay curves show single-exponential decay lifetimes of 0.917 and 0.988 ns, respectively, in the toluene solution and neat films. Based on the PLQY and decay lifetimes, the radiative and the non-radiative rate (k_r and k_{nr}) of **PTPC** can be calculated from the following equations: $k_r = \text{PLQY}/\tau$, and $k_{nr} = 1/\tau - k_r$. The k_r s of **PTPC** are calculated to be 10.04 and $8.37 \times 10^8 \text{ s}^{-1}$ in the toluene solution and neat film, which are higher than those of the most reported deep-blue emitters. The k_{nr} s are calculated to be 0.87 and $1.75 \times 10^8 \text{ s}^{-1}$.

Electrochemical and electrical properties

The HOMO energy levels of **PTPC** were obtained *via* the cyclic voltammetry (CV) measurements in anhydrous CH_2Cl_2 solution with 0.1 M tetra-*n*-butylammonium hexafluorophosphate. As displayed in Fig. S9 (ESI[†]), the onset oxidation potential of **PTPC** is 0.82 V, corresponding to a HOMO level of -5.41 eV . The LUMO energy level can be calculated to be -2.21 eV from the equation of $\text{LUMO} = (\text{HOMO} + E_g) \text{ eV}$. Meanwhile, single-carrier devices were fabricated to assess their charge carrier transporting properties. The structures of the hole-only device (HOD) and electron-only device (EOD) are indium tin oxide (ITO)/1,4,5,8,9,11-hexaazatriphenylene-hexacarbonitrile (HATCN, 20 nm)/**PTPC** (80 nm)/HATCN (20 nm)/Al (150 nm) and ITO/(1,3,5-tri(phenyl-2-benzimidazolyl)benzene) (TPBi, 10 nm)/**PTPC** (80 nm)/LiF (1 nm)/Al (150 nm), respectively. It can be seen from the current density–voltage plots that current densities of the HOD and EOD are significantly increased with the increase of applied voltages (Fig. S10, ESI[†]), illustrating that the material **PTPC** has relatively balanced bipolar transporting capacities. The reasonable HOMO/LUMO energy levels and bipolar transporting properties of **PTPC** can ensure effective carrier injection, transportation and recombination, which is quite important for developing high efficiency and low roll-off OLEDs.

EL performance of non-doped OLEDs

To further verify the HLCT properties and EL performance of **PTPC**, a non-doped device (**D1**) with the simple structure of ITO/HATCN (20 nm)/1,1-bis[4-*N,N*-di(4-tolyl)amino]phenyl]cyclohexane (TAPC, 50 nm)/4,4',4''-tri(*N*-carbazolyl)triphenylamine (TCTA, 5 nm)/**PTPC** (10 nm)/TPBi (40 nm)/LiF (1 nm)/Al (150 nm) was fabricated. Here, ITO was the anode, HATCN was utilized as the hole-injection layer. TPAC served as the hole-transporting layer. TCTA was used as the exciton-blocking layer. TPBi served as electron-transporting and hole-blocking layers. LiF and Al served as the electron-injection layer and cathode, respectively. The energy level alignment of the device and

chemical structure of the used organic materials are shown in Fig. S11 (ESI[†]). The device and EL performance of the device are shown in Fig. 4 and are summarized in Table S4 (ESI[†]). The device **D1** exhibits a low turn-on voltage (V_{on}) of 3.1 V, corresponding to the energy of S_1 state of **PTPC**, which illustrates the highly efficient injection, transportation and recombination of holes and electrons. The luminance of 1000 and 7118 (maximum) cd m^{-2} are obtained at applied voltages of 4.7 and 7.2 V. Significantly, the device **D1** not only realizes high forward-viewing efficiencies of 6.78% (at 54.3 cd m^{-2}), 2.66 cd A^{-1} , and 2.60 lm W^{-1} for EQE, current efficiency, and power efficiency without using light out-coupling enhancement techniques, but also shows a stable deep-blue emission with a peak at 411 nm, FWHM of 55 nm and CIE coordinates of (0.156, 0.059), which is close to the standard blue CIE color coordinates of (0.15, 0.06), defined by the International Telecommunication Union (ITU). Significantly, the EQEs still remain as high as 6.73% and 6.18% at the luminance of 100 and 1000 cd m^{-2} , exhibiting a very low efficiency roll-off. As mentioned above, the possibility of the TADF process can be ruled out due to the nanosecond-scale single-exponential decay lifetime and large ΔE_{ST} , while the good linear correlation between luminance and current density implies that the TTA process is not the main emission mechanism for the device **D1** (Fig. S11c, ESI[†]). As seen in Fig. S3 and S12 (ESI[†]), the excited states of S_1 , S_2 , T_7 and T_8 show HLCT transition character, which ensures the existence of the high-lying reverse intersystem crossing (hRISC) process. Furthermore, a large ΔE_{ST} of 0.711 eV and small energy differences of 0.0015 and 0.0506 eV between S_1 and T_7 as well as S_2 and T_8 can effectively enhance the hRISC rate, making more singlet excitons be generated *via* two hRISC channels. It is notable that such high efficiency achieved at high luminance is one of the highest values in non-doped OLEDs with $\text{CIE}_y \leq 0.07$ (Table 1).

In addition, the non-doped devices (**D2**, **D3** and **D4**) were fabricated by increasing the thickness of the emissive layer and changing the electron-transporting layer (Fig. S13, ESI[†]). It can be seen from Fig. 5 and Fig. S14 (ESI[†]) that the devices not only emit stable deep-blue light, but also show high EQEs of 6.34%, 6.71% and 6.67% for **D2**, **D3** and **D4**, respectively. These results signify that the deep-blue emitter **PTPC** has overriding characteristics of reproducibility and performance stability. Achieving such high EL performance can be attributed to the high PLQY, appropriate HOMO/LUMO energy levels, bipolar transporting properties, and HLCT excited state.

ASE properties

As mentioned above, due to the twisted, rigid and bulky molecular backbone of **PTPC**, which leads to a large spatial overlap between the HOMO and LUMO energy levels and efficiently weakens the intermolecular π - π stacking interaction and k_{nr} , realizing excellent stability (T_d : 520 °C), a high PLQY (Film: 82.7%) and a large radiative rate from S_1 to S_0 transition (Film: $8.37 \times 10^8 \text{ s}^{-1}$). Encouraged by these results, we further

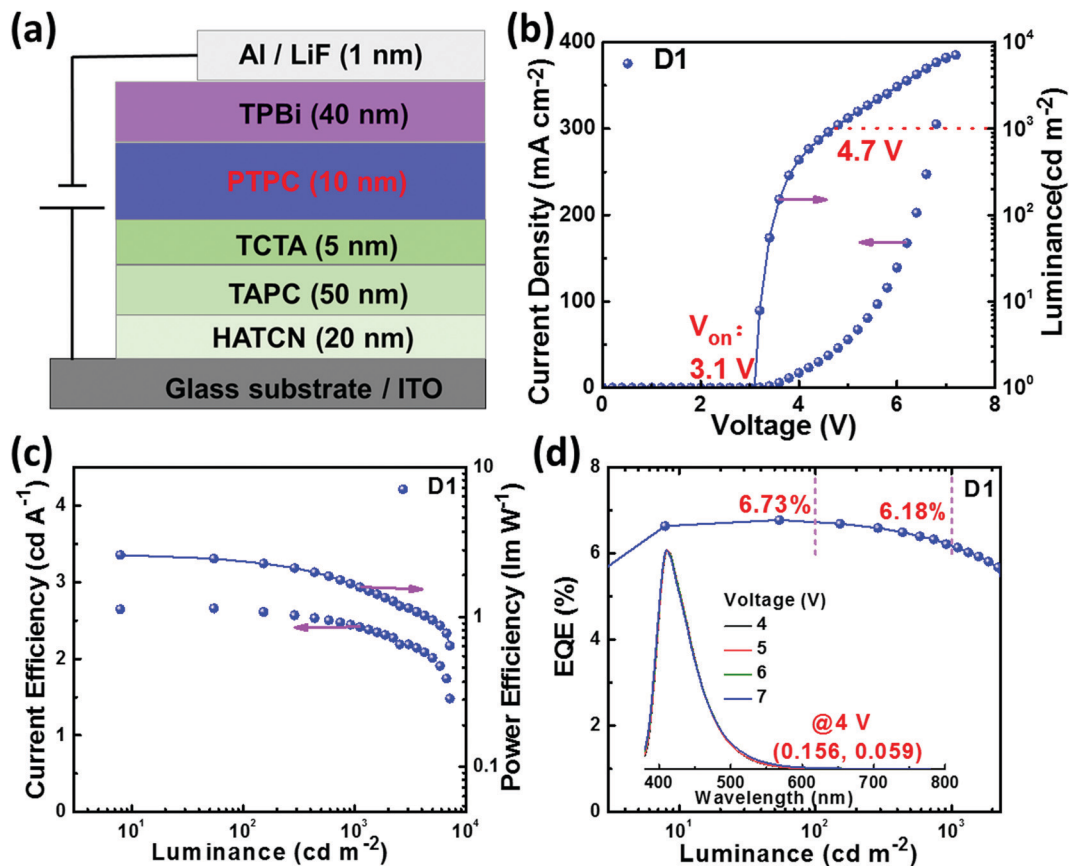


Fig. 4 The device structure diagram and EL performance of **D1**. (a) Device structure diagram. (b) Current density and luminance versus voltage curves. (c) Current efficiency and power efficiency versus luminance curves. (d) EQE versus luminance curves. The inset shows the normalized EL spectra at different voltages.

Table 1 The summary of reported non-doped OLEDs with CIE_y ≤ 0.07

Ref.	EQE (%)		CIE(x,y)
	EQE _{max} ^a	EQE ₁₀₀₀ ^b	
This work	6.78	6.18	(0.157, 0.059)
Ref. 1 ³⁵	6.4	5.58	(0.151, 0.066)
Ref. 2 ³⁶	5.08	4.44	(0.16, 0.06)
Ref. 3 ³⁷	7.1	5.3	(0.16, 0.06)
Ref. 4 ³⁸	6.57	—	(0.17, 0.07)
Ref. 5 ³⁹	1.97	—	(0.16, 0.06)
Ref. 6 ⁴⁰	8.9	—	(0.150, 0.060)
Ref. 7 ⁴¹	3.02	—	(0.167, 0.056)
Ref. 8 ⁴²	3.38	—	(0.154, 0.063)
Ref. 9 ⁴³	5.29	—	(0.155, 0.058)
Ref. 10 ²¹	5.74	4.8	(0.152, 0.054)
Ref. 11 ⁴⁴	4.6	4.04	(0.154, 0.058)
Ref. 12 ⁴⁵	3.19	—	(0.156, 0.054)
Ref. 13 ⁴⁶	6.33	6.32	(0.151, 0.066)
Ref. 14 ³⁰	4.78	—	(0.159, 0.060)
Ref. 15 ⁴⁷	6.73	4.01	(0.156, 0.055)

^a and ^b represent the maximum EQE and the value at 1000 cd m⁻².

determined the ASE properties of **PTPC** with the hope to obtain a substantial threshold. As shown in Fig. 6a, b and Fig. S15–S19 (ESI[†]), with the gradual increase of excitation power, the emission of all of the studied films (2, 4, 6, 8, 10 and 100 wt%-**PTPC** blended in mCP host) exhibited spectral narrowing.

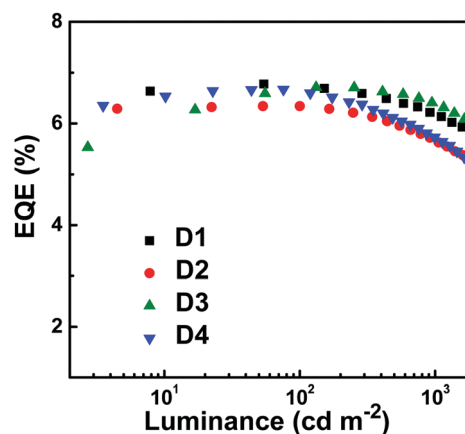


Fig. 5 EQE versus luminance curves of the deep-blue devices (**D1**, **D2**, **D3** and **D4**).

The ASE effect can be well-evidenced by a clear change in the intensity slope with an increasing excitation energy density. As shown in Fig. S20a (ESI[†]), the ASE peaks can be adjustable from 419 to 431 nm by changing the doping concentration of **PTPC**. The optimized result was achieved by the sample with a doping concentration of 8 wt%, where the FWHM reduces from 41.0 to

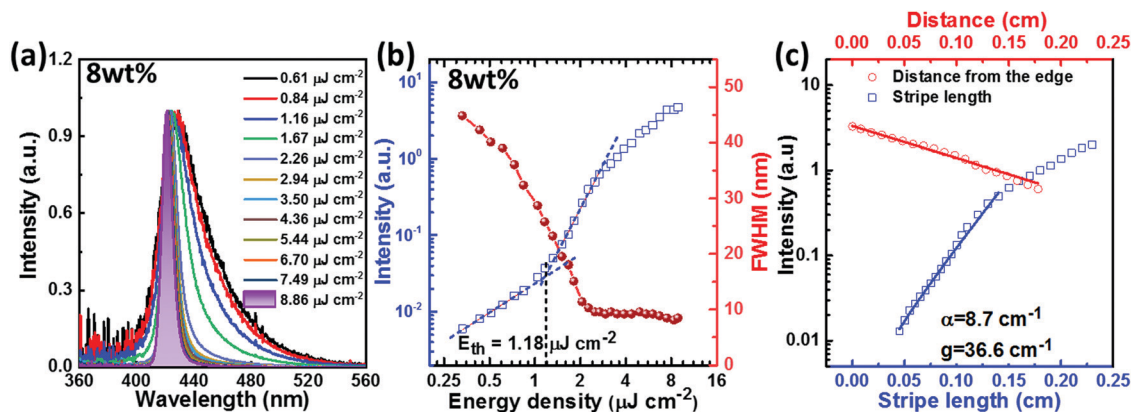


Fig. 6 (a) PL spectra of 8 wt% PTPC doped in mCP films at different excitation powers below and above the ASE threshold. (b) PL peak intensity and FWHM as a function of the excitation energy density of 8 wt% PTPC doped in mCP films. (c) Emission intensity variation with respect to both the stripe beam length and the distance from the sample edge.

8.1 nm when the excitation power was increased from 0.61 to 8.86 $\mu\text{J cm}^{-2}$. The ASE threshold was determined to be 1.18 $\mu\text{J cm}^{-2}$ from the change in the slope of PL peak intensity versus excitation energy density curve (Fig. 6b), which is much lower than that of the common D-A molecules.^{48–52} According to the dependence of the PL peak intensity on the stripe length, a gain coefficient (g) of 36.6 cm^{-1} can be estimated. And the loss coefficient (α) of 8.7 cm^{-1} is obtained by fitting the PL peak intensity as a function of the distance between the sample edge and the excitation region.

Most notably, the neat film not only shows a low threshold of 2.13 $\mu\text{J cm}^{-2}$, but also exhibits more excellent photo-stability (with the 7% reduction) than the 8 wt%-PTPC blended mCP film (with the 82% reduction) after continuously pumping 10 000 pulses at an energy density of 8 $\mu\text{J cm}^{-2}$ (Fig. S20b, ESI[†]), which should be due to the stable, rigid and bulky D- π -A type molecular skeleton. These fantastic results illustrate that the HLCT emitter PTPC with 1,2,4,5-TPB as the π -bridge should be a promising candidate for developing highly efficient and stable deep-blue OLEDs and OSSLS. Based on this molecular system, suppressing the internal conversion process from higher triplet states (T_n) to T_1 by increasing the T_n - T_1 energy gap could be an effective avenue to further enhance the EL performance.

Conclusions

In summary, a high-performance large-size deep-blue bipolar HLCT emitter PTPC exhibiting AEE characteristic is successfully designed and characterized by incorporating PI and tPCz units into the 1,2,4,5-TPB core. Combining twisted and rigid molecular backbones with bulky weak D/A groups can effectively limit the π -conjugation length and suppress the aggregation quenching problem, guaranteeing good stability, a high PLQY and a fast radiative rate. Thus, it could not only achieve deep-blue light with an EL peak of 411 nm and a high EQE of 6.78% in a non-doped OLED with PTPC as the emitter, but also exhibit excellent ASE performance with a low threshold of

1.18 $\mu\text{J cm}^{-2}$, which is one of the best values among those of the reported deep-blue D-A type materials. Such results provide a rational design strategy to construct efficient and stable deep-blue optoelectronic materials for future applications in OLEDs and OSSLS.

Conflicts of interest

There are no conflicts to declare.

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