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On the fundamentals of organic mixed ionic/ electronic conductors

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The first Telluride Science meeting (formerly TSRC) on organic mixed ionic and electronic conductors (OMIECs), Oct 3–7, 2022, brought together researchers across the field to understand the fundamental processes and identify out-standing questions related to this exciting class of materials. OMIECs are organic materials that promote the transport of mobile electronic charge carriers while simultaneously supporting ionic transport and ionic–electronic coupling. These properties open up broad areas of applications from energy to bioelectronics. Devices include batteries, supercapacitors, actuators, electrochromic displays, and organic electrochemical transistors (OECTs). They possess the key strengths of traditional organic electronic materials, such as synthetic tunability and low-temperature processing. Despite the recent advances in devices and applications achieved with such materials, many challenges and gaps in understanding remain. These topics hold the key to designing next-generation materials and devices that continue to push the limits of performance and stability and facilitate novel functionality. This perspective aims to summarize the current understanding, conversations, and debates that made this TSRC particularly engaging, enabling new directions and searching for missing pieces of the OMIEC puzzle.

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The materials – search for model systems

As in other fields within materials science, it is important to identify model systems that will enable accessible and repeatable fundamental studies. At their most basic level, OMIECs support electronic transport through an electronic transport pathway (*i.e.*, along and between conjugated polymer backbones or between nearby redox centers), and ionic transport

occurs through free volume or an ion transport pathway. The ions are needed to ensure charge neutrality when the conjugated or redox system is doped (chemically or electrochemically).

The opportunity and the challenge with organic systems is that such an arrangement can be achieved through many routes: blends, composites, single component systems with varied side chains, block co-polymers, *etc.* A single model system has eluded the field as a whole, and different questions demand different model systems. Poly(3-hexylthiophene)

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(P3HT) is the most studied system in this area. Many fundamental studies have focused on electrochemistry in organic electrolytes, such as TBAPF₆ in acetonitrile. Controlled electrochemical doping led to highly conducting films.¹ P3HT, in combination with salts (*e.g.*, LiClO₄) and poly(ethyleneoxide) (PEO), finds an audience in energy applications or other areas where dry mixed conduction is sought.²

Poly(ethylenedioxythiophene):poly(styrenesulfonate) (PEDOT:PSS) is the most studied OMIEC, a combination of a doped conducting polymer and a polyelectrolyte, and can be regarded as a polyelectrolyte complex. PEDOT:PSS with aqueous electrolytes, including NaCl or KCl, has been well studied owing to their commercial availability and ubiquity in the community of researchers interested in aqueous operation or the role of hydration, *e.g.*, for OECTs or actuators. More recently, conjugated backbones with ionic side groups (so-called conjugated polyelectrolytes) and oligoether side chains fill yet another need and allow operation in aqueous electrolytes. Additional materials strategies from side-chain free polymers and radical polymers are also gaining popularity.

Even so, single component systems are still largely represented by a few chemical structures. While polythiophene backbones are the p-type work-horse of the community, naphthalene diimide (NDI) acceptor units combined with thiophene donors can be regarded as the n-type model system. The electrochemistry of this polymer (*e.g.*, NDI-T) is quite complex but points to a backbone with localized charges on the NDI units.³ Another n-type mixed conductor, the ladder-type poly(benzimidazobenzophenanthroline) (BBL), stands as one of the few side-chain free mixed conductors and has shown high performance in OECTs. Pure p-type redox polymers bear,

for example, triarylamine or carbazole redox units.⁴ On the small molecules front, the films of oligoether tethered fullerenes,⁵ thiophene-flanked diketopyrrolopyrrole units,⁶ polycyclic systems consisting of naphthalene bis-isatin dimer core with rhodanine terminals,⁷ or thiophene-EDOT based cores⁸ are among the structures that can be electrochemically doped and de-doped in aqueous media, while they are yet to be commonly used by different research teams. All these systems can be modified with additives, dopants, and through processing and post-processing.

No one system can answer all the key questions, and no one system is completely ideal (*i.e.*, using PEDOT:PSS in commercial formulations means some additives may not be known to the researcher, and small batch novel synthesized polymers can have variable levels of impurities, defects, and distributions of molecular weights). Nevertheless, some parallels and generalities can be gleaned, putting these diverse materials on a broader spectrum of understanding that benefits the community at large.

The operating environment dictates properties

The composition, structure, and properties of OMIECs are complex and interrelated. Compositionally, all species play a critical role. Structurally, OMIECs often form semi- or *para*-crystalline domains in an amorphous matrix. Multicomponent OMIECs can undergo phase separation, such as the segregation into conjugated polymer-rich and polyelectrolyte-rich domains



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in PEDOT:PSS.^{9,10} Ions and additives affect the degree of crystallinity, coherence length within the crystalline domains, phase separation, and phase purity, amongst other structural parameters.¹¹ Harsh processing techniques can enhance crystallinity and increase the ratio of conjugated polymer to polyelectrolyte in multicomponent OMIECs.¹²

OMIECs have applications in a variety of operating environments. The permeable nature of OMIECs makes them sensitive to operating environments that often include contact with, or immersion in, electrolytes when operated in electrochemical devices. Here again, all species play a critical role, and the role of ion and solvent transport between the OMIEC and its environment must be considered. Contacting electrolytes (e.g., solid polymer, aqueous, organic, or ionic liquid) interact with OMIECs in different ways and to different degrees, depending on characteristics of the solvent (polarity, miscibility) and the ion (size, solvation, polarizability).

The passive swelling (just due to contact with the electrolyte, no electrochemical potential applied) and active swelling (driven by electrochemical potential) of OMIECs must be considered to understand the morphology of the film during operation. During passive swelling, ions and solvent penetrate and swell the OMIEC. The choice of water *versus* acetonitrile electrolytes leads to different degrees of swelling, e.g., determined by *in situ* spectroscopic ellipsometry coupled with cyclic voltammetry.^{13,14} The degree of swelling strongly depends on the miscibility of the electrolyte and the OMIEC, which is often determined by the OMIEC side chain chemistry (oligoether, alkyl, tethered ion) and architecture (length, branching).^{15–18} The miscibility of solvent and ions in the OMIEC can vary, resulting in differential solvent and ion uptake.¹⁹ Further, solvated ions may shed or retain solvation shells when injected into OMIEC. The degree of swelling may differ between amorphous and crystalline domains²⁰ or between conjugated polymer-rich and polyelectrolyte-rich phases.²¹ Ion and solvent uptake into the crystalline domains expands and modifies the crystal structure, in some cases enhancing or disrupting crystallinity dependent on both the OMIEC, electrolyte, and electronic charge density.^{20,22–27}

If the as-prepared OMIEC contains mobile ions that differ from those in the contacting electrolyte, ion exchange will occur, such as ion-proton exchange in PEDOT:PSS.²¹ In some systems, larger concentrations of fixed ionic (polyelectrolyte) charge in the OMIEC introduces Donnan exclusion, limiting ion uptake to oppositely charged counterions.²¹ In such cases, di- and poly-valent counterions can chelate with the fixed charge groups, becoming trapped and inducing ionic crosslinks.²⁸ But in most high-performing OMIECs, Donnan exclusion is suppressed or absent, and cation and anion transport, along with solvent transport, contribute to OMIEC swelling.^{21,29}

A change in the electrochemical potential of the OMIEC can actively modulate how much it swells. The electrochemical oxidation or reduction of the OMIEC modulates the electronic carrier and the counterbalancing ion (dopant) density. This accumulated charge, in turn, can modulate optoelectronic properties – which is core to electroactive functionality of

OMIECs in a variety of applications, e.g. electrochemical displays but also electrochemically switchable nanoantennas as metasurfaces.³⁰ Dopant modulation can occur by both ion uptake and expulsion.^{14,29} This modulation of electrostatics also changes the miscibility of the OMIEC and electrolyte surroundings, bringing in or expelling solvent and charge-balanced populations of anions and cations, beyond just the dopant transport.²¹

The degree of active swelling can vary greatly depending on the electrolyte choice and the OMIEC chemical structure. Microscopically, electrochemical cycling can induce crystallite expansion and contraction.^{20,24,25} Extreme active swelling can lead to dissolution/delamination,³¹ but can also be leveraged in reversible swelling/gel formation.³² The effects of active swelling on composition, structure, and properties are often not fully reversible.³³

Additionally, under operating conditions, impurities can play a large role. Dissolved oxygen can limit the electrochemical stability window and accelerate OMIEC degradation.^{34,35} Further, OMIECs with shallow HOMO levels can be spontaneously doped (oxidized).³⁶ Residual catalysts within the OMIEC can determine its electronic conductivity.³⁷ The contacting electrodes also affect the OMIEC, with electrode material (along with oxygen and potential) dictating the potential stability window and degradation kinetics of the OMIEC.³⁴

All this demands the use of multimodal *in situ* characterization in various operating conditions to capture and interrogate the complexities. Structure, composition, and properties must be mapped across relevant environments and potentials. This requires simultaneous (or parallel) probing of electronic (stored charge, electrical conductivity), compositional (individual anions, cations, and solvent), structural (micro-, mesoscale, macroscopic), spectral, and mechanical properties of OMIECs in application-relevant electrolyte environments under controlled potentials. Some popular *in situ* characterization tools are electrochemical UV-VIS-IR absorption spectroscopy, quartz crystal microbalance with dissipation monitoring (QCM-D), X-ray characterization, X-ray photoelectron spectroscopy (XPS), spectroscopic ellipsometry, and scan-probe microscopies.³⁸

Where are the ions in the structure?

A large effort has been made to understand where ions reside within the volume of OMIECs during electrochemical doping/dedoping. However, getting a holistic picture of ions has proven challenging due to the heterogeneous microstructure of OMIECs and the experimental challenges of working with (especially aqueous) electrolytes *in situ*. To understand ion arrangement in OMIECs, the driving force for ion intercalation has to be considered, favourable microstructures have to be found that feature efficient ionic–electronic coupling, and finally, it has to be understood how ions are distributed at device length scales where electric fields can drive their motion.

The ability for ions to penetrate OMIECs is highly dependent on the mixing enthalpy.³⁹ It has been shown that polar



sidechains are one powerful way to facilitate bulk ion penetration when using atomic ions (*e.g.*, Na⁺, Cl⁻, *etc.*).⁴⁰ Volumetric doping may also be achieved by modifying the charge distribution on ionic species.⁴¹ There is also evidence to suggest that charge neutral ion pairs can diffuse into OMIECs⁴² and ions can remain in OMIEC films following doping/dedoping cycles^{21,29} indicating mixing between ions and OMIECs is thermodynamically favorable.

We also note that the number of charge-compensated ion-electron/hole pairs scales with the applied potential.^{43,44} This indicates that the doping level is thermodynamically tied to the electric potential, where the energy for changing the doping level may be tied to the entropy for different carrier densities⁴⁵ voltage-dependent enthalpy of mixing between ions and OMIECs,^{39,46} and/or the chemical/quantum capacitance of the conjugated polymer.^{47,48}

Clarifying where ions reside within the microstructure of OMIECs is essential to understand the coupling between ionic and electronic charges. From an energetic standpoint, we would expect electrons/holes to first populate the lower energy sites of the microstructure within the ordered regions.⁴⁹ This is also supported by electrochemical cyclic voltammetry of different amorphous and semicrystalline P3HT films (regioregular *vs.* regio-random, as-cast *versus* well-crystallized).^{1,50} However, during chemical doping, reactants must diffuse through the amorphous region of the film, leading to doping/dedoping of the amorphous region first.^{51,52} Similarly, ion conduction in OMIECs requires free volume for ions, which are often hydrated, to move. Thus, there is likely to be a balance between doping of the lowest energy sites first *vs.* doping the sites accessible to ions.

There are many reports stating that the degree of swelling during doping/dedoping is largest in the amorphous parts of the microstructure.^{53,54} However, at higher doping densities, ions can disrupt the crystals, leading to reversible^{26,51} and irreversible^{51,55,56} phase changes. In some cases, ion pathways form within the microstructure *via* self-assembly, minimizing changes to the microstructure during doping/dedoping.²²

One must also consider the rearrangement of ions across device length scales. While typical models for OECT operation use C^* as the governing quantity to solve for carrier densities,⁵⁷ this model ignores the ion drift under the electric field. Ions are charged and will also move under the same applied potential driving hole current across the OECT channel. Recently, lateral ion drift along the channel has been taken into account to accurately capture experimentally measured potential drops along the channel.^{58,59} Similarly, ionic and electronic drift and diffusion have been modeled in PEDOT:PSS by treating the OMIEC as a composite of a pure electrolyte phase and pure electronic phase with a local electric field at their interface.⁶⁰ While this model works well for PEDOT:PSS, it might be challenging to extend to homopolymer OMIECs which may not have a distinct electrolyte phase.

The role of swelling and ion motion

Another common thread through many of our discussions was the role of swelling in OMIECs. There is broad agreement that

swelling is the factor that distinguishes OMIEC applications from dry organic semiconductor applications. Even then, there is no broad consensus on how swelling should be defined. The literature distinguishes “passive” (swelling at the open circuit potential) from “active” (applied bias-induced swelling) swelling. But many newer, low threshold voltage materials are partially doped in ambient so that demarcation becomes ambiguous. Unfortunately, quantifying swelling at the level needed to understand what is occurring on the microscale is quite difficult. Bulk swelling is typically measured *via* electrochemical quartz crystal microbalance with dissipation monitoring (EQCM-D). EQCM-D provides great insight into total mass, but mass does not distinguish exactly which species are moving. It could be a cation, anion, or solvent moving or some complex combination of the three that give rise to the total mass change.²⁹ For this reason, there is a need for more tools that identify which species are moving under what conditions. Ion motion can be studied indirectly through moving front experiments, though this may not be appropriate for systems in which doping occurs, also, *via* ion expulsion. X-ray fluorescence experiments have also provided valuable insights,²¹ yet there remains a need for operando studies or adoption of ambient XPS techniques to shed light on ionic composition during operation.

Even once the identity of mobile species is resolved, real questions remain about where the swelling occurs. There is a common observation that bulk (QCM-D) swelling is larger than the observed swelling of the semicrystalline lattice.²⁵ This leads to speculation that most swelling occurs in the amorphous regions, which could disrupt the percolative electronic network.²⁰ New emerging techniques could start proving or disproving this hypothesis.⁶¹ However, that level of swelling, which can range from 20 to hundreds of percent, is hard to understand from a mechanical standpoint. When a film is swelling more than 100%, a persistent question is what holds the film (often decorated with highly hydrophilic side chains) together? We discussed ionic *vs.* covalent cross-linking, hydrogen bonding, or even electrostatic effects of the polaron, among other options, and agree that this issue requires further investigation.

Controlling mixed conduction

Molecular design trends

The transition from organic field effect transistors (OFETs) to OECTs saw the utilization of many similar materials, whereby the need to facilitate ion transport in water was often met by replacing alkyl side chains with oligoether ones or with the addition of ionic units.^{62–65} However, emerging studies are showing that mixed conduction can be achieved without side chains altogether.^{66,67} So, the logical question arises: are side chains (ion-free or bearing) necessary?

It has been shown that side-chain engineering modulates the mixed conduction capability of polymer films, typically portrayed as a matter of balancing ionic and electronic



ions are needed for conductivity modulation, thermal annealing has so far been considered detrimental, as shown by Flagg, *et al.* in a p-type polymer. By improving crystallinity, the resistance of the active layer to ion penetration is increased. This leads to a lower OECT performance by decreasing mobility.⁶³ However, a recent report showed that operation with a large anion (PF₆ or TFSI) is possible after annealing the film up to 300 °C.⁸⁵ The positive or negative effects of annealing differ from one backbone to another, and seem to be directly tied to their glass transition and melting temperature, rather than to the nature of the side chains themselves.

Many devices comprise OMIEC thin films (on a substrate) with a thickness of typically less than 100 nanometers, which ensures that ions can rapidly reach the complete volume of the material, essential for, *e.g.*, achieving a high switching speed. For some devices such as fiber based OECTs, the material (*e.g.* highly aligned PEDOT:PSS microfibers⁸⁶) may act as both the OMIEC as well as the *de facto* substrate and must hence feature adequate mechanical robustness. The OMIEC material then tends to have dimensions of at least several micrometers, which considerably reduces the speed of ion ingress, but could be mitigated through the use of porous materials.⁸⁷

Emergent phenomena owing to high charge densities

OMIECs display a finite potential window of high electrical conductivity, with the conductance decreasing dramatically beyond a critical gate voltage-controlled charge carrier density. While this phenomenon is reversible and not due to chemical or electrochemical degradation of the OMIEC's π -system, the precise cause is not fully understood yet.

Measurements of *in situ* conductance coupled with cyclic voltammetry show bell-shaped conductance profiles for redox polymer systems.⁴ For P3HT in TBAPF₆/acetonitrile, for example, a plateau-like conductance is found over a large electrochemical potential range of 0.4 up to 0.9 V *vs.* Fc/Fc⁺ where polaron and bipolaron species are formed.¹ A possible explanation for the behavior may be associated with band filling,^{88–90} while others have proposed that energetic/structural disorder at high carrier densities leads to decreased charge carrier mobility.^{1,91–93} For a further discussion on plateau-like conductivity and the decrease at very high potentials, we refer to electrochemical literature.⁹⁴

In BBL, the maximum electrical conductivity coincides with a change in the polarity of the Seebeck coefficient as the carrier density approaches 1 electron/repeating unit.⁹⁵ This is explained by the formation of multiply charged species and the opening of a hard Coulomb gap around the Fermi energy caused by double occupation and Coulomb interactions occurring at high doping levels. It was also found that the soft Coulomb gap previously proposed to explain the drop in electrical conductivity is of minor importance,⁹⁶ particularly for a rigid polymer like BBL.

Recent work has also shown that the decrease in conductance at high charge carrier density is associated with electron–

cation interactions occurring at the electrolyte and organic semiconductor interface that alter the potential energy landscape in which electronic charges move.⁹⁷ In fullerene (C₆₀) single crystals gated through ionic liquids, the degree of conductance suppression is considerably affected by the size of the ions. Small ions support activated transport and a complete conductance suppression at around 1 electron per C₆₀, while larger ions decrease the extent of suppression and allow exploring transport up to 1.7 s per C₆₀.

Although the decrease in electrical conductivity may pose challenges for conventional logic gates and sensors, it also leads to a unique transistor current response with a Gaussian shape. This response can simulate the activation and deactivation of sodium channels present in biological neurons, creating conductance-based organic electrochemical neurons that accurately replicate crucial biological neural features.⁹⁸ Therefore, these neurons can serve as event-based neuromorphic sensing and processing components, displaying exceptional biorealistic capabilities.

Modeling and simulation

Both quantum chemistry and molecular dynamics (MD) simulations have been used to interrogate the transport mechanisms and structure–function relationships of OMIECs. Recent work⁶¹ shows how grazing-incidence resonant X-ray diffraction (GIRXRD) complemented by periodic DFT calculations allows the determination of the mean position of the ion relative to the polymer backbone in the crystalline regions, which somewhat surprisingly is found near the lamellae midpoint (and not close to the polymer backbone). Multiscale methods combining MD and Kinetic Monte Carlo (KMC) have been instrumental as a tool to describe charge transport and percolation in conjugated materials and will continue to be key in OMIEC research. A multiscale transport model based on information from the structural and electrical characterization of PEDOT:PSS blends was also recently reported to explain the exponential-like conductivity scaling at different PEDOT:PSS ratios.⁹⁹ Charge transport from one PEDOT-rich grain to the next happens thanks to isolated/dispersed PEDOT chains in a PSS-rich matrix, with PSS-conjugated groups enhancing tunneling efficiency.

A coarse-grained multiscale model was used to systematically investigate the effects of side chain polarity in a series of random OMIEC copolymers.¹⁹ The highest charge mobility (and doping) is associated with the highest fraction of polar side chains. The results where a similar fraction of polar and apolar side-chains were mixed point to a significant drop in charge mobility due to a shorter polaron delocalization length arising from the increased torsional disorder along the polymer backbone.

Multiscale models also shed light on time-dependent ion-charge interactions in mixed-conducting materials. Recent work¹⁰⁰ has incorporated DFT calculations and MD simulations to capture the rapid charge fluctuation on the polymer backbone due to the presence of solvent and ions. While this



approach does not consider charge transfer across polymer chains, it captures both ion-induced instantaneous polarization and slow (~ 1 ns) charge reorganization due to conformational changes in polymer chains. Interactions across polymer chains dominate charge dynamics, which are expected to play a major role in bulk polymer aggregates or crystallites. While promising, this study highlights how representing the interaction between ionic and electronic charge carriers remains the most challenging task for any computational model of OMIECs. Although some methodological barriers still exist for reliably reproducing amorphous and semi-crystalline morphologies and ion-polaron coupling, ample opportunities exist for using already mature methodologies to study other structure-function and morphology-function relationships in OMIECs. Such efforts are made urgent by the extreme complexity and potential variety of OMIEC materials (ionic species, polymer backbone and connectivity, side chain polarity, length, and frequency). Despite the number of unknowns in the field, we believe a systematic OMIEC optimization effort guided by computational screening will soon be within reach as the community gradually uncovers relationships and correlations.

Spanning different systems

Despite some generalities across systems and common tools to study them, different systems and applications are governed by very different phenomena. For example, electronic transport is mechanistically different in non-conjugated *versus* conjugated redox polymers. Specifically for non-conjugated redox polymers, the electronic charge is localized on pendant groups where electrons can hop from site to site.¹⁰¹ In contrast, conjugated polymers distribute charge through the polymer's pi-conjugated backbone.

For electron conduction in non-conjugated redox polymers and small radical molecules, it is important to consider the proximity of the redox sites. As the site-to-site distance decreases, the conductivity increases. For example, a 1000-fold difference in the solid-state conductivity occurs when changing the hopping distance from 5.7 to 11.6 Å in certain crystals of small radical molecules.¹⁰² For non-conjugated redox polymers, the flexibility of the polymer chain must be considered in addition to redox site proximity.¹⁰³ As the chain's flexibility increases, the diffusivity of the polymer increases, and there are more opportunities for redox sites to come within closer proximity through thermal motion. For example, the highly flexible and low- T_g radical-containing polymer poly(4-glycididloxy-2,2,6,6-tetramethylpiperidine-1-oxyl) exhibited the highest solid-state conductivity within its class of 28 S m⁻¹. In the case of carbazole redox polymers (polyvinylcarbazole), conductivities as high as 2×10^{-2} S cm⁻¹ could be obtained upon chemical doping.⁴ In this case, the simultaneous cross-linking and doping of the starting materials into bicarbazole redox dimers could be proven by a combination of UV-vis-NIR spectroscopy and *in situ* and *ex situ* conductivity measurements.

In contrast to dry conditions, non-conjugated redox polymers exhibit mixed ion–electron–solvent transport in electrochemical

cells. The charge transport is affected by polymer–solvent and polymer–ion interactions, swelling, and the ion's solvation shell.^{104–106} For example, EQCM-D experiments revealed that both anions and cations could participate in the polymer's redox behavior, which should be considered in future designs.¹⁰⁴ Further, electrolyte choice can yield 1000 \times changes in the charge storage of non-conjugated redox polymers,¹⁰⁵ indicating a rich future design space for electrolytes.

The combination of ac-impedance spectroscopy and dc measurements of mixed conductors (PEDOT:PSS and conjugated polyelectrolytes) can help to decouple electronic and ionic charge transport as a function of relative humidity and corresponding water uptake.^{15,107} The influence of relative humidity was particularly prominent for the ion conductivity, which reached 10^{-2} S cm⁻¹ in the wet state. The electronic conductivity is less sensitive to humidity but requires the existence of electronic percolation pathways throughout the films, which is directly interrelated with the morphology upon exposure to humidity. It is interesting to note that the humidity. Response of the polyelectrolyte phases could be used to tune mechanical properties and make actuators.¹⁰⁸

When translating OMIECs into dry systems, the advantages conferred by solvent-assisted ion transport are compromised, and instead, the segmental mobility of the ion-conducting chain becomes a determining factor in the intrinsic limitations of ion transport. In the case of side-chain-based OMIECs, it becomes essential to carefully consider the structure of the side chains, ensuring optimal ion solvation and promoting fast dynamics to facilitate ion motion. Recent joint computational and experimental studies have provided valuable insights into the intricacies of Li-salt-doped OMIECs based on oligoethylene glycol-substituted thiophene derivatives.^{109–112} These reports on dry OMIECs highlight the challenge of effectively separating the dynamics of the side chains responsible for ion transport from the more rigid conjugated core imparting semicrystallinity for efficient electronic transport.

The rate at which an OMIEC transports ions is affected by whether or not the ions are contained within a solvent shell, as is typically the case for PEDOT:PSS and oligoether side chain containing polymers in aqueous electrolytes.^{56,63,113,114} In other cases, for example, for P3HT with polymer electrolytes, the ion moves without a solvent shell, where solubility interactions dominate.¹¹⁵ It is typically, but not always, observed that a solvent shell improves ion transport. Nevertheless, a single perspective unifying the two extremes and providing a predictive structure–property relationship remains challenging.

Conclusions and perspectives

While our focus has been on conjugated polymers and organic small molecules, it is clear that the library of OMIECs continues to expand by the day. Small molecules, for example, have found an important place in precision molecular design and the extreme case of poly-crystalline morphologies. MOFs, COFs, and porous conjugated polymers, while their integration into



devices presents barriers currently, afford three-dimensional frameworks where porosity can be controlled at the molecular scale, and chemical functionality can be finely dialed in. As for electrolytes, ionic liquids, as first investigated in electrolyte-gated transistors, are now numerous, and biocompatible versions with various ionic units may introduce new characteristics to the same OMIEC devices. Broadening the library of mixed conductors beyond polymers – perhaps even introducing hybrid systems – will thus open new applications for OMIEC devices.

The quest for model systems is useful but may be futile given the breadth of conditions and target applications under the OMIEC tent. Nevertheless, they prove useful for fundamental studies. The OMIEC must be considered a complete system, including the extrinsic or intrinsic ions, additives, and solvents, as appropriate. In addition, given the volatility of the solvents or the extrinsic nature of some electrolytes, multi-modal and operando characterization is needed to make direct device-material connections. Many differences in molecular design and film microstructures are not considered or well controlled, adding further challenges. Moving forward, questions remain about the exact location, interactions, and compositions (ionic) within OMIEC materials and how to model them properly. Can such understanding unlock designs for ideal materials? More stable formulations? Or allow for less common and highly non-linear responses toward novel applications? It's likely, but the community needs to tackle these questions collaboratively to get there.

Conflicts of interest

There are no conflicts to declare.

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