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Atomic defects in monolayer ordered double transition metal carbide ($\text{Mo}_2\text{TiC}_2\text{T}_x$) MXene and CO_2 adsorption†

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Transition metal carbides (MXenes) with formula $\text{M}_{n+1}\text{C}_n\text{T}_x$ ($n = 2$ and 3) have been emerging as a new family of two-dimensional (2D) materials that have great potential in electronic applications and CO_2 conversion catalysts. It has been already found that the electronic and electrochemical properties of $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes can be tuned by replacing the two outer titanium layers with molybdenum layers. Similar to other 2D materials, intrinsic defects can be formed in the synthesized MXene flakes and the formation of defects can influence the performance of these materials. Herein, we systematically study the effect of the different types of structural defects on the structural stability, electronic behavior, and electrochemical properties of ordered $\text{Mo}_2\text{TiC}_2\text{T}_x$ terminated with the specific surface functional groups of fluorine, oxygen, and hydroxide. The calculated defect formation energies imply that the formation of defects is dependent on the surface terminations, where the O-terminated MXenes demand more energy than the F- and OH-terminated MXenes. We found that defect formation is more feasible in the outer molybdenum layers than in the inner titanium layer. Our results predicted that the CO_2 molecule adsorbs on the defective surfaces through a spontaneous and exothermic process that is critical to its capture, while the perfect surface weakly attracts the molecule through a nonspontaneous and endothermic process. Thus, our study predicts that the electronic and electrochemical properties of $\text{Mo}_2\text{TiC}_2\text{T}_x$ can be tuned by forming specific defects and these MXenes could be promising materials for CO_2 adsorption and conversion.

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1 Introduction

MXenes are 2D materials that have shown great promise in energy storage applications,¹ including lithium-ion batteries^{2,3} and supercapacitors,⁴ and catalysis.⁵ Quasi-two-dimensional (2D) $\text{Mo}_2\text{TiC}_2\text{T}_x$ is an ordered double transition metal MXene, synthesized in 2015,⁶ in which the outer titanium layers of $\text{Ti}_3\text{C}_2\text{T}_x$ (the first MXene discovered in 2011⁷) are replaced with molybdenum layers. In this MXene, the titanium layer is sandwiched by the molybdenum layers and carbon atoms occupy the octahedral sites between these two transition metals, as shown in Fig. 1a. It was shown that the electrochemical and electronic properties of the ordered $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXene are different from those of $\text{Ti}_3\text{C}_2\text{T}_x$ MXene.^{6,8,9} While $\text{Ti}_3\text{C}_2\text{T}_x$ is a metallic-like conductor,^{8–10} it was confirmed

that $\text{Mo}_2\text{TiC}_2\text{T}_x$ exhibits semiconductor-like transport behavior.^{6,10} It was predicted that the ordered $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXene terminated by the F group would be more semiconducting than that terminated by the OH group, while –O termination would provide a metallic behavior.¹⁰

Defect formation was observed in 2D materials, such as graphene,^{14,20,21} BN,^{22–24} and MoS_2 .^{13,25} It has been well documented that the presence of defects affects the electrical, electrochemical, electronic, optoelectronic, and mechanical properties of these 2D materials.^{13,14,20,21,23} Intrinsic defects such as atomic vacancies and vacancy clusters were also observed using scanning tunneling microscopy (STM) images in synthesized MXene flakes,^{16,17} which the synthesis conditions might affect the defect concentrations of MXene flakes.^{16,18,19} Therefore, the formation of different kinds of defects is unavoidable in MXenes. Similar to other 2D materials, the presence of intrinsic defects might alter the electronic and electrochemical properties of monolayer $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes.

Structural defects in 2D materials could be unfavorable. For example, the presence of defects leads to weakening of the mechanical strength^{21,26} and the electronic performance of graphene.²⁷ However, certain types of defects could be favorable and provide applicable functionalities by adjusting the local properties.

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Fig. 1 Fully relaxed structural models of different types of defects considered in monolayer ordered $\text{Mo}_2\text{TiC}_2\text{T}_x$. (a) Top and side view of the perfect MXene. Top views of (b) V_{T_x} , (c) V_{Mo} , (d) V_{Mo_2} , (e) V_{T_x2} , (f) V_{MoTi} , and (g) $\text{V}_{\text{Mo}_2\text{Ti}}$.

For example, the mechanical properties of graphene could be also increased by controlled defect creation²¹ or the properties of MoS_2 could be tailored by controlled defect formation.¹³ Compared to these 2D materials, while intrinsic defects were previously observed in MXene flakes,^{16,17} the influence of defects on the electronic and electrochemical properties of MXenes has not extensively been investigated. From the rare investigations,^{11,12,16} it was confirmed that defects have unfamiliar effects on the MXene properties compared to other 2D materials depending on the type of defects, reconstruction of the structure, and local chemical bonding.^{11,12,16} It was predicted that titanium vacancies in $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes could not alter the metallic conductivity considerably; however, they affect the surface morphology and terminations.¹⁶ The effects of carbon vacancies in Ti_2CT_x indicated the improvement of electronic conductivity and the enhancement of mechanical flexibility.¹¹ It was also investigated that the intrinsic defects in Ti_2XT_x ($\text{X} = \text{C}, \text{N}$) MXenes could change the electronic properties from metallic to semiconducting or from semiconducting to metallic depending on the type of defect and the system.¹²

Previous studies indicate that various MXenes with different transition metals and functional groups might exhibit quite different behavior toward defects because of the comparatively complex structures and variety of chemical bonding between atoms. Despite tuning the electronic and electrochemical properties of $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes by manipulating the outer transition

metal layers,⁶ which are important for electronics, energy storage, sensors, and other applications, the presence of defects might significantly change the properties of ordered $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes. Therefore, a detailed atomistic understanding of the influence of different types of defects on the electronic and electrochemical properties of ordered $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes is important to explore their real-device applications, in which understanding these properties is crucial.

In this study, we performed first-principles calculations based on density functional theory (DFT) to explore the effects of different types of defects on the structural stability, electronic behavior, and electrochemical properties of ordered $\text{Mo}_2\text{TiC}_2\text{T}_x$ ($\text{T} = \text{O}, \text{F}, \text{and OH}$). The formation of defects on the top and bottom molybdenum sublayers is more feasible because of the direct contact with the etchant solution; however, the presence of defects in the titanium layer is also possible, which are inherited from the MAX phase. Moreover, cluster-defects were observed within the same sublayer in $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes.¹⁶ Fig. 1 shows the optimized defect structures of the seven different types of defects considered in this study, including a single T_x vacancy (V_{T_x}) and two adjacent T_x vacancies (V_{T_x2}) on both sides of the flake, a single molybdenum vacancy (V_{Mo}), two adjacent molybdenum vacancies (V_{Mo_2}) within the same sublayer, a vacancy complex of a single molybdenum vacancy and one nearby titanium vacancy (V_{MoTi}), a vacancy complex of two adjacent molybdenum

vacancies forming within the same sublayer and one nearby titanium vacancy ($V_{\text{Mo}_2\text{Ti}}$), and two molybdenum vacancies forming within two different sublayers (V_{MoMo}). The top view of V_{MoMo} is the same as V_{Mo} and is not shown in Fig. 1. The structural stability of all types of defects is investigated through their formation energies. The electronic properties of these defected $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes are also analyzed to highlight the influence of these effects on the conductance of MXenes terminated with different functional groups. The effects of these defects on the electrochemical properties of MXenes are also studied by evaluating the CO_2 adsorption performance on $\text{Mo}_2\text{TiC}_2\text{O}_2$.

2 Computational methods

First-principles calculations based on density functional theory (DFT) are performed using the Vienna ab initio Simulation Package (VASP).²⁸ We employed the generalized gradient approximation with the Perdew–Burke–Ernzerhof (PBE)²⁹ functional to describe the exchange and correlation of electrons. The projected augmented-wave (PAW)³⁰ potential is applied to take the interaction between the ion-core and valence electrons into consideration and a plane wave basis set with a sufficient energy cutoff of 520 eV is employed to consider valence electrons. A $4 \times 4 \times 1$ supercell is used to model 2D $\text{Mo}_2\text{TiC}_2\text{T}_x$ ($T = \text{F}, \text{O}, \text{and OH}$) MXene monolayers to provide at least a 10 Å distance between the point defects. A large vacuum of at least 20 Å is used to avoid any undesirable interaction between an MXene sheet and its periodically repeated images along the c -axis. The Brillouin zone is sampled with $5 \times 5 \times 1$ Monkhorst–Pack³¹ k -points for structural optimizations and with a $10 \times 10 \times 1$ mesh for electronic structure analysis. The atomic positions and lattice constants are fully optimized using the conjugate gradient method. The geometries of the MXenes are relaxed until the forces acting on each atom become less than $0.001 \text{ eV \AA}^{-1}$. The criterion for energy convergence is 10^{-6} eV per cell. The partial occupancies are determined using the Methfessel–Paxton smearing scheme with a smearing width of 0.1 eV.³² We employed the spin-dependent GGA plus Hubbard U (GGA + U)³³ methodology to apply the correlation effects in transition metals of Mo and Ti with a value of $U = 4 \text{ eV}$, which has been applied for $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes¹⁰ and extensively applied for their oxides.^{34,35} We performed spin-polarized DFT calculations with the antiferromagnetic (AFM) configurations predicted by Anasori *et al.*¹⁰ for the perfect $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes chemically terminated by different terminations. We also studied the probable effect of the defects on the AFM configurations of the MXenes. However, our results show that the considered defects cannot change the AFM configurations predicted for the perfect MXenes.

We calculated the formation energies of defects as $E_f = E_{\text{MXene_defect}} - E_{\text{MXene_perfect}} + \sum_i n_i \mu_i$, where $E_{\text{MXene_defect}}$ and $E_{\text{MXene_perfect}}$ are the total energies of the MXene supercell with specific defects and related perfect structure without defects, respectively, n_i is the number of removed atoms, and μ_i is the chemical potential of the removed atom(s). Since it is not straightforward to obtain the exact value of μ_i , which is

dependent on pressure and temperature,³⁶ we calculated E_f with the chemical potential between -2.0 and 0.0 eV .

The adsorption energy of the CO_2 molecule is calculated as $E_{\text{ads}} = E_{\text{MXene+mol}} - (E_{\text{MXene}} + E_{\text{mol}})$ by taking van der Waals (vdW) interactions into account by employing the dispersion-corrected DFT-D2 scheme,³⁷ which is essential for the accurate description of the interaction between CO_2 and a surface.^{38–40} $E_{\text{MXene+mol}}$ is the total energy of the adsorbate–substrate system; E_{MXene} is the energy of the MXene monolayer; and E_{mol} is the energy of the isolated CO_2 molecule. We investigated the CO_2 adsorption on the MXene by placing the molecule close to different sites of the monolayer parallel and perpendicular in different possible directions to find the largest adsorption energy. Negative (positive) values of the adsorption energy indicate an exothermic (endothermic) process.

3 Results and discussion

3.1 Formation energy of defects

We calculated the formation energy of different types of defects considered (see Fig. 1) in both non-terminated and terminated ordered MXenes to better understanding the energetic costs of the formation of defects. We first evaluated these energies for non-terminated Mo_2TiC_2 to gain more information about the feasibility of forming defects during the first step of exfoliation when the flakes are not terminated. Then, those energies are presented for terminated MXenes with specific functional groups. Since a defect formation energy is expected to be positive, it costs (consumes) energy to make a defect.

Our results for the non-terminated MXene (Fig. 2a) indicate that the formation energies of V_{Mo} are lower than the other types of defects, varying from 6.65 to 8.65 eV with the chemical potential between -2.0 and 0.0 eV . The formation of V_{Mo_2} consumes approximately the same amount of energy as V_{MoMo} and almost twice as much energy as V_{Mo} , which indicates a linear relationship between the formation energy of defects and the number of V_{Mo} . This linear relationship was also found for titanium vacancy clusters with up to six titanium vacancies from the outer sublayer of the Ti_3C_2 MXene.¹⁶ Moreover, V_{MoTi} and $V_{\text{Mo}_2\text{Ti}}$ defect formation is more energy consuming than that of other types of defects. The results show that defect formation in the titanium sublayer is energetically more difficult than within the outer molybdenum sublayers. Our obtained results can be confirmed by the study of Sang *et al.*,¹⁶ where they calculated the formation energy of a single titanium vacancy from the outer (2.85 eV) and inner (6.49 eV) sublayer of a $4 \times 5 \times 1$ supercell of the Ti_3C_2 MXene.

The effect of the surface functions on the calculated defect formation energies (Fig. 2b–d) indicates relatively the same behavior as determined for the pristine Mo_2TiC_2 MXene. The formation of V_{Mo} , V_{Mo_2} , V_{MoTi} , $V_{\text{Mo}_2\text{Ti}}$, and V_{MoMo} on the MXene terminated by oxygen (Fig. 2c) is more energy consuming than that for the pristine non-terminated MXene (Fig. 2a), while the energetic costs of these defects on the MXenes terminated by fluorine and hydroxide (Fig. 2b and d) are comparable with



Fig. 2 Formation energies of all types of defects as a function of the chemical potential of the removed atom(s) in the range of $-2 \text{ eV} < \mu < 0 \text{ eV}$ in (a) Mo_2TiC_2 , (b) $\text{Mo}_2\text{TiC}_2\text{F}_2$, (c) $\text{Mo}_2\text{TiC}_2\text{O}_2$, (d) and $\text{Mo}_2\text{TiC}_2(\text{OH})_2$.

those on the pristine MXene (Fig. 2a). This mainly originates from the effect of the surface functions on the bonding strength of the MXenes. It was comprehensively predicted that the bond strength between the atoms in the MXenes terminated by F and OH is relatively similar since the degree of covalency in both surface functions is the same (one electron), but the binding between atoms is stiffer in the MXenes functionalized by O atoms than those terminated by F or OH.⁸ This is mainly because O atoms share two electrons to form a covalent bond with the transition metals. Therefore, the F and O atoms make effectively single and double bonds with the surface, respectively, and that is why O atoms make stronger bonds with the surface. F and OH groups receive similarly one electron from the surface and that is why the bond strength of M–F and M–OH bonds is almost of the same order. The higher formation energy of a titanium vacancy on $\text{Ti}_{n+1}\text{C}_n\text{T}_x$ ($n = 1, 2$) MXenes terminated by oxygen than that for pristine and –F and –OH terminated MXenes was also predicted.^{12,16} Moreover, the calculated formation energies demonstrate that fluorine and hydroxide defect formation are comparable and are much more feasible than oxygen defect formation. It is worth mentioning that the formation of functional group atoms is considered on both sides of the flake because of the direct contact of the flakes with the etchant solution.

To further explore the effect of the replacement of the outer titanium layers in $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes by molybdenum layers on the defect formation energies, we compared our calculated formation energies of V_{Mo} on a $4 \times 4 \times 1$ supercell of $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes with the calculated formation energies of V_{Ti} from the outer titanium sublayers of a $4 \times 5 \times 1$ supercell of $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes.¹⁶ We scaled linearly the calculated formation energies of V_{Ti} to provide the same vacancy concentration considered in our study since a linear relationship between the formation energies of V_{Ti} clusters and the number of V_{Ti} was predicted. The scaled formation energies of V_{Ti} in the pristine Ti_3C_2 MXene ($\text{Ti}_3\text{C}_2\text{-}V_{\text{Ti}}$: 3.56 eV) and terminated $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes ($\text{Ti}_3\text{C}_2\text{F}_2\text{-}V_{\text{Ti}}$: 4.53, $\text{Ti}_3\text{C}_2\text{O}_2\text{-}V_{\text{Ti}}$: 9.68, and $\text{Ti}_3\text{C}_2(\text{OH})_2\text{-}V_{\text{Ti}}$: 4.34) indicate that the replacement of the outer titanium layers in $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes by molybdenum layers makes the formation of defects more difficult in the non-terminated Mo_2TiC_2 MXene and MXenes terminated by fluorine and hydroxide; however, the formation of defects in the MXene terminated by oxygen functional groups is comparable. This is because of the fact that Mo–Mo bonds are stronger than Ti–Ti bonds.⁴¹ Thus, based on our study, the energetic costs of the formation of defects become higher by manipulating the outer titanium layers in $\text{Ti}_3\text{C}_2\text{T}_x$ MXenes by molybdenum layers

because of the intrinsic bond strengths in these two MXene systems and V_{Mo} is the most probable defect to form during the synthesis of MXene monolayers.

3.2 Electronic properties

Before evaluating the electronic properties of defected $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes, we present the electronic properties of perfect $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes with different terminations. The total densities of states (DOS) of the perfect MXenes are displayed in Fig. 3. Similar to all predicted pristine mono-M element MXenes (M is a transition metal), which are metallic,⁸ the results indicate that the pristine ordered Mo_2TiC_2 MXene is also metallic (Fig. 3a). It was predicted that some of the mono-M element MXenes with surface terminations could exhibit semiconducting behavior depending upon the participating components and surface terminations.⁸ Here, we consider a specific surface termination (*e.g.*, pure -F, -O or -OH) and our results show that ordered $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes terminated by F and OH become semiconducting (Fig. 3b and d) and the band-gap of the OH-terminated MXene is smaller than that of the F-terminated MXene, while the O-terminated MXene exhibits metallic behavior (Fig. 3c). Our results are in reasonable agreement with an earlier study by Anasori *et al.*,¹⁰ where they also investigated the electronic structure of MXenes with mixed terminations since it was observed from EDX analysis that the synthesized surface has a mixture of -F, -O, and -OH terminations. Their results indicated that the synthesized $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes have a

semiconducting behavior, except for the MXene terminated by high concentrations of oxygen.

The effect of different types of defects on the electronic properties of non-terminated and terminated ordered MXenes is evaluated. Our calculated results of the total DOS of the non-terminated MXene (Fig. S1, ESI†) indicate that the defective Mo_2TiC_2 MXene is still metallic. Fig. 4 presents the DOS of the different types of defected MXenes under investigation, while for all different surface terminated MXenes (F, O and OH), it is very clear that the V_{T_x} vacancy makes these MXenes again metallic as we can observe states due to vacancies crossing the Fermi energy levels in these structures (Fig. 4a, e and i). For V_{MO} and V_{MoMo} vacancies, we do not observe any significant changes in the electronic DOS of the structures; for F and OH terminations, the MXenes are semiconducting, while for O terminations the structure retains its metallic nature. From the other defected MXenes, we can conclude that there are very minor changes in the electronic states, making the structures more semiconducting due to shifting of the bands away from the Fermi energy level.

3.3 CO_2 adsorption on $\text{Mo}_2\text{TiC}_2\text{O}_2$ defects

The interactions of different molecules with defective structures such as TiC_2 ,¹⁵ TiO_2 ,⁴² graphene and silicene,⁴³ BCN,⁴⁴ and MoS_2 ^{45–49} demonstrated that the molecular adsorption on the surfaces could be enhanced in the presence of defects. Defects act as electron trap sites and materials with defects have been

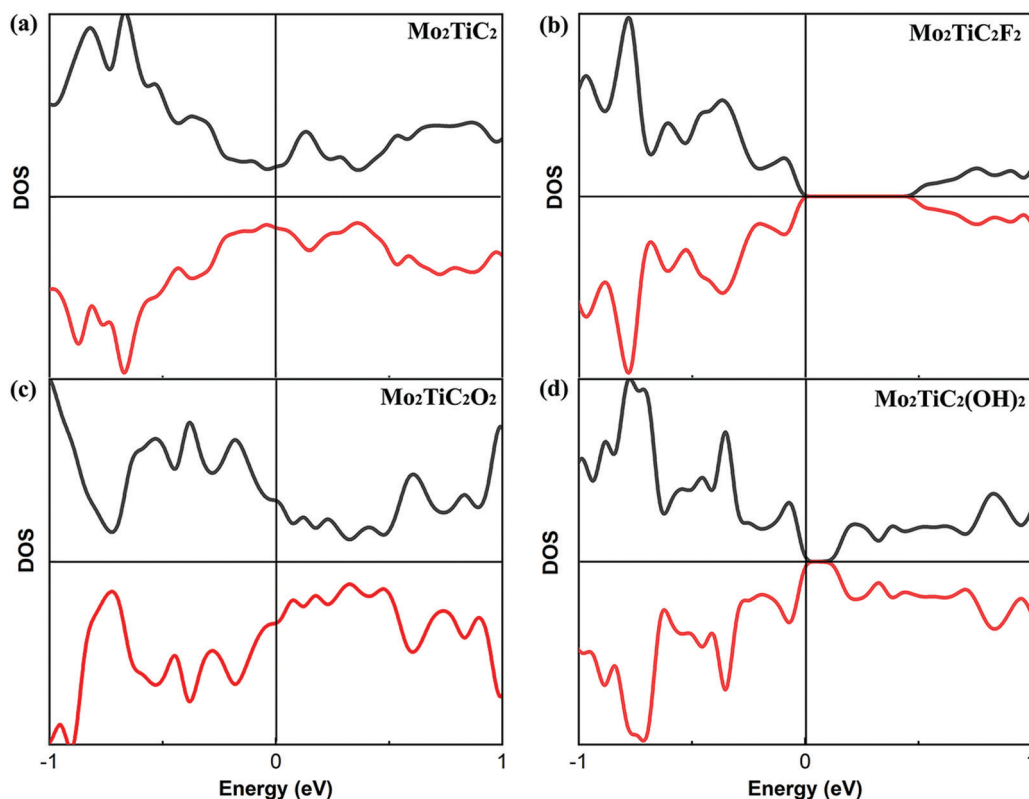


Fig. 3 Calculated total density of states (DOS) of perfect MXenes. (a) Pristine MXene, (b) $\text{Mo}_2\text{TiC}_2\text{F}_2$, (c) $\text{Mo}_2\text{TiC}_2\text{O}_2$, and (d) $\text{Mo}_2\text{TiC}_2(\text{OH})_2$. Black and red lines show spin-up and -down states, respectively.

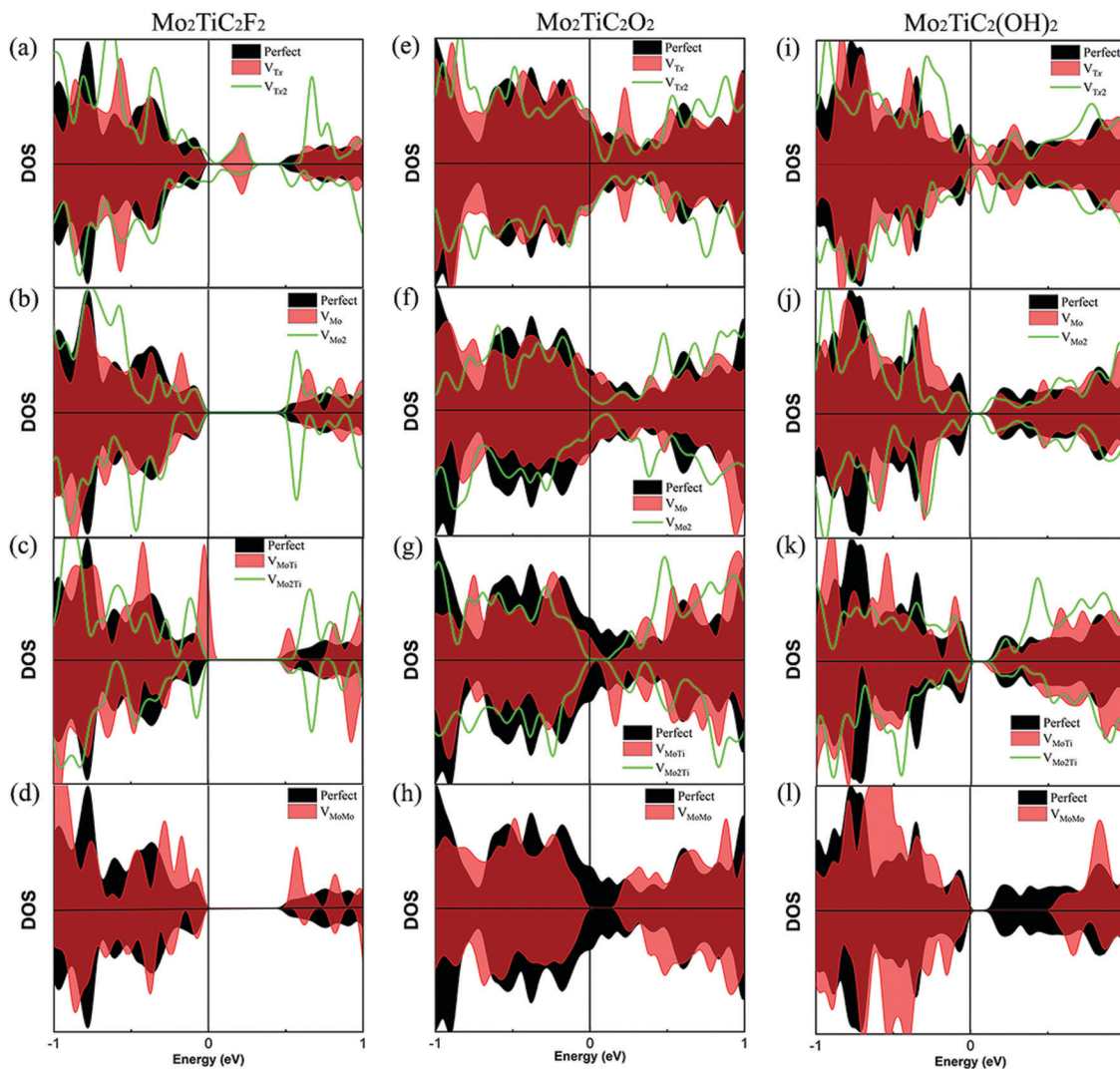


Fig. 4 Effect of different types of defects on the DOS of MXenes with different terminations. (a–d) $\text{Mo}_2\text{TiC}_2\text{F}_2$, (e–h) $\text{Mo}_2\text{TiC}_2\text{O}_2$, and (i–l) $\text{Mo}_2\text{TiC}_2(\text{OH})_2$. The type of defect is shown in the different panels.

found to have stronger adsorption energies for CO_2 molecules. Very recently, we theoretically studied the adsorption of CO_2 on BCN and graphene and found that a CO_2 molecule binds with an adsorption energy of $-18.91 \text{ kJ mol}^{-1}$ and $-13.19 \text{ kJ mol}^{-1}$ for these two 2D materials, respectively.⁵⁰ We also found that B-rich BCN samples exhibit significantly enhanced adsorption energies of CO_2 . Earlier, we investigated the effect of Stone-Wales (SW) defects on the adsorption energy of CO_2 on BCN and found that SW defects act as strong adsorption sites for CO_2 .⁴⁴ In this section, the influence of the considered defects on the reactivity and electrochemical transport properties of $\text{Mo}_2\text{TiC}_2\text{T}_x$ MXenes is investigated by modeling the interaction of a CO_2 molecule on the monolayer. The synthesized MXenes have a mixture of $-\text{O}$, $-\text{F}$, and $-\text{OH}$ terminations and the type of surface terminations and the percentage of each type depend on the etching and delamination conditions. Here, we only consider the MXene terminated by oxygen surface functions since the MXenes are more likely functionalized with a higher percentage of oxygen

than fluorine and hydroxide.^{6,10} We first evaluate the capability of a $3 \times 3 \times 1$ supercell of the clean $\text{Mo}_2\text{TiC}_2\text{O}_2$ MXene for CO_2 adsorption. Then, we discuss the reactivity of defected $\text{Mo}_2\text{TiC}_2\text{O}_2$ MXenes with a CO_2 molecule.

Our calculated results indicate that the CO_2 molecule is physisorbed on the perfect $\text{Mo}_2\text{TiC}_2\text{O}_2$ MXene, with a nonspontaneous reaction energy of 0.21 eV (see Fig. 5a). The carbon atom of CO_2 interacts with the oxygen surface functions and the molecule is placed at a distance of 2.89 Å while it is not significantly tilted. Moreover, a comparative nonspontaneous Gibbs free reaction energy (at 298.15 K) of 0.23 eV for the physisorption of a CO_2 molecule on the $\text{Mo}_3\text{C}_2\text{O}_2$ MXene was calculated.⁵¹ Thus, we can realize that Mo-containing MXenes are weakly reactive for CO_2 capture.

We further evaluated the proficiency of defected $\text{Mo}_2\text{TiC}_2\text{O}_2$ MXenes for CO_2 adsorption. Before that, we analyzed the electron localization function (ELF) over the surfaces, which provides a qualitative understanding of the empirical concept

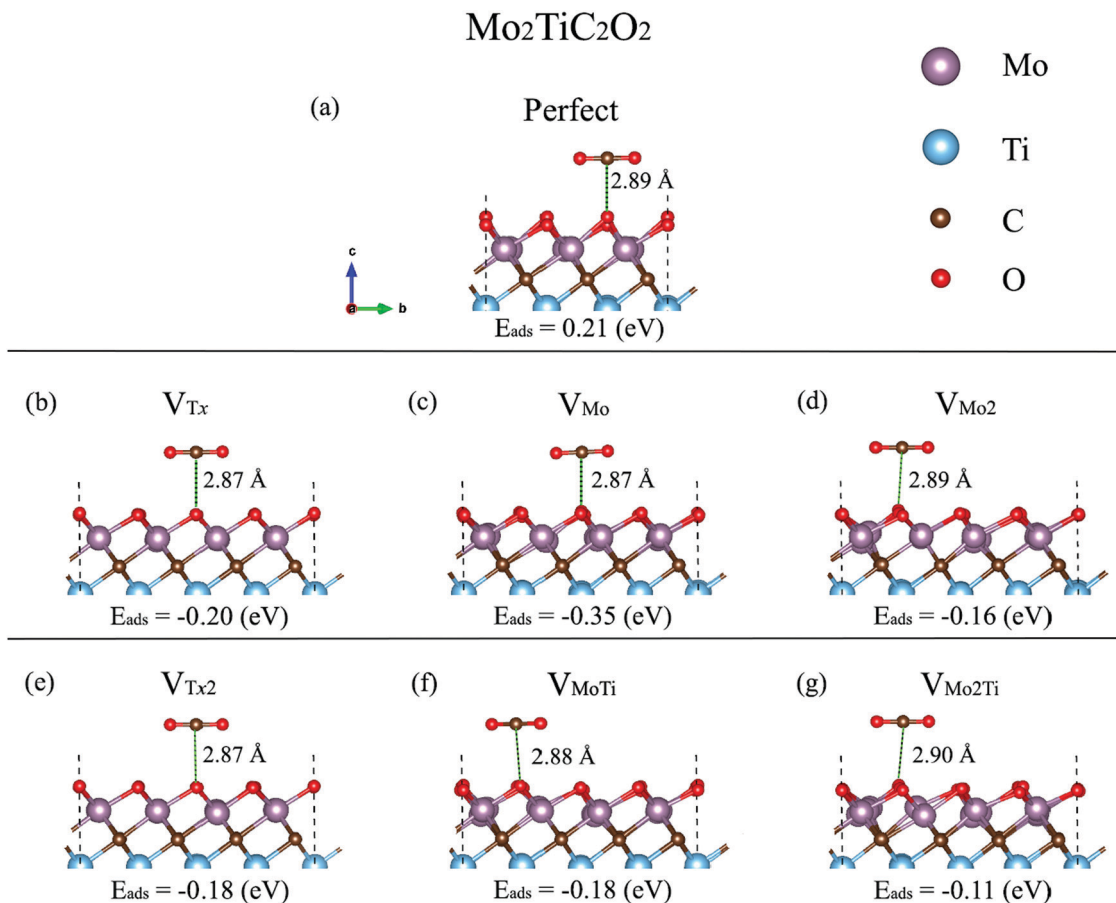


Fig. 5 Relaxed adsorption configurations of a CO_2 molecule on the perfect and defected $\text{Mo}_2\text{TiC}_2\text{O}_2$ MXene. (a) Perfect MXene, (b) MXene- V_{Tx} , (c) MXene- V_{Mo} , (d) MXene- V_{Mo_2} , (e) MXene- $V_{\text{Tx}2}$, (f) MXene- V_{MoTi} , and (g) MXene- $V_{\text{Mo}_2\text{Ti}}$.

of electron localization and the behavior of the electrons over the defective surfaces, as shown in Fig. 6. The calculated ELF for the non-terminated MXene and MXenes terminated by fluorine and hydroxide are shown in Fig. S2–S4 (ESI[†]), respectively. The ELF is defined in the range between 0 and 1, where the upper limit of

ELF = 1 exhibits perfect localization and ELF = 0.5 corresponds to a uniform electron gas. Here, we showed ELFs in the range between 0 and 0.5 since perfect localization was not detected over the surfaces. As it is shown in Fig. 6, electrons are generally localized on the surface functional groups over the MXene.

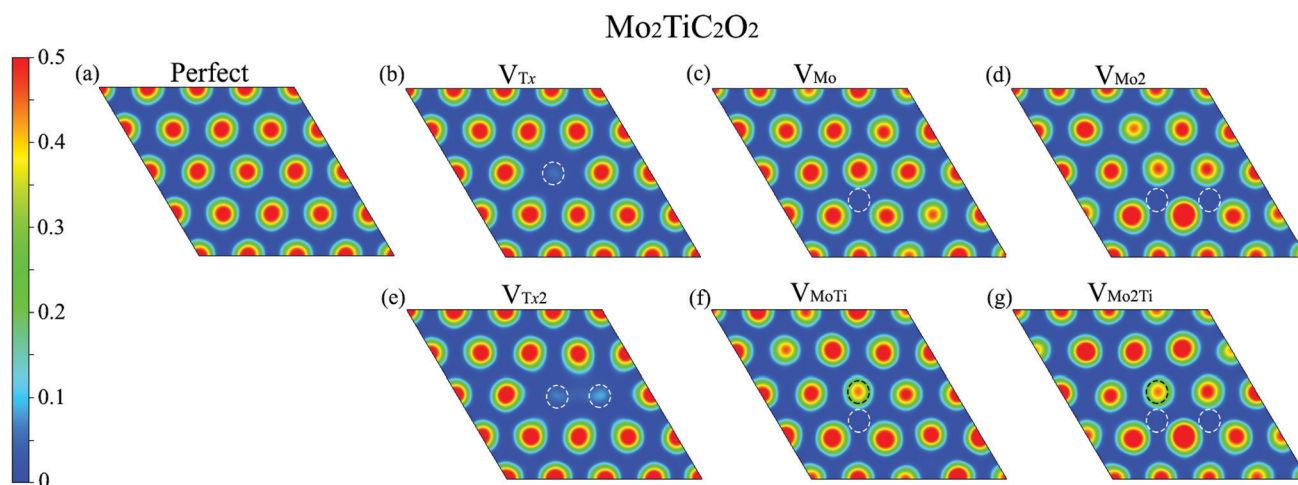


Fig. 6 Electron localization function (ELF) on the plane perpendicular to the c -axis at a close distance on top of surface terminations of the defected $\text{Mo}_2\text{TiC}_2\text{O}_2$ MXenes. Dashed circle lines highlight the place of the removed atoms.

In the case of V_{T_x} and $V_{T_{x2}}$, electrons are completely delocalized on the place of the removed oxygen atom(s) and freely move to other oxygen atoms on the surface. However, in other types of defects, electrons are partially delocalized on the oxygen atoms close to the removed atoms and move to other places. Based on the ELF analysis, we placed the CO_2 molecule close to different sites of the monolayer, in which electrons are more localized, parallel and perpendicular in different possible directions to find the largest adsorption energy in defective MXenes.

Our calculated adsorption energies (Fig. 5) imply that the molecule spontaneously adsorbs on the defective surfaces with adsorption energies between -0.11 and -0.35 eV, hence indicating that the molecule is trapped in the defect. The molecule is placed at distances between 2.87 and 2.90 Å and is not considerably tilted, relatively similar to the case of the perfect MXene. The adsorption energies on the defective surfaces including V_{T_x} and $V_{T_{x2}}$ are comparative (-0.20 and -0.18 eV, respectively). The molecule is highly attracted by the defect formed by a single Mo atom with a spontaneous reaction energy of -0.35 eV, while our results show a spontaneous reaction energy of -0.16 eV for MXene- V_{Mo_2} . Moreover, the results indicate that a Ti vacancy decreases the proficiency of MXene- V_{Mo} and MXene- V_{Mo_2} in CO_2 adsorption from -0.35 to -0.18 eV in MXene- V_{Mo} and from -0.16 to -0.11 eV in MXene- V_{Mo_2} . We also performed Bader charge analyses of the perfect and defected MXene- CO_2 systems to quantify the electron transfer from the MXenes to the CO_2 molecule, but we observed that the charge transfer to the CO_2 molecule is quite negligible since the molecule is not chemisorbed on the surfaces. The electron density difference isosurface plot for the case of the MXene- V_{T_x} surface, as an example, is shown in Fig. S5 (ESI†).

4 Conclusion

In the present study, we have investigated atomic defects in ordered $Mo_2TiC_2T_x$ terminated with specific surface functional groups of fluorine, oxygen or hydroxide. Different atomic vacancy defects in the top and sublayer of the MXenes have been studied. Through defect formation energy calculations, we conclude that the defects are dependent on the surface terminations, where the O-terminated MXenes demand more energy than the F- and OH-terminated MXenes; also the defect formation is more feasible in the outer molybdenum layers than in the inner titanium layer. Meanwhile, DOS calculations indicate that there are minor changes in the electronic states with a shift in states away from the Fermi energy level, making the system more semiconducting.

Our results of the CO_2 molecule adsorption indicate the potential catalytic properties of the $Mo_2TiC_2T_x$ MXenes and the catalytic properties of the MXenes can be enhanced by forming defects in the monolayer. While a CO_2 molecule adsorbs on the perfect O-terminated MXene through a nonspontaneous and endothermic process, defect formation on the monolayer leads to direct CO_2 interactions with the surfaces through a spontaneous and exothermic process that is crucial to its capture. Our study shows that defect formation on the monolayer cannot activate the CO_2 molecule since the bond angle $\angle O-C-O$ is not considerably

changed. Our study unveils a route to the possibility of enhancing the catalytic properties of $Mo_2TiC_2T_x$ MXenes by the formation of specific defects in the monolayer.

Conflicts of interest

There are no conflicts to declare.

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References

- 1 B. Anasori, M. R. Lukatskaya and Y. Gogotsi, *Nat. Rev. Mater.*, 2017, **2**, 16098.
- 2 M. Zhao, M. Torelli, C. E. Ren, M. Ghidui, Z. Ling, B. Anasori, M. W. Barsoum and Y. Gogotsi, *Nano Energy*, 2016, **30**, 603–613.
- 3 M. Naguib, J. Halim, J. Lu, K. M. Cook, L. Hultman, Y. Gogotsi and M. W. Barsoum, *J. Am. Chem. Soc.*, 2013, **135**, 15966–15969.
- 4 Y. Peng, B. Akuzum, N. Kurra, M. Zhao, M. Alhabeab, B. Anasori, E. C. Kumbur, H. N. Alshareef, M. Ger and Y. Gogotsi, *Energy Environ. Sci.*, 2016, **9**, 2847–2854.
- 5 G. Fan, X. Li, Y. Ma, Y. Zhang, J. Wu, B. Xu, T. Sun, D. Gao and J. Bi, *New J. Chem.*, 2017, **41**, 2793–2799.
- 6 B. Anasori, Y. Xie, M. Beidaghi, J. Lu, B. C. Hosler, L. Hultman, P. R. C. Kent, Y. Gogotsi and M. W. Barsoum, *ACS Nano*, 2015, **9**, 9507–9516.
- 7 M. Naguib, M. Kurtoglu, V. Presser, J. Lu, J. Niu, M. Heon, L. Hultman, Y. Gogotsi and M. W. Barsoum, *Adv. Mater.*, 2011, **23**, 4248–4253.
- 8 M. Khazaei, A. Ranjbar, M. Arai and S. Yunoki, *J. Mater. Chem. C*, 2017, **5**, 2488–2503.
- 9 R. Khaledialidusti, B. Anasori and A. Barnoush, *Phys. Chem. Chem. Phys.*, 2020, **22**(6), 3414–3424.
- 10 B. Anasori, C. Shi, E. J. Moon, Y. Xie, C. A. Voigt, P. R. C. Kent, S. J. May, S. J. L. Billinge, M. W. Barsoum and Y. Gogotsi, *Nanoscale Horiz.*, 2016, **1**, 227–234.
- 11 T. Hu, X. Wang and X. Wang, *Phys. Chem. Chem. Phys.*, 2017, **19**, 31773–31780.
- 12 A. Bandyopadhyay, D. Ghosh and S. K. Pati, *Phys. Chem. Chem. Phys.*, 2018, **20**, 4012–4019.
- 13 W. Zhou, X. Zou, S. Najmaei, Z. Liu, Y. Shi, J. Kong and J. Lou, *Nano Lett.*, 2013, **13**, 2615–2622.

- 14 F. Banhart, J. Kotakoski and A. V. Krasheninnikov, *ACS Nano*, 2011, **5**, 26–41.
- 15 W. Cheong, L. Zheng, F. Ren, G. Ying, X. Cao, D. Wang, Q. Peng, G. Wang and C. Chen, *J. Am. Chem. Soc.*, 2019, **141**, 4086–4093.
- 16 X. Sang, Y. Xie, M. Lin, M. Alhabeab, K. L. Van Aken, Y. Gogotsi, P. R. C. Kent, K. Xiao and R. R. Unocic, *ACS Nano*, 2016, **10**, 9193–9200.
- 17 L. H. Karlsson, J. Birch, J. Halim, M. W. Barsoum and P. O. Å. Persson, *Nano Lett.*, 2015, **15**, 4955–4960.
- 18 J. Halim, S. Kota, M. R. Lukatskaya, M. Naguib, M. Zhao, E. J. Moon, J. Pitcock, J. Nanda, S. J. May, Y. Gogotsi and M. W. Barsoum, *Adv. Funct. Mater.*, 2016, **2**, 3118–3127.
- 19 A. Lipatov, M. Alhabeab, M. R. Lukatskaya, A. Boson, Y. Gogotsi and A. Sinitiskii, *Adv. Electron. Mater.*, 2016, **2**, 1600255.
- 20 L. Vicarelli, S. J. Heerema, C. Dekker and H. W. Zandbergen, *ACS Nano*, 2015, **9**, 3428–3435.
- 21 G. López-polín, C. Gómez-navarro, V. Parente, M. I. Katsnelson, F. Pérez-murano and J. Gómez-herrero, *Nat. Phys.*, 2015, **11**, 26.
- 22 D. Wong, J. V. Jr, L. Ju, J. Lee, S. Kahn, H. Tsai, C. Germany, T. Taniguchi, K. Watanabe, A. Zettl, F. Wang and M. F. Crommie, *Nat. Nanotechnol.*, 2015, **10**, 949–953.
- 23 Y. Liu, X. Zou and B. I. Yakobson, *ACS Nano*, 2012, **6**, 7053–7058.
- 24 R. Bourrellier, S. Meuret, A. Tararan, O. Ste, M. Kociak, L. H. G. Tizei and A. Zobelli, *Nano Lett.*, 2016, **16**, 4317–4321.
- 25 M. Amani, D. Lien, D. Kiriya, J. Xiao, A. Azcatl, J. Noh, S. R. Madhupathy, R. Addou, M. Dubey, K. Cho, R. M. Wallace, S. Lee, J. He, X. Zhang, E. Yablonovitch and A. Javey, *Science*, 2015, **350**, 1065–1068.
- 26 Y. Wei, J. Wu, H. Yin, X. Shi, R. Yang and M. Dresselhaus, *Nat. Mater.*, 2012, **11**, 759–763.
- 27 O. V. Yazyev and S. G. Louie, *Nat. Mater.*, 2010, **9**, 806–809.
- 28 G. Kresse and J. Furthmüller, *Comput. Mater. Sci.*, 1996, **6**, 15–50.
- 29 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, 3865–3868.
- 30 G. Kresse and D. Joubert, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 11–19.
- 31 H. J. Monkhorst and J. D. Pack, *Phys. Rev. B: Solid State*, 1976, 5188–5192.
- 32 M. Methfessel and A. T. Paxton, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1989, **40**, 3616–3621.
- 33 S. L. Dudarev, G. A. Botton, S. Y. Savrasov, C. J. Humphreys and A. P. Sutton, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1998, **57**, 1505–1509.
- 34 M. Nolan, S. D. Elliott, J. S. Mulley, R. A. Bennett and M. Basham, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2008, **77**, 235424.
- 35 G. Hautier, S. P. Ong, A. Jain, C. J. Moore and G. Ceder, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, **85**, 155208.
- 36 W.-X. Li, C. Stampfl and M. Scheffler, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2003, **68**, 165412.
- 37 S. Grimme, *J. Comput. Chem.*, 2006, **27**, 1787–1799.
- 38 R. Khaledialidusti, A. K. Mishra and A. Barnoush, *ACS Omega*, 2019, **4**(14), 15935–15946.
- 39 R. Khaledialidusti, A. K. Mishra and A. Barnoush, *J. Chem. Phys.*, 2018, **149**, 224702.
- 40 A. K. Mishra, A. Roldan and N. H. De Leeuw, *J. Phys. Chem. C*, 2016, **120**, 2198–2214.
- 41 M. Khazaei, A. Ranjbar, D. Bogdanovski and R. Dronskowski, *Phys. Chem. Chem. Phys.*, 2018, **20**, 8579–8592.
- 42 O. Bikondoa, C. L. Pang, R. Ithnin, C. A. Muryn, H. Onishi and G. Thornton, *Nat. Mater.*, 2006, **5**, 189–192.
- 43 H. H. Gürel, V. O. O. Zçelik and S. Ciraci, *J. Phys. Chem. C*, 2014, **118**, 27574–27582.
- 44 N. Kumar, K. Moses, K. Pramoda, S. N. Shirodkar, A. K. Mishra, U. V. Waghmare, A. Sundaresana and C. N. R. Rao, *J. Mater. Chem. A*, 2013, **1**, 5806–5821.
- 45 S. Cho, S. J. Kim, Y. Lee, J. Kim, W. Jung, H. Yoo, J. Kim, H. Jung, B. Engineering, A. Technology, A. Development and D. Development, *ACS Nano*, 2015, **9**, 9314–9321.
- 46 L. Qi, Y. Wang, L. Shen and Y. Wu, *Appl. Phys. Lett.*, 2016, **1089**, 063103.
- 47 B. Akdim, R. Pachter and S. Mou, *Nanotechnology*, 2016, **27**, 185701.
- 48 H. Li, M. Huang and G. Cao, *Phys. Chem. Chem. Phys.*, 2016, **18**, 15110–15117.
- 49 C. Gonzalez, B. Biel and Y. Dappe, *Phys. Chem. Chem. Phys.*, 2017, **19**, 9485.
- 50 A. Kumar and S. Mishra, *J. Mol. Graphics Modell.*, 2019, **93**, 107446.
- 51 N. Li, X. Chen, W. Ong, D. R. Macfarlane, X. Zhao, A. K. Cheetham and C. Sun, *ACS Nano*, 2017, **11**, 10825–10833.