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**PAPER**

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Unravelling the origin of dual photoluminescence in  $\text{Au}_2\text{Cu}_6$   
clusters by triplet sensitization and photon upconversion

## PAPER

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[View Journal](#) | [View Issue](#)Cite this: *J. Mater. Chem. C*, 2022, **10**, 4597Unravelling the origin of dual photoluminescence in  $\text{Au}_2\text{Cu}_6$  clusters by triplet sensitization and photon upconversion†

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Ligand-protected, noble-metal clusters are promising as luminescent materials and photosensitizers because of their diverse chemical compositions and structures, as well as their easily tunable electronic and photophysical properties with single-atom precision. Many studies have shown that the photoluminescence (PL) properties of these clusters can be significantly enhanced by alloying the metal core or by increasing the rigidity of the ligand environment. However, the fundamental nature of their PL (whether fluorescent or phosphorescent) and the elementary process for enhancing their PL have not been fully clarified. In this study, we established that  $\text{Au}_2\text{Cu}_6(\text{S-Adm})_6(\text{PPh}_3)_2$  (S-Adm = 1-adamantanethiolate) clusters exhibit not only fluorescence but also phosphorescence at room temperature. Further, we discovered that  $\text{Au}_2\text{Cu}_6$  serves as a triplet sensitizer, and through the analysis of the photon-upconversion phenomenon, we determined the rate constants and quantum yields of all radiative and non-radiative processes involving the excited singlet and triplet states of the clusters. The temperature dependence of the PL and the theoretical calculations indicated that thermally activated intersystem crossing occurs through the spin–vibronic coupling mechanism mediated by higher excited states ( $\text{S}_2$  and/or  $\text{T}_2$ ) where direct spin–orbit coupling is possible.

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## Introduction

Ligand-protected noble-metal clusters have attracted significant attention because of their diverse chemical compositions and structures, as well as their potential applications as photosensitizers in photocatalytic reactions,<sup>1,2</sup> singlet oxygen generation,<sup>3,4</sup> solar cells,<sup>5,6</sup> and photon upconversion.<sup>7</sup> To realize these applications, it is important to profoundly understand the underlying photophysics of the clusters, including their photo-absorption and photoluminescence (PL) properties. As is well known, dual PL and multiple PL decay are frequently observed in ligand-protected gold, silver, and their alloy clusters at room temperature;<sup>8–12</sup> however, their origin remains unclear. The exact origin of the PL is difficult to identify because of the presence of multiple emission sites in the clusters, such as the metal core and surface staple,<sup>7–12</sup> as well as the contribution of electronic states with different spin multiplicities (*i.e.*, excited singlet and triplet states) to these emission sites.<sup>7,13–15</sup> Therefore, it would be desirable to determine the rate constants and quantum yields of all radiative and non-radiative processes

involving the excited singlet and triplet states of the cluster, so that the complex behavior of the excited-state relaxation processes of the clusters can be accurately understood.

As shown in Scheme 1, the  $\text{Au}_2\text{Cu}_6(\text{S-Adm})_6(\text{PPh}_3)_2$  (where S-Adm = 1-adamantanethiolate, hereinafter referred to as  $\text{Au}_2\text{Cu}_6$ ) cluster, first synthesized by Zhu *et al.*,<sup>16</sup> possesses a structure in which the  $\text{Cu}_6(\text{S-Adm})_6$  ring surrounds a  $\text{PPh}_3$ -terminated Au dimer. This cluster exhibits a strong emission peak around 660 nm in dichloromethane (DCM), which is assigned to the ligand-to-metal charge-transfer (LMCT) fluorescence transition.<sup>16,17</sup> Additionally, to

Scheme 1 Chemical structure of  $\text{Au}_2\text{Cu}_6$ .

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tailor the PL characteristics, the excited-state charge-transfer property of the cluster can be perturbed using phosphine ligands that functionalize the phenyl groups with electron-donating or electron-withdrawing groups.<sup>16</sup> Such flexibility in controlling the electronic and photophysical properties suggests that Au<sub>2</sub>Cu<sub>6</sub>-based clusters are promising as luminescent materials and photosensitizers.

In this study, we established that the Au<sub>2</sub>Cu<sub>6</sub> cluster exhibits not only fluorescence but also phosphorescence at room temperature and can sensitize molecular triplets. Thus far, various triplet sensitizers, such as heavy-metal (*e.g.*, Os, Ir, Ru, and Pt) complexes<sup>18–23</sup> and semiconductor nanocrystals,<sup>24–37</sup> have been developed and applied to triplet fusion upconversion (TF-UC). Recently, we presented atomically precise metal clusters as a new addition to the library of triplet sensitizers.<sup>7</sup> Herein, we report yet another example of triplet sensitization and TF-UC by metal clusters. Through a detailed analysis of the TF-UC phenomenon occurring in the Au<sub>2</sub>Cu<sub>6</sub> sensitizer-erylene emitter pair, we successfully determined the rate constants and quantum yields of all radiative and non-radiative processes involving the excited singlet and triplet states of the Au<sub>2</sub>Cu<sub>6</sub> cluster. The temperature dependence of the dual PL revealed that thermally activated intersystem crossing (ISC) occurs in this cluster, and the ISC mechanism is discussed based on the results of density functional theory (DFT) and time-dependent DFT (TD-DFT) calculations.

## Experimental

### Chemicals

The following special-grade reagents were purchased from FUJIFILM Wako Pure Chemicals Co., Ltd: hydrogen tetrachloroaurate(III) tetrahydrate (HAuCl<sub>4</sub>·4H<sub>2</sub>O), copper(I) chloride (CuCl), triphenylphosphine (PPh<sub>3</sub>), anthracene, DCM, and tetrahydrofuran (THF). Platinum(II) octaethylporphyrin (PtOEP), *N,N'*-bis(2,5-di-*tert*-butylphenyl)-3,4,9,10-perylene-dicarboximide (PDI), and 1-adamantanethiol (HS-Adm) were purchased from Sigma Aldrich. Sodium borohydride (NaBH<sub>4</sub>), perylene, and 9,10-diphenylanthracene (DPA) were purchased from Tokyo Chemical Industry Co., Ltd. All the chemicals were used as received without further purification. Ultrapure water ( $\geq 18.2$  M $\Omega$  cm) was produced using a Direct-Q UV distillation system.

### Synthesis

Au<sub>2</sub>Cu<sub>6</sub>(S-Adm)<sub>6</sub>(PPh<sub>3</sub>)<sub>2</sub> (Scheme 1) was synthesized following a method previously reported by Zhu and coworkers.<sup>16</sup> Briefly, 19.8 mg of CuCl (0.2 mmol) was dispersed in a 1 : 1 acetonitrile and methanol mixed solution, after which 136 mg (0.8 mmol) of 1-adamantanethiol was added. The mixture was stirred until a clear colorless solution was obtained. Thereafter, a solution of 198 mg (0.4 mmol) of Au(PPh<sub>3</sub>)Cl dissolved in 5 mL of toluene and 1 mL of 60 mg mL<sup>-1</sup> of NaBH<sub>4</sub> aq. were added simultaneously. After stirring for 60 h, the solution was centrifuged to obtain a purple precipitate. The precipitate solution was

washed with toluene three or four times and subsequently extracted with DCM to obtain pure Au<sub>2</sub>Cu<sub>6</sub>(S-Adm)<sub>6</sub>(PPh<sub>3</sub>)<sub>2</sub>.

The solution containing the product was diluted with a large amount of DCM ( $\sim 10^{-8}$  M) and subjected to electrospray ionization mass spectrometry (JEOL, JMS-T100LP AccuTOF LC-Plus). The mass spectrum was monitored in the positive-ion mode. A NaI methanol solution was employed for the calibration. Fig. S1 (ESI<sup>†</sup>) displays the positive-ion mode electrospray ionization mass spectrum of the product. The peaks around 2303 *m/z* correspond to Au<sub>2</sub>Cu<sub>6</sub>(S-Adm)<sub>6</sub>(PPh<sub>3</sub>)<sub>2</sub>. The isotope pattern of this peak is consistent with that obtained theoretically.

### Absorption and PL measurements

Ultraviolet-visible (UV-vis) absorption spectra were recorded on a Lambda 650 spectrometer (PerkinElmer). The PL spectra were recorded using a SpectraPro 2300i polychromator coupled with a liquid-nitrogen-cooled charge-coupled device camera (Spec-10:100B/LN, Roper Scientific). For the PL spectrum measurements, a linearly polarized continuous-wave (CW) diode laser with a wavelength of 532 or 640 nm (Edmund Optics) was employed as the excitation source. The temperature dependence of the PL spectrum and lifetime was measured using a variable temperature cell holder (CoolSpeK, Unisoku Co., Ltd). In measuring the PL lifetime of the clusters, a 478 or 634 nm-picosecond pulsed laser (pulse width: 40–80 ps, P1L048X or P1L063X, Advanced Laser Diode System) was employed as the excitation source. All the signals, including the synchronization of the pulse laser and an avalanche photodiode (APD, SPCM-AQRH-61, PerkinElmer), were fed into a time-correlated single-photon-counting (TC-SPC) card (TimeHarp 260, PicoQuant) operated in time-tagged and time-resolved modes.<sup>38,39</sup> Data acquisition and fluorescence decay analyses were performed using the SymPhoTime 64 (PicoQuant) software. To determine the fluorescence lifetime, single or double exponentials were fitted to the fluorescence curves by maximum likelihood estimation. For the excitation spectrum measurement, white light from an Xe arc lamp through an optical fiber was supplied into a monochromator (CM110, Spectral Products) and spectrally divided into monochromatic light. The solution sample was irradiated with the monochromatic light, thus obtained, while scanning the wavelength of the monochromatized light, and the emission from the sample was detected by the APD. A silicon photodiode (SM1PD2A-CAL-SP, Thorlabs Japan Inc.) with an amplifier (PDA200C, Thorlabs Japan Inc.) was used to record the light intensity at each wavelength and to correct the signal intensity of the excitation spectrum. To monitor only the emission band of interest, appropriate filters were placed in front of the APD. The experimental setup for the upconverted emission decay measurements has been thoroughly described elsewhere.<sup>7</sup> Briefly, a 532 nm-CW laser, periodically blocked at 500 Hz by an optical chopper (MC2000B controller, MC1F10HP blade, Thorlabs), was employed as the excitation source. Based on the scattered light decay measurements under this condition, the time resolution was estimated to be 4.5  $\mu$ s, which was sufficiently shorter than the decay time of the delayed (or upconverted) fluorescence of the



perylene emitter (*i.e.*, several hundred  $\mu\text{s}$ ). The upconverted emission signal was detected by an APD (APD410A/M, Thorlabs) and monitored synchronously by an oscilloscope (GDS1062A, Chip One Stop). The decay profiles were averaged 256 times and stored in the oscilloscope.

### Electrochemical measurement

To examine the electrochemical properties of the cluster, differential pulse-voltammetry measurements were conducted in an acetonitrile solution containing tetrabutyl ammonium perchlorate using ECstat-301 (EC Frontier Co., Ltd). For the experiment, glassy carbon was employed as the working electrode,  $\text{Ag}/\text{Ag}^+$  as the reference electrode, and a Pt wire as the counter electrode, with a scan rate of  $50 \text{ mV s}^{-1}$ . A ferrocenium/ferrocene ( $\text{Fc}/\text{Fc}^+$ ) redox couple was used as an internal reference, and the measured potentials were converted to normal hydrogen electrode (NHE) potentials by  $+0.69 \text{ V}$  or the energy level of the highest occupied molecular orbital (HOMO) relative to the vacuum level by  $E_{\text{HOMO}} (\text{V, vs. NHE}) = -E_{\text{HOMO}} (\text{eV, vs. vacuum level}) - 4.44$ .<sup>40</sup> For perylene, DPA, and anthracene, the HOMO and lowest unoccupied molecular orbital (LUMO) energies were obtained from the literature.<sup>41,42</sup>

### Transient absorption spectroscopy of $\text{Au}_2\text{Cu}_6$

Subnanosecond transient absorption (TA) spectroscopy measurements of the  $\text{Au}_2\text{Cu}_6(\text{S-Adm})_6(\text{PPh}_3)_2$  cluster with and without perylene in deaerated THF were performed using a picosecond TA spectroscopy system (picoTAS, UNISOKU Co., Ltd) based on the randomly interleaved pulse-train method.<sup>43</sup> The pump source was a picosecond Nd:YAG laser (355 nm, EKSPLA PL-2210A, 1 kHz, full width at half maximum (FWHM) = 25 ps) equipped with an optical parametric generator (EKSPLA PG403, 410–700 nm, 50  $\mu\text{J}$  per pulse @500 nm), and the probe light source was a supercontinuum radiation source (SMHP-20.2, Leukos, 20 MHz, FWHM = 50–100 ps, 410–2000 nm). The temporal resolution of the system was estimated to be 80–100 ps with a 10–90% rise time. The TA measurements were performed using a quartz cell with an optical path length of 2 mm.

### Quantum chemical calculations

DFT and TD-DFT calculations were implemented with the Gaussian 16 program (ES64L-G16, RevB.01).<sup>44</sup> To reduce the central processing unit cost, all the S-Adm ligands in the clusters were replaced with 1-*tert*-butylthiolate (S-CMe<sub>3</sub>) ligands. The geometric structure of  $\text{Au}_2\text{Cu}_6(\text{S-CMe}_3)_6(\text{PPh}_3)_2$  in the  $S_0$ ,  $S_1$ , or  $T_1$  state was optimized using the BP86 functional<sup>45</sup> with the basis sets of def2-SV(P) for Cu and Au atoms<sup>46</sup> and 6-31G(d,p) for H, C, P, and S atoms;<sup>47</sup> the SDD pseudopotentials with scalar relativistic effects were included.<sup>48</sup> Furthermore, harmonic vibrational frequency analysis was performed for all the optimized structures, and it was confirmed that there were no imaginary frequencies. Using these optimized structures, the vertical transition energies of  $S_0 \rightarrow S_n$  ( $n = 1\text{--}200$ ),  $S_1 \rightarrow S_0$ , and  $T_1 \rightarrow S_0$  were calculated using the TD-B3LYP functional,<sup>49–51</sup> with a polarizable continuum model (solvent = THF). In this

calculation, the basis sets of def2-SV(P) for the Cu and Au atoms and 6-31G(d,p) for H, C, P, and S atoms were employed, and the SDD pseudopotentials were also implemented. Note that the theoretical levels chosen in this study have previously been shown to yield relatively accurate results in reproducing the electronic structure of ligand-protected noble metal clusters.<sup>7,52</sup> The vertical transition energies from  $S_2$  to  $S_0$  at the  $S_1$  equilibrium geometry and the transition energy from  $T_2$  to  $S_0$  at the  $T_1$  equilibrium geometry were also calculated at the same theoretical level. The molecular structures and frontier molecular orbitals were constructed using GaussView 6.0 and Avogadro 1.2.0, respectively.

## Results and discussion

### Photoabsorption and emission properties

The UV-vis absorption spectra of  $\text{Au}_2\text{Cu}_6$  are shown in Fig. 1. Two relatively sharp peaks around 520 nm and 590 nm are observed, which are consistent with the features of the reported absorption spectra of  $\text{Au}_2\text{Cu}_6$  in DCM.<sup>16</sup> Conversely, in the reported emission spectrum of  $\text{Au}_2\text{Cu}_6$  in DCM, only a single LMCT fluorescence band (band I) around 660 nm is observed.<sup>16,17</sup>



Fig. 1 Absorption and emission spectra of  $\text{Au}_2\text{Cu}_6$  (10  $\mu\text{M}$ ) in (a) dichloromethane (DCM) and (b) tetrahydrofuran (THF). (c) Excitation spectra obtained by analysis at the emission wavelength of  $\lambda \leq 700 \text{ nm}$  (band I, solid line) or  $\lambda \geq 800 \text{ nm}$  (band II, dashed line). (d) Photoluminescence decay curves of  $\text{Au}_2\text{Cu}_6$  in deaerated DCM and THF ( $\lambda_{\text{ex}} = 478 \text{ nm}$ ), where both bands I and II were monitored. The inset shows the decay and rise of the selectively observed band (bands I [ $\lambda \leq 700 \text{ nm}$ ] and II [ $\lambda \geq 800 \text{ nm}$ ]) emissions for  $\text{Au}_2\text{Cu}_6$  in deaerated THF.

However, Fig. 1a and b show another emission band (band II) at relatively long wavelengths in DCM and THF. The excitation spectra were measured using appropriate filters to selectively monitor only the emission of band I or II, as shown in Fig. 1c. The excitation spectra were virtually identical and matched the absorption spectra, confirming that both emission bands originated from  $\text{Au}_2\text{Cu}_6$ . Band II was not observed in the literatures presumably because the detector used was not sensitive to photons with wavelengths longer than 750 nm. As shown in Fig. 1d, two decay components having very different lifetimes were observed in the PL decay curve. The fast decay component originated from band I and the slow component from band II, as confirmed by the selective observation of the PL decay curves (Fig. S2 (ESI<sup>†</sup>)) and the inset of Fig. 1d). Notably, the emission decay of band I and the emission rise of band II showed practically identical single exponential kinetics, both having a lifetime of 9 ns. As listed in Table 1, the emission lifetimes of band I were 19.0 and 9.0 ns in DCM and THF, respectively, while those of band II were 5.2 and 4.9  $\mu\text{s}$ , respectively. Judging from these results, the band-I emission corresponds to fluorescence,<sup>16,17</sup> while the newly observed emission (band II) can be tentatively assigned to phosphorescence. This suggests the formation of an excited triplet state in the  $\text{Au}_2\text{Cu}_6$  cluster. Although the fluorescence quantum yields ( $\Phi_f$ ) of  $\text{Au}_2\text{Cu}_6$  in DCM and THF were both quite small, there was a pronounced solvent effect on the fluorescence radiative rate constant ( $k_f$ ), with a 4.9-fold increase in DCM compared to that in THF. This solvent effect is attributable to the strong LMCT nature of the fluorescence transition,<sup>16,17</sup> as mentioned above. The emission quantum yields of band II ( $\Phi_{\text{PL}}$ ), which is considered to be phosphorescence, were determined by relative estimation (see ESI<sup>†</sup>) to be 0.17 (DCM) and 0.21 (THF). However, notably, the value of  $\Phi_{\text{PL}}$  corresponds to the product of the ISC yield ( $\Phi_{\text{ISC}}$ ) and the phosphorescence quantum yield ( $\Phi_p$ ), i.e.,  $\Phi_{\text{PL}} = \Phi_{\text{ISC}} \cdot \Phi_p$ . The determination of these quantum yields will be discussed later. As shown in Fig. S3 (ESI<sup>†</sup>), dynamic light scattering measurements suggest that the photostability of  $\text{Au}_2\text{Cu}_6$  in DCM is low and that the clusters dissociate rapidly upon laser irradiation. Therefore, all the results shown below were obtained using THF as the solvent, in which the photostability of the  $\text{Au}_2\text{Cu}_6$  cluster is significantly higher than that in DCM.

### Quenching experiments and triplet sensitization

As shown in Fig. 2, the energy of the  $S_1$  state of perylene (2.84 eV) was sufficiently higher than the excitation photon energy (1.96 eV). Furthermore, the LUMO energy of perylene was higher than that of the cluster by more than 1 eV (Fig. S4, ESI<sup>†</sup>).

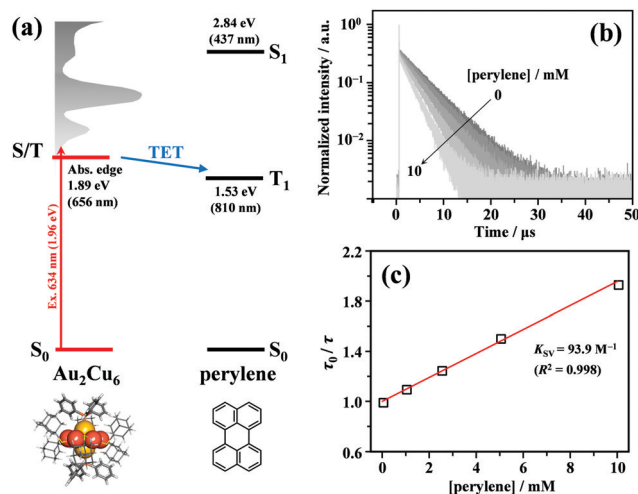


Fig. 2 (a) Energy level diagram of the  $\text{Au}_2\text{Cu}_6$  and perylene pair. (b) PL decay curves of  $\text{Au}_2\text{Cu}_6$  (10  $\mu\text{M}$ ) with [perylene] = 0–10 mM in deaerated THF, where both band-I and -II were monitored. To prevent the high concentration of perylene from absorbing the excitation light, an excitation wavelength of  $\lambda_{\text{ex}} = 634 \text{ nm}$  was used in this quenching experiment. (c) Stern–Volmer plot of phosphorescence decay times (i.e., PL lifetimes of band-II) at different perylene concentrations.  $K_{\text{SV}}$  and  $R^2$  represent the Stern–Volmer constant and coefficient of determination.

Therefore, neither singlet energy transfer nor photoinduced electron transfer could occur. In contrast, the  $T_1$  state of perylene (1.53 eV) was lower in energy than the absorption edge of the cluster (1.89 eV); therefore, triplet energy transfer (TET) could occur in this pair. Energetically, there is also the possibility of hole transfer from the cluster to perylene (Fig. S4b, ESI<sup>†</sup>); however, no evidence suggesting this process was obtained in the present experiments. This is probably due to the fact that the HOMO localized in the  $\text{Au}_2$  and  $\text{Cu}_6\text{S}_6$  portions is almost completely sterically shielded by the six adamantane and two  $\text{PPh}_3$  ligands, hindering hole transfer. As shown in Fig. 2b, the lifetime of the long-lived component (band II) gradually decreased with an increase in the perylene concentration from 0 to 10 mM. The Stern–Volmer plot of the observed lifetimes is shown in Fig. 2c. The plot was well reproduced with a linear fit, and the slope yielded a Stern–Volmer constant ( $K_{\text{SV}}$ ) of  $93.9 \text{ M}^{-1}$ . To confirm that this quenching was caused by TET, TA spectroscopy measurements were performed with and without perylene. Without perylene, the  $\text{Au}_2\text{Cu}_6$  cluster exhibited an absorption band at 450–490 nm, attributed to the triplet–triplet (T–T) absorption, and ground-state bleaching at 520 and 590 nm. Both bands exhibited decay times of approximately 3  $\mu\text{s}$ .

Table 1 Photoabsorption and photoluminescence parameters of  $\text{Au}_2\text{Cu}_6$  in deaerated dichloromethane (DCM) and tetrahydrofuran (THF).<sup>a</sup>

Solvent	$\lambda_{\text{abs}}/\text{nm}$ ( $\epsilon/\text{M}^{-1}\text{cm}^{-1}$ )	$\lambda_{\text{em}}/\text{nm}$ band-I band-II	$\tau_f/\text{ns}$ band-I	$\Phi_f$ band-I	$k_f/\text{s}^{-1}$	$\tau_p/\mu\text{s}$ band-II	$\Phi_{\text{PL}}$ band-II
DCM	593 (4268)	684 (1.813) 814 (1.523)	19.0	$5.3 \times 10^{-2}$	$2.8 \times 10^6$	5.2	0.17
THF	591 (4483)	674 (1.840) 811 (1.529)	9.0	$5.1 \times 10^{-3}$	$5.7 \times 10^5$	4.9	0.21

<sup>a</sup> Note that all the data were obtained by preparing solution samples in a light-shielded environment and measuring them immediately afterwards.

<sup>b</sup> The values in parentheses are in eV.



Fig. 3 (a) Transient absorption (TA) spectra at the time delay of 20  $\mu$ s and (b) decay time profiles of the TA signal at 490 nm, measured for  $\text{Au}_2\text{Cu}_6$  (270  $\mu\text{M}$ ) with 0 and 10 mM perylene in deaerated THF at 298 K (excitation wavelength,  $\lambda_{\text{ex}}$  = 532 nm).

(Figs. S5a and b, ESI<sup>†</sup>), which is shorter than the phosphorescence lifetime of 4.9  $\mu$ s. This difference is attributed to the higher cluster concentration (20 $\times$ ) required for the TA measurement compared to that for the emission-lifetime measurement, owing to the weak TA signal. At such a high concentration, T-T annihilation (TTA) may occur between the clusters, shortening their lifetime. With perylene (10 mM), the decay time of the TA band at the 450–490 nm region was significantly lengthened, as shown in Fig. S5c (ESI<sup>†</sup>), and a new TA peak around 490 nm was observed with a time delay of 20  $\mu$ s, as shown in Fig. 3a. Since this peak was not observed without perylene, it can be attributed to the T-T absorption of perylene.<sup>53,54</sup> The decay time of this T-T absorption was determined to be 7.4  $\mu$ s, which was considerably shorter than the triplet lifetime of perylene in deaerated THF ( $\sim 0.7$  ms) because of the TTA between the perylene molecules and the quenching by the cluster sensitizer. These results demonstrate that the band-II emission was phosphorescence and that its quenching was caused by TET from the cluster in the excited triplet state to perylene in the  $S_0$  state (*i.e.*, the triplet sensitization of perylene [ $E_T$  = 1.53 eV] by  $\text{Au}_2\text{Cu}_6$ ). From the Stern–Volmer analysis (see ESI<sup>†</sup>), the rate constant of this TET process ( $k_{\text{TET}}$ ) was determined to be  $1.9 \times 10^7 \text{ s}^{-1}$ . When 9,10-diphenylanthracene (DPA) ( $E_T$  = 1.77 eV)<sup>55</sup> was employed as the triplet quencher, the TET rate constant was dramatically reduced to  $2.3 \times 10^5 \text{ s}^{-1}$  (Fig. S6 and Table S1, ESI<sup>†</sup>). When anthracene ( $E_T$  = 1.84 eV)<sup>56</sup> was used, no phosphorescence quenching (*i.e.*, TET) occurred. Therefore, we inferred that the  $T_1$  state of  $\text{Au}_2\text{Cu}_6$  is located around 1.7–1.8 eV, within 0.2 eV below the absorption edge of  $S_1$  (1.89 eV). The triplet perylene sensitization ability of  $\text{Au}_2\text{Cu}_6$  suggests that TF-UC can occur in this sensitizer-emitter pair, which will be discussed in the next subsection.

### TF-UC in the $\text{Au}_2\text{Cu}_6$ –perylene pair

Fig. 4a shows the PL spectrum obtained when a solution of  $\text{Au}_2\text{Cu}_6$  (11  $\mu\text{M}$ ) with 5 mM perylene was excited with a 532 nm-CW laser. Along with the fluorescence and phosphorescence of the cluster at 600–900 nm, perylene fluorescence was observed in the wavelength region shorter than the excitation wavelength. As shown in Fig. 4b, the time profile of this fluorescence signal exhibited sub-millisecond order decay. Since the upconverted emission intensity ( $I$ ) is proportional to the square of the triplet emitter ( $^3\text{E}^*$ ) concentration, its kinetics,  $I(t)$ , is given by

$$I(t) = I(0) \left( \frac{1 - \beta}{\exp(t/\tau_T) - \beta} \right)^2 \quad (1)$$

where  $\tau_T$  is the lifetime of  $^3\text{E}^*$ , and  $\beta$  is a dimensionless parameter between 0 and 1.<sup>57</sup> Details are provided in ESI<sup>†</sup>. By fitting the decay curve using eqn (1),  $\tau_T$  = 0.48 ms and  $\beta$  = 0.83 (Fig. S7, Tables S2 and S3, ESI<sup>†</sup>) were obtained. Notably, the value of  $\tau_T$  is five orders of magnitude higher than that of the prompt perylene fluorescence (5.0 ns in THF<sup>7</sup>), indicating that a delayed fluorescence *via* TTA occurred between the triplet perylene molecules, *i.e.*, TF-UC, triggered by the photoexcitation of the  $\text{Au}_2\text{Cu}_6$  sensitizer.

As depicted in Fig. 5, when a relatively high concentration of  $^3\text{E}^*$  is generated *via* the sensitizer-emitter TET process, TTA can occur efficiently between the  $^3\text{E}^*$  molecules because of their



Fig. 4 (a) PL spectrum obtained for the deaerated THF solution of  $\text{Au}_2\text{Cu}_6$  (11  $\mu\text{M}$ ) and perylene (5 mM), and (b) decay curve of the upconverted emission of perylene obtained by monitoring only the  $\lambda \leq 492$  nm region using short-pass filters. The fitting of the decay curve was performed using eqn (1). The 532 nm-CW laser was used in these measurements. Note that the upconverted emission of perylene was also observed when using a 640 nm-CW laser.



Fig. 5 Schematic diagram of TF-UC occurring in the system of a sensitizer–emitter pair.

very long lifetimes. This process causes one emitter molecule to return to the  $S_0$  state and the other to transition to the  $S_1$  state (*i.e.*, the UC state,  $^1E^*$ ), from which UC fluorescence is generated.<sup>20,58,59</sup> The TF-UC efficiency ( $\Phi_{UC}$ ), conventionally determined by relative estimation using eqn (S1) (ESI<sup>†</sup>), corresponds to the number of observed UC photons divided by the number of photons absorbed by the sensitizer ( $\#^1S^*$ ). Therefore,  $\Phi_{UC}$  is not necessarily a quantity unique to the TF-UC system, as it is strongly affected by several loss factors depending on the sample and experimental conditions, as mentioned below. Conversely, the formation yield of the UC state ( $\Phi_{UCs}$ ), *i.e.*, the number of upconverted states generated ( $\#^1E^*$ ) per  $\#^1S^*$ , is unique to the system; thus, the following relationship holds:<sup>60</sup>

$$\Phi_{UCs} = \Phi_{ISC} \Phi_{TET} \Phi_{TTA} \quad (2)$$

Note that the theoretical maximum value of  $\Phi_{UCs}$  is 0.50 (50%) because  $\Phi_{TTA} \leq 0.50$ . As can be seen in eqn (2),  $\Phi_{UCs}$  is proportional to the  $\Phi_{ISC}$  of the sensitizer, and the  $\Phi_{ISC}$  of the cluster can be estimated by evaluating  $\Phi_{UCs}$ .<sup>7</sup> Experimentally, the value of  $\Phi_{UCs}$  is determined by appropriately correcting for the loss of UC photons due to the reabsorption of UC photons by the sensitizer and emitter molecules (*i.e.*, optical outcoupling yield,  $\Phi_{out}$ )<sup>61</sup> and quenching of  $^3E^*$  by the sensitizer (*i.e.*, emitter quenching yield,  $\Phi_q$ ) against  $\Phi_{UC}$ .<sup>57</sup> This is expressed in the following equation:

$$\Phi_{UCs} = \frac{\Phi_{UC}}{\Phi_{out}(1 - \Phi_q)\Phi_f} \quad (3)$$

where  $\Phi_f$  represents the fluorescence quantum yield of the emitter.<sup>60</sup> Therefore, if the values of  $\Phi_{TET}$  and  $\Phi_{TTA}$  are known, the  $\Phi_{ISC}$  value can be obtained using eqn (2). The  $\Phi_{TET}$  can be determined by Stern–Volmer analysis, as shown in Fig. 2c, Fig. S6, and Table S1 (ESI<sup>†</sup>). However, since the  $\Phi_{TTA}$  strongly depends on the emitter concentration, solvent type, intensity of the excitation laser, and so on, great care must be taken when using the values obtained from the literature. Moreover, since TF-UC measurements are usually performed at high emitter concentrations ( $\geq 1$  mM), the value of  $\Phi_f$  is most likely to be

considerably different from that of the dilute solutions of the emitter.<sup>62</sup> To avoid uncertainty in the values of  $\Phi_{TTA}$  and  $\Phi_f$ , which are quantities related only to the emitter, we, herein, propose the following relative method.

Using a representative metal porphyrin sensitizer (*e.g.*, PtOEP) having a known  $\Phi_{ISC}$  (denoted herein as  $\Phi_{ISC}^r$ ) and the same emitter as a reference system, we evaluate its  $\Phi_{UCs}$  at the same emitter concentration, in the same solvent, and under the same measurement conditions (*e.g.*, optical setup and excitation wavelength). Note that yields marked with “r” on the right shoulder indicate those of the reference system thereafter. Applying the ratio in eqn (2) to the reference and target sample,  $\Phi_{TTA}$  can be cancelled out as follows:

$$\frac{\Phi_{UCs}}{\Phi_{UCs}^r} = \frac{\Phi_{ISC} \Phi_{TET}}{\Phi_{ISC}^r \Phi_{TET}^r} \quad (4)$$

Using eqn (3), the left-hand side of eqn (4) can be expressed and experimentally determined by the following equation:

$$\frac{\Phi_{UCs}}{\Phi_{UCs}^r} = \frac{\Phi_{UC}/\Phi_{out}(1 - \Phi_q)}{\Phi_{UC}^r/\Phi_{out}^r(1 - \Phi_q^r)} \quad (5)$$

By combining eqn (4) and (5), we can estimate the  $\Phi_{ISC}$  of the sensitizer of interest by the following equation:

$$\Phi_{ISC} = \frac{\Phi_{ISC}^r \Phi_{TET}^r}{\Phi_{TET}} \cdot \frac{\Phi_{UC}/\Phi_{out}(1 - \Phi_q)}{\Phi_{UC}^r/\Phi_{out}^r(1 - \Phi_q^r)} \quad (6)$$

Notably, this equation does not include the parameters of the emitter,  $\Phi_{TTA}$  and  $\Phi_f$ .

The value of  $\Phi_{TET}$  at a given perylene concentration can be estimated using the  $K_{sv}$  value ( $93.9 \text{ M}^{-1}$ ) obtained from the Stern–Volmer plot shown in Fig. 2c. For example,  $\Phi_{TET} = 0.32$  under the emitter concentration condition ( $[\text{perylene}] = 5 \text{ mM}$ ) in Fig. 4a, where the prominent UC emission was observed. In this study, PtOEP with  $\Phi_{ISC} \approx 1.0$  was used as the reference sensitizer,<sup>63</sup> and a series of parameters that need to be substituted into eqn (6) were evaluated for the pair of PtOEP sensitizer and perylene emitter in deaerated THF (Table 2). The parameters of the  $\text{Au}_2\text{Cu}_6$ –perylene system were evaluated following the same process under the same experimental conditions, and the  $\Phi_{ISC}$  of the cluster was finally determined. As summarized in Tables 2 and 3, through this method, the

Table 2 Measured and corrected upconversion quantum yields and their related parameters

$\lambda_{ex}^a/\text{nm}$	Sensitizer	Emitter	$[E]/\text{mM}$	$\Phi_{TET}$	$\Phi_{UC}^b$	$\Phi_{out}^c$	$1 - \Phi_q^d$	$\Phi_{UCs}^e$
532	$\text{Au}_2\text{Cu}_6$	Perylene	5.0	0.32	0.0018	0.60	0.65	0.0049
	PtOEP			0.99	0.0052	0.64	0.27	0.032

<sup>a</sup> The 532 nm excitation light intensity,  $I_{ex}$  ( $\text{W cm}^{-2}$ ), was set to the intensity region where the  $\Phi_{UC}$  becomes constant ( $I_{ex} = 8.04$  and  $3.09 \text{ W cm}^{-2}$  for the reference and sample systems, respectively). <sup>b</sup> PDI ( $\Phi_f = 0.97$  in toluene) was used as the reference dye, and the values were calculated by the relative method using eqn (S1) (ESI). <sup>c</sup> Calculated using eqn (S8) (ESI). <sup>d</sup> Calculated using eqn (S7) (ESI). <sup>e</sup> Calculated using eqn (3).



**Table 3** Quantum yields and rate constants for the radiative and nonradiative processes of the  $\text{Au}_2\text{Cu}_6(\text{S-Adm})_6(\text{PPh}_3)_2$  cluster in deaerated THF at room temperature

$T/\text{K}$	$\Phi_f$	$\Phi_{\text{ISC}}^a$	$\Phi_{\text{IC}}^b$	$\Phi_p$	$\Phi'_{\text{ISC}}^c$	$k_f^d/\text{s}^{-1}$	$k_{\text{ISC}}^e/\text{s}^{-1}$	$k'_{\text{IC}}^f/\text{s}^{-1}$	$k_p^g/\text{s}^{-1}$	$k'_{\text{ISC}}^h/\text{s}^{-1}$
293	$5.1 \times 10^{-3}$	0.47	0.52	0.44	0.56	$5.7 \times 10^{5c}$ (1.75 $\mu\text{s}$ )	$5.3 \times 10^7$	$5.8 \times 10^7$	$9.0 \times 10^4$ (11.1 $\mu\text{s}$ )	$1.1 \times 10^5$

<sup>a</sup> Intersystem crossing (ISC) quantum yield from  $S_1$  to  $T_1$  ( $\Phi_{\text{ISC}}$ ) was calculated using eqn (6). <sup>b</sup> Internal conversion (IC) quantum yield ( $\Phi_{\text{IC}}$ ) was calculated by  $1 - (\Phi_f + \Phi_{\text{ISC}})$ . <sup>c</sup> ISC quantum yield from  $T_1$  to  $S_0$  ( $\Phi'_{\text{ISC}}$ ) was calculated by  $1 - \Phi_p$ . <sup>d</sup> Radiative rate constant from  $S_1$  to  $S_0$  (*i.e.*, fluorescence transition) was calculated by  $\Phi_f \tau_f^{-1}$ . The value in parentheses represents the radiative lifetime of  $S_1$ . <sup>e</sup> Calculated by  $\Phi_{\text{ISC}} \tau_f^{-1}$ . <sup>f</sup> Calculated by  $\Phi_{\text{IC}} \tau_f^{-1}$ . <sup>g</sup> Radiative rate constant from  $T_1$  to  $S_0$  (*i.e.*, phosphorescence transition) was calculated by  $\Phi_p \tau_p^{-1}$ . <sup>h</sup> Calculated by  $\Phi'_{\text{ISC}} \tau_p^{-1}$ .

$\Phi_{\text{UCS}}$  for the  $\text{Au}_2\text{Cu}_6$ -perylene system was estimated to be 0.0049 (0.49%), resulting in an estimated  $\Phi_{\text{ISC}}$  of 0.47 for the  $\text{Au}_2\text{Cu}_6$  cluster. Thus, the phosphorescence quantum yield,  $\Phi_p$ , was determined to be 0.44 from the relationship,  $\Phi_{\text{PL}} = \Phi_{\text{ISC}} \cdot \Phi_p = 0.21$  (see Table 1), and the ISC quantum yield from  $T_1$  to  $S_0$  ( $\Phi'_{\text{ISC}}$ ) was 0.56. All the quantum yields and rate constants determined for the radiative and non-radiative processes are summarized in Table 3. The results show that the  $\text{Au}_2\text{Cu}_6$  cluster has relatively high quantum yields of ISC and phosphorescence at room temperature and, in contrast, very small fluorescence quantum yields. Notably,  $\text{Au}_2\text{Cu}_6$  possesses a very long fluorescence radiative lifetime (1.75  $\mu\text{s}$ ) and a small ISC rate of  $5.3 \times 10^7 \text{ s}^{-1}$ , which are strikingly different from those of many transition-metal complexes that exhibit ultrafast ISC.<sup>64–66</sup>

### Temperature effects on photophysical processes

To gain a profound understanding of the radiative and non-radiative processes in  $\text{Au}_2\text{Cu}_6$ , we measured the temperature ( $T$ ) dependence of the PL spectrum (Fig. 6a) and lifetime (Fig. 6b).



**Fig. 6** Temperature dependence of (a) PL; (b) PL decay; and (c) quantum yields of fluorescence ( $\Phi_f$ ), phosphorescence ( $\Phi_p$ ), internal conversion ( $\Phi_{\text{IC}}$ ), and ISC ( $\Phi_{\text{ISC}}$ ) of  $\text{Au}_2\text{Cu}_6$  (20  $\mu\text{M}$ ) in deaerated THF. Note that both band-I and -II emissions were monitored in the lifetime measurements. (d) Arrhenius-type plot of the estimated ISC rate constant versus the inverse of the absolute temperature for  $\text{Au}_2\text{Cu}_6$  in deaerated THF. The solid line represents the best fit of the data to eqn (7).

The results revealed the  $T$ -dependence of the quantum yield of all the radiative and non-radiation processes (Fig. 6c and Tables S4, S5, ESI†). The fluorescence intensity, lifetime, and quantum yield increased with decreasing temperature, accompanied by a concomitant decrease in the phosphorescence intensity, suggesting that the ISC process is suppressed at relatively low temperatures. Indeed, the  $\Phi_{\text{ISC}}$  and rate constant ( $k_{\text{ISC}}$ ), which were estimated from the  $T$ -dependence of  $\Phi_p$  and  $\Phi_{\text{PL}}$  (Table S5, ESI†) by assuming that the singlet and triplet radiative rate constants ( $k_f$  and  $k_p$ , respectively) were temperature-independent,<sup>67</sup> exhibited a remarkable decline with decreasing temperature. This fact indicates that the ISC proceeds *via* a thermally activated pathway. Thus, we assume that  $k_{\text{ISC}}$  can be expressed by an Arrhenius-type equation (eqn (7)).<sup>68,69</sup>

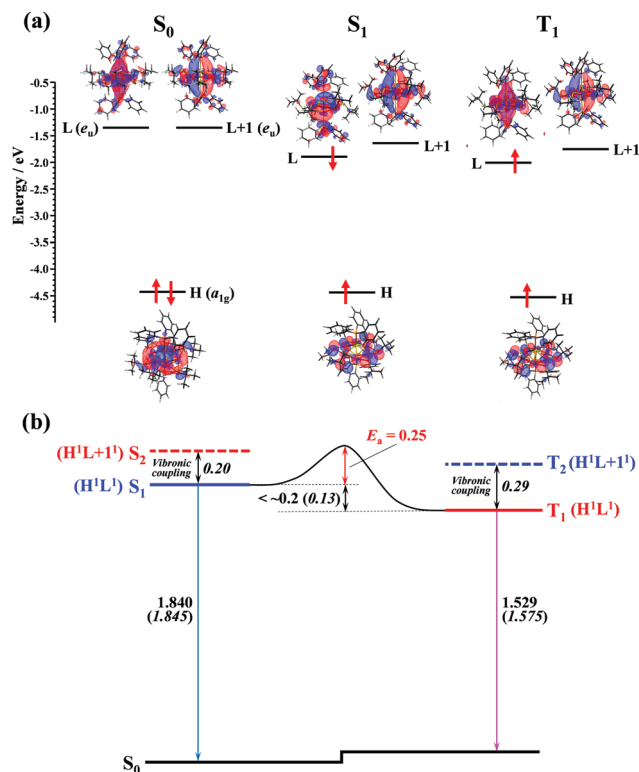
$$k_{\text{ISC}} = k_0 + \frac{A}{\sqrt{k_B T}} \exp\left(-\frac{E_a}{k_B T}\right) \quad (7)$$

where  $k_0$  and  $E_a$  represent the temperature-independent rate constant and the activation energy for the ISC process, respectively. The Arrhenius-type plot shown in Fig. 6d was fitted well by eqn (7), yielding  $A = 1.34 \times 10^{13} \text{ cm}^{1/2} \text{ s}^{-1}$  and  $E_a = 1994 \text{ cm}^{-1}$  (0.25 eV). Note that  $k_0$  is very small ( $< 10^3 \text{ s}^{-1}$ ) and can be neglected compared to the second term. The rate constants of the internal conversion (IC) process also decreased with a decrease in the temperature from 293 to 193 K (Table S5, ESI†), although the decrease was not as pronounced as that in the case of the ISC. However, an abrupt decrease was observed at temperatures below 193 K ( $k_B T = \sim 134 \text{ cm}^{-1}$ ), suggesting that the IC process may be highly promoted by low-frequency vibrational modes ( $< 140 \text{ cm}^{-1}$ ).

### Theoretical consideration

To gain insights into the mechanism of the thermally activated ISC with the small rate constant (*ca.*  $5 \times 10^7 \text{ s}^{-1}$ ), DFT and TD-DFT calculations for the ground and excited states of the  $\text{Au}_2\text{Cu}_6(\text{S-CMe}_3)_6(\text{PPh}_3)_2$  cluster were performed. In the  $S_0$  structure with  $D_{3d}$  symmetry, the LUMO (L) and LUMO+1 (L+1), which are mainly derived from  $\pi$ -like orbitals spreading along the P-Au-Au-P bond axis, were degenerate and orthogonal to each other (Fig. 7a). Therefore,  $S_1$  and  $S_2$  or  $T_1$  and  $T_2$  are degenerate in the  $S_0$  structure (Table S6, ESI†). However, in the equilibrium structure of the  $S_1$  and  $T_1$  states (Fig. S9, ESI†), the symmetry lowering (Fig. S10, ESI†) induced the weak splitting between L and L+1 (*i.e.*, the  $S_1$  and  $S_2$  [or  $T_1$  and  $T_2$ ] states split energetically), as shown in Fig. 7. The calculated vertical transition energies from  $S_1$  to  $S_0$  (1.845 eV) and  $T_1$  to  $S_0$  (1.575 eV) agree well with the energies of the fluorescence





**Fig. 7** (a) Computed energy-level diagram for the optimized structures of  $S_0$ ,  $S_1$ , and  $T_1$  (lower panel). The electronic configuration of each electronic state is shown with upward ( $\alpha$ -spin) and downward ( $\beta$ -spin) arrows in the notations HOMO (H), LUMO (L), and LUMO+1 (L+1). (b) Schematic representation of the ISC mechanism of the  $\text{Au}_2\text{Cu}_6$  cluster. All values are in eV. Numbers in *italics* are the calculated values, while other numbers are estimated from the experiments.

(1.840 eV) and phosphorescence (1.529 eV) emission peaks (Table 1 and Fig. 7b). The shape of the HOMO (H), localized mainly on the Cu<sub>6</sub>S<sub>6</sub> ring, is the same between the S<sub>1</sub> and T<sub>1</sub> states, and the S<sub>1</sub> and T<sub>1</sub> states have the same H<sup>1</sup>L<sup>1</sup> electronic configuration. In the S<sub>1</sub> state, although the shape of the LUMO is somewhat different from that of T<sub>1</sub>, there is no orthogonality between the LUMO (precisely, singly occupied molecular orbital on the high energy side) of S<sub>1</sub> and that of T<sub>1</sub>. According to the El-Sayed rules,<sup>70</sup> therefore, the direct spin-orbit coupling (SOC) between S<sub>1</sub> and T<sub>1</sub> is considered to be inactive. In fact, the temperature-independent term (*k*<sub>0</sub>) corresponding to the S<sub>1</sub> → T<sub>1</sub> direct tunneling was negligibly small. Conversely, the S<sub>2</sub> and T<sub>2</sub> states have the H<sup>1</sup>L + 1<sup>1</sup> electronic configuration (Fig. 7b and Table S6, ESI<sup>†</sup>), thereby allowing the direct SOC between S<sub>2</sub> and T<sub>1</sub>, as well as between S<sub>1</sub> and T<sub>2</sub>. It is also theoretically suggested that the S<sub>2</sub> and T<sub>2</sub> states are energetically close to the S<sub>1</sub> and T<sub>1</sub> states (<0.3 eV). Therefore, ISC should be partially allowed *via* the spin-vibronic coupling, mediated by the S<sub>2</sub> and T<sub>2</sub> states.<sup>70</sup> Since the SOC constant of the Au atom is considerably larger than those of the light elements (S, P, C, and H atoms) that constitute the organic ligands, the L and L+1, which are mainly derived from the Au dimer moiety, are expected to contribute significantly to the S<sub>2</sub>-T<sub>1</sub> and S<sub>1</sub>-T<sub>2</sub> SOC matrix elements. Currently, there is no experimental evidence of the strong

coupling (*i.e.*, avoided crossing) between the  $S_1$  and  $T_1$  potential energy surfaces. This may be responsible for the activation barrier in the ISC process. The reverse ISC from  $T_1$  to  $S_1$  is unlikely to occur because the activation energy is as large as 0.5 eV.<sup>71</sup> Consequently, no thermally activated delayed fluorescence was observed in this cluster.

## Conclusions

This study demonstrates that  $\text{Au}_2\text{Cu}_6$  clusters exhibit both fluorescence and phosphorescence at room temperature. Additionally, the quantum yields and rate constants of all radiative and non-radiative processes in this cluster were determined by analyzing the TF-UC phenomenon induced when  $\text{Au}_2\text{Cu}_6$  is employed as a triplet sensitizer. Temperature-dependent fluorescence/phosphorescence measurements and theoretical calculations suggested that thermally activated ISC occurs *via* the spin–vibronic coupling mechanism mediated by the SOC-allowed  $S_2$  and  $T_2$  states. Metal clusters composed of heavy-metal elements, such as Au and Ag atoms, possess considerably larger SOC constants than those of organic molecules and relatively small S–T energy gaps, which inevitably leads to considerably stronger S–T mixing. Consequently, it is highly probable that excited states with a triplet nature are involved in various excited-state relaxation processes, including PL phenomena. Therefore, the approach reported herein, based on triplet sensitization by metal clusters, will be very useful not only for elucidating the origin of PL in clusters but also for expanding their application as photosensitizers.

## Author contributions

Daichi Arima: experiment, data acquisition, analysis, and drafted the initial manuscript. Yoshiaki Niihori: experiment, revision, and funding acquisition. Masaaki Mitsui: supervision, writing – revision and editing, theoretical calculations, and funding acquisition.

## Conflicts of interest

There are no conflicts to declare.

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