

## PAPER

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11, 8146Organic ammonium iodide salts as passivation  
for buried interface enables efficient and stable  
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The interfaces between a halide perovskite absorber film and its charge transporting layers have been identified as one of the key bottlenecks for achieving high performance and stability. Herein we report a benzothieno[3,2-*b*]benzothiophene (BTBT) derivative and a carbazole derivative to passivate the interface between the sputtered NiO<sub>x</sub> and the halide perovskite in p-i-n perovskite solar cells (PSCs). The treatment of the NiO<sub>x</sub> with the BTBT derivative significantly enhances the open-circuit voltage (*V*<sub>oc</sub>), boosting the champion power conversion efficiency (PCE) from 16.6% to 18.6%. We carry out an in-depth analysis of the interface employing a lift-off procedure revealing that the BTBT molecules reside at the bottom interface and successfully suppress a (redox) reaction at this interface. As a result of the improved interface quality, the stability of the PSCs is superior under maximum power point (MPP) conditions and in ambient conditions, retaining 89% of the initial PCE after 250 hours at 25% relative humidity and 95% after 1 hour under continuous MPP tracking (MPPT).

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## 10th Anniversary Statement

Over the past ten years, the Journal of Materials Chemistry C has continued along the path of the Journal of Materials Chemistry as a prime destination for studies that combine fundamental advances in materials chemistry with device aspects. This also holds for the hybrid perovskite field, with many inspiring and impactful papers being published in the journal. The topic of the current manuscript, interfacial engineering of perovskite solar cells, has received significant attention over the past years. With the current manuscript, we show that designed large organic ammonium cations can be employed as interlayers at the hole-transporting layer/perovskite interface to enhance the power conversion efficiency (PCE) and stability of inverted architecture perovskite solar cells (PSCs). While inverted architecture PSCs generally reach lower PCEs than their regular architecture counterparts, we are convinced that interfacial engineering can greatly reduce this gap and that designed large organic molecules possess significant potential as interlayer materials.

## 1. Introduction

Perovskite solar cells have emerged as a promising photovoltaic technology because of their excellent performance and the possibility to be processed at a low cost. Over a decade of intense research on PSCs has led the power conversion efficiency (PCE) to reach a certified 25.7% and 31.3% for single junction and perovskite/silicon (Si) tandem solar cells, respectively.<sup>1</sup> Despite the high performance, multiple bottlenecks need to be overcome towards the commercialization of PSCs. Firstly, as PSCs are mainly solution-processed devices, the performance is still significantly inferior on a larger scale compared to other photovoltaic (PV) technologies. Secondly, the perovskite material has the tendency to degrade in outdoor conditions, such as high temperatures, UV exposure, and

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ambient atmosphere.<sup>2,3</sup> Consequently, the high initial performance declines rather quickly over time in normal operating conditions. Many approaches are being investigated to improve the stability of PSCs, such as compositional<sup>3</sup> and interfacial engineering<sup>4,5</sup> of the perovskite layer, development of more stable transport layers, and process engineering of all the layers.<sup>6</sup>

Typically, PSCs are fabricated in one of two architectures, *viz.* n-i-p and p-i-n (inverted). In the n-i-p configuration, the electron transporting layer (ETL) is processed first, followed by the perovskite layer and the hole transporting layer (HTL). Whereas in p-i-n, the HTL is processed first, followed by the perovskite layer and the ETL. The majority of the research has been conducted on n-i-p cells, and the highest PCEs have also been obtained in this architecture. However, these n-i-p cells frequently use TiO<sub>2</sub> as ETL, which is prone to induce degradation under UV light,<sup>7</sup> and they generally employ doped organic semiconductors as HTL, such as spiro-OMeTAD, which are unstable at high temperatures<sup>8</sup> or polymers like poly(triaryl)-amine (PTAA) which are costly, suffer from batch to batch variation and need to be doped with hygroscopic dopants. Therefore, research attention has lately been shifting to p-i-n cells since they are generally more stable and can employ cheap transport materials such as inorganic metal oxides (*e.g.*, NiO<sub>x</sub>), which can be deposited as HTL in this configuration. Additionally, these inorganic materials can be processed with large-scale compatible deposition techniques such as sputtering<sup>9</sup> and thermal evaporation. Nevertheless, the PCE of p-i-n PSCs is still lagging compared to n-i-p PSCs, certainly in the case when sputtered NiO<sub>x</sub><sup>10–12</sup> is employed as HTL, where mainly the inferior open-circuit voltage,  $V_{oc}$ , is at the origin of the lower PCE.<sup>12</sup> It is argued that the lower  $V_{oc}$  for NiO<sub>x</sub>-based HTLs originates from an inferior energy band alignment of the valence band maxima relative to the perovskite layer. Therefore, doping of the solution-processed NiO<sub>x</sub> layer with Li, Mg, and Cu has been an approach to improve the alignment.<sup>13,14</sup> However, Caleb *et al.*<sup>12</sup> showed the formation of a hole blocking PbI<sub>2–x</sub>Br<sub>x</sub> layer at the NiO<sub>x</sub>/perovskite interface due to redox reactions. Excess formamidinium iodide (FAI) in the halide perovskite precursor was employed, which suppressed

the formation of the hole blocking PbI<sub>2–x</sub>Br<sub>x</sub>, resulting in a significantly enhanced  $V_{oc}$ , close to the  $V_{oc}$  when PTAA is used as HTL. Liu *et al.*,<sup>15</sup> introduced phenethylammonium iodide (PEAI) to modify the NiO<sub>x</sub>/perovskite interface, which simultaneously contributed to enhanced crystallinity and stability of the halide perovskite layer, passivating interface defects, forming a two-dimensional PEA<sub>2</sub>PbI<sub>4</sub> perovskite layer, and superior interface contact properties. Other organic molecules, such as guanidinium (Gua), ethylammonium (EA), *n*-butylammonium (BA) and dimethylammonium (DEA) have also been employed at the HTL/perovskite interface and it was found that the binding affinity of the molecules to the substrates is one of the key determinants to form a 2D phase at the buried interface.<sup>16</sup> In the work by Chen *et al.*<sup>16</sup> only Gua was able to form a 2D phase (Gua<sub>2</sub>PbI<sub>4</sub>) and was therefore used as passivation layer. However, the influence on the photovoltaic performance of PSCs was not shown for the other organic molecules, which don't seem to form a 2D phase according to the XRD results. Therefore, it is not clear whether the capability of an organic molecule to form a 2D phase at the bottom interface is essential to act as a passivation layer.

In this work, two large organic ammonium iodide salts, a carbazole (Cz) derivative and a benzothieno[3,2-*b*]benzothiophene (BTBT) derivative are used as an interlayer between sputtered NiO<sub>x</sub> and halide perovskite, as shown in Fig. 1. The BTBT derivative was recently used by some of the authors to form quasi-2D cesium lead iodide perovskites with significantly enhanced stability compared to butylammonium-based ones.<sup>17</sup> In the current work, the potential of this BTBT derivative as a passivator at HTL/perovskite interfaces is evaluated. Photoluminescence quantum yield (PLQY) measurements show that both interlayer materials have the capability to passivate the interfacial defects, with  $V_{oc}$  enhancements of 40 mV and 100 mV for the Cz and BTBT interlayers, respectively, compared to the control devices. External quantum efficiency (EQE) measurements reveal that the BTBT interlayer results in a superior charge carrier extraction at the NiO<sub>x</sub>/perovskite interface compared to the Cz-based one and consequently gives rise to a higher short-circuit current density ( $J_{sc}$ ), which is potentially



Fig. 1 (a) Molecular structure of BTBT and Cz. (b) Perovskite solar cell stack depiction. The control sample does not contain any modification between NiO<sub>x</sub> and the perovskite layer.

related to the well-known high charge carrier mobility of BTBT derivatives.<sup>18–21</sup>

## 2. Results and discussion

The BTBT or Cz ammonium iodide salts were spin-coated from dimethylformamide (DMF) on top of the HTL, followed by a drying step. Next, the 3D halide perovskite was deposited *via* spin coating and annealed (details are present in the ESI†). The sample without the interlayer is referred to as control, whereas the samples with the carbazole and BTBT-based interlayer are referred to as Cz and BTBT, respectively.

### A. PL and TRPL

To get insights into the impact of the interlayers on the non-radiative interface recombination and on the hole extraction, we performed steady-state (SSPL) and time-resolved photoluminescence (TRPL) measurements on glass samples with and without HTL (NiO<sub>x</sub>). As shown in Fig. 2a, the films with both interlayers deposited directly on glass (without HTL) give rise to a very similar SSPL peak which is more intense than for the control film, indicating enhanced radiative recombination by the passivation of interfacial defects. The average lifetime of

the carriers,  $\tau_{\text{avg}}$ , was extracted from the TRPL data using a biexponential fit:

$$\tau_{\text{avg}} = \frac{A_1 \cdot \tau_1 + A_2 \cdot \tau_2}{A_1 + A_2}$$

The values for  $A_1$ ,  $A_2$ ,  $\tau_1$  and  $\tau_2$  are displayed in the ESI† (Table S1). Based on this fitting, an extended average lifetime of the photo-generated carriers is revealed, going from 196 ns for the control to 285 ns and 317 ns for the Cz and BTBT interlayers, respectively, indicating the suppression of non-radiative recombination with the interlayers. The TRPL measurements conducted on the substrates with NiO<sub>x</sub>, as shown in Fig. 2c and d, reveal a very slightly accelerated decay in the first tens of ns with both of the interlayers. Such accelerated decay could be either due to enhanced interfacial trap recombination or due to improved hole extraction. Given the improved lifetimes of the photo-generated carriers in the presence of the interlayers, as determined from the measurements on the glass substrates, it is suggested that the faster decay on NiO<sub>x</sub> is due to improved hole extraction. Nevertheless, the carriers have a very similar lifetimes (13.5–14.5 ns) with and without the interlayers, so the influence of the interlayers on the  $J_{\text{sc}}$  of solar



Fig. 2 (a) Steady-state PL and (b) time-resolved PL measurements for perovskite films on glass, with and without the interlayers. (c and d) Raw and fitted data of time-resolved PL measurements for perovskite films with and without interlayers on the HTL (NiO<sub>x</sub>).

cell devices is expected to be rather limited based on these measurements.

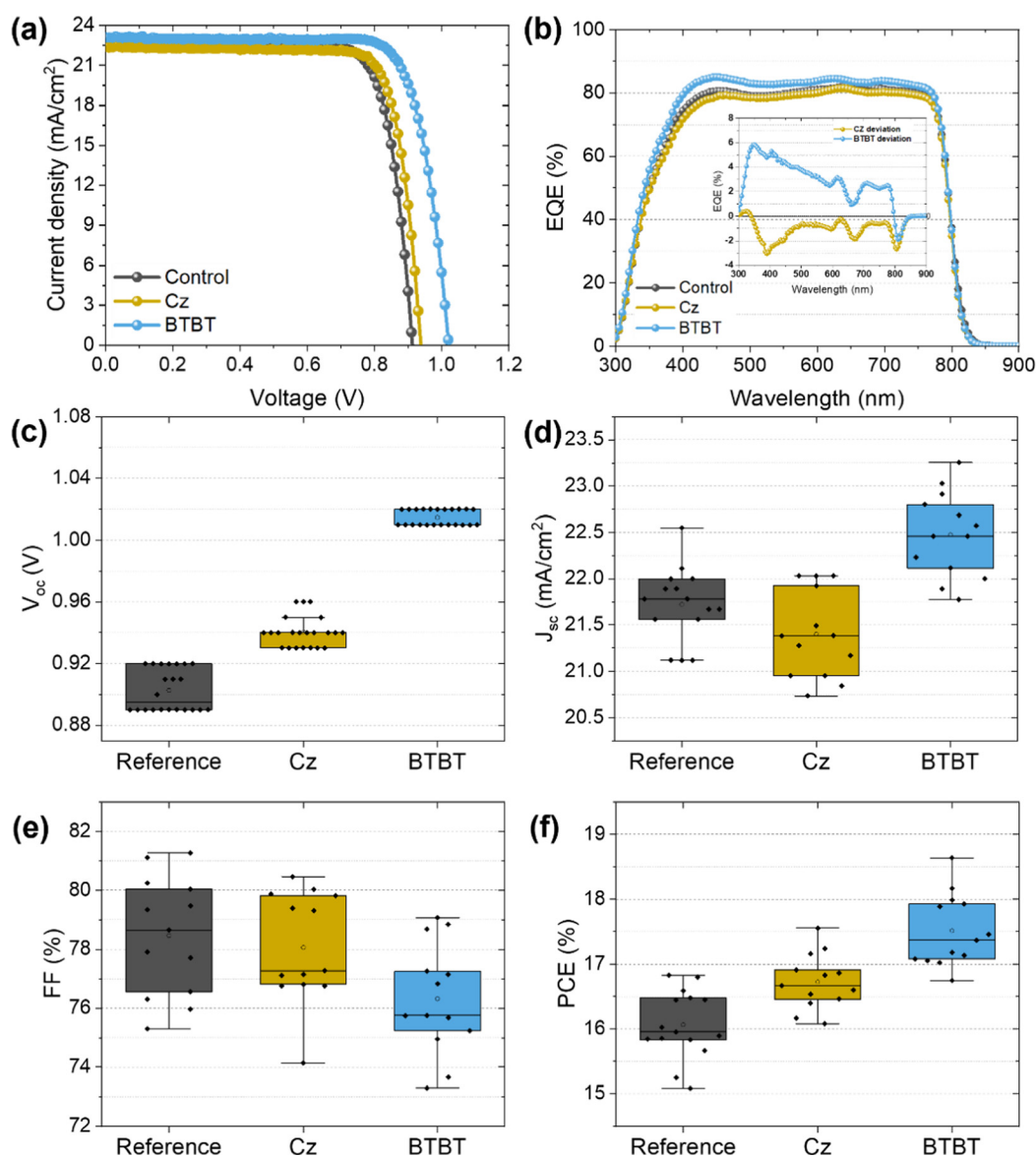
## B. JV and EQE

The photovoltaic performance is shown in Fig. 3 and summarized in Table 1. As seen from the *JV* curves in Fig. 3a, the interlayers mainly have an influence on the  $V_{oc}$ , which is to be expected for the passivation of interface defects. The improvement with the BTBT interlayer is quite significant compared to that based on Cz, exceeding the  $V_{oc}$  of the control sample ( $\sim 0.91$  V) with 100 mV compared to 30 mV using Cz. Small differences in  $J_{sc}$  were noticed as well, which are more clearly visualized by the EQE measurements shown in Fig. 3c. The inset figure represents the difference between the EQE signal of the cell with an interlayer and the control cell as a function of

**Table 1** Solar cell device parameters extracted from the *JV*-curves of the champion devices. The average values and deviation for the PCE is noted within the brackets

Interlayer	$J_{sc}$ [mA cm <sup>-2</sup> ]	$V_{oc}$ [V]	FF [%]	PCE [%]
N.A. (Control)	22.65	0.92	79.47	16.58 (16.0 ± 0.3)
Cz	22.43	0.94	80.46	16.94 (16.6 ± 0.2)
BTBT	23.10	1.02	79.07	18.64 (17.5 ± 0.4)

wavelength, such that the influence of the interlayers on the EQE at each wavelength is more apparent. As smaller wavelengths are mainly absorbed near the bottom interface of the perovskite film, where these interlayers are present according to HAXPES measurements (discussed in next section), the influence of the interlayers on the EQE signal is expected to be most significant for wavelengths in the range of 300–500 nm.



**Fig. 3** (a) *JV*-curves of the champion PSCs with and without the interlayers and (b) the corresponding EQE output. The statistical distribution of the measured (c–f)  $V_{oc}$ ,  $J_{sc}$ , FF and PCE.

Indeed, Fig. 3b reveals that the BTBT interlayer improves the EQE signal of the control film by approximately 5% (absolute) in the small wavelength region, while the Cz interlayer seems to slightly suppress the signal. This corresponds well with the measured  $J_{sc}$ , with values of  $22.65 \text{ mA cm}^{-2}$ ,  $22.43 \text{ mA cm}^{-2}$  and  $23.1 \text{ mA cm}^{-2}$  for the control film, Cz interlayer, and BTBT interlayer, respectively. These results suggest that the BTBT interlayer improves the charge carrier extraction at the interface, as was suggested by the TRPL measurements, which may either be the consequence of improved energy level alignment or the well-known superior charge-carrier mobility of BTBT molecules.<sup>15,18–20,22</sup> Since a lift-off procedure (*vide infra*) is only possible with PTAA as a sacrificial layer, PSCs with PTAA as HTL were prepared as well, for which the  $J$ - $V$  curves are shown in Fig. S2 (ESI†). For this HTL, an improvement in  $V_{oc}$  is achieved as well using both interlayers. In comparison with  $\text{NiO}_x$  as HTL, the improvement with the BTBT interlayer is however more limited and quite similar to that achieved with Cz as the interlayer. This suggests that the influence of the interlayers on the interface with  $\text{NiO}_x$  is more complex than for PTAA and likely related to the redox reactions which can occur at this interface.<sup>12</sup> For both HTLs, the champion efficiencies were achieved when using BTBT as an interlayer, achieving 18.6% and 20.6% for  $\text{NiO}_x$  and PTAA, respectively.

### C. XRD, XPS and HAXPES

To gain more insight into the interlayer formation, an analysis method is required which can probe this buried interface. As surface-sensitive techniques like XPS are not suitable in the normal solar cell configuration, a lift-off process<sup>23</sup> is applied (peeling off the perovskite film from its original substrate and attaching it in reverse orientation to a new substrate, as illustrated in Fig. S1, ESI†). The main requirement for this approach is the use of an (HTL) material that can be dissolved in a perovskite-compatible solvent. Unfortunately, the preferred HTL,  $\text{NiO}_x$ , does not allow for the use of such an approach, meaning that an alternative HTL like PTAA, which is usually dissolved in chlorobenzene or toluene, has to be used for this purpose.<sup>23,24</sup> The lift-off process is described in more detail in the ESI† (Fig. S1).

XRD measurements on the lift-off samples (Fig. 4a) do not reveal the formation of 2D halide perovskites, as no reflections are observed at angles below  $10^\circ$ . However, it is clear that the presence of the interlayers results in different ratios of the main perovskite reflection at  $\sim 14^\circ 2\theta$  and the main  $\text{PbI}_2$  reflection at  $12.8^\circ 2\theta$ . For the control film without interlayer, the main  $\text{PbI}_2$  reflection is more intense than the main perovskite reflection, suggesting that the  $\text{PbI}_2$  content is mainly present at the HTL/perovskite interface in the inverted cell configuration. With the Cz interlayer, the two reflections are of a similar magnitude, while for the BTBT interlayer, the intensity of the  $\text{PbI}_2$  reflection is significantly diminished. However, it is not clear whether the  $\text{PbI}_2$  content is present at the interface or in the bulk. In case the  $\text{PbI}_2$  content would be concentrated at the interface, the XRD results could, at least in part, explain the superior performance of the BTBT interlayer since  $\text{PbI}_2$  is known to act

as a hole blocking layer.<sup>11</sup> To verify whether the  $\text{PbI}_2$  content as determined by XRD is present at the interface or in the bulk, both XPS and HAXPES measurements were performed on lift-off samples, denoted as “Ref top” and “BTBT top” in Fig. 4b, and on a pure  $\text{PbI}_2$  film, as a reference, which was spin coated on a separate substrate. Typically, the information depth of standard XPS is limited to the very first few nanometers, making it extremely sensitive to ultrathin surface layers, while the harder X-rays of HAXPES allow to significantly increase the kinetic energy of photoelectrons thereby enabling information depths up to about 100 nm,<sup>24,25</sup> thus representing bulk properties. In order to establish the surface and bulk properties of the perovskite film w./w.o. the BTBT interlayer, the I-3d<sub>5/2</sub> core level was characterized using high energy resolution. Focusing first on the XPS results presented in Fig. 4c, a single peak can be recognized at a binding energy of 619.2 eV for the  $\text{PbI}_2$  reference. In contrast, the perovskite absorber w./w.o. BTBT reveals a main I-3d<sub>5/2</sub> core line at 619.4 eV (with no additional peak or shoulder at 619.2 eV suggesting no substantial contribution by  $\text{PbI}_2$ ) followed by a chemically shifted component at 621.4 eV which is significant for the pure perovskite reference, but much weaker when adding BTBT before perovskite deposition. The difference in peak position of the main I-3d<sub>5/2</sub> core line for both perovskites (w/wo BTBT layer) as compared to the  $\text{PbI}_2$  reference is confirmed by the more bulk sensitive HAXPES measurements which, however, do not show the additional chemically shifted component at 621.4 eV. By combining both experiments (XPS, HAXPES), we can conclude that (i) there is no substantial amount of  $\text{PbI}_2$  accumulated at the interface between the pure perovskite and its HTL (suggesting that the  $\text{PbI}_2$  content, revealed by the XRD measurements, is located more deeply in the perovskite bulk), but that (ii) there exists an interface layer of unknown origin (visible in the surface sensitive XPS, but not in the bulk-sensitive HAXPES measurements) which is strongly suppressed in the presence of the BTBT interlayer. The additional intensity observed in the I-3d<sub>5/2</sub> core level spectra at 621.4 eV suggests the presence of iodide in a different (less negative) oxidation state as compared to the perovskite and  $\text{PbI}_2$  reference. Higher binding energies have been reported in the presence of oxygen<sup>26</sup> which, in a simple picture, can be assigned to the stronger electronegativity of oxygen (3.44) as compared to iodine (2.66). Although a clear assignment of the measured peak position (621.4 eV) of this new spectral component is not possible at the moment, it can be speculated that the peak is the result of an oxidation reaction (*e.g.* with oxygen from the ITO that was used during the preparation of the lift-off films). Such an oxidation reaction is expected to occur as well when the perovskite film is deposited on top of  $\text{NiO}_x$  during the processing of the PSCs, which will be suppressed in the presence of BTBT. We can also speculate that the existence of an additional interlayer (in the absence of BTBT) might negatively influence the energy band alignment between the perovskite and the HTL, resulting in a reduced  $V_{oc}$ . The chemical identification of this new component and its impact on energy level alignment will require a dedicated study





Fig. 4 (a) XRD measurements performed on lift-off samples with and without the interlayers. (b) Illustration of the four different configurations used for XPS and HAXPES measurements. (c) XPS and HAXPES measurements performed on lift-off samples concentrating on the I-3d<sub>5/2</sub> core level. (d) HAXPES measurement performed on all configurations concentrating on the S-1s core level.

which is beyond the scope of the present work. Finally, as these Cz and BTBT ammonium salts can be dissolved in the same solvents that are used for the perovskite solution (DMF and DMSO), the interlayer could potentially be fully or partially dissolved during the spin coating step of the 3D halide perovskite on top of this layer. As a result, the interlayer molecules could hypothetically enter into the perovskite bulk structure. To assess whether this is the case, additional HAXPES measurements were performed on four different device stack configurations using PTAA as HTL, as shown in Fig. 4b. “Ref bottom” and “BTBT bottom” samples contain the normal configuration as used in the PSCs, meaning that they haven’t been subjected to a lift-off process. These samples are prepared on silicon substrates instead of glass/ITO substrates to facilitate the preparation of small samples by breaking, as required for the XPS and HAXPES measurements, and to prevent charging effects during these measurements. As mentioned earlier, “Ref top” and “BTBT top” samples are obtained by the lift-off process. For this experiment, we focused on the sulfur 1s core level since this element is only present in the BTBT molecules

and not in the 3D halide perovskite. Since Cz only contains elements that are the same as those in the 3D perovskite (H, N and C), we limited this study to BTBT. The HAXPES results for the four configurations are shown in Fig. 4d, presenting the S-1s core level region. As expected, for both configurations without the BTBT interlayer, no signal corresponding to sulfur is present. For the lift-off configuration with the BTBT interlayer, a very clear S-1s signal can be noticed which does not exist for the normal configuration with the BTBT interlayer. Based on these results, it can be concluded that the presence of the BTBT molecules is limited to the HTL/perovskite bottom interface and therefore, these interlayers passivate this interface, improving the optoelectronic properties and device performance as described above. It must be noted that while PTAA was used as the HTL instead of NiO<sub>x</sub> to perform this analysis, we expect the interaction of the interlayer molecules, *via* their ammonium (NH<sub>3</sub><sup>+</sup>) tethering group, to be stronger with the polar NiO<sub>x</sub> surface than with the more apolar PTAA surface. As such, the absence of BTBT molecules present in the bulk of the 3D halide perovskite as determined for the case of PTAA as

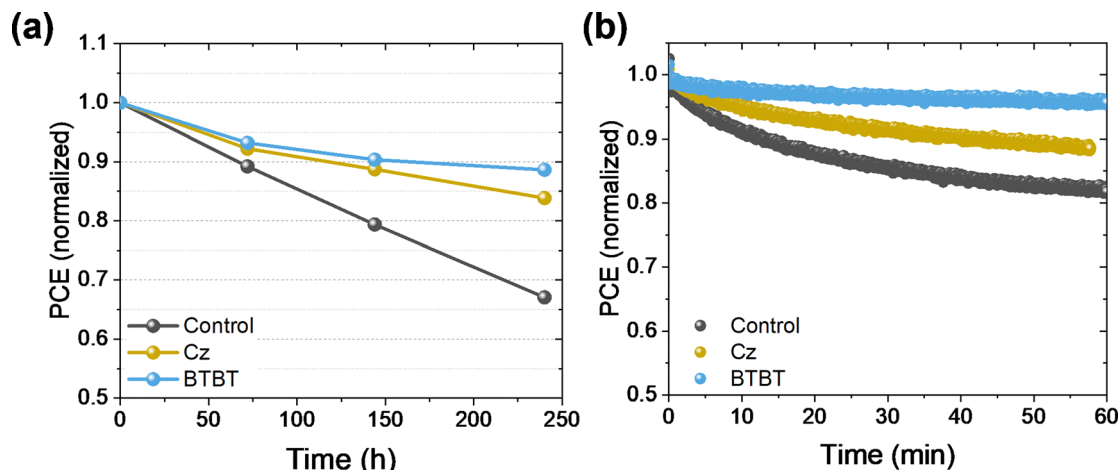


Fig. 5 Stability test on PSCs without encapsulation in (a) ambient conditions (25% RH) and (b) under continuous illumination and MPPT conditions.

HTL, is a strong indicator that the same will be true for the case of  $\text{NiO}_x$  as HTL.

#### D. Stability

The PSCs were exposed to a relative humidity (RH) of 25% in the dark for 250 hours, and the PCE was measured at intervals of 70–100 hours to track the decline of the performance, as shown in Fig. 5. For both interlayers, a clear improvement in the stability is obtained, retaining 84% and 89% of the initial performance for Cz and BTBT, respectively, while the control cell (*i.e.* without the Cz or BTBT interlayer) only retained 67% of the initial performance. The critical factor for these cells was the FF, as the  $J_{sc}$  and  $V_{oc}$  remained rather unchanged over the course of the tests (Fig. S3, ESI†), suggesting that the perovskite layer itself was not the most sensitive material to the stress conditions, but rather one of the interfaces. Given the positive effect of the interlayers on the stability, we hypothesize that the stabilization of the  $\text{NiO}_x$ /perovskite interface is at least partly responsible for the improved stability. Additionally, samples with and without the interlayers were exposed to continuous illumination for an hour while held under maximum power point (MPP) conditions. Over this time, the interlayers significantly improve the stability under operational conditions, as the control only retained 82% of its initial performance while 88% and 95% was retained for Cz and BTBT, respectively. This also shows that the BTBT interlayer has a profound impact on improving the stability at the interface.

### 3. Conclusion and outlook

In summary, we have developed two new interface passivation materials based on benzothieno[3,2-*b*] benzothiophene (BTBT) and carbazole (Cz). These materials were introduced at the  $\text{NiO}_x$ /perovskite interface in p-i-n perovskite solar cells, a key bottleneck for the device's performance and stability. The PSCs with the BTBT interlayer showed higher PCE and stability than the control and Cz-based devices. To get mechanistic insights, we performed XPS, HAXPS, PLQY, and TRPL measurements.

These results show that passivation of the  $\text{NiO}_x$ /perovskite interface *via* the BTBT derivative suppresses a (redox) reaction at this interface and significantly enhances the open-circuit voltage ( $V_{oc}$ ) by suppressing the non-radiative recombination. These results provide a design rationale for new passivation materials and self-assembled monolayers. We believe that the BTBT core can be combined with phosphonic acid functionalities to create new self-assembled monolayers in further work, given the success of Cz-based phosphonic acid SAMs.<sup>27–30</sup>

### Author contributions

Stijn Lammar: investigation, validation, visualization, writing – original draft, conceptualization Wouter Van Gompel: conceptualization, investigation, supervision, writing – review & editing Stijn Lenaers: investigation, writing – review & editing Martijn Mertens: investigation Hans-Gerd Boyen: investigation, supervision, writing – review & editing Derese Desta: investigation Afshin Hadipour: supervision, writing – review & editing Laurence Lutsen: project administration, supervision, funding acquisition Dirk Vanderzande: conceptualization, supervision Anurag Krishna: writing – review & editing, visualization Yaser Abdulraheem: funding acquisition Tom Aernouts: project administration, supervision, funding acquisition Jef Poortmans: supervision.

### Conflicts of interest

There are no conflicts to declare.

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## References

- 1 M. A. Green, E. D. Dunlop, J. Hohl-Ebinger, M. Yoshita, N. Kopidakis and X. Hao, *Prog. Photovoltaics*, 2022, **30**, 3–12.
- 2 C. C. Boyd, R. Checharoen, T. Leijtens and M. D. McGehee, *Chem. Rev.*, 2019, **119**, 3418–3451.
- 3 Z. Wang, Z. Shi, T. Li, Y. Chen and W. Huang, *Angew. Chem., Int. Ed.*, 2017, **56**, 1190–1212.
- 4 M. Ye, C. He, J. Iocozzia, X. Liu, X. Cui, X. Meng, M. Rager, X. Hong, X. Liu and Z. Lin, *J. Phys. D: Appl. Phys.*, 2017, **50**, 373002.
- 5 Z. Guo, Z. Wu, Y. Chen, S. Wang and W. Huang, *J. Mater. Chem. C*, 2022, **10**, 13611–13645.
- 6 H. Xiang, P. Liu, W. Wang, R. Ran, W. Zhou and Z. Shao, *Chem. Eng. J.*, 2021, **420**, 127599.
- 7 T. Leijtens, G. E. Eperon, S. Pathak, A. Abate, M. M. Lee and H. J. Snaith, *Nat. Commun.*, 2013, **4**, 2885.
- 8 W. Song, L. Rakocevic, R. Thiruvallur Eachambadi, W. Qiu, J. P. Bastos, R. Gehlhaar, Y. Kuang, A. Hadipour, T. Aernouts and J. Poortmans, *ACS Appl. Mater. Interfaces*, 2021, **13**, 44294–44301.
- 9 H. Lee, Y.-T. Huang, M. W. Horn and S.-P. Feng, *Sci. Rep.*, 2018, **8**, 5590.
- 10 X. Zheng, Z. Song, Z. Chen, S. S. Bista, P. Gui, N. Shrestha, C. Chen, C. Li, X. Yin, R. A. Awni, H. Lei, C. Tao, R. J. Ellingson, Y. Yan and G. Fang, *J. Mater. Chem. C*, 2020, **8**, 1972–1980.
- 11 N. Pant, A. Kulkarni, M. Yanagida, Y. Shirai, T. Miyasaka and K. Miyano, *Adv. Mater. Interfaces*, 2020, **7**, 1901748.
- 12 C. C. Boyd, R. C. Shallcross, T. Moot, R. Kerner, L. Bertoluzzi, A. Onno, S. Kavadiya, C. Chosy, E. J. Wolf, J. Werner, J. A. Raiford, C. de Paula, A. F. Palmstrom, Z. J. Yu, J. J. Berry, S. F. Bent, Z. C. Holman, J. M. Luther, E. L. Ratcliff, N. R. Armstrong and M. D. McGehee, *Joule*, 2020, **4**, 1759–1775.
- 13 L. Xu, X. Chen, J. Jin, W. Liu, B. Dong, X. Bai, H. Song and P. Reiss, *Nano Energy*, 2019, **63**, 103860.
- 14 W. Chen, Y. Wu, J. Fan, A. B. Djurišić, F. Liu, H. W. Tam, A. Ng, C. Surya, W. K. Chan, D. Wang and Z. He, *Adv. Energy Mater.*, 2018, **8**, 1703519.
- 15 H. Zhu, Y. Liu, F. T. Eickemeyer, L. Pan, D. Ren, M. A. Ruiz-Preciado, B. Carlsen, B. Yang, X. Dong, Z. Wang, H. Liu, S. Wang, S. M. Zakeeruddin, A. Hagfeldt, M. I. Dar, X. Li and M. Grätzel, *Adv. Mater.*, 2020, **32**, 1907757.
- 16 B. Chen, H. Chen, Y. Hou, J. Xu, S. Teale, K. Bertens, H. Chen, A. Proppe, Q. Zhou, D. Yu, K. Xu, M. Vafaie, Y. Liu, Y. Dong, E. H. Jung, C. Zheng, T. Zhu, Z. Ning and E. H. Sargent, *Adv. Mater.*, 2021, **33**, 2103394.
- 17 P. Denis, M. Mertens, W. T. M. Van Gompel, A. Maufort, S. Mertens, Z. Wei, M. Van Landeghem, S. Gielen, B. Ruttens, D. Deduytsche, C. Detarvernier, L. Lutsen, F. Grozema, K. Vandewal and D. Vanderzande, *Adv. Opt. Mater.*, 2022, **10**, 2200788.
- 18 K. Takimiya, I. Osaka, T. Mori and M. Nakano, *Acc. Chem. Res.*, 2014, **47**, 1493–1502.
- 19 Y. Kiyota, T. Kadoya, K. Yamamoto, K. Iijima, T. Higashino, T. Kawamoto, K. Takimiya and T. Mori, *J. Am. Chem. Soc.*, 2016, **138**, 3920–3925.
- 20 R. Janneck, N. Pilet, S. P. Bommanaboyena, B. Watts, P. Heremans, J. Genoe and C. Rolin, *Adv. Mater.*, 2017, **29**, 1703864.
- 21 C. Niebel, Y. Kim, C. Ruzié, J. Karpinska, B. Chattopadhyay, G. Schweicher, A. Richard, V. Lemaure, Y. Olivier, J. Cornil, A. R. Kennedy, Y. Diao, W. Lee, S. Mannsfeld, Z. Bao and Y. H. Geerts, *J. Mater. Chem. C*, 2015, **3**, 674–685.
- 22 W. T. M. Van Gompel, R. Herckens, P. H. Denis, M. Mertens, M. C. Gélvez-Rueda, K. Van Hecke, B. Ruttens, J. D'Haen, F. C. Grozema, L. Lutsen and D. Vanderzande, *J. Mater. Chem. C*, 2020, **8**, 7181–7188.
- 23 X. Yang, D. Luo, Y. Xiang, L. Zhao, M. Anaya, Y. Shen, J. Wu, W. Yang, Y. H. Chiang, Y. Tu, R. Su, Q. Hu, H. Yu, G. Shao, W. Huang, T. P. Russell, Q. Gong, S. D. Stranks, W. Zhang and R. Zhu, *Adv. Mater.*, 2021, **33**, 1–10.
- 24 J. Li, L. Zuo, H. Wu, B. Niu, S. Shan, G. Wu and H. Chen, *Adv. Funct. Mater.*, 2021, **2104036**, 1–10.
- 25 O. Sambalova, E. Billeter, J. Mann, T. Miyayama, D. Burnat, A. Heel, D. Bleiner and A. Borgschulte, *Surf. Interface Anal.*, 2020, **52**, 811–817.
- 26 NIST X-ray Photoelectron Spectroscopy (XPS) Database, [https://srdata.nist.gov/xps/main\\_search\\_menu.aspx](https://srdata.nist.gov/xps/main_search_menu.aspx), (accessed December 2022).
- 27 S. Chen, C. Zuo, B. Xu and L. Ding, *J. Semicond.*, 2021, **42**, 120203.
- 28 A. Magomedov, A. Al-Ashouri, E. Kasparavičius, S. Strazdaite, G. Niaura, M. Jošt, T. Malinauskas, S. Albrecht and V. Getautis, *Adv. Energy Mater.*, 2018, **8**, 1801892.
- 29 S. Y. Kim, S. J. Cho, S. E. Byeon, X. He and H. J. Yoon, *Adv. Energy Mater.*, 2020, **10**(44), 2002606.
- 30 J. Dagar, M. Fenske, A. Al-Ashouri, C. Schultz, B. Li, H. Köbler, R. Munir, G. Parmasivam, J. Li, I. Levine, A. Merdasa, L. Kegelmann, H. Näsström, J. A. Marquez, T. Unold, D. M. Többs, R. Schlattmann, B. Stegemann, A. Abate, S. Albrecht and E. Unger, *ACS Appl. Mater. Interfaces*, 2021, **13**, 13022–13033.