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Tailoring dielectric properties using designed polymer-grafted ZnO nanoparticles in silicone rubber†

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Polymer grafts were used to tailor the interphases between ZnO nanoparticles (NPs) and silicone matrices. The final electrical properties of the nanocomposites were tuned by the grafted interphases, by controlling the inter-particle distance and the NP-morphology. The nanocomposites can be used in electrical applications where control of the resistivity is desired. Hansen's solubility parameters were used to select a semi-compatible polymer for grafting to obtain anisotropic NP morphologies in silicone, and the grafted NPs self-assembled into various morphologies inside the silicone matrices. The morphologies in the semi-compatible nanocomposites could be tuned by steering the graft length of poly(*n*-butyl methacrylate) via entropic matrix-graft wetting using surface-initiated atom-transfer radical polymerization. Image analysis models were developed to calculate the radius of primary NPs, the fraction of aggregates, the dispersion, and the face-to-face distance of NPs. The dielectric properties of the nanocomposites were related to the morphology and the face-to-face distance of the NPs. The dielectric losses, above 100 Hz, for nanocomposites with grafted NPs were approximately one decade lower than those of pristine NPs. The isotropic nanocomposites increased the resistivity up to 100 times compared to that of neat silicone rubber, due to the trapping of charge carriers by the interphase of dispersed NPs and nanoclusters. On the other hand, the resistivity of anisotropic nanocomposites decreased 10–100 times when the inter-particle distance in continuous agglomerates was close to the hopping distance of charge carriers. The electrical breakdown strength increased for compatible isotropic nanocomposites, and the temperature dependence of the resistivity and the activation energy were ~50% lower in the nanocomposites with grafted NPs. These flexible dielectric nanocomposites are promising candidates for low-loss high-voltage transmission cable accessories, mobile electronic devices, wearables and sensors.

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1. Introduction

A sustainable society requires renewable generation of electric power and reduced energy consumption *via* improved energy efficiency. More efficient electrical insulation systems need a range of mechanical and electrical properties. Tailored materials are of interest for specific applications, leading to

components with smaller size, and better reliability and efficiency. This is of current importance, for example for integrating distant renewable energy resources. Polymeric materials are being used as the main insulation of power cables, cable accessories and cable joints. They are also found as insulation in generators and motors. Further, flexible polymer nanocomposites are being considered as insulation materials for mobile electronic devices, so-called wearables, since they are known to be durable and elastic, and to have potentially tunable properties.^{1–3} Nanoparticles (NPs) are especially interesting as they might exhibit different or additional features than their corresponding bulk phase.⁴ Nanocomposites containing SiO₂ nanoparticles (NPs), grown *in situ* in polyester-amide,⁵ or dispersed in polyester⁶ or polyamide-imide,⁷ show excellent resistance against partial discharges and are used as insulation for electrical motors. The addition of SiO₂, TiO₂ or MgO NPs to polyethylene (PE) resulted in cable-insulation materials with higher electric breakdown strength^{8,9} or reduced space-charge

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(PDMS or silicone rubber) commonly used as outdoor insulation of high-voltage components and cable joints, but of growing interest for soft electronics. We characterized interparticle distances of annealed grafted ZnO NPs, various NP-morphologies in silicone matrices and how they depend on the N_M/N_g and σ . Finally, we investigated how the morphology of NPs influenced the dielectric properties (*e.g.* resistivity, electric constant, and losses) below, at and above the glass transition temperature of the grafts ($T_g \approx 20^\circ\text{C}$).⁴¹

2.1 Materials

2.2 Synthesis

2.2.1 Silanization and immobilization of initiators onto ZnO NPs to produce ZnO–Br. The ZnO NPs were dried under vacuum at 190 °C, dispersed in a mixture of deionized water and ethanol (1 : 1 by vol.) and silanized using APTES prior to the immobilization of ATRP-initiators (α -BiB). The silanization procedure to obtain ZnO–APTES, and the subsequent immobilization procedure to obtain ZnO–Br, were both performed according to previous work^{42,43} and are described in more detail in the ESI.† The average amount of Br-groups was 0.021 μ mol Br per mg of ZnO–Br, according to the mass loss detected by TGA.

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$$A = \sum_{j=1}^n \pi r_j^2 \frac{P(r_j)}{n} \quad (2)$$

where $P(r_j)$ is the radius-dependent probability of each curve, r_j is the NP radius of bucket j , and n is the total number of buckets in the histogram. The total fraction of agglomerates is thus defined as $\theta_{\text{aggl},B} = 1 - A_{\text{prim}}/A_{\text{tot}}$. As in method A, $\theta_{\text{aggl},B} = 0\%$ means no agglomeration, while a value of 100% means that all the primary NPs are agglomerated. Homogeneously distributed clusters close to the size of large primary NPs could potentially be considered as primary NPs with this method. The method is free from arbitrary model parameters.

(C) The relative cluster size was obtained by calculating the “average radius” ($\langle r \rangle$) in the micrographs and dividing by the “mean radius” of primary NPs ($\langle r_{0B} \rangle$, from method B) giving the relative average radius $r_{\text{rel},B} = \langle r \rangle / \langle r_{0B} \rangle$. If $r_{\text{rel},B}$ is close to or equal to 1, less agglomeration is present, while a higher $r_{\text{rel},B}$ value indicates more agglomeration. This strategy is robust, relatively easy to implement and does not use any arbitrary model parameters. However, in the case of homogeneously sized nanoclusters and very few primary NPs, the $r_{\text{rel},B}$ value underestimates the ratio of agglomerates. Therefore, it is valuable to compare $r_{\text{rel},B}$ with $r_{\text{rel},A} = \langle r \rangle / \langle r_{0A} \rangle$. Method C is considered more reliable than the methods A and B when most clusters are slightly larger than or slightly smaller than the r_{cutoff} .

(D) The face-to-face (or centre-to-centre) 2D-distances between NPs were compared to the corresponding theoretical distances for randomly positioned, dispersed spheres. The latter distances were calculated on 2D-slices cut from Monte-Carlo generated 3D-geometries of composites with non-intersecting, randomly positioned spherical fillers.⁴² The theoretical (2D) centre-to-centre distance for a system with equally sized particles with diameter 1 is described as a function of volume filler fraction ϕ as

$$D_{\text{cci}}^1(\phi) = \left(\frac{a_i}{1 + b_i \times \phi} \right)^{c_i} + d_i \quad (3)$$

where the subscript i corresponds to the i^{th} neighbour. The values used were $a = 131.57$, $b = 73\,978$, $c = 0.5622$, $d = 0.5353$ for the nearest neighbour and $a = 509.02$, $b = 30\,746$, $c = 0.4994$, $d = 0.01020$ for the 51st nearest neighbor.⁴² To obtain distances for systems with particle diameters other than 1, $D_{\text{cci}}^1(\phi)$ was multiplied by the diameter of the primary NPs. When method B is used to calculate the radius of the primary NPs, the theoretical centre-to-centre distance thus becomes $\check{D}_{\text{cci},B} = 2\langle r_{0B} \rangle D_{\text{cci}}^1(\phi)$ and the theoretical face-to-face distance is $\check{D}_{\text{ffi},B} = \check{D}_{\text{cci}} - 2\langle r_{0B} \rangle$.

If the primary NPs are clearly distinguishable, the ratio $\Delta_{\text{cci},B} = \langle D_{\text{cci}} \rangle / \check{D}_{\text{cci},B}$ of the distance $\langle D_{\text{cci}} \rangle$ and the analytical distance $\check{D}_{\text{cci},B}$ will decrease with decreasing dispersion, but if the primary NPs within the agglomerates are not treated as separated, $\Delta_{\text{cci},B}$ will instead increase with decreasing dispersion. This can be handled by introducing the cut-off radius (r_{cutoff}) from method A and treating the observed NPs with a radius larger than r_{cutoff} as clusters of adjacent NPs on a honey comb

lattice.¹² All the distances and the average radius are then recalculated. With this adjustment, a value of the ratio $\Delta_{\text{cci},D}$ less than 1 corresponds to a poorer dispersion, while a value of $\Delta_{\text{cci},D}$ greater than 1 corresponds to a better dispersion. The same applies for the new face-to-face ratio $\Delta_{\text{ffi},B}$. When experimental distances are compared with theoretical, it is important to use the volume filler fraction (ϕ) calculated from the micrograph to avoid systematic errors introduced due to visibility differences of NPs in cross-sections, deposited or microtomed samples.

2.4 Instrumentation and characterization methods

Fourier Transform Infrared spectroscopy (FTIR, 4000–600 cm^{-1}) was carried out with a Perkin-Elmer Spectrum 2000 FTIR instrument equipped with a MKII Golden Gate, single reflection ATR system (Specac Ltd, London, UK). The ATR-crystal was a MKII heated diamond 45° ATR top-plate.

A Kaiser-test kit, containing phenol (~80% in ethanol), potassium cyanide in H_2O :pyridine, and ninhydrin (6% in ethanol), was used to determine the concentration of primary amines on ZnO-APTES using a Cary UV-Vis Spectrophotometer with CaryWinUV software detecting the absorbance at 576 nm, which is associated with the amine-ninhydrin-complex concentration.^{44,45}

Thermogravimetric analysis (TGA) was performed on a Mettler Toledo TGA/DSC1, using STARe software to process the data. The samples were heated from 50 to 700 °C in ceramic cups at a rate of 10 °C min^{-1} in N_2 (flow rate 50 mL min^{-1}).

Transmission electron microscopy (TEM) of ZnO NPs was performed using a Hitachi HT7700 TEM at 100 kV accelerating voltage. Diluted suspensions (<0.1 mg mL^{-1}) of unmodified or modified NPs in EtOH or THF were applied on ultrathin carbon-coated copper grids (Ted Pella, Inc.) and examined in the microscope after drying.

A Hitachi S-4800 field emission scanning and transmission electron microscope (SEM/STEM) was used to study cross-sections of composites prepared by freeze-fracturing in liquid nitrogen and dispersions of pristine and modified NPs applied onto ultrathin carbon coated copper grids (TED Pella, Inc.).

DMF-SEC was carried out using a TOSOH EcoSEC HLC-8320GPC system equipped with an EcoSEC RI detector, three columns (PSS PFG 5 μm ; Microguard, 100 Å, and 300 Å) columns (M_w resolving range: 0.1–300 kDa) from PSS GmbH, using DMF (0.2 mL min^{-1}) with 0.01 M LiBr as the mobile phase at 50 °C. A conventional calibration method was used utilizing poly(methyl methacrylate) standards. Corrections for flow rate fluctuations were applied using toluene as an internal standard. PSS WinGPCUnity software version 7.2 was used to process the data.

2.4.1 Dielectric characterization. All the samples were conditioned at 23 °C and 0% RH for at least two days prior to dielectric characterization measurements. The contact pressure was 4 kPa for the electrode and 23 kPa for the guard. Soft and smooth rubber films ensured good contact to the electrodes without significant deformation. Relative permittivity and dielectric loss-tangent were measured using an insulation



diagnostic analyser (Megger IDAX300) with its output voltage connected to a large planar electrode below the sample, and its input connected to a guarded brass electrode with a diameter of 20 mm and a guard gap of 0.15 mm, above the sample. Dry samples were tested inside a climate chamber (Weiss Technik) set to 20 °C or 50 °C at 0% of RH. The voltage was sinusoidal at 200 V (peak) with frequencies sweeping from 10 kHz to 0.01 Hz. This type of measurement includes the effects of polarization, capacitance, and conduction. Resistivity was measured inside a climate chamber (CTS C-20/350) at −10 °C, 23 °C, and 50 °C. The samples were positioned on top of a high-voltage electrode connected to a Heinzinger PNC5 high-voltage DC power-supply outside the chamber. The brass electrode used for permittivity measurements was positioned on top of the film and connected to a Keithley 6512 electrometer. The polarization process took place under a constant DC electric field for 12 min. The resistivity was calculated based on the final polarization current. Electrical breakdown tests were performed on film samples between two spherical electrodes (6 mm in diameter), immersed in silicone oil. An AC voltage at 50 Hz was applied between the electrodes, starting at zero and increasing at 1 kV s^{−1} peak value, in the style of method 'A' in ASTM D149-97a. The electrodes were adjusted to provide sufficient force to prevent the sample from sliding between the electrodes. The sample thickness was measured with a Mitotoyo ABS Digimatic Indicator ID-SX (6 mm in diameter) at each tested site and the measured value was used to divide the recorded breakdown voltage.

3. Results and discussion

The major advantage, as well as challenge, of NPs is the control of the chemistry of their large specific surface area. The manufacturing of pure NPs with a defined size and surface chemistry on a large scale is still a significant challenge. Today, a range of commercial ZnO NPs of various purities and surface reactivities are available, but the presence of organic remnants may affect any subsequent surface modification and thus the final properties of nanocomposites. Therefore, two types of commercially available ZnO NPs were compared (see the section "Comparison of commercial ZnO NPs" in the ESI†) and it was found that the organic content of ZnO NPs produced by flame-spray pyrolysis⁴⁶ was 10 times less than that of ZnO NPs produced by salt-precipitation, which also contained residual salt-precursors⁴⁷ detected by FTIR and TGA (Fig. S1 and S2†). Further, the resistivity of a silicone rubber with ZnO NPs containing salt-precursors and pure ZnO NPs was $\sim 10^{10}$ Ω m and $\sim 10^{14}$ Ω m, respectively (Fig. S3†). The chemically pure, non-porous and highly crystalline ZnO NPs produced by flame-spray pyrolysis were therefore selected for this study. Micrometer sized ZnO particles (and bulk ZnO) are known to be semi-conductors with properties depending on doping elements and grain-boundaries.^{48,49} However, flame-sprayed, highly pure ZnO NPs neither contain doping elements nor large grain-boundaries. Low filling fractions of pure metal-oxide NPs have significantly reduced the conductivity by the trapping of charge carriers.^{11,12} The trapping has mainly been attributed to "deep

traps" on the NP-interface;⁵⁰ the induction of dipoles;⁵¹ and the scattering or charge differences by the NP-polymer interphase.^{8,52} In this study, we designed the interphase of flame-sprayed pristine ZnO NPs⁴⁶ by silanization and controlled polymer grafting in order to tune the dielectric properties of polymer nanocomposites.

3.1 Detection of accessible amines of ZnO-APTES

ZnO NPs were silanized by aminopropyltriethoxysilane (APTES) (Fig. S4a†). The accessible primary amines of the silane coating were monitored by spectrophotometry (at 570 nm) through the formation of the Ruhemanns-Blue complex with ninhydrin, using the Kaiser-test (Fig. S4b†).^{44,45} The concentration of primary amines was shown to be directly related to the final graft density, as the ninhydrin-complex detects the most accessible amines.⁴² The amine density was calculated using Lambert-Beer's law as previously described by Malmström *et al.*⁴² and was found to be 0.08 ± 0.01 amines nm^{−2}.

3.2 Characterization of ZnO-g-PBMA and free PBMA

Prior to the controlled grafting of PBMA from ZnO NPs, ATRP-initiators were immobilized on the primary amines of ZnO-APTES by the formation of an amide-bond (ZnO-Br, Fig. 1) as described in previous reports (see the ESI†).^{42,43} In this study, the SI-ATRP grafting was conducted from the end group of ZnO-Br in the presence of a sacrificial initiator to allow fast and easy characterization of the polymerization (Fig. 1).

The PBMA-grafts were cleaved and the degrafted NPs characterized by FTIR (Fig. 2). FTIR-spectra of ZnO-APTES, ZnO-Br and grafted ZnO NPs showed the absorbance peaks of amines, amides, and carboxyl groups or their absence in pristine and degrafted ZnO NPs at 1800–1500 cm^{−1} (Fig. 2 enlarged image) and the characteristic absorbance of C-H around 2930 cm^{−1}, which confirms the successful surface modification of the NPs (Fig. 2).

According to TGA, the concentration of immobilized initiators on ZnO-Br was 0.021 μ mol Br per mg of ZnO-Br, which is equal to 0.8 initiators nm^{−2}. The concentration of initiators was ten times higher than the determined amount of available amines and can be explained by the greater accessibility of the initiators compared to that of the bulky ninhydrin-complex; the initiators are able to reach primary amines deeper inside the "gooey-layer" of silane, as suggested by Malmström *et al.*⁴²

The free polymers were characterized by DMF-SEC and their average molar masses were compared to those of the cleaved chains (Table 1, Fig. S5 and S6†). The narrow dispersities indicate good control (Table 1). Free PBMA-chains were slightly longer and exhibited a slightly more narrow size distribution than the cleaved PBMA-grafts, in agreement with grafting from cellulose substrates.⁵³

According to the thermograms, the content of volatile organics increased with each surface modification, and the mass-losses are displayed next to the sample names in Fig. S7.† The average graft density of ~ 0.08 grafts nm^{−2} was confirmed from the correlation between the molar mass of cleaved grafts and the organic mass losses, and this is equal to a grafting



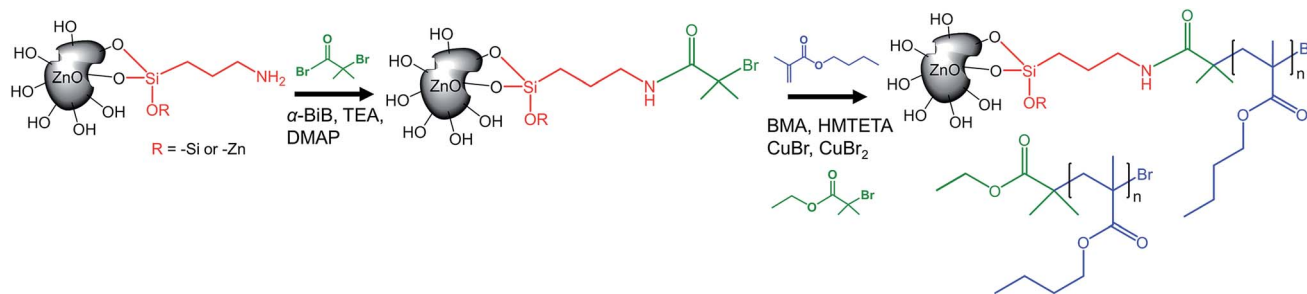


Fig. 1 The immobilization of the ATRP-initiator (green), and the surface-initiated grafting of BMA (blue) from ZnO-APTES and ZnO-Br, respectively. The free chains of PBMA are simultaneously polymerized from the sacrificial initiator by ATRP in parallel with the grafting reaction.



Fig. 2 FTIR spectra of ZnO, ZnO-APTES, ZnO-Br, ZnO-g-PBMA60k and degrafted ZnO between 4000 and 600 cm^{-1} . Spectra have been vertically scaled and offset for clarity. The enlarged spectrum shows detailed overlapping spectra between 1800 and 1500 cm^{-1} without offset.

efficiency of $\sim 10\%$. The calculated molar mass of grafts (M_g) and the degree of polymerization of the grafts (N_g) in Table 1 were obtained from the grafting density and the organic content according to eqn (S1) in the ESI.†

Finally, the band gaps of pristine ZnO, ZnO-g-PBMA10k and ZnO-g-PBMA60k (Table S2†) were calculated according to eqn (S2) in the ESI,† using the absorption peak obtained by UV-Vis spectroscopy, Fig. S8.† The tiny red-shift of grafted ZnO NPs indicated that the band gaps did not significantly change with surface grafting and thus remained constant at ~ 3.3 eV.

3.3 Dispersion of ZnO, ZnO-APTES and ZnO-g-PBMA

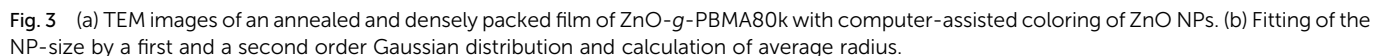
It is challenging to determine the distribution of primary NPs and the fraction of agglomerates from SEM and TEM images, or

dynamic light scattering, especially from NPs with a broad size distribution. The size of smaller agglomerates (nanoclusters) may be similar to that of larger primary NPs, which creates the illusion of well-dispersed NPs at low magnification, when in fact small agglomerates are present. Another difficulty is in distinguishing between agglomerates and clusters of individually separated NPs. To improve and to simplify the distinction between single NPs and agglomerates, we developed four methods to assess the radius and the dispersion of NPs considering different aspects. Methods A and B determine the respective fractions of agglomerates ($\theta_{\text{aggl},A}$ and $\theta_{\text{aggl},B}$) by using either the user-defined cut-off radius (r_{cutoff}) or the “mean radius” (r_{0B}) at the highest peak of the distribution histogram. Method A is reliable for detecting both homogeneously

Table 1 Comparison of molar masses and dispersities of free PBMA-chains, cleaved PBMA-grafts and calculated average molar masses of PBMA and correlating degree of polymerization

Name of NPs	Free PBMA		Cleaved grafts		Calculated	
	M_w (g mol^{-1})	D_M	M_w (g mol^{-1})	D_M	M_g (g mol^{-1})	N_g
ZnO-g-PBMA10k	13 000	1.18	9000	1.09	10 000	70
ZnO-g-PBMA50k	54 000	1.12	50 000	1.19	47 000	331
ZnO-g-PBMA60k	68 000	1.17	60 000	1.25	61 000	429
ZnO-g-PBMA80k	101 000	1.22	74 000	1.30	83 000	584





bution of non-annealed NPs. TEM and STEM images of non-annealed ZnO, ZnO-APTES and ZnO-g-PBMA were studied and compared by methods A-D and provided a better insight into the spontaneous agglomeration caused by the long-range (polar) and short-range (vdW) interactions of the samples (Fig. 4). The images at different magnifications (1.5k to 40k), were transformed into high-contrast images, and the fractions of agglomerates, the radii, the relative radii, the face-to-face distances, and the deviation ratios (Table 2), as well as the size distributions of primary NPs and the agglomerate ratio were calculated by methods A-D (Fig. 5).

	Method A	Method B	Method C				Method D	
	Fraction of agglomerates (%)		Average radius (nm)	Mean radius (nm)	Relative average radius		Face-to-face distance (nm)	Face-to-face deviation ratio
Name of the NP	$\theta_{\text{agglo,A}}$	$\theta_{\text{agglo,B}}$	$\langle r \rangle$	$\langle r_{0B} \rangle$	$r_{\text{rel,A}}$	$r_{\text{rel,B}}$	$D_{\text{ff1,D}}$	$\Delta_{\text{ff1,D}}$
ZnO	90	90	60–70	30–40	2.0	2.0	—	0.4–0.5
ZnO-APTES	80–90	0–50	70–90	70–90	2.4	1.1	—	0.1–0.3
ZnO-g-PBMA10k	30–70	70–80	40–50	20–40	1.3	1.5	6–8*	0.3–0.6
ZnO-g-PBMA60k	0–50	0–60	30–60	20–50	1.2	1.1	21–22*	0.8–11
ZnO-g-PBMA80k	0–30	0–40	30–40	30–40	1.1	1.1	26–28*	1.1–14



Fig. 4 Transmission images of agglomerates and primary particles of (a) pristine ZnO NPs, (b) ZnO-APTES, (c) ZnO-g-PBMA10k and (d) ZnO-g-PBMA80k. The inset pictures of annealed samples are 200 nm wide.

For pristine ZnO NPs, methods A and B successfully determined similar fractions of agglomerates ($\sim 90\%$, Table 2) despite the heterogeneous geometry of the NPs (Fig. 4a and 5a).

The large amount of agglomerates indicated strong NP-NP attraction (polarity). The average radius ($\langle r \rangle$) was approximately twice as large as the radius of the primary ZnO NPs.

The ZnO-APTES formed homogeneously distributed nanoclusters with agglomerated cores classified as 80–90% agglomerates (Fig. 4b) and only a few primary NPs (Fig. 5b). The formation of nanoclusters was detected by method A (as $\theta_{\text{aggl},A} \gg \theta_{\text{aggl},B}$) and confirmed by method C (since $r_{\text{rel},A} \gg r_{\text{rel},B}$), and was probably an effect of insufficient shielding of short-range core-core attractions during drying (Table 2). The agglomerates and distribution of ZnO-g-PBMA10k appeared very differently compared to ZnO-APTES (Fig. 4c and 5c). Both the radii ($\langle r_{0B} \rangle$ and $\langle r \rangle$) were close to the radius of the primary NPs (~ 33.5 nm), and thus described individual NPs separated from the closest neighbours (Table 2). The fractions of agglomeration were different ($\theta_{\text{aggl},A} < \theta_{\text{aggl},B}$), while the relative radii were similar ($\langle r_{\text{rel},A} \approx \langle r_{\text{rel},B} \rangle$), indicating a heterogeneous distribution of nanoclusters. The large variations in face-to-face distances of primary NPs within and between nanoclusters suggested that some core-core attraction between ZnO-g-PBMA10k remained during drying, although the grafts prevented complete core-core agglomeration.

Multiple images of annealed and non-annealed NPs with long grafts (ZnO-g-PBMA60k and ZnO-g-PBMA80k) displayed

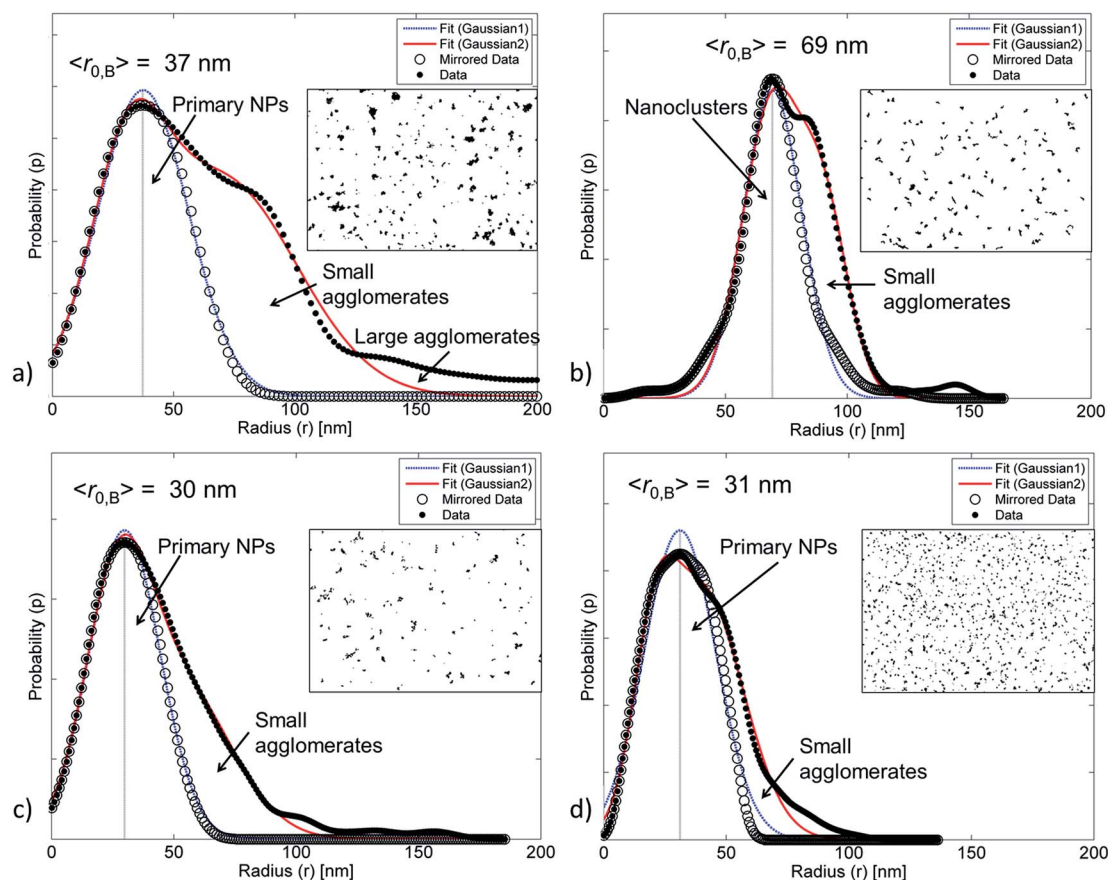


Fig. 5 The size distribution and the mean radius calculated from transmission images of (a) ZnO NPs, (b) ZnO-APTES, (c) ZnO-g-PBMA10k and (d) ZnO-g-PBMA80k. The arrows show the areas under the curves, representing the fraction of primary NPs, small and large agglomerates. The inset images are 12 μm wide cut-outs of the full images for simplified comparison.



Name of the nanocomposite	Silicone matrix	N_M	ZnO NP content		N_M/N_g
			vol%	wt%	
<i>1.5ZnO-g-PBMA10k</i>	PDMS	847	1.5	8.4	12
<i>2.5ZnO-g-PBMA10k</i>	Silicone rubber	550	2.5	14	8
<i>1.5ZnO-g-PBMA60k</i>	PDMS	847	1.5	8.4	2
<i>2.5ZnO-g-PBMA80k</i>	Silicone rubber	550	2.5	14	0.9

The nanocomposite films containing pristine ZnO NPs (*1.5ZnO*) were homogeneous and white translucent, while *1.5ZnO-APTES* and *1.5ZnO-g-PBMA10k* were more transparent

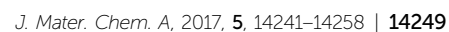




Fig. 7 Photo of neat PDMS and nanocomposites containing 1.5 vol% of ZnO NPs, ZnO-APTES, ZnO-g-PBMA10k or ZnO-g-PBMA60k.

and heterogeneous. The 1.5ZnO-g-PBMA60k samples contained regions of NPs separated by almost completely transparent regions providing the highest transmittance among the nanocomposites (Fig. 7 and S9a†). The increase in transparency of coated ZnO NPs can be explained by a combination of matching of the refractive indices⁴³ and regions containing few NPs.

Cross-sections of 1.5ZnO revealed randomized isotropic distributions of NPs and small nanoclusters incorporated in the PDMS-matrix (Fig. S10†), as an effect of the enthalpic compatibility between ZnO-NPs and the silicone matrix, according to the HSPs (Fig. 6). Cross-sections of 1.5ZnO-g-PBMA10k , with N_M/N_g equal to 12 (Table 1) expectedly formed phase-separated rounded agglomerates of grafted NPs, as the short grafts were

entropically unable to be wetted by the much longer polymer matrix. More surprisingly, minor regions of distributed NPs were also present (Fig. 8a and b), and might be an effect of the semi-efficient shielding of the enthalpic interactions between the ZnO NPs and the silicone matrix, in agreement with the heterogeneous dispersions of ZnO-g-PBMA10k (Fig. 4c and 5c). Cross-sections of 1.5ZnO-g-PBMA60k ($N_M/N_g = 2$) formed continuous superstructures of well-separated film-forming NPs (Fig. 8c and d and S11†). The long grafts appeared to entangle mainly with the other PBMA-grafts and partly with the matrix, as $N_M > N_g$, and formed continuous non-spherical agglomerates. This indicates that the initially well-dispersed ZnO-g-PBMA60k (in THF) self-assembled into superstructures of NPs during the thermal treatment of the films prior to curing. The interphase of grafts and matrix appeared to be seamless and the ZnO-cores were clearly separated by the grafts in agreement with annealed NPs (Fig. 3).

3.4.2 Nanocomposites of silicone rubber with 2.5 vol% NPs. The neat Sylgard® 184 (*silicone rubber*) and nanocomposites containing 2.5 vol% ZnO NPs (2.5ZnO), ZnO-g-PBMA10k (2.5ZnO-g-PBMA10k) and ZnO-g-PBMA80k (2.5ZnO-g-PBMA80k) were reproducibly prepared as white or white-translucent films (Fig. S9b†) similar to the nanocomposites with 1.5 vol% NPs, but they exhibited different NP morphologies (Fig. 9).

Cross-sections of 2.5ZnO-g-PBMA10k ($N_M/N_g \approx 8$) showed the anticipated formation of large rounded and elongated agglomerates of closely packed NPs (6–8 nm). Some of the agglomerates percolated and formed superstructures of several hundred microns in size containing islands of neat silicone

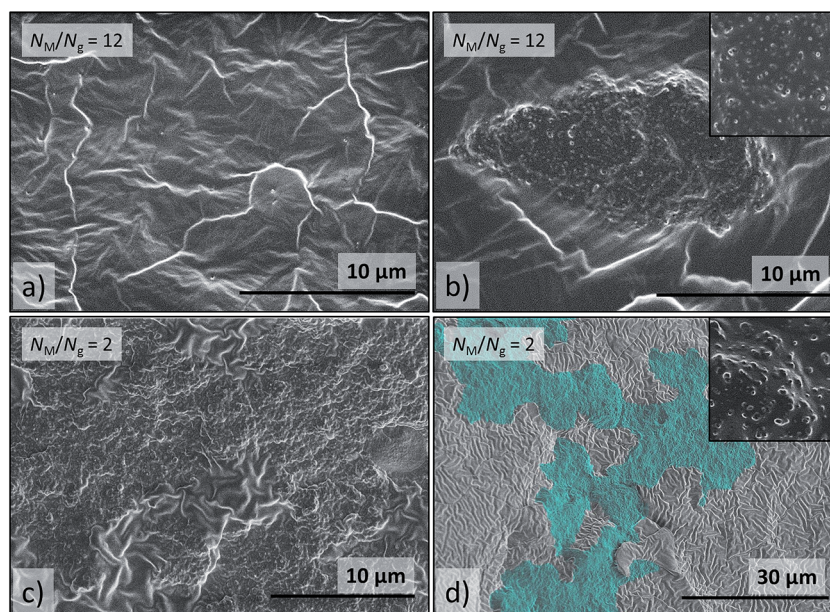


Fig. 8 SEM images of cross-sections of (a) and (b) 1.5ZnO-g-PBMA10k ($N_M/N_g = 12$) and (c) and (d) 1.5ZnO-g-PBMA60k ($N_M/N_g = 2$). (a) Region with low concentration of dispersed ZnO-g-PBMA10k, (b) a large spherical agglomerate of separated ZnO-g-PBMA10k, (c) continuous non-spherical agglomerate of ZnO-g-PBMA60k, (d) superstructure of ZnO-g-PBMA60k digitally highlighted in blue using Adobe Photoshop® for clarification. The original image is available in the ESI.† The inset pictures are 2 μm wide and show well-separated grafted NP cores of the agglomerates.

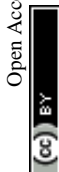




Fig. 9 SEM-images of cross-sections of (a) and (b) 2.5ZnO-g-PBMA10k ($N_M/N_g \approx 8$) and (c) and (d) 2.5ZnO-g-PBMA80k ($N_M/N_g \approx 0.9$). (a) Agglomerate of densely packed ZnO-g-PBMA10k containing areas with no NPs, (b) percolation of two large agglomerates of ZnO-g-PBMA10k, (c) nanoclusters and strings of ZnO-g-PBMA80k, (d) well-dispersed ZnO-g-PBMA80k in silicone rubber. The inset pictures in (b) and (d) are 2 μ m wide and show clustered ZnO-g-PBMA10k and well-separated ZnO-g-PBMA80k, respectively.

rubber (Fig. 9), which probably was an effect of improved wetting by the shorter polymer chains of the silicone matrix ($N_M/N_g < 10$), in contrast to 1.5ZnO-g-PBMA10k ($N_M/N_g > 10$).

Interestingly, cross-sections of 2.5ZnO-g-PBMA80k ($N_M/N_g \approx 0.9$) revealed regions of well-dispersed NPs and the presence of small agglomerates forming short strings (Fig. 9). It appears that the matching of N_g and N_M significantly improved the wetting of the semi-compatible grafts by the silicone matrix, and that this allowed the NPs to entangle with, and disperse in, the matrix, in a manner similar to that of homopolymer nanocomposites.^{32,33} The interconnected strings of grafted NPs appear to be effectively separated by the long brushes, in accordance with Fig. 3 and 4.

3.4.3 Phase diagram of the morphologies of nanocomposites. The morphologies of nanocomposites with grafted NPs (1.5ZnO-g-PBMA10k, 2.5ZnO-g-PBMA10k, 1.5ZnO-g-PBMA60k and 2.5ZnO-g-PBMA80k) were compared with previously reported data for annealed homopolymer nanocomposites^{33,56–60} in a phase diagram inspired by and partly reproduced from Kumar *et al.*³⁴ with data from Green *et al.*⁵⁶ (\square), Archer *et al.*⁵⁷ (\circ), Bates *et al.*⁵⁸ (Δ), Kumar *et al.*³³ (\diamond), and Jestin *et al.*^{59,60} (\star) (see Fig. 10). The data from the current study are represented by filled pentagons (\triangle). Each morphology observed in this study fits the corresponding phases in the diagram, where the product of σ and the square root of the graft length ($\sqrt{N_g}$) is plotted against the N_M/N_g ratio. This indicates that the entropical contributions are dominant for semi-compatible multipolymer nanocomposites ($RED \approx 1$). In the case of 2.5ZnO-g-PBMA80k, mainly well-dispersed NPs were observed together with the minor formations of short strings, indicated by a red pentagon with blue edges (Fig. 10). The formation of strings was similar to that reported

by Jestin *et al.*⁶⁰ The 2.5ZnO-g-PBMA10k is shown by a green pentagon with black edges, representing phase-separated continuous agglomerates. It should be noted that both the NPs and the self-assembled structures in the present study are significantly larger in size than the grafted SiO₂ or Fe₂O₃ NPs in the cited studies.



Fig. 10 A phase-diagram inspired by and partly reproduced from a review by Kumar *et al.*³³ with data from Green *et al.*⁴⁸ (\square), Archer *et al.*⁴⁹ (\circ), Bates *et al.*⁵⁰ (\diamond), Kumar *et al.*³² (\diamond), Jestin *et al.*^{51,52} (\star), and data from the current study (\triangle). The symbols are color coded where red data points represent well-dispersed samples, the black represent phase-separated structures, the blue represent strings, the green represent connected sheets and superstructures, and the purple points symbolize samples with small clusters of NPs. The edges of the pentagons are colored if a minor additional phase was also present.

3.5 Dielectric properties of silicone nanocomposites

The dielectric properties of the sample films were characterized multiple times under dry conditions at temperatures below, at and above the T_g of the PBMA-grafts. The dielectric losses, the relative permittivity and the resistivity were linked to the NP-morphologies and the face-to-face distance of the NPs, while the breakdown strength was dependent on the NP-matrix compatibility (Fig. S12†).^{3,61} The breakdown strength was improved by the addition of compatible, highly pure ZnO NPs (despite the relatively large filling content), while it decreased with increasing graft-length for the semi-compatible PBMA-grafted NPs. Importantly, all breakdown voltage levels were well above the E-fields used in high-voltage cable accessories, joints, and rotating machines.

The focus of this study is the correlation of the dielectric losses and the resistivity with the NP-morphologies, the surface grafting and the inter-particle distances, which are described in the following sections.

3.5.1 Dielectric spectroscopy of PDMS samples with 1.5 vol% NPs. The dielectric properties of PDMS, 1.5ZnO, 1.5ZnO-g-PBMA10k and 1.5ZnO-g-PBMA60k were evaluated at 20 °C and 50 °C (Fig. 11).

For PDMS, the real part of the relative permittivity ($\epsilon'_{r,PDMS} \approx 3$) exhibited no significant frequency dependence and was found to be slightly lower at 50 °C,⁶² due to a reduction in the number of aligned dipoles in PDMS. With the addition of pristine ZnO NPs (1.5ZnO), the ϵ'_r increased at 20 °C as an effect of the higher ϵ'_r of the NPs ($\epsilon'_{r,ZnO} \approx 8.5$).⁶³ The ϵ'_r was about the same for all the nanocomposites and may be an effect of the absence of impurities, doping agent and grain-boundaries of the highly pure, flame-sprayed ZnO NPs⁴⁶ (Fig. S1 and S2†). However, the ZnO-g-PBMAs were slightly frequency dependent.

The dielectric loss ($\tan \delta = \epsilon''_r / \epsilon'_r$, where ϵ''_r is the imaginary component of ϵ'_r) of PDMS was low and made it possible to investigate the effect of added nanofillers (Fig. 11b and d). The $\tan \delta$ of PDMS increased with temperature at low frequencies and the slope originating from DC conduction was shifted towards higher frequencies at 50 °C. The addition of pristine ZnO NPs (1.5ZnO) led to a higher and constant $\tan \delta$ above 0.1 Hz, while $\tan \delta$ at 50 °C decreased at higher frequencies.

At both 20 and 50 °C, the addition of grafted NPs introduced a frequency-dependent $\tan \delta$, which resulted in lower losses at higher frequencies than with pristine ZnO NPs. This effect is attributed to the absence of polar groups (–OH and moisture) at the interface of the ZnO-g-PBMAs. At 20 °C, the



Fig. 11 Dielectric spectroscopy data of neat PDMS, 1.5ZnO, 1.5ZnO-g-PBMA10k and 1.5ZnO-g-PBMA60k at 20 °C (a and b) and at 50 °C (c and d). Fig. (a) and (c) show the relative permittivity, ϵ'_r versus frequency, and (b) and (d) the loss tangent, $\tan \delta$ versus frequency.



1.5ZnO-g-PBMA60k suffered greater losses at low frequencies than *1.5ZnO-g-PBMA10k*. This may be attributed to the slower orientation of entangled long PBMA-grafts in the superstructures of *1.5ZnO-g-PBMA60k*. However, at 50 °C and at low frequencies, $\tan \delta$ decreased and *1.5ZnO-g-PBMA60k* exhibited the lowest losses of all the samples (Fig. 11d). The non-linear behaviour of the $\tan \delta$ may be an effect of the glass transition of the entangled long PBMA-grafts, and the lower DC conductivity can be explained by the face-to-face distance of 21–22 nm (Table 2), which is longer than the hopping distance of charge carriers (<10 nm) according to Wang *et al.*⁶⁴

3.5.2 Dielectric spectroscopy of silicone rubber samples with 2.5 vol% NPs. To evaluate the effect of NPs in a commercial grade of silicone rubber, Sylgard® 184 was used. It was observed that the neat silicone rubber (*silicone rubber*) had a significantly higher dielectric loss than PDMS (Fig. 11b and d and 12b and d), which is probably a result of the reinforcing silica filler. Therefore, a higher filling ratio of NPs was used than in the PDMS-based nanocomposites. Dielectric spectroscopy data of *silicone rubber*, *2.5ZnO*, *2.5ZnO-g-PBMA10k*, and *2.5ZnO-g-PBMA80k* at 20 °C and 50 °C are shown in Fig. 12.

The ϵ'_r of *silicone rubber* was observed to be similar to that of PDMS (Fig. 11a and c and 12a and c). However, the increase in

ϵ'_r below 0.1 Hz, observed for *silicone rubber*, and also for *2.5ZnO-g-PBMA10k* (Fig. 12a and c), has previously been explained by a “quasi-DC” conduction double-layer effect,^{52,65} due to the formation of small “capacitors”, or dipoles, as the charges are located on opposite sides of the NPs (*i.e.* fumed silica or *ZnO-g-PBMA10k*). The slope of $\tan \delta$ in *silicone rubber* and *2.5ZnO-g-PBMA10k* was close to -1 below ~ 50 Hz. The *2.5ZnO-g-PBMA10k* exhibited the highest $\tan \delta$ below ~ 50 Hz, and an increase in the DC conductivity, caused by the superstructures with face-to-face distances close to the hopping distance (<10 nm) (Fig. 12b and d).⁶⁴ However, above ~ 50 Hz, the $\tan \delta$ can be significantly reduced in agreement with what was seen for *1.5ZnO-g-PBMA10k* and *1.5ZnO-g-PBMA60k*.

The isotropic morphology of NPs in *2.5ZnO* and *2.5ZnO-g-PBMA80k* effectively reduced the “quasi-DC” effect, since the dispersed NP-interphases immobilize the charge carriers (Fig. 12c and d). For *2.5ZnO*, $\tan \delta$ was constant over the investigated frequency range and showed the highest value at 10 kHz and among the lowest at 0.01 Hz (less DC conduction). The *2.5ZnO-g-PBMA80k* (containing dispersed grafted NPs) exhibited a low $\tan \delta$ at both low and high frequencies.

3.5.3 Relationship between the relative dielectric losses and the NP morphology. The $\tan \delta$ values of *1.5ZnO-g-PBMA10k*,

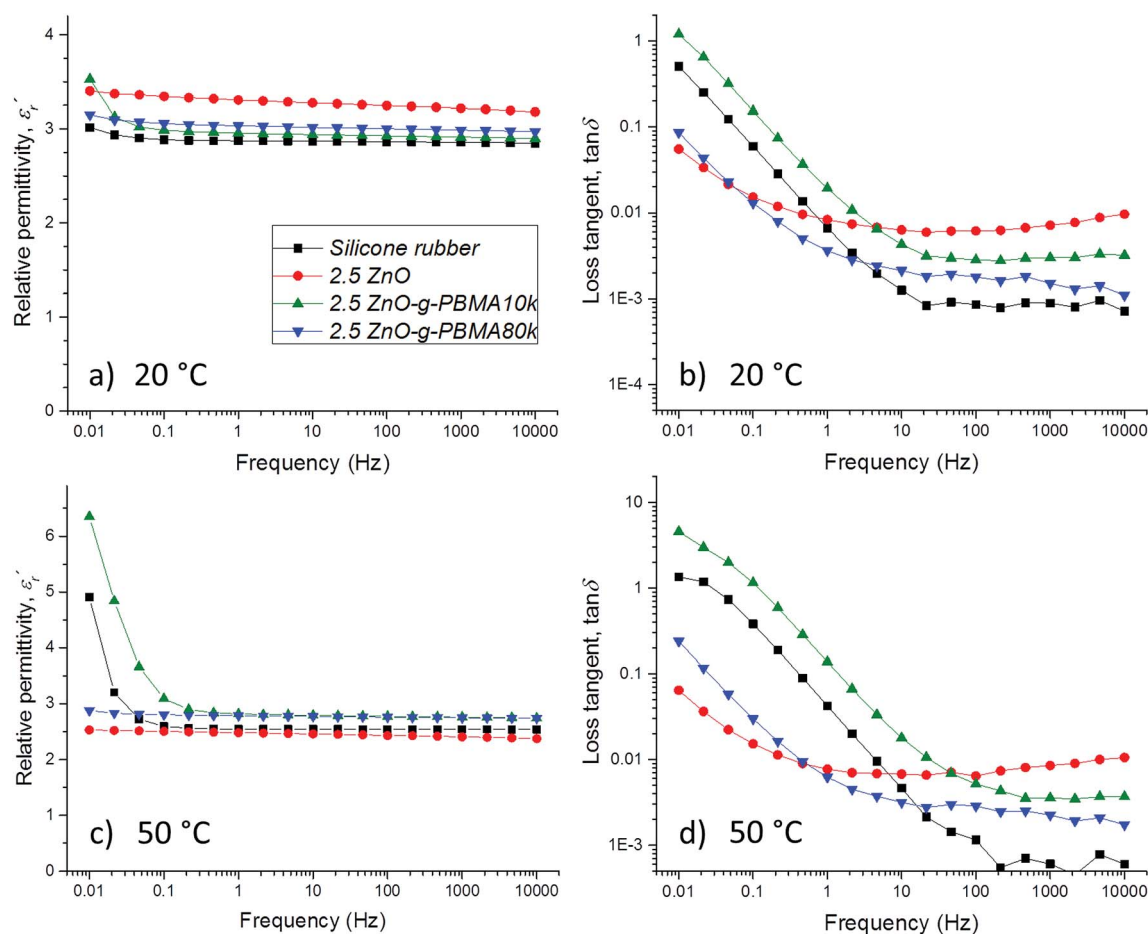
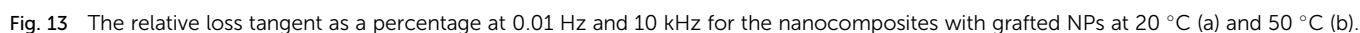
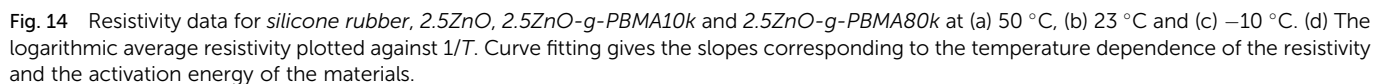


Fig. 12 Dielectric spectroscopy data of *silicone rubber*, *2.5ZnO*, *2.5ZnO-g-PBMA10k* and *2.5ZnO-g-PBMA80k* at 20 °C (a and b), and 50 °C (c and d). Fig. (a) and (c) display the relative permittivity, ϵ'_r versus frequency, and (b) and (d) the loss tangent, $\tan \delta$ versus frequency.



a percentage, was related to the morphology and the face-to-face distance of the NPs at 20 °C and 50 °C (Fig. 13). All four nanocomposites with grafted NPs exhibited lower losses at 10 kHz at both temperatures, due to the less polar interphases of ZnO-*g*-PBMA than pristine ZnO NPs. The relative $\tan \delta$ at 10 kHz decreased at 50 °C for all four samples. The *2.5ZnO-g-PBMA80k* exhibited 80–90% lower relative losses at both frequencies and



temperatures. At low frequencies, the interphase of the well-dispersed grafted NPs efficiently trapped charge carriers, and at high frequencies the dielectric loss was lower, due to a reduced amount of polar groups and moisture.

At 0.01 Hz, the *2.5ZnO-g-PBMA10k* with connected structures exhibited large relative losses at both 20 °C and 50 °C, in contrast to *1.5ZnO-g-PBMA60k* with connected structures and the phase-separated *1.5ZnO-g-PBMA10k*, both of which exhibited lower relative losses at 50 °C. The large relative loss tangent (240%) and increased DC-conductivity of *2.5ZnO-g-PBMA10k* at 50 °C were an effect of the morphologies and short face-to-face distance of the ZnO-cores (<10 nm) (Fig. 13). The interesting dielectric properties obtained at low frequencies were complemented by DC resistivity measurements to gain more insight into the conduction mechanisms.

3.5.4 Resistivity measurements of silicone rubber containing 2.5 vol% NPs. The resistivities of *silicone rubber*, *2.5ZnO*, *2.5ZnO-g-PBMA10k*, and *2.5ZnO-g-PBMA80k* at different electric fields (0.2–8 kV mm⁻¹) are shown in Fig. 14. The samples were tested above (50 °C), around (23 °C) and below (−10 °C) the *T_g* of PBMA. The scatter in data could, however, be up to one decade at high resistivities (>10¹⁶ Ω m), due to the low signal-to-noise ratio.

It was concluded that all the evaluated materials essentially exhibited a resistivity independent of the applied electric field within the evaluated interval, and the resistivity increased with decreasing temperature. The isotropic nanocomposites with dispersed NPs (*2.5ZnO* and *2.5ZnO-g-PBMA80k*) had a higher resistivity than the *silicone rubber* (Fig. 14), indicating a trapping of charge carriers by the interphase of the well-dispersed NPs, despite the high filling ratio. These findings are in agreement with the reported higher resistivities of LDPE-nanocomposites with ZnO NPs.¹¹ The enhanced resistivity indicates that ZnO-g-PBMA80k was still readily available for trapping of charge carriers.

The lower resistivity of *2.5ZnO-g-PBMA10k* (Fig. 14) is most likely an effect of the connected structures with the short face-to-face distances (6–8 nm), in the range for efficient hopping of charge carriers (Fig. 13).⁶⁴ This may be compared with the longer face-to-face distance of 26–28 nm for the *2.5ZnO-g-PBMA80k*. We verified that the observed reduction in resistivity was not caused by ionic contamination (Fig. S3†) by reproducing the synthesis and evaluating the effect of different cleaning approaches. The pure interphases of NPs were confirmed by the greater resistivity of *2.5ZnO-g-PBMA80k*.

Since it is considered important to obtain stable dielectric properties over a large temperature interval for high voltage dielectrics and electronic applications, the temperature dependence of the resistivity was calculated using the Arrhenius equation. The temperature dependence of the resistivity and the activation energy (*E_a*) were obtained by curve fitting of the logarithmic resistivities *versus* *T*⁻¹, where the slopes correspond to the *E_a/R* (Fig. 14d and S13a†). The temperature dependence of the resistivity decreased for the nanocomposites compared to *silicone rubber*. The lowest *E_a* was obtained for *2.5ZnO-g-PBMA80k* followed by *2.5ZnO-g-PBMA10k* and *2.5ZnO* (Fig. 14d and S13b†).

4. Conclusion

For the first time, Hansen's solubility parameters (HSPs) have been directly applied to predict the compatibility between the polymer grafts on NPs and polymer matrices. By using HSPs, we designed nanocomposites with semi-compatible polymer grafts and matrices in order to obtain various NP morphologies, ranging from spherical agglomerates *via* continuous superstructures to well-dispersed NPs. We discovered that the morphology of the NPs in a semi-compatible nanocomposite is strongly affected by the entropic wetting of the grafts by the matrix (*N_M/N_g* ratio) similar to homopolymer nanocomposites. This indicates that the morphology of ZnO NPs and thus, the dielectric properties of the nanocomposites can be tuned by controlling the graft length and the *N_M/N_g* ratio. Further, the resistivity of anisotropic nanocomposites containing connected morphologies is dependent on the tunable face-to-face distance of the NP-cores, as the resistivity decreased when the distance was close to the hopping distance (<10 nm). Image analysis models were developed to improve the characterization of NPs and their distribution, by calculating the face-to-face distance, the radius of the primary NPs, and the fraction of agglomerates.

More specifically, we obtained nanocomposites of commercial silicone with low dielectric loss where the resistivity increased up to 100 times. The interphase of well-dispersed grafted NPs provides a unique combination of high resistivity, by the trapping of charge carriers, and low dielectric loss at higher frequencies, due to the absence of polar groups and moisture. The well-dispersed pristine ZnO NPs improved the electrical breakdown strength of silicone due to their large and compatible NP-matrix interphase. Further, the thermal dependence of the DC resistivity was reduced by up to 50% for nanocomposites compared to neat silicone rubber.

We believe that these tunable properties of silicone based nanocomposites with semi-compatible polymer grafts are highly versatile and promising materials for low-loss high-voltage transmission cable accessories, as well as for flexible-dielectric materials for mobile electronic devices, wearables and sensors.

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